Quantum Computation Introduction

Truong-Son Van

Fulbright University Vietnam Reading Group

September 26, 2022

Logistics

Helpful videos:

```
https://www.tsvan.xyz/reading.html
Follow John Preskill's notes
http:
//theory.caltech.edu/~preskill/ph219/ph219_2021-22.html
```

Info (dates, speakers, references) is hosted on:

- Preskill's class: https://www.youtube.com/playlist?list= PLOojjrEqIyPy-1RRD8cTD_1F1hflo89Iu
- UC Berkeley Vazirani's class: https://www.youtube.com/ playlist?list=PLXEJgM3ycgQW5ysL69uaEdPoof4it6seB

Basics: Complex Vector Space

A *complex vector space* is a non-empty set \mathbb{V} , whose elements we call vectors, with three operations

- **2** Negation $-: \mathbb{V} \to \mathbb{V}$
- **3** Scalar multiplication: $\cdot : \mathbb{C} \times \mathbb{V} \to \mathbb{V}$

and a distinguished element called zero vector $0 \in \mathbb{V}$. The operations above obey usual rules with the scalar multiplication obeys rules for complex numbers.

Example 1

 \mathbb{C}^n , space of complex polynomials $\mathbb{C}[x]$, space of complex-valued square integrable functions $L^2(\mathbb{R};\mathbb{C})$.

(Un)fortunately, we do have to work with complex numbers...

Basics: Inner Product

Definition 2

Given a vector space \mathbb{V} , an *inner product* (dot product or scalar product) is a function

$$\langle \cdot, \cdot \rangle : \mathbb{V} \times \mathbb{V} \to \mathbb{C}$$

that satisfies



Basics: Hilbert Space

Definition 3

A Hilbert space is a complete vector space with an inner product.

Example 4

- lacktriangledown \mathbb{C}^n wiht the usual inner product

$$\langle f, g \rangle = \int \overline{f}g$$

Fact: finite dimensional Hilbert spaces are isomorphic to \mathbb{R}^n . This concept is more useful in infinite dimensional spaces.

Basics: Hilbert Space

Definition 5

Given a Hilbert space \mathcal{H} , a hermitian operator is a mapping $A:\mathcal{H}\to\mathcal{H}$ such that $A=A^\dagger$, where $A^\dagger:\mathrm{Dom}(A^\dagger)\subset\mathcal{H}\to\mathcal{H}$ is defined by the following relation

$$\langle Ax, y \rangle = \langle x, A^{\dagger}y \rangle$$
.

- For a bounded operator or in finite dimensions, hermitian is the same with self-adjoint. Issues come up when we deal with unbounded operators in infinite dimensions.
- We mostly deal with finite dimensions, however. So this is not an issue. Just get used to the terminology of physicists.

Basics: Quantum Mechanics

We approach this from the axiomatic point of view.

Axiom 1 (States)

A state is a complete description of a physical system. In quantum mechanics, a state is a **ray** in a Hilbert space.

Question

What is a ray?

Let \sim be an equivalence relation on ${\cal H}$ defined by

$$x \sim y \iff x = \alpha y, \alpha \in \mathbb{C}$$
.

In QM, a ray, donoted by $|\varphi\rangle$, is an element in the equivalence class \mathcal{H}/\sim .

When we write $|\varphi\rangle$, we typically think of representation φ of this equivalence class.

Notations (abuse):

$$\alpha |\varphi\rangle := |\alpha\varphi\rangle \equiv |\varphi\rangle , \qquad (1)$$

$$|\varphi\rangle + |\theta\rangle := |\varphi + \theta\rangle$$
 (2)

Be careful! Despite (1),

$$\alpha |\varphi\rangle + |\theta\rangle \neq |\varphi + \theta\rangle$$
.

Given $\varphi \in \mathcal{H}$, we denote $\langle \varphi |$ to be a linear functional (dual vector) mapping from $\mathcal{H} \to \mathbb{C}$, such that

$$\langle \varphi | (\theta) = \langle \varphi, \theta \rangle$$
.

By abuse of notation again, because we think of $|\theta\rangle$ to be the representation θ of the equivalence class, we write

$$\langle \varphi | \theta \rangle := \langle \varphi | (\theta) = \langle \varphi, \theta \rangle.$$

This is called Dirac notation.

Given $\varphi \in \mathcal{H}$, we denote $\langle \varphi |$ to be a linear functional (dual vector) mapping from $\mathcal{H} \to \mathbb{C}$, such that

$$\langle \varphi | (\theta) = \langle \varphi, \theta \rangle$$
.

By abuse of notation again, because we think of $|\theta\rangle$ to be the representation θ of the equivalence class, we write

$$\langle \varphi | \theta \rangle := \langle \varphi | (\theta) = \langle \varphi, \theta \rangle.$$

This is called Dirac notation. Incredibly annoying to mathematicians!



Axiom 2 (Observables)

An observable is a property of a physical system that in principle can be measured. In QM, an observable is a hermitian operator.

Facts:

- Hermitian operators have real eigenvalues
- ② A "compact" hermitian operator A can be decomposed into sums orthogonal projectors to eigenvectors.

$$A=\sum_n a_n E_n$$

where a_n 's are eigenvalues and E_n are the orthogonal projectors to eigenvectors.

Axiom 3 (Measurement)

A measurement is a process in which information about the state of a physical system is acquired by an observer. In QM, the measurement of an observable A prepares an eigenstate of A, and the observer learns the value of the corresponding eigenvalue.

As QM is inherently probabilistic, the probabilities for the outcomes are determined by "Born" rule:

$$\mathbb{P}(a_n) = ||E_n|\varphi\rangle||^2 = \langle \varphi|E_n|\varphi\rangle.$$

The expectation value of the measurement outcome:

Axiom 3 (Measurement)

A measurement is a process in which information about the state of a physical system is acquired by an observer. In QM, the measurement of an observable A prepares an eigenstate of A, and the observer learns the value of the corresponding eigenvalue.

As QM is inherently probabilistic, the probabilities for the outcomes are determined by "Born" rule:

$$\mathbb{P}(a_n) = ||E_n|\varphi\rangle||^2 = \langle \varphi|E_n|\varphi\rangle.$$

The expectation value of the measurement outcome:

$$\langle A \rangle = \sum a_n \mathbb{P}(a_n) = \langle \varphi | A | \varphi \rangle$$
.

- 4 ロト 4 個 ト 4 恵 ト 4 恵 ト - 恵 - 夕 Q ()

Axiom 4 (Time evolution)

Time evolution is determined by the Schrödinger equation:

$$rac{d}{dt}\ket{\Psi(t)}=-iH(t)\ket{\Psi(t)}$$

Axiom 5 (Composite systems)

The Hilber space of composite system AB is the tensor product of A and B:

$$\mathcal{H}_{AB} = \mathcal{H}_{A} \otimes \mathcal{H}_{B}$$

of dimension d_Ad_B and orthonormal basis $\{|e_i\rangle_A\otimes|k_j\rangle_B\}$.

↓□▶ ↓□▶ ↓□▶ ↓□▶ ↓□ ♥ ♀○

Rule: if
$$|a\rangle=\sum a_i\,|e_i\rangle\in\mathcal{H}_A$$
 and $b=\sum b_j\,|k_j\rangle\in\mathcal{H}_B$, then $|a\rangle\otimes|b\rangle:=\sum_{i,j}a_ib_j\,|e_i\rangle\otimes|k_j\rangle$.

Rule: if
$$|a\rangle=\sum a_i\,|e_i\rangle\in\mathcal{H}_A$$
 and $b=\sum b_j\,|k_j\rangle\in\mathcal{H}_B$, then $|a\rangle\otimes|b\rangle:=\sum_{i,j}a_ib_j\,|e_i\rangle\otimes|k_j\rangle$.

The rest of QM just flows from here, up to our imagination.

Qubits

The building block of classical information is the bit, which takes values in $\{0,1\}$. The building block of quantum information is the qubit.

Qubits

The building block of classical information is the bit, which takes values in $\{0,1\}$. The building block of quantum information is the qubit.

Question: why qubits? other than they live in the simplest non-trivial quantum system?

Qubits

$$\mathcal{H} = \operatorname{span}\{|0\rangle, |1\rangle\}, \dim(\mathcal{H}) = 2.$$

Think

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \;, \quad |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \;.$$

All the possible *unitary* transformations are the Pauli operators:

$$I\,,\quad \sigma_1=\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}\,,\quad \sigma_2=\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}\,,\quad \sigma_3=\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}\,.$$

Then

$$|\Psi\rangle = a|0\rangle + b|1\rangle$$
.

Measure σ_3 :

$$\mathbb{P}(|0\rangle) = |a|^2 \qquad \mathbb{P}(|1\rangle) = |b|^2$$
.

