### Answers to theory questions LB

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#### 1 Pros and cons of LB

#### 1.1 Pros

- Explicit update rule, no system of equations to solve.
  - Done by doing a trick. Introduce a new distribution function and see that the sum over states give the same global quantities.
- Easy to parallelize.
  - Each cell only depend locally on the neighbouring cells in the streaming step and only on itself in the collision.
- Simulating mesoscale
  - Unlike NS, LB looks at distributions of "bunches" of particles. May catch behaviour too fine for NS to detect.
- Can handle multiphase flows relatively easy.
  - Usually done by modifying the collision operator.

#### 1.2 Cons

- Explicit update rule
  - Need small timesteps.
  - Finite propagation speed of information.
- Only using a square grid can be limiting.

## 2 Convective-/viscous term

Consider the boltzmann equation,

$$\frac{\delta f}{\delta t} + \vec{v} \cdot \nabla f = -\Delta (f - f^{eq}), \tag{1}$$

and the BGK approximation

$$\Delta(f - f^{eq}) = -\frac{1}{\tau}(f - f^{eq}) \tag{2}$$

#### 3 Linear LBM to non-linear NSE

In the LBM we have a linear update rule,

$$f_i(\vec{x} + \vec{c}_i dt, t + dt) = f_i(\vec{x}, t) - \frac{1}{\tau} (f_i(\vec{x}, t) - f_i^{eq}), \tag{3}$$

which gives us the updated distribution of particles with velocity  $\vec{c_i}$  at position  $\vec{x} + \vec{c_i}dt$  and time t + dt. i = 0 : Q - 1 where Q is the size of our velocity space.

#### 3.1 Heuristic argument

The NSE solves a global system on a macroscale. To get the macroscopic velocities from the LBM velocities we need to integrate the velocity density distribution over the velocity set. We hence introduce nonlinearity.

#### 4 LBM memory usage

## 5 Conservation of mass and momentum in the collision rule

Collision rule:

$$f_i^*(\vec{x}, t) = f_i(\vec{x}, t) - \frac{1}{\tau} (f_i(\vec{x}, t) - f_i^{eq}(\vec{x}, t))$$
(4)

Approximation to the equilibrium state which holds when the distribution is close to equilibrium:

$$f_i^{eq}(\vec{x}, t) \approx w_i \rho \left[ 1 + \frac{\vec{c}_i \cdot \vec{u}}{c_s^2} + \frac{(\vec{c}_i \cdot \vec{u})^2}{c_s^4} - \frac{\vec{u} \cdot \vec{u}}{c_s^2} \right].$$
 (5)

Properties on  $w_i$  (there are more, but only these are needed here):

$$\sum_{i=0}^{Q-1} w_i c_i^{\alpha} c_i^{\beta} = c_s^2 \delta_{\alpha\beta}; \quad \alpha = x, y, z; \quad \beta = x, y, z.$$
 (6)

$$\sum_{i=0}^{Q-1} w_i = 1,\tag{7}$$

$$\sum_{i=0}^{Q-1} w_i c_i^{\alpha} = 0, \quad \alpha = x, y, z,$$
 (8)

#### 5.1 Conservation of mass

Before the collision step we have the density

$$\rho = \sum_{i=0}^{Q-1} f_i(\vec{x}, t). \tag{9}$$

After the collision the density is

$$\rho^* = \sum_{i=0}^{Q-1} f_i^*(\vec{x}, t). \tag{10}$$

Lets first look at the sum over the equilibrium distribution:

$$\sum_{i=0}^{Q-1} f_i^{eq}(\vec{x}, t) = \sum_{i=0}^{Q-1} \left( w_i \rho \left[ 1 + \frac{\vec{c}_i \cdot \vec{u}}{c_s^2} + \frac{(\vec{c}_i \cdot \vec{u})^2}{2c_s^4} - \frac{\vec{u} \cdot \vec{u}}{2c_s^2} \right] \right). \tag{11}$$

By virtue of eq. 6, 7 and 8 eq. 11 becomes after some calculation

$$\sum_{i=0}^{Q-1} f_i^{eq}(\vec{x}, t) = \rho. \tag{12}$$

Hence the after collison density becomes

$$\rho^* = \rho,\tag{13}$$

and we have mass conservation.

#### 5.2 conservation of momentum

Momentum before collision:

$$\rho \vec{u} = \sum_{i=0}^{Q-1} f_i(\vec{x}, t) \vec{c}_i. \tag{14}$$

After the collision the momentum is

$$\rho^* \vec{u}^* = \sum_{i=0}^{Q-1} f_i^* (\vec{x}, t) \vec{c_i}. \tag{15}$$

Lets first look at the sum over the equilibrium distribution:

$$\sum_{i=0}^{Q-1} w_i \vec{c_i} f_i^{eq}(\vec{x}, t) = \sum_{i=0}^{Q-1} \left( \vec{c_i} w_i \rho \left[ 1 + \frac{\vec{c_i} \cdot \vec{u}}{c_s^2} + \frac{(\vec{c_i} \cdot \vec{u})^2}{2c_s^4} - \frac{\vec{u} \cdot \vec{u}}{2c_s^2} \right] \right).$$
 (16)

By using eq. 8 the first and last term inside the square bracket becomes zero. The second term inside the square bracket becomes

$$\sum_{i=0}^{Q-1} \left( \frac{w_i \vec{c}_i (\vec{c}_i \cdot \vec{u}_i)}{2c_s^2} \right) = \vec{u}. \tag{17}$$

By expanding the third term and using eq. 6 we get

$$\sum_{i=0}^{Q-1} \left( w_i \vec{c}_i \frac{(\vec{c}_i \cdot \vec{u}_i)^2}{2c_s^4} \right) = \sum_{i=0}^{Q-1} \left( w_i \vec{c}_i \frac{\vec{c}_i \cdot \vec{c}_i + 2\vec{c}_i \cdot \vec{u} + \vec{u} \cdot \vec{u}}{2c_s^4} \right)$$
(18)

$$=\sum_{i=0}^{Q-1} \frac{w_i \vec{c}_i(\vec{c}_i \cdot \vec{c}_i)}{2c_s^4} + \frac{\vec{u}}{c_s^2}.$$
 (19)

# 6 Conservation of mass and momentum in the streaming rule