

# Qudit entanglers using quantum optimal control and Rydberg interaction on nuclear-spin

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# Quantum computing beyond qubits

#### Universal gate set for quantum computing

- Qubits : SU(2) + one entangling gate such as CNOT
- Qudits : SU(d) + one entangling gate

# Quantum computing beyond qubits

#### Universal gate set for qudits:

- Gates are harder to implement for qudits
  - 1. Constructive analytical approach [1][2]

# Quantum computing beyond qubits

#### Universal gate set for qudits:

- Gates are harder to implement for qudits
  - 1. Constructive analytical approach:
  - 2. Numerical approach: Quantum control

### High fidelity entangling qudit gates using quantum control

[1] M. Ringbauer, M. Meth, L. Postler, R. Stricker, R. Blatt, P. Schindler, and T. Monz, arXiv:2109.06903 (2021)

[2] M. S. Blok, V. V. Ramasesh, T. Schuster, K. O'Brien, J. M. Kreikebaum, D. Dahlen, A. Morvan, B. Yoshida, N. Y. Yao, and I. Siddiqi, Phys. Rev. X 11, 021010 (2021)

[3] SO, A. Mitra, M. J. Martin, and I. H. Deutsch, Phys. Rev. A 104, L060401 (2021).

# Nuclear Spin Qudit: Group-II Atoms

## **Strontium:** 87**Sr**<sup>[1][2]</sup>

• The ground space of  $^{87}$ Sr, J=0 and F=I+J=I

F: Total angular momentum

I: Nuclear spin

J: Electron angular momentum

• I=9/2 gives d=10 (qudecimal) [d=2F+1]

$$\frac{}{|-9/2\rangle} \quad \overline{|-7/2\rangle} \quad \overline{|-5/2\rangle} \quad \overline{|-3/2\rangle} \quad \overline{|-1/2\rangle} \quad \overline{|1/2\rangle} \quad \overline{|3/2\rangle} \quad \overline{|5/2\rangle} \quad \overline{|7/2\rangle} \quad \overline{|m_I = 9/2\rangle}$$

[1] Boyd, M. M. (2007). High precision spectroscopy of strontium in an optical lattice: Towards a new standard for frequency and time, JILA Ph. D Thesis.

[2] Martin, M. J. (2013). "Quantum metrology and many-body physics: pushing the frontier of the optical lattice clock." JILA Ph. D Thesis.

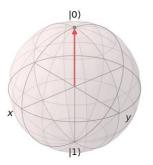
## Control Hamiltonian for quantum control

- We use the quantum control technique of GRAPE to generate qudit entanglers
- We have a controllable two qudit Hamiltonian,

$$H(t) = H^{(1)}(t) + H^{(2)}(t) + H_{\text{Ent}}$$

The single atom Hamiltonian is given as,

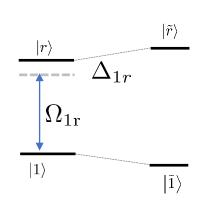
$$H^{(\alpha)}(t) = -\boldsymbol{\mu}.\boldsymbol{B}_{\mathrm{rf}}(t) = \Omega_{\mathrm{rf}} \left[ \cos \phi(t) F_x^{\alpha} + \sin \phi(t) F_y^{\alpha} \right] \equiv H_{\alpha}[\phi(t)]$$



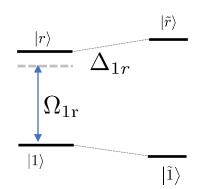
**Larmor Precession** 

## Rydberg blockade and Rydberg dressing<sup>[1][2]</sup>

$$H(t) = H_1[\phi(t)] + H_2[\phi(t)] + H_{Ent}$$



Atom 1

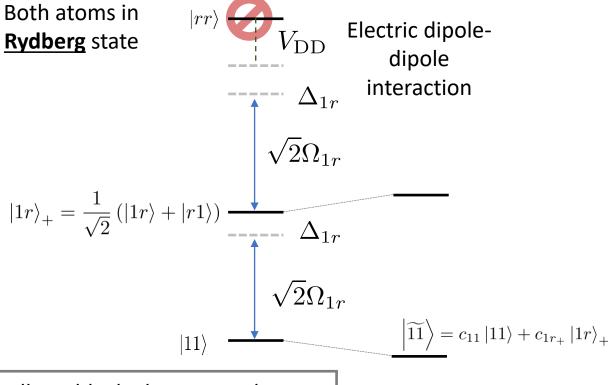


Atom 2

 $|1\rangle$ : ground state

 $|r\rangle$ : highly excited state

Rydberg blockade



Rydberg blockade occurs when

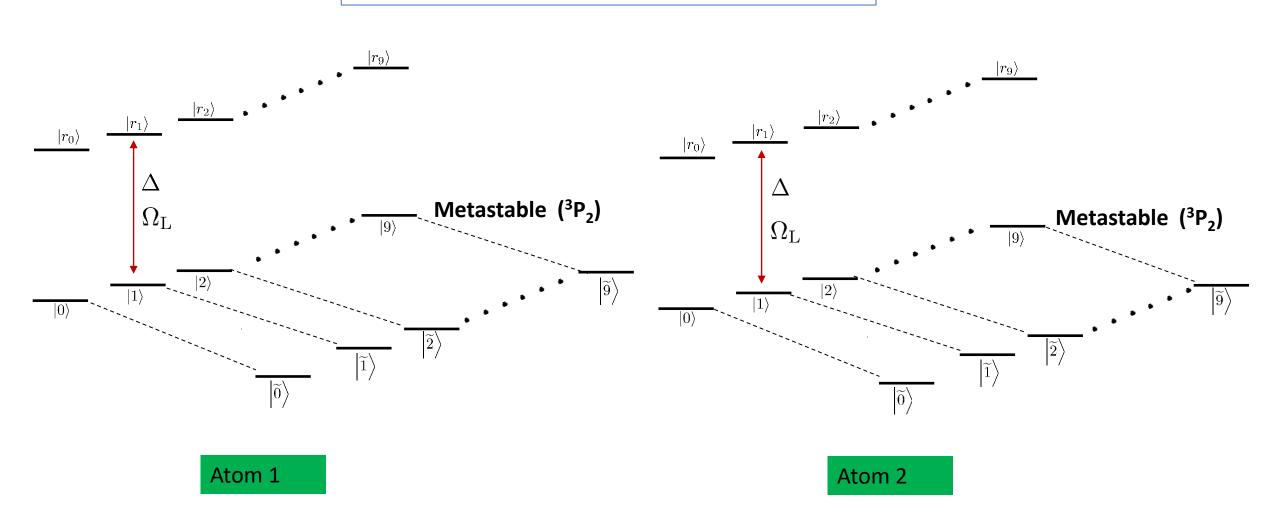
 $|V_{\rm DD}| \gg \Omega_{1r} \quad |V_{\rm DD}| \gg |\Delta_{1r}|$ 

[1] A. Mitra, M. J. Martin, G. W. Biedermann, A. M. Marino, P. M. Poggi, and I. H. Deutsch, Phys. Rev. A 101, 030301 (R) (2020).

[2] J. E. Johnson and S. L. Rolston, Phys. Rev. A 82, 033412 (2010).

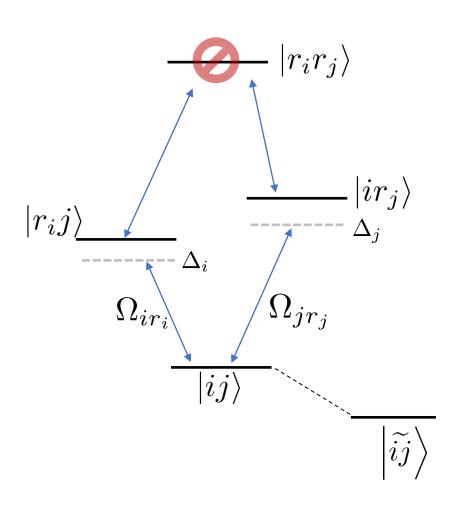
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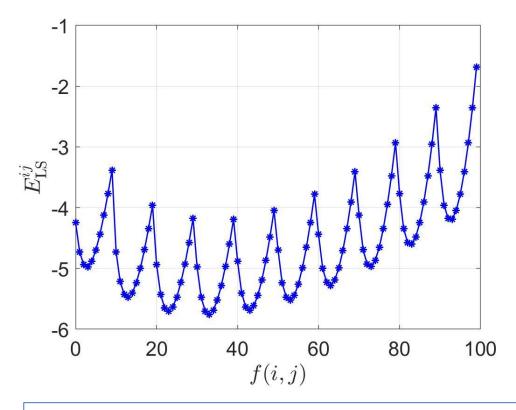


$$\left|\widetilde{ij}\right\rangle = C_{ij}\left|ij\right\rangle + C_{r_ij}\left|r_ij\right\rangle + C_{ir_j}\left|ir_j\right\rangle$$

$$H_{\mathrm{Ent}} = \sum_{ij} E^{ij} \left| \widetilde{ij} \right\rangle \left\langle \widetilde{ij} \right|$$

# Understanding the entangling Hamiltonian

$$H_{\mathrm{Ent}} = \sum_{ij} E^{ij} \left| \widetilde{ij} \right\rangle \left\langle \widetilde{ij} \right|$$



$$f(i,j) = 10i + j; 0 \le i, j < 10$$

Tensor light shift kind of effect

$$H(t) = H_1[\phi(t)] + H_2[\phi(t)] + H_{\rm Ent}$$

# Quantum optimal control

Partial isometry

$$\begin{pmatrix}
U_{11} & U_{12} & \cdots & U_{1d} \\
U_{21} & U_{22} & \cdots & U_{2d} \\
\vdots & \vdots & \ddots & \vdots \\
U_{d1} & U_{d2} & \cdots & U_{dd}
\end{pmatrix}$$
Unitary Transformation:  $\{|e_i\rangle\} \to \{|f_i\rangle\}$ 

$$V_{\text{tar}} = \sum_{i=1}^{k} |f_i\rangle \langle e_i| + U_{\perp}$$

- Cost function:  $\mathcal{F}[c(t)] = \frac{1}{k^2} \left| \sum_{i=1}^k \langle f_i | V[c(t)] | e_i \rangle \right|^2$
- Divide the total time T into N equal timesteps each of width  $\Delta t$

• Finding the optimal control waveforms involves maximizing  $\mathcal{F}[\boldsymbol{c}(t)]$ 

## Quantum optimal control

Partial isometry

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- Divide the total time T into N equal timesteps each of width  $\Delta t$
- Number of free parameter:  $N_{\min}(k,d) = 2kd k^2 1$ .

# Encoding a Qutrit entangler in Qudit

☐ We are interested in the CPhase which generalizes CZ gate

CPhase 
$$|j\rangle |k\rangle = \omega^{jk} |j\rangle |k\rangle; \quad \omega = \exp(2\pi i/d)$$

$$|0\rangle = \left| m_I = \frac{-9}{2} \right\rangle \quad |00\rangle \to |00\rangle \qquad |12\rangle \to e^{\frac{4\pi i}{3}} |12\rangle$$

$$|1\rangle = \left| m_I = \frac{-7}{2} \right\rangle \quad |02\rangle \to |02\rangle \qquad |20\rangle \to |20\rangle$$

$$|2\rangle = \left| m_I = \frac{-5}{2} \right\rangle \quad |10\rangle \to |10\rangle \qquad |22\rangle \to e^{\frac{4\pi i}{3}} |21\rangle$$

$$|2\rangle = \left| m_I = \frac{-5}{2} \right\rangle \quad |11\rangle \to e^{\frac{2\pi i}{3}} |11\rangle$$

$$H(t) = H_1[\phi(t)] + H_2[\phi(t)] + H_{\text{Ent}}$$

$$H_{\alpha}[\phi(t)] = -\mu. \boldsymbol{B}_{\text{rf}}(t) = \Omega_{\text{rf}} [\cos \phi(t) F_x + \sin \phi(t) F_y]$$

We control the system by changing the phase in a piecewise constant way with the cost function

$$\mathcal{F}[c(t)] = \frac{1}{k^2} \left| \sum_{i=1}^k \langle f_i | V[c(t)] | e_i \rangle \right|^2$$

# Encoding a Qutrit entangler in Qudit

$$H(t) = H_1[\phi(t)] + H_2[\phi(t)] + H_{\rm Ent}$$

 $N_{min}$ = 643 (minimum number of parameters)

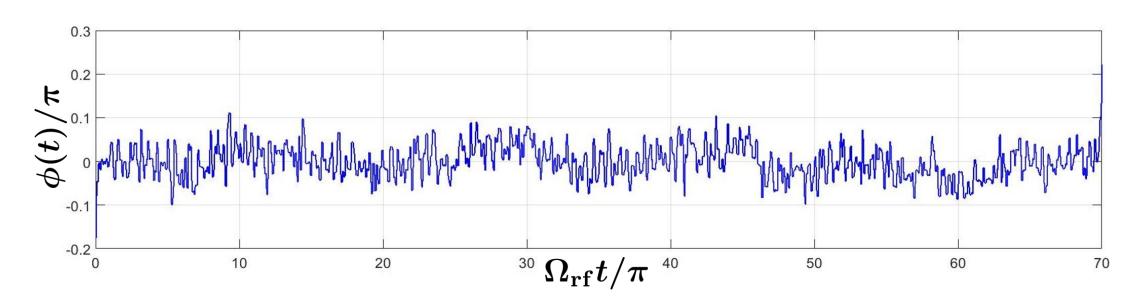
N: Total Number of controls

N=1500

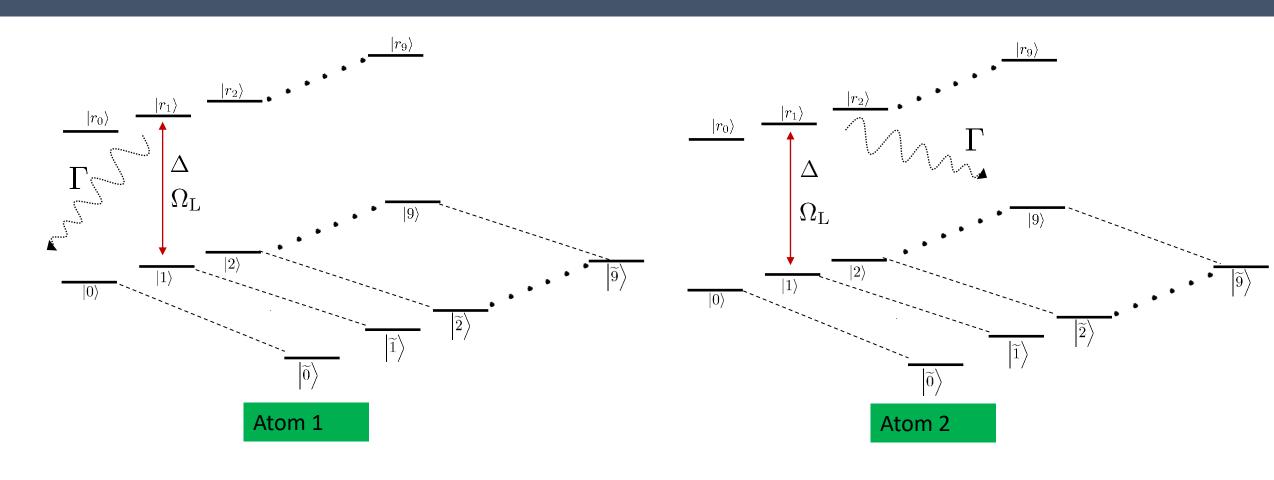
T: Total time for the Simulation

 $\Omega_{\rm rf}$ T=70 $\pi$ 

Fidelity=0.999



# Decoherence: Effective Hamiltonian for qudits

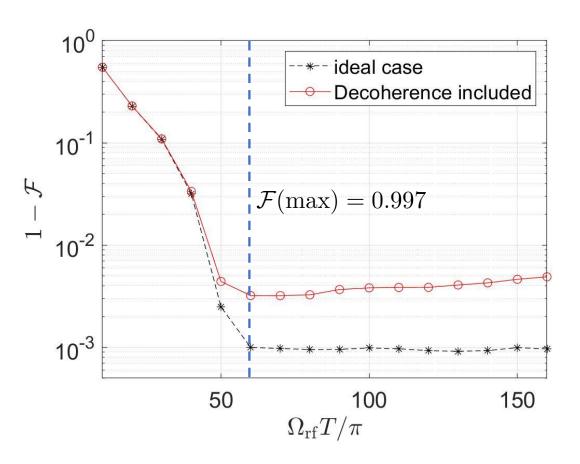


$$H_{\mathrm{Ent}}^{\mathrm{eff}} = \sum_{ij} \left( E^{ij} - i \gamma_{\mathrm{decay}}^{ij} / 2 \right) \left| i \tilde{j} \right\rangle \left\langle i \tilde{j} \right|$$

# Qutrit Cphase gate

$$H(t) = H_1[\phi(t)] + H_2[\phi(t)] + H_{\text{Ent}}$$

$$\mathcal{F}[c(t)] = \frac{1}{k^2} \left| \sum_{i=1}^k \langle f_i | V_{\text{eff}}[c(t)] | e_i \rangle \right|^2$$



# Fidelity: Cphase gate

$$H(t) = H_1[\phi(t)] + H_2[\phi(t)] + H_{\text{Ent}}$$

$$\mathcal{F}[c(t)] = \frac{1}{k^2} \left| \sum_{i=1}^k \langle f_i | V_{\text{eff}}[c(t)] | e_i \rangle \right|^2$$

d	$\mathcal{F}(\max)$	$\Omega_{ m rf}T_{ m opt}$
2	0.998	$24\pi$
3	0.997	$60\pi$
5	0.9900	$162\pi$
7	0.9700	$270\pi$

- This values can be further optimized
- A larger qudit stores much more information than a qubit system

# Summary:

- A quantum control approach for creating CPhase gate for qudits with few  $\mu s$ .
- Decoherence induced by finite lifetime of the Rydberg states is studied in detail.

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■ The preprint to appear in arXiv soon.



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### Thanks for the attention

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