

# Lecture Notes on Index Theory

- David Alvaro  
Dept of Physics UCSD  
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## 3.5. Index theory

Consider a smooth 2D vector field  $\mathbf{v}(\phi)$ . The angle that vector  $\mathbf{v}$  makes with respect to  $\hat{\phi}_1$  and  $\hat{\phi}_2$  axes is a scalar field:

$$\theta(\phi) = \tan^{-1} \left( \frac{v_2(\phi)}{v_1(\phi)} \right)$$

As long as  $v$  has finite length [doesn't vanish or blow up] the angle  $\theta$  is well defined. Expect that we can integrate  $\nabla \theta$  over a closed curve  $C$  in phase space to get

$$\oint_C d\phi \cdot \nabla \theta = 0$$

[ $\theta$  at all points is independent of path taken]

$\sum \nabla \theta$  is gradient of a scalar field

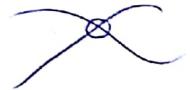
However this can fail if  $\mathbf{v}(\phi)$  vanishes or diverges at one or more pts. in the interior of  $C$ .

Define:  $W_C(V) = \frac{1}{2\pi} \oint_C d\phi \cdot \nabla \theta$

then  $W_C(V) \in \mathbb{Z}$  is an integer valued function of  $C$  which is the change in  $\theta$  around the curve  $C$ . This must be an integer because  $\theta$  is well defined only upto multiples of  $2\pi$ . Differential changes of  $\theta$  are in general well defined.

If  $\mathbf{v}(\phi)$  is finite  $\rightarrow$  neither diverges nor vanishes anywhere in  $C$  then  $W_C(V) = 0$ .

Assuming that  $\mathbf{v}$  never diverges, any singularities in  $\theta$  must be arising from pts where  $\mathbf{v} \rightarrow 0$ , which in general occurs at isolated points. since two variables



The index of a two-dimensional vector field  $\mathbf{V}(\varphi)$  at a point  $\varphi$  is the integer-valued winding of  $\mathbf{V}$  about that point.

$$\text{ind } C \mathbf{V} \varphi = \lim_{\rho_0 \rightarrow 0} \frac{1}{2\pi} \oint_{C(P_0)} d\varphi \cdot \nabla \varphi$$

$$= \lim_{\rho_0 \rightarrow 0} \frac{1}{2\pi} \oint_{C(P_0)} d\varphi \cdot \left( \frac{\nabla V_2 - V_2 \nabla V_1}{V_1^2 + V_2^2} \right)$$

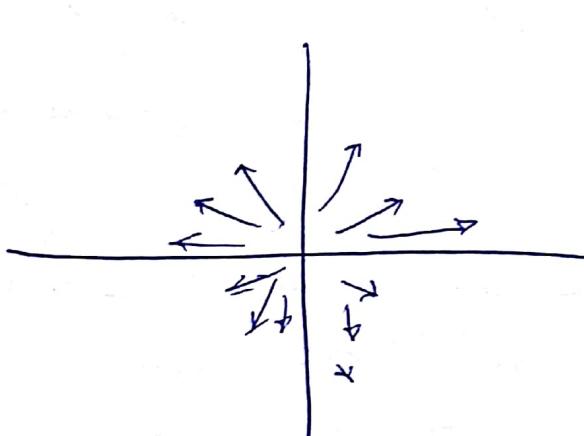
The index of a closed curve  $C$  is given by the sum of the indices at all the singularities enclosed by the curve.

$$\left[ \because \theta = \tan^{-1} \left( \frac{V_2}{V_1} \right) \right]$$

$$d\theta = \frac{1}{1 + \left( \frac{V_2}{V_1} \right)^2} \times \left( \frac{V_1 dV_2 - V_2 dV_1}{V_1^2 + V_2^2} \right)$$

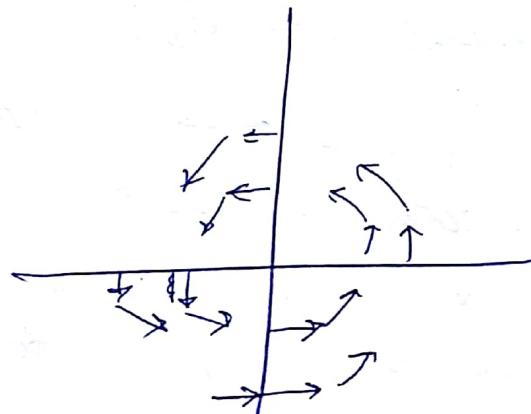
$$W_C \mathbf{V} \varphi = \sum_{\varphi_i \in \text{int}(C)} \text{ind } C \mathbf{V} \varphi_i$$

\* Technically should weight the index of each enclosed singularity by the signed number of times the curve  $C$  encloses that singularity. For simplicity and clarity, assume that curve  $C$  is homeomorphic to the circle  $S^1$ .

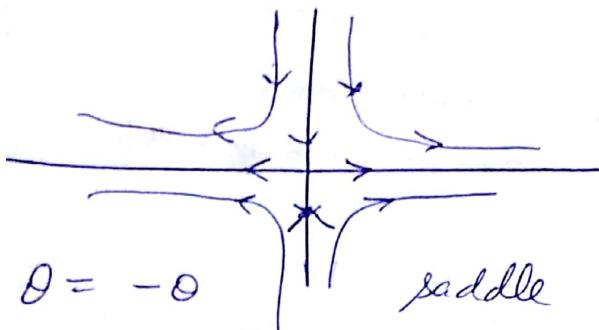


Index

Unstable  
Node.



Periodic cycle  
(purely imaginary)



$\theta = -\phi$  saddle

Index -1

$$1^{\text{st}} \quad v = c u_s g_s \rightarrow \theta = \phi \\ \text{diagram}$$



saddle

$\theta = \frac{\pi}{2} - \phi$

$$v = c(-y_s) u_s$$

$$\rightarrow \theta = \phi + \frac{\pi}{2}$$

+1

$$2^{\text{nd}} \quad v = c(u_s - y_s)$$

$$\rightarrow \theta = -\phi$$

$$v = c(y_s) u_s$$

$$\rightarrow \theta = -\phi + \frac{\pi}{2}$$

-1

$$v = (u_s - y_s^2, 2uy_s)$$

$$\theta = 2\phi$$

$$+2$$

→ single fixed pt  
at  $(0, 0)$  index +2

$$v = (u_s + y_s^2, u_s + 2uy_s)$$

↓  
two fixed pts:  
 $(u_s, y_s) = \{(0, 1), (0, -1)\}$

$$v = \begin{pmatrix} 2u_s - y_s \\ 1+2y_s \\ 2u_s \end{pmatrix}$$

$$m_{0,1} = \begin{pmatrix} 0 & -2 \\ 2 & 0 \end{pmatrix} \quad m_{0,-1} = \begin{pmatrix} 0 & 2 \\ 2 & 0 \end{pmatrix}$$

$$, \text{ each } +1 \quad +1$$

$$, \therefore w_{C \cap V} = +1 + (-1) = 0$$

### Properties of index/winding no.

- The index  $\varphi_0$  (vs of  $N=2$  vector field  $v$ ) at  $\varphi_0$  is winding no.

field  $v$  at pt  $\varphi_0$  is winding no. of  $v$  about that point.

- The winding number  $w_{C \cap V}$  of a curve  $C$  is the sum of indices of the singularities enclosed by that curve. Smooth deformations of  $C$  do not change the winding no. One must stretch  $C$  over a fixed pt singularity in order to change  $w_{C \cap V}$ .

- Uniformly rotating each vector in the vector field by angle  $\beta$  has the effect of sending  $\theta \rightarrow \theta + \beta$ . This leaves all indices and winding nos invariant.

- Nodes and spirals whether stable or unstable have index of +1 (as do the special cases of centers, stars and degenerate nodes)
- Saddle points have index -1.
- Clearly any closed orbit must lie on a dome of index +1.

### 3.6.1 Gauss-Bonnet Thm

Deep result in mathematics: connects local geometry of 2D manifold to global topological structure.

Content of theorem:

$$\int_M dA \kappa = 2\pi \chi(M) = 2\pi \sum_i \text{ind}(V_i)$$

where  $M$  is a 2D manifold (a topological space locally homeomorphic to  $\mathbb{R}^2$ )  
 $\kappa$  is the local gaussian curvature of  $M$ :

$$\kappa = (1/R_1 R_2)^2 \text{ where } R_1, R_2 \text{ are}$$

$\chi(M)$ : Euler characteristic of  $M$  principal radii of curvature.

given by  $\chi(M) = 2 - 2g$  where  $g$  is the genus of  $M$ , which

$\nabla(\varphi)$  is any smooth vector field, is the no. of holes (or handles) on  $M$ ,  $\varphi_i$  are the singularities of  $M$ .

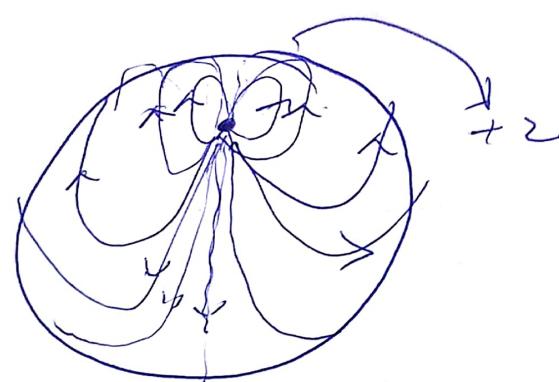
points of the vector field, which are fixed pts of the dynamics  $\dot{\varphi} = \nabla(\varphi)$

Consider  $M = S^2$  unit 2-sphere.

$$\text{at any pt } R_1=R_2=1 \rightarrow \kappa = 1 \quad \int_M dA \kappa = 4\pi$$

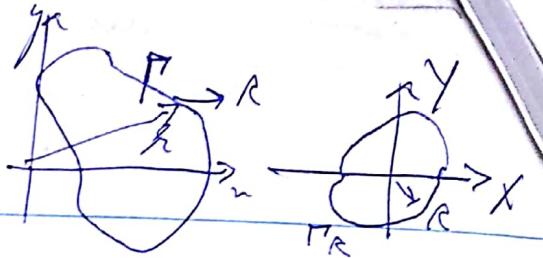
$$4\pi = (2\pi) \chi(M) \rightarrow \chi(M) = \sum_i \text{ind}(V_i) = 2 - 2g = 2$$

Any smooth vector field on  $S^2$  must have a singularity or singularities:  $\sum_i \text{ind}(V_i) = \chi(S^2) = 2$  for sphere



$$x = i \quad y = j$$

$$\frac{dy}{ds} = \frac{xy - yx}{x^2 + y^2}$$



$$IR = \frac{1}{2\pi} \oint_C \frac{d\theta}{ds} ds = \frac{1}{2\pi} \int_s \frac{xy - yx}{x^2 + y^2} ds$$

$$IR = \frac{1}{2\pi} \oint_C \frac{x dy - y dx}{x^2 + y^2}$$

$r(s) = (x(s), y(s))$  describes  $\Gamma_R$

$R(s) = (x, y)$  can be regarded as a position vector on a plane with  $x, y$  axis describes another curve  $\Gamma_R$

$\Gamma_R$  is closed since  $R$  returns to its original value after a complete cycle.

$\Gamma_R$  encircles origin  $IR$  times, anticlockwise

Theorem 3.1 Suppose on and inside  $\Gamma$ :  $x, y$  and their 1st derivatives are cont and  $x$  and  $y$  are not simultaneously 0. (No cusp or self-intersection) then  $IR$  is 0

green's theorem

$$IR = \frac{1}{2\pi} \oint_C \left[ \frac{\partial}{\partial x} \left( \frac{x}{x^2 + y^2} \right) + \frac{\partial}{\partial y} \left( \frac{y}{x^2 + y^2} \right) \right] dxdy$$

↳ region enclosed by  $\Gamma_R$

integrals are identically 0.

$$(u_n = -v_y)$$

green's theorem

$$\oint_C f dx + g dy = \iint_D \left( \frac{\partial g}{\partial x} - \frac{\partial f}{\partial y} \right) dA$$

## Index Theory

Strang NL  
lec 8 YouTube

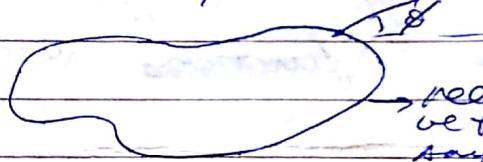
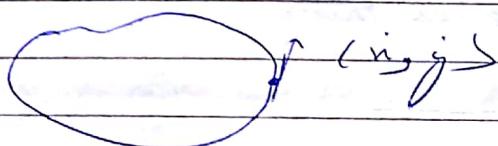
Aim: Provides "global" info about phase portrait.

Linearizing: local methods → However depends on info in <sup>and depth</sup> <sub>neighborhood of pt</sub>

Topology: complex - winding No. Physics: Gauss.

Index of a closed curve  $c$  → (does not intersect itself, <sup>and depth</sup> space through <sub>closed loop</sub>)  
 $c = \text{simple closed curve, not necessarily a closed trajectory}$ 

$$\phi = \tan^{-1} [y/x]$$

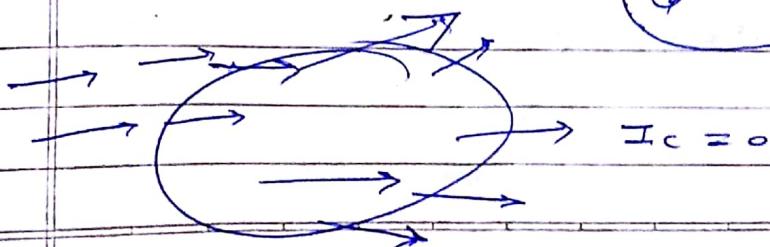
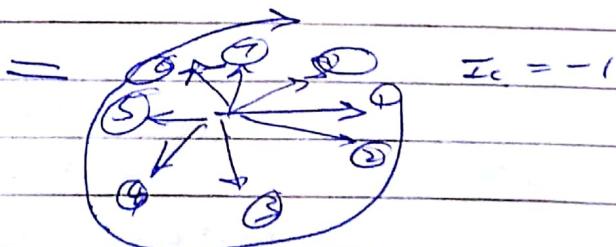
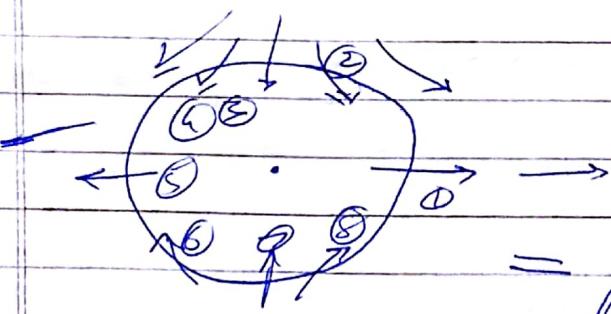
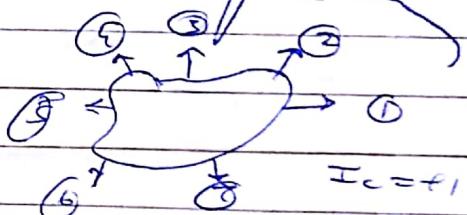


as  $n$  goes around  $c$  once, counterclockwise,  
 $\phi$  changes continuously.

If  $x, y$  are continuous, let  $\sum \phi I_c = \text{net charge in } \phi$   
around  $c$ . Then  $I_c = \frac{1}{2\pi} \sum \phi I_c$  is index of  $c$  w.r.t.  
vector field  $(x, y)$

$$\sum \phi I_c = 2\pi$$

$$I_c = +1$$

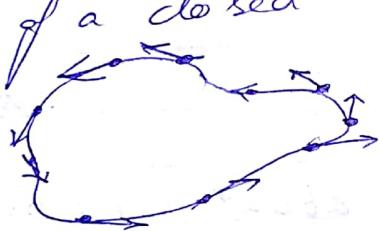


Teacher's Sign.: \_\_\_\_\_

... and stable or unstable

### Properties:

↳ Index of a closed trajectories



Any closed trajectory has index

$$I_C = +1$$

Questions?

Berry phase curvature related to index theory



Brouwer's Fixed Point and Sperner's Lemma  $\rightarrow$  Topology  
triangulation and index theory

Green's functions, divergence, Stokes' theorem

## Properties of the Index

- Suppose  $C$  can be deformed continuously into  $C'$  without passing through a fixed pt. Then  $\int_C \phi \, I_C = \int_{C'} \phi \, I_{C'}$
- (Deforming  $C$  into  $C'$  continuous then  $\int_C \phi \, I_C = \int_{C'} \phi \, I_{C'} =$  change continuously. But  $I_C$  is an integer Hence it can't change without jumping)
- If  $C$  doesn't enclose any fixed points, then  $I_C = 0$  by property ① → shrink  $C$  to a tiny circle without changing index. But  $\phi$  is essentially constant on such a circle, because all vectors point in nearly same direction thanks to the assumed smoothness of the vector field. Hence  $\int_C \phi \, I_C = 0 \quad I_C = 0$

3. If we reverse all arrows in vector field by changing  $t \rightarrow -t$ , the index is unchanged.
- $$\dot{x} = \frac{dx}{dt} = f_{\text{ext}} \quad \text{let } dt = -dt' \rightarrow \dot{x} = -f_{\text{ext}}$$
- All angles change from  $\phi$  to  $\pi + \phi$   $\int_C \phi \, I_C$  stays the same
- Suppose closed curve  $C$  is actually a trajectory for system, i.e.  $C$  is a closed orbit. Then  $I_C = +1$  because vector field is everywhere tangent to  $C$  because  $C$  is a trajectory

Index of a pt.

Let  $n^*$  be an isolated fixed pt.

$I_C$  is independent of  $C$  hence only dependent on  $n^*$ .



④ Index not related to stability

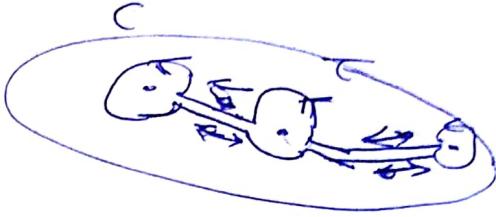
Spirals, centers, degenerate nodes and stars all have  $I=+$

⑤ saddle pt. is truly different.

Thm : If a closed curve  $C$  surrounds  $n$  isolated  
fixed pts.  $n_1, n_2, \dots, n_n$ . Then

$$I_C = I_1 + I_2 + \dots + I_n$$

$I_k$  is index of  $n_k$  for  $k = 1, \dots, n$



Thm Any closed orbit in the phase plane must enclose  
fixed pts whose indices sum to +1.

Corollary : Always at least one fixed pt inside  
any closed orbit in the phase plane.

If there is only one fixed point inside it, it cannot  
be a saddle point.

In every case the charge in  $\phi$  must be a multiple of  $2\pi$ .

$$\oint \phi \, I_F = 2\pi I_F$$

3.4

where  $I_F$  is defined and is an integer.

Index of  $\Gamma$  with  $\phi$  vector field  $(x, y)$

( $\Gamma$  is always counter clockwise)

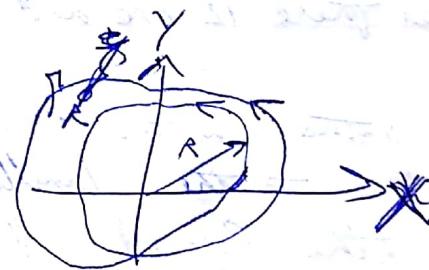
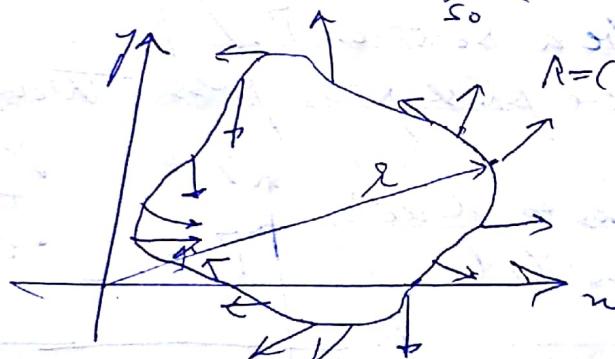
$$I_F = \frac{1}{2\pi} \oint ds = \frac{1}{2\pi} \int_{S_0}^{S_1} \frac{dx - y \, dy}{x^2 + y^2} ds$$

$$\frac{d \tan \phi}{ds} = \frac{dy}{dx}$$

$$\frac{d\phi}{ds} = \frac{xy' - yx'}{x^2 + y^2}$$

from 3.4

$$I_F = \frac{1}{2\pi} \int_{S_0}^{S_1} \frac{(d\phi)}{(ds)} ds = \frac{1}{2\pi} \int_{S_0}^{S_1} \frac{xy' - yx'}{x^2 + y^2} ds$$



As  $x(s)$  describes  $\Gamma$ ,  $R(s) = (x, y)$  describes another curve  $\Gamma_R$  on  $x$  &  $y$  plane.

$\Gamma_R$  is closed since it returns to original value after complete cycle.

$\Gamma_R$  encircles origin  $I_F$  times.

$$I_F = \frac{1}{2\pi} \oint_{\Gamma_R} \frac{xdy - ydx}{x^2 + y^2}$$

$$\left( \frac{dy}{dt} - \frac{dx}{dt} \right) ds$$

$$(dy - dx) \frac{ds}{dt}$$

Then suppose  $\Gamma$  lies in a simply connected region on which  $x, y$  and their 1st derivatives are continuous and  $x$  and  $y$  are not simultaneously 0 (No singularity)

then  $I_F \neq 0$

Proof: Green's theorem  $\oint P \, dx + Q \, dy = \iint_D \left( \frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dx \, dy$

$$\text{We can write } dx = x_n \frac{du}{dt} + x_y \frac{dy}{dt} \text{ similarly } dy = y_n \frac{du}{dt} + y_x \frac{dy}{dt}$$

$$\text{then } I_F = \oint_{\Gamma} \left( \frac{x_y - y_x}{x^2 + y^2} du + \frac{x_n - y_n}{x^2 + y^2} dy \right)$$

## Geometrical Aspects of Plane Autonomous systems

Intro 3

Intro

where the linear approximation is 0,

Jordan, Smith

index of a point provides supporting inform<sup>n</sup> on its nature and complexity in strongly nonlinear cases.

Phase diagram doesn't give info about behaviour of paths at infinity beyond boundaries (Global paths).

### 3.1 The index of a point.

Given the system:  $\dot{x} = X(x, y)$ ,  $\dot{y} = Y(x, y)$  (3.15)

let  $\Gamma$  be any smooth closed curve consisting of ordinary points of  $\Gamma$  (3.15). Let  $\Omega$  be a point on  $\Gamma$ .

Then there is one and only one phase path passing through  $\Omega$ .

The paths (without implication to direction) belong to the family described by the eqn:

$$\frac{dy}{dx} = \frac{Y(x, y)}{X(x, y)}$$

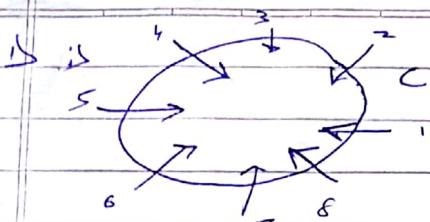
$s = (x, y)$  is tangential to phase path through the point and points in direction of increasing  $t$ .

$$\tan \phi = Y/X$$

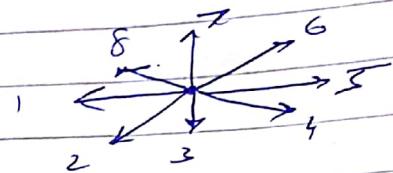
The curve  $\Gamma$  is traversed in the counterclockwise direction and variation of  $\phi$  is followed along it.  $\phi$  will not return to original value but differ by  $2\pi n$  after a full cycle.

Example.

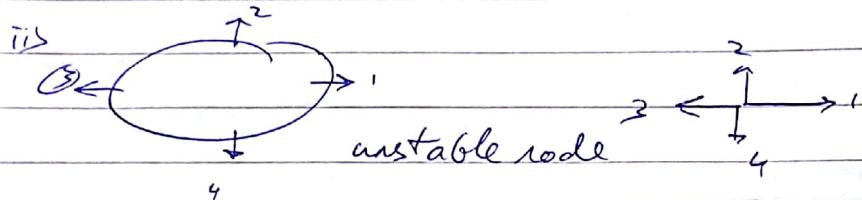
## Example



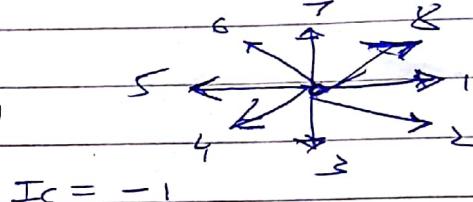
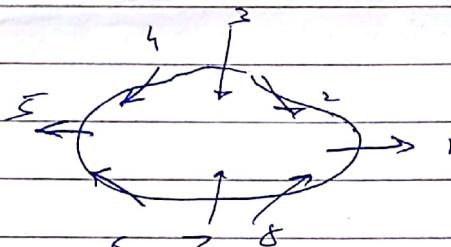
i) stable node



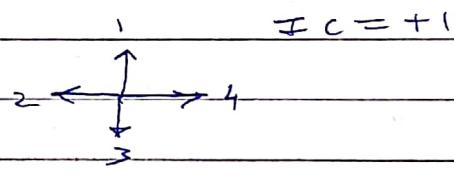
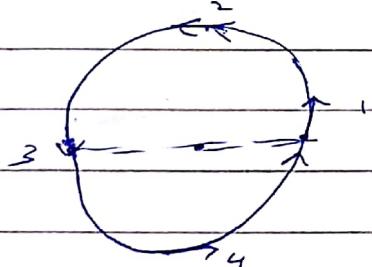
$I_C = +1$  anticlockwise



ii) saddle



iv) Center



$$83 \quad \begin{aligned} i &= u^2 - y^2 \\ j &= 2uy \end{aligned}$$
$$\begin{aligned} n &= \cos\theta & y &= \sin\theta & i &= u^2 - y^2 \\ \phi &= \tan^{-1} \frac{\cos(2\theta)}{\cos(\theta)} \end{aligned}$$

On unit circle

$$u^2 + y^2 = 1$$

$(1, 0)$	$(0, 1)$	$(-1, 0)$	$(0, -1)$
$\phi = 0$	$u = 0$	$i = 0$	$i = 0$
$j = \pm 1$	$j = -1$	$j = 1$	$j = -1$
$\frac{1}{\sqrt{2}}(1, 1) \rightarrow j = \frac{1}{\sqrt{2}}(1, 0)$	$\boxed{I_C = 0}$		

Teacher's Sign.: \_\_\_\_\_

$$\dot{x} = 2n^2 - 1$$
$$\dot{y} = 2ny$$

$$\tan \phi = \left( \frac{\sin 2\theta}{\cos 2\theta} \right) = \tan 2\theta$$
$$I_C = +2\phi$$

