Application of Tensor Analysis in Physics

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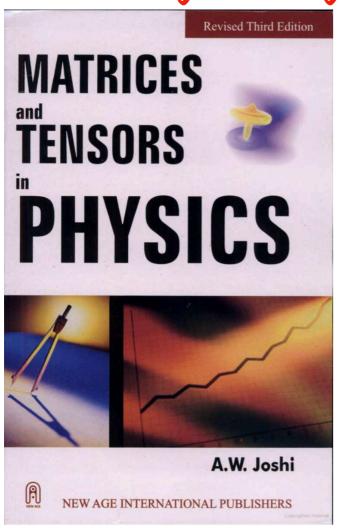
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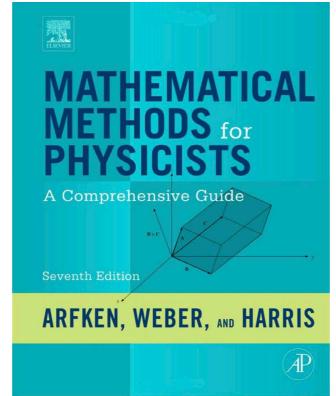
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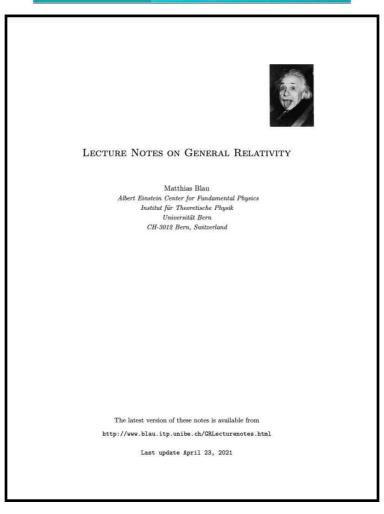


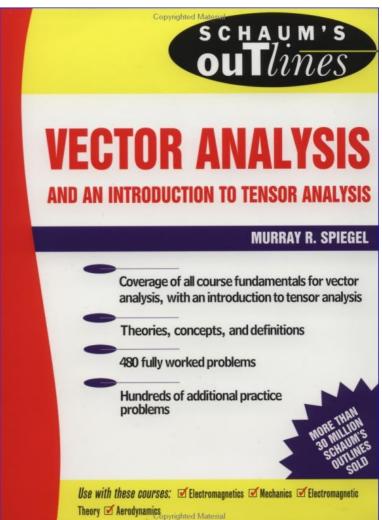
Teaching material for

School of Basic Science, Indian Institute of Technology, (Dandi, India Useful references and study materials









Lecture contents

- Coordinate transformation
- Tensor of rank zero (scalar)
- Tensor of rank one (vector): Contra variant and Co variant
- Tensor of rank two: Contra variant, Co variant and Mixed
- Tensor of higher ranks: Contra variant, Co variant and Mixed
- Various examples: Metric tensor, reciprocal tensor, associated, relative, absolute, symmetric and antisymmetric, permutation tensor
- Tensor calculus and its applications in Physics
- Assignments for students

Coordinate transformation

Forward transformation/mapping

$$x^1$$

$$\bar{x}^{\alpha} = \bar{x}^{\alpha}(x^1, x^2, \cdots, x^N)$$

$$\forall \quad \alpha = 1, 2, \cdots, N$$

$$x^2$$

$$d\bar{x}^{\alpha} = \frac{\partial \bar{x}^{\alpha}}{\partial x^{\beta}} dx^{\beta} = J^{\alpha}_{\beta}(\bar{x}, x) dx^{\beta} \qquad \forall \quad \alpha, \beta = 1, 2, \cdots, N$$

$$\forall \quad \alpha, \beta = 1, 2, \cdots, N$$

$$J^{\alpha}_{\beta}(\bar{x}, x) = \frac{\partial \bar{x}^{\alpha}}{\partial x^{\beta}} \qquad \forall \quad \alpha, \beta = 1, 2, \cdots, N$$

$$\forall \quad \alpha, \beta = 1, 2, \cdots, \Lambda$$

Inverse transformation/mapping

$$r^{N}$$

$$x^{\beta} = x^{\beta}(\bar{x}^1, \bar{x}^2, \cdots, \bar{x}^N) \qquad \forall \quad \beta = 1, 2, \cdots, N$$

$$dx^{\beta} = \frac{\partial x^{\beta}}{\partial \bar{x}^{\alpha}} d\bar{x}^{\alpha} = \left[J^{-1}(\bar{x}, x) \right]_{\alpha}^{\beta} d\bar{x}^{\alpha} \qquad \forall \quad \alpha, \beta = 1, 2, \cdots, N$$

$$\bar{x}^{\Lambda}$$

Inverse transformation matrix:

$$\left[J^{-1}(\bar{x},x)\right]_{\alpha}^{\beta} = \frac{\partial x^{\beta}}{\partial \bar{x}^{\alpha}} \quad \forall \quad \alpha,\beta = 1,2,\cdots,N$$

$$\forall \quad \alpha, \beta = 1, 2, \cdots, N$$

Tensor of rank zero (Scalar)

$$\mathbf{A} \ \phi = \phi(x^{\beta}) \qquad \bar{x}^{\alpha} = \bar{x}^{\alpha}(x^{\beta}) \qquad \forall \quad \alpha, \beta = 1, 2, \cdots, N \qquad \bar{\phi} = \bar{\phi}(\bar{x}^{\alpha}) \quad \mathbf{B}$$

Under the above mentioned coordinate transformation if we found

$$\phi = \bar{\phi}$$

Which means the function is invariant under the above mentioned coordinate transformation

Tensor of rank zero

Scalar

Simple example- Rotation about a specific axis/ arbitrary direction in 3D:

$$\bar{x}^{\alpha} = R^{\alpha}_{\beta} x^{\beta} \quad \forall \quad \alpha, \beta = 1, 2, 3$$

$$\bar{x}^{\alpha} = R^{\alpha}_{\beta} x^{\beta} \quad \forall \quad \alpha, \beta = 1, 2, 3$$

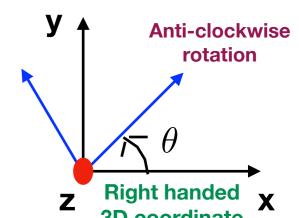
$$d\bar{x}^{\alpha} = R^{\alpha}_{\beta} dx^{\beta} \quad \forall \quad \alpha, \beta = 1, 2, 3$$

Rotation matrix:
$$R^{\alpha}_{\beta} = \frac{\partial \bar{x}^{\alpha}}{\partial x^{\beta}}$$
 $\forall \alpha, \beta = 1, 2, 3$ R: Orthogonal Matrix

$$\left[R^T R\right]_{\gamma}^{\alpha} = \left[R^{-1} R\right]_{\gamma}^{\alpha} = \frac{\partial x^{\alpha}}{\partial \bar{x}^{\beta}} \frac{\partial \bar{x}^{\beta}}{\partial x^{\gamma}} = \frac{\partial x^{\alpha}}{\partial x^{\gamma}} = \delta_{\gamma}^{\alpha} \qquad \forall \quad \alpha, \beta = 1, 2, 3$$

$$\bar{x}_{\alpha}\bar{x}^{\alpha} = (R_{\alpha}^{\gamma}x_{\gamma})\left(R_{\beta}^{\alpha}x^{\beta}\right) = \left(R_{\alpha}^{\gamma}R_{\beta}^{\alpha}\right)x_{\gamma}x^{\beta} = \delta_{\beta}^{\gamma}x_{\gamma}x^{\beta} = x_{\beta}x^{\beta} \quad \forall \quad \alpha, \beta, \gamma = 1, 2, 3$$

Length of a vector is unchanged under rotation (Scalar)



Tensor of rank one (Vector): Contra variant & Co variant

Contra variant vector:

$$\bar{A} \stackrel{\bar{x}^{\alpha} = \bar{x}^{\alpha}(x^{\beta})}{\forall \beta = 1, 2, \cdots, N} \stackrel{\bar{x}^{\alpha} = \bar{x}^{\alpha}(x^{\beta})}{\bar{A}^{\alpha}} \stackrel{\forall \alpha, \beta = 1, 2, \cdots, N}{\forall \alpha = 1, 2, \cdots, N} \bar{A}^{\alpha}(x^{\alpha}) \equiv \bar{A}^{\alpha} \\ \bar{A}^{\alpha} = \frac{\partial \bar{x}^{\alpha}}{\partial x^{\beta}} A^{\beta}$$

Transformation rule of a Contra variant vector or Contra variant tensor of rank one. Simple example: Velocity transform as a contra variant vector.

$$\bar{v}^{\alpha} = \frac{d\bar{x}^{\alpha}}{dt} = \frac{\partial \bar{x}^{\alpha}}{\partial x^{\beta}} \frac{dx^{\beta}}{dt} = \frac{\partial \bar{x}^{\alpha}}{\partial x^{\beta}} v^{\beta}$$

Co variant vector:

$$\bar{A} \stackrel{\bar{A}_{\beta}(x^{\beta}) \equiv A_{\beta}}{\forall \beta \equiv 1, 2, \cdots, N} \bar{A}_{\alpha}(\bar{x}^{\alpha}) \equiv \bar{A}_{\alpha} \\ \bar{A}_{\alpha} = \frac{\partial x^{\beta}}{\partial \bar{x}^{\alpha}} A_{\beta}$$

$$\bar{A}_{\alpha} = \frac{\partial x^{\beta}}{\partial \bar{x}^{\alpha}} A_{\beta}$$

Transformation rule of a Co variant vector or Co variant tensor of rank one. Simple example: Components of gradient of a scalar function transform as a co variant vector.

$$(\bar{\nabla}\phi)_i = \bar{\partial}_i \phi = \frac{\partial \phi}{\partial \bar{x}^i} = \frac{\partial x^j}{\partial \bar{x}^i} \ \frac{\partial \phi}{\partial x^j} = \frac{\partial x^j}{\partial \bar{x}^i} \ \partial_j \phi = \frac{\partial x^j}{\partial \bar{x}^i} \ (\nabla\phi)_j$$

Tensor of rank two: Contra variant, Co variant & Mixed

Contra variant tensor:

$$\mathbf{A} \quad \overline{A^{\beta_1\beta_2}(x^\beta) \equiv A^{\beta_1\beta_2}} \quad \overline{x}^\alpha = \overline{x}^\alpha(x^\beta) \quad \forall \quad \alpha, \beta = 1, 2, \cdots, N \\ \forall \quad \beta, \beta_1, \beta_2 = 1, 2, \cdots, N \quad \overline{A^{\alpha_1\alpha_2}(x^\alpha)} \equiv \overline{A^{\alpha_1\alpha_2}} \\ \overline{A}^{\alpha_1\alpha_2} = \frac{\partial \overline{x}^{\alpha_1}}{\partial x^{\beta_1}} \frac{\partial \overline{x}^{\alpha_2}}{\partial x^{\beta_2}} A^{\beta_1\beta_2} \quad \overline{A^{\beta_1\beta_2}}$$

Transformation rule of a Contra variant tensor of rank two.

Co variant tensor:

$$\bar{A} \stackrel{A_{\beta_1\beta_2}(x^\beta) \equiv A_{\beta_1\beta_2}}{\forall \beta,\beta_1,\beta_2=1,2,\cdots,N} \bar{A}_{\alpha_1\alpha_2} = \frac{\partial x^{\beta_1}}{\partial \bar{x}^{\alpha_1}} \frac{\partial x^{\beta_2}}{\partial \bar{x}^{\alpha_2}} A_{\beta_1\beta_2} \qquad \bar{A}_{\alpha_1\alpha_2=1,2,\cdots,N} \mathbf{B}$$

Transformation rule of a Co variant tensor of rank two.

Mixed tensor:

$$\bar{A}^{\alpha_1}_{\beta_2} \qquad \bar{A}^{\alpha_1}_{\alpha_2} = \frac{\partial \bar{x}^{\alpha_1}}{\partial x^{\beta_1}} \frac{\partial x^{\beta_2}}{\partial \bar{x}^{\alpha_2}} A^{\beta_1}_{\beta_2} \qquad \bar{A}^{\alpha_1}_{\alpha_2} = \frac{\partial \bar{x}^{\alpha_1}}{\partial x^{\beta_1}} \frac{\partial x^{\beta_2}}{\partial \bar{x}^{\alpha_2}} A^{\beta_1}_{\beta_2}$$

Transformation rule of a mixed tensor of rank two.

Simple example: Krönecker Delta function transform as a mixed tensor of rank two.
$$\bar{\delta}_k^j = \frac{\partial \bar{x}^j}{\partial \bar{x}^k} = \frac{\partial \bar{x}^j}{\partial x^m} \frac{\partial x^n}{\partial \bar{x}^k} \frac{\partial x^m}{\partial x^n} = \frac{\partial \bar{x}^j}{\partial x^m} \frac{\partial x^n}{\partial \bar{x}^k} \ \delta_n^m$$

Tensor of higher ranks: Contra variant, Co variant & Mixed

Contra variant tensor:

Transformation rule of a Contra variant tensor higher rank. Co variant tensor:

$$A \underbrace{A_{\beta_1\beta_2\cdots\beta_k}}_{\bar{A}_{\alpha_1\alpha_2\cdots\alpha_m}} = \frac{\partial x^{\beta_1}}{\partial \bar{x}^{\alpha_1}} \frac{\partial x^{\beta_2}}{\partial \bar{x}^{\alpha_2}} \cdots \frac{\partial x^{\beta_k}}{\partial \bar{x}^{\alpha_m}} A_{\beta_1\beta_2\cdots\beta_k} = \prod_{j=1,2,\cdots,m} \prod_{p=1,2,\cdots,k} \frac{\partial x^{\beta_p}}{\partial \bar{x}^{\alpha_j}} A_{\beta_1\beta_2\cdots\beta_k}$$

Transformation rule of a Co variant tensor of rank two.

Mixed tensor:

$$\mathbf{A} \begin{array}{c} \mathbf{A} \begin{array}{c} \mathbf{A} \\ \mathbf{A} \\ \boldsymbol{\gamma_{1} \gamma_{2} \cdots \gamma_{q}} \end{array} & \bar{x}^{\alpha} = \bar{x}^{\alpha}(x^{\beta}) \quad \forall \quad \alpha, \beta = 1, 2, \cdots, N \\ \bar{A}_{\lambda_{1} \lambda_{2} \cdots \lambda_{p}}^{\alpha_{1} \alpha_{2} \cdots \alpha_{m}} = \left(\frac{\partial \bar{x}^{\alpha_{1}}}{\partial x^{\beta_{1}}} \frac{\partial \bar{x}^{\alpha_{2}}}{\partial x^{\beta_{2}}} \cdots \frac{\partial \bar{x}^{\alpha_{m}}}{\partial x^{\beta_{k}}} \right) \left(\frac{\partial x^{\gamma_{1}}}{\partial \bar{x}^{\lambda_{1}}} \frac{\partial x^{\gamma_{2}}}{\partial \bar{x}^{\lambda_{2}}} \cdots \frac{\partial x^{\gamma_{q}}}{\partial \bar{x}^{\lambda_{p}}} \right) A_{\gamma_{1} \gamma_{2} \cdots \gamma_{q}}^{\beta_{1} \beta_{2} \cdots \beta_{k}} \\ = \prod_{j=1,2,\cdots,m} \prod_{i=1,2,\cdots,k} \prod_{a=1,2,\cdots,q} \prod_{b=1,2,\cdots,p} \frac{\partial \bar{x}^{\alpha_{j}}}{\partial x^{\beta_{i}}} \frac{\partial x^{\gamma_{a}}}{\partial \bar{x}^{\lambda_{b}}} A_{\gamma_{1} \gamma_{2} \cdots \gamma_{q}}^{\beta_{1} \beta_{2} \cdots \beta_{k}} \end{array}$$

Transformation rule of a mixed tensor of rank two.

Various examples

Metric tensor:

$$ds^2 = g_{\alpha\beta}dx^{\alpha}dx^{\beta} = g^{\alpha\beta}dx_{\alpha}dx_{\beta}$$

$$d\bar{s}^2 = \bar{g}_{\gamma\lambda}d\bar{x}^{\gamma}d\bar{x}^{\lambda} = \bar{g}^{\gamma\lambda}d\bar{x}_{\gamma}d\bar{x}_{\lambda}$$

$$ds^{2} = d\bar{s}^{2} \longrightarrow g_{\alpha\beta} = \bar{g}_{\gamma\lambda}J_{\alpha}^{\gamma}J_{\beta}^{\lambda} = \bar{g}_{\gamma\lambda} \frac{\partial \bar{x}^{\gamma}}{\partial x^{\alpha}} \frac{\partial \bar{x}^{\lambda}}{\partial x^{\beta}}$$

$$ds^{2} = d\bar{s}^{2} \longrightarrow g^{\alpha\beta} = \bar{g}^{\gamma\lambda} \left[J^{-1} \right]_{\alpha}^{\gamma} \left[J^{-1} \right]_{\beta}^{\lambda} = \bar{g}^{\gamma\lambda} \frac{\partial x^{\alpha}}{\partial \bar{x}^{\gamma}} \frac{\partial x^{\beta}}{\partial \bar{x}^{\lambda}}$$

Receprocal tensor:

$$g^{\beta\lambda} = \frac{\text{Cofactor of } g_{\beta\lambda}}{g}$$
 where $g = \text{Det}(g_{\alpha\beta}) \neq 0$

where
$$g = \text{Det}(g_{\alpha\beta}) \neq 0$$

$$g_{\alpha\beta}g^{\beta\lambda} = \delta^{\lambda}_{\alpha}$$

Associated tensor:

$$A^{\alpha}_{*\beta} = g^{\gamma\alpha} A_{\gamma\beta}$$

$$A^{\alpha}_{*\beta} = g^{\gamma\alpha} A_{\gamma\beta} \qquad A^{\alpha\beta} = g^{\gamma\alpha} g^{\lambda\beta} A_{\gamma\lambda}$$

Relative and absolute tensor:

$$\bar{A}_{\lambda_{1}\lambda_{2}\cdots\lambda_{p}}^{\alpha_{1}\alpha_{2}\cdots\alpha_{m}} = \left|\frac{\partial x}{\partial \bar{x}}\right|^{w} \left(\frac{\partial \bar{x}^{\alpha_{1}}}{\partial x^{\beta_{1}}} \frac{\partial \bar{x}^{\alpha_{2}}}{\partial x^{\beta_{2}}} \cdots \frac{\partial \bar{x}^{\alpha_{m}}}{\partial x^{\beta_{k}}}\right) \left(\frac{\partial x^{\gamma_{1}}}{\partial \bar{x}^{\lambda_{1}}} \frac{\partial x^{\gamma_{2}}}{\partial \bar{x}^{\lambda_{2}}} \cdots \frac{\partial x^{\gamma_{q}}}{\partial \bar{x}^{\lambda_{p}}}\right) A_{\gamma_{1}\gamma_{2}\cdots\gamma_{q}}^{\beta_{1}\beta_{2}\cdots\beta_{k}}$$

$$= \left|\frac{\partial x}{\partial \bar{x}}\right|^{w} \prod_{j=1,2,\cdots,m} \prod_{i=1,2,\cdots,k} \prod_{a=1,2,\cdots,q} \prod_{b=1,2,\cdots,p} \frac{\partial \bar{x}^{\alpha_{j}}}{\partial x^{\beta_{i}}} \frac{\partial x^{\gamma_{a}}}{\partial \bar{x}^{\lambda_{b}}} A_{\gamma_{1}\gamma_{2}\cdots\gamma_{q}}^{\beta_{1}\beta_{2}\cdots\beta_{k}}$$

If the weight factor

 $w = 0 \longrightarrow \text{Absolute tensor}$

 $w = 1 \longrightarrow (\text{Relative}) \text{ tensor density}$

Various examples

Symmetric tensor:

$$A_{\alpha\beta} = A_{\beta\alpha}$$

symmetric under index exchange $\alpha \longleftrightarrow \beta$

Anti-symmetric tensor:

$$A_{\alpha\beta} = -A_{\beta\alpha}$$

 $A_{\alpha\beta} = -A_{\beta\alpha}$ anti – symmetric under index exchange $\alpha \longleftrightarrow \beta$

Simple example: A rank two tensor can be expressed as a sum of an symmetric and anti-symmetric tensor

$$A_{\alpha\beta} = C_{\alpha\beta} + D_{\alpha\beta}$$

$$C_{\alpha\beta} = \frac{1}{2} (A_{\alpha\beta} + A_{\beta\alpha}) = \frac{1}{2} (A_{\beta\alpha} + A_{\alpha\beta}) = C_{\beta\alpha}$$
 symmetric under $\alpha \leftrightarrow \beta$

$$D_{\alpha\beta} = \frac{1}{2} \left(A_{\alpha\beta} - A_{\beta\alpha} \right) = -\frac{1}{2} \left(A_{\beta\alpha} - A_{\alpha\beta} \right) = -D_{\beta\alpha} \quad \text{anti-symmetric under} \quad \alpha \leftrightarrow \beta$$

Permutation tensor: Levi-Civita

$$arepsilon_{a_1 a_2 a_3 \dots a_n} = egin{cases} +1 & ext{if } (a_1, a_2, a_3, \dots, a_n) ext{ is an even permutation of } (1, 2, 3, \dots, n) \ -1 & ext{if } (a_1, a_2, a_3, \dots, a_n) ext{ is an odd permutation of } (1, 2, 3, \dots, n) \ 0 & ext{otherwise} \end{cases}$$

$$\varepsilon_{i_1\ldots i_n}\varepsilon^{j_1\ldots j_n}=\delta^{j_1\ldots j_n}_{i_1\ldots i_n}$$

$$arepsilon_{i_1\ldots i_k}{}_{i_{k+1}\ldots i_n}arepsilon^{i_1\ldots i_k}{}_{i_{k+1}\ldots j_n}=\delta_{i_1\ldots i_k}^{i_1\ldots i_k}{}_{i_{k+1}\ldots i_n}^{j_{k+1}\ldots j_n}=k!\;\delta_{i_{k+1}\ldots i_n}^{j_{k+1}\ldots j_n}$$

$$arepsilon_{i_1\ldots i_n}arepsilon^{i_1\ldots i_n}=n! \ arepsilon_{i_1i_2\ldots i_n}arepsilon_{j_1j_2\ldots j_n}=egin{bmatrix} \delta_{i_1j_1} & \delta_{i_1j_2} & \ldots & \delta_{i_1j_n} \ \delta_{i_2j_1} & \delta_{i_2j_2} & \ldots & \delta_{i_2j_n} \ dots & dots & dots & dots \ \delta_{i_nj_1} & \delta_{i_nj_2} & \ldots & \delta_{i_nj_n} \ \end{pmatrix}.$$

$$arepsilon_{a_1 a_2 a_3 \ldots a_n} = \prod_{1 \leq i < j \leq n} \operatorname{sgn}(a_j - a_i)$$

$$arepsilon_{ijk}arepsilon^{imn}=\delta_{j}{}^{m}\delta_{k}{}^{n}-\delta_{j}{}^{n}\delta_{k}{}^{m}$$

$$arepsilon_{jmn}arepsilon^{imn}=2{\delta_j}^i$$

$$arepsilon_{ijk}arepsilon^{ijk}=6.$$

Useful for vector calculus

Tensor calculus and its application in Physics

Vector calculus:

$$\vec{A}.\vec{B} = A_i B^j \delta^i_j = A_i B^i = A_i B_i$$

In 3D Euclidean geometry

$$ds^{2} = \delta^{ij} dx_{i} dx_{j} = dx_{i} dx_{i} = dx^{2} + dy^{2} + dz^{2}$$

$$(\vec{A} \times \vec{B})_{i} = \epsilon_{ijk} A^{j} B^{k} = \epsilon_{ijk} A_{j} B_{k}$$

$$(\vec{\nabla} \phi)_{i} = \partial_{i} \phi$$

$$\nabla^{2} \phi = \partial_{i} \partial^{i} \phi = \partial_{i} \partial_{i} \phi$$

$$(\vec{\nabla} \times \vec{A})_{i} = \epsilon_{ijk} \partial_{j} A_{k}$$

$$(\nabla^{2} \vec{A})_{i} = \partial_{j} \partial_{j} A_{i}$$

$$\vec{\nabla} \phi . \vec{\nabla} \psi = \partial_{i} \phi \ \partial_{i} \psi$$

$$\vec{\nabla} \phi . \vec{\nabla} \phi = \partial_{i} \phi \ \partial_{i} \phi = (\partial_{i} \phi)^{2}$$

$$A_{\mu}B^{\mu} = -A_0B_0 + A_iB_i$$

 $A_{\mu} = (-A_0, \vec{A}), \quad B^{\mu} = (B_0, \vec{B})$

In 3+1 D Minkowski flat geometry

$$ds^{2} = \eta^{\alpha\beta} dx_{\alpha} dx_{\beta} = -dt^{2} + dx^{2} + dy^{2} + dz^{2} \quad (in \ c = 1 \ unit)$$

$$\partial_{\mu} = \left(\partial_{t}, \vec{\nabla}\right), \quad \partial^{\mu} = \eta^{\mu\nu} \partial_{\nu} = \left(\partial_{t}, -\vec{\nabla}\right), \quad \Box = \eta^{\mu\nu} \partial_{\mu} \partial_{\nu} = \partial_{\mu} \partial^{\mu} = \partial_{t}^{2} - \nabla^{2}$$

$$\eta_{\mu\nu} = \operatorname{diag}\left(-1, 1, 1, 1\right) = \eta^{\mu\nu}$$

$$\eta_{\mu\nu} \eta^{\mu\nu} = (-1)^{2} + 3 \times (1)^{2} = 4 \quad \text{in } 3 + 1 \text{ D Minkowski flat space}$$

$$\eta_{\mu\nu} \eta^{\nu\alpha} = \delta^{\alpha}_{\mu}$$

$$\partial_{\mu} \phi = \left(\partial_{t}, \vec{\nabla}\right) \phi$$

$$\partial^{\mu} \phi = \left(\partial_{t}, -\vec{\nabla}\right) \phi$$

$$\Box \phi = \left(\partial_{t}^{2} - \nabla^{2}\right) \phi$$

$$\Box \psi - (\partial_t - \nabla) \psi$$

$$\Box A^{\mu} = (\partial_t^2 - \nabla^2) A^{\mu}$$

$$\Box C^{\mu\nu} = (\partial_t^2 - \nabla^2) C^{\mu\nu}$$

$$(\partial_{\mu}\phi)(\partial^{\mu}\phi) = (\partial_t\phi)^2 - (\partial_i\phi)^2$$

$$(\partial_{\mu}\phi)(\partial^{\mu}\psi) = (\partial_t\phi)(\partial_t\psi) - (\partial_i\phi)(\partial_i\psi)$$

Tensor calculus and its application in Physics

Electromagnetic theory in Minkowski flat space:

$$F_{\mu\nu} = \partial_{[\mu}A_{\nu]} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu} \qquad A_{\mu} = \left(-\Phi, \vec{A}\right)$$

$$F_{\mu\nu} = -F_{\nu\mu} \qquad \text{anti - symmetric under} \qquad \mu \leftrightarrow \nu$$

$$\eta_{\mu\nu} = \operatorname{diag}\left(-1, 1, 1, 1\right) = \eta^{\mu\nu} \qquad F^{\mu\nu} = \partial^{[\mu}A^{\nu]} = \eta^{\mu\alpha}\eta^{\nu\beta}F_{\alpha\beta}$$

$$\eta_{\mu\nu}\eta^{\mu\nu} = (-1)^2 + 3 \times (1)^2 = 4 \quad \text{in } 3 + 1 \text{ D Minkowski flat space}$$

$$\eta_{\mu\nu}\eta^{\nu\alpha} = \delta^{\alpha}_{\mu}$$

$$A^{\mu} = \eta^{\mu\nu}A_{\nu} = \left(\Phi, \vec{A}\right) \qquad \Box F^{\mu\nu} = 0 \qquad \to \text{ EM wave eqns}$$

$$\partial_{\mu} = \left(\partial_{t}, \vec{\nabla}\right), \quad \partial^{\mu} = \eta^{\mu\nu}\partial_{\nu} = \left(\partial_{t}, -\vec{\nabla}\right), \quad \Box = \eta^{\mu\nu}\partial_{\mu}\partial_{\nu} = \partial_{\mu}\partial^{\mu} = \partial_{t}^{2} - \nabla^{2}$$

$$\partial_{\mu}F^{\mu\nu} = 0 = \partial^{\mu}F_{\mu\nu} \quad \text{Gauss law of electric field \& Ampere's law in magnetism} \blacktriangleleft$$

$$G_{\mu\nu} = \frac{1}{2!}\epsilon_{\mu\nu\alpha\beta}F^{\alpha\beta} \qquad \text{Electromagnetic dual tensor}$$
Sourceless

 $\partial_{\mu}G^{\mu\nu} = 0 = \partial^{\mu}G_{\mu\nu} \longrightarrow \text{Bianchi identity}$

→ Faraday's law of EM induction & Gauss law of magnetism

Sourceless

Maxwell's

eans

Assignments for students

- Q1. Show that the elements of 3D rotation matrix (consider rotation about x,y,z and arbitrary axis) transform like a tensor. Apart from rotation matrix if we consider a general matrix then in that case is it possible to show that the elements transform like a tensor?
- Q2. Show that the vector product is unique to 3-D space, that is, only in three dimensions can we establish a one-to-one correspondence between the components of an antisymmetric tensor (second-rank) and the components of a vector.
- Q3. Show that if all the components of any tensor of any rank vanish in one particular coordinate system, they vanish in all coordinate systems.
- Q4. Prove that a necessary and sufficient condition that a tensor of rank R become an invariant by repeated contraction in that R be even and that the number of contra variant and co variant indices be equal to R/2.

Thanks for your time....