Quantum Computation Introduction

Truong-Son Van

Fulbright University Vietnam Reading Group

September 26, 2022

Logistics

Info (dates, speakers, references) is hosted on:
`https://www.tsvan.xyz/reading.html`
Follow John Preskill's notes
`http://theory.caltech.edu/~preskill/ph219/ph219_2021-22.html`
Helpful videos:

1. Preskill's class: `https://www.youtube.com/playlist?list=PL0ojjrEqIyPy-1RRD8cTD_lF1hflo89Iu`
2. UC Berkeley Vazirani's class: `https://www.youtube.com/playlist?list=PLXEJgM3ycgQW5ysL69uaEdPoof4it6seB`

Basics: Complex Vector Space

A *complex vector space* is a non-empty set $\mathbb{V}$, whose elements we call vectors, with three operations

1. Addition $+ : \mathbb{V} \times \mathbb{V} \to \mathbb{V}$
2. Negation $- : \mathbb{V} \to \mathbb{V}$
3. Scalar multiplication: $\cdot : \mathbb{C} \times \mathbb{V} \to \mathbb{V}$

and a distinguished element called *zero vector* $0 \in \mathbb{V}$. The operations above obey usual rules with the scalar multiplication obeys rules for complex numbers.

Example 1

$\mathbb{C}^n$, space of complex polynomials $\mathbb{C}[x]$, space of complex-valued square integrable functions $L^2(\mathbb{R}; \mathbb{C})$.

(Un)fortunately, we do have to work with complex numbers...

Basics: Inner Product

Definition 2

Given a vector space $\mathbb{V}$, an *inner product* (dot product or scalar product) is a function

$$\langle \cdot, \cdot \rangle : \mathbb{V} \times \mathbb{V} \to \mathbb{C}$$

that satisfies

1. $\langle V, V \rangle \geq 0$ and $\langle V, V \rangle = 0$ iff $V = 0$.
2. $\langle V_1 + V_2, V_3 + V_4 \rangle = \langle V_1, V_3 \rangle + \langle V_2, V_4 \rangle$.
3. $\langle cV_1, V_2 \rangle = c\langle V_1, V_2 \rangle$ and $\langle V_1, cV_2 \rangle = \overline{c}\langle V_1, V_2 \rangle$.
4. $\langle V_1, V_2 \rangle = \overline{\langle V_2, V_1 \rangle}$.

Basics: Hilbert Space

Definition 3

A Hilbert space is a complete vector space with an inner product.

Example 4

1. $\mathbb{C}^n$ wiht the usual inner product
2. $L^2(\mathbb{R}; \mathbb{C})$ where

$$\langle f, g \rangle = \int \overline{f} g$$

Fact: finite dimensional Hilbert spaces are isomorphic to $\mathbb{R}^n$. This concept is more useful in infinite dimensional spaces.

Basics: Hilbert Space

Definition 5

Given a Hilbert space $\mathcal{H}$, a hermitian operator is a mapping $A : \mathcal{H} \to \mathcal{H}$ such that $A = A^\dagger$, where $A^\dagger : \mathrm{Dom}(A^\dagger) \subset \mathcal{H} \to \mathcal{H}$ is defined by the following relation

$$\langle Ax, y \rangle = \langle x, A^\dagger y \rangle .$$

- For a bounded operator or in finite dimensions, hermitian is the same with self-adjoint. Issues come up when we deal with unbounded operators in infinite dimensions.
- We mostly deal with finite dimensions, however. So this is not an issue. Just get used to the terminology of physicists.

Basics: Quantum Mechanics

Basics: Quantum Mechanics (đều)

We approach this from the axiomatic point of view.

Axiom 1 (States)

*A state is a complete description of a physical system. In quantum mechanics, a state is a **ray** in a Hilbert space.*

Basics: Quantum Mechanics (đều)

> **Question**
> *What is a ray?*

Let $\sim$ be an equivalence relation on $\mathcal{H}$ defined by

$$x \sim y \iff x = \alpha y, \alpha \in \mathbb{C}.$$

In QM, a ray, donoted by $|\varphi\rangle$, is an element in the equivalence class $\mathcal{H}/\sim$.

When we write $|\varphi\rangle$, we typically think of representation φ of this equivalence class.

Basics: Quantum Mechanics (đều)

Notations (abuse):

$$\alpha \left|\varphi\right\rangle := \left|\alpha\varphi\right\rangle \equiv \left|\varphi\right\rangle , \tag{1}$$

$$\left|\varphi\right\rangle + \left|\theta\right\rangle := \left|\varphi + \theta\right\rangle . \tag{2}$$

Be careful! Despite (1),

$$\alpha \left|\varphi\right\rangle + \left|\theta\right\rangle \neq \left|\varphi + \theta\right\rangle .$$

Basics: Quantum Mechanics (đều)

Given $\varphi \in \mathcal{H}$, we denote $\langle\varphi|$ to be a linear functional (dual vector) mapping from $\mathcal{H} \to \mathbb{C}$, such that

$$\langle\varphi|\,(\theta) = \langle\varphi, \theta\rangle.$$

By abuse of notation again, because we think of $|\theta\rangle$ to be the representation θ of the equivalence class, we write

$$\langle\varphi|\theta\rangle := \langle\varphi|\,(\theta) = \langle\varphi, \theta\rangle.$$

This is called Dirac notation.

Basics: Quantum Mechanics (dểu)

Given $\varphi \in \mathcal{H}$, we denote $\langle\varphi|$ to be a linear functional (dual vector) mapping from $\mathcal{H} \to \mathbb{C}$, such that

$$\langle\varphi|\left(\theta\right) = \langle\varphi, \theta\rangle.$$

By abuse of notation again, because we think of $|\theta\rangle$ to be the representation θ of the equivalence class, we write

$$\langle\varphi|\theta\rangle := \langle\varphi|\left(\theta\right) = \langle\varphi, \theta\rangle.$$

This is called Dirac notation.
Incredibly annoying to mathematicians!

Basics: Quantum Mechanics (đều)

Axiom 2 (Observables)

An observable is a property of a physical system that in principle can be measured. In QM, an observable is a hermitian operator.

Facts:

1. Hermitian operators have real eigenvalues
2. A "compact" hermitian operator A can be decomposed into sums orthogonal projectors to eigenvectors.

$$A = \sum_n a_n E_n$$

where a_n's are eigenvalues and E_n are the orthogonal projectors to eigenvectors.

Basics: Quantum Mechanics (đều)

Axiom 3 (Measurement)

A measurement is a process in which information about the state of a physical system is acquired by an observer. In QM, the measurement of an observable A prepares an eigenstate of A, and the observer learns the value of the corresponding eigenvalue.

As QM is inherently probabilistic, the probabilities for the outcomes are determined by "Born" rule:

$$\mathbb{P}(a_n) = ||E_n |\varphi\rangle||^2 = \langle\varphi| E_n |\varphi\rangle .$$

The expectation value of the measurement outcome:

Basics: Quantum Mechanics (đều)

Axiom 3 (Measurement)

A measurement is a process in which information about the state of a physical system is acquired by an observer. In QM, the measurement of an observable A prepares an eigenstate of A, and the observer learns the value of the corresponding eigenvalue.

As QM is inherently probabilistic, the probabilities for the outcomes are determined by "Born" rule:

$$\mathbb{P}(a_n) = ||E_n \, |\varphi\rangle \, ||^2 = \langle\varphi| \, E_n \, |\varphi\rangle \ .$$

The expectation value of the measurement outcome:

$$\langle A \rangle = \sum a_n \mathbb{P}(a_n) = \langle\varphi| \, A \, |\varphi\rangle \ .$$

Basics: Quantum Mechanics (đều)

Axiom 4 (Time evolution)

Time evolution is determined by the Schrödinger equation:

$$\frac{d}{dt}\left|\Psi(t)\right\rangle = -iH(t)\left|\Psi(t)\right\rangle$$

Axiom 5 (Composite systems)

The Hilber space of composite system AB is the tensor product of A and B:

$$\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$$

of dimension $d_A d_B$ and orthonormal basis $\{\left|e_i\right\rangle_A \otimes \left|k_j\right\rangle_B\}$.

Basics: Quantum Mechanics (đều)

Rule: if $|a\rangle = \sum a_i |e_i\rangle \in \mathcal{H}_A$ and $b = \sum b_j |k_j\rangle \in \mathcal{H}_B$, then

$$|a\rangle \otimes |b\rangle := \sum_{i,j} a_i b_j |e_i\rangle \otimes |k_j\rangle \ .$$

Basics: Quantum Mechanics (đều)

Rule: if $|a\rangle = \sum a_i |e_i\rangle \in \mathcal{H}_A$ and $b = \sum b_j |k_j\rangle \in \mathcal{H}_B$, then

$$|a\rangle \otimes |b\rangle := \sum_{i,j} a_i b_j |e_i\rangle \otimes |k_j\rangle \ .$$

The rest of QM just flows from here, up to our imagination.

Qubits

The building block of classical information is the bit, which takes values in $\{0, 1\}$. The building block of quantum information is the qubit.

Qubits

The building block of classical information is the bit, which takes values in $\{0, 1\}$. The building block of quantum information is the qubit.

Question: why qubits? other than they live in the simplest non-trivial quantum system?

Qubits

$\mathcal{H} = \operatorname{span}\{|0\rangle, |1\rangle\}$, $\dim(\mathcal{H}) = 2$.
Think

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

All the possible *unitary* transformations are the the Pauli operators:

$$I, \quad \sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

Then, a qubit looks like

$$|\varphi\rangle = \alpha |0\rangle + \beta |1\rangle.$$

As a convention, we would like to normalize the coefficients α and β so that $||\,|\varphi\rangle\,||^2 = 1$.

Qubits

Measure the effect of σ_3. Recall Axiom 3.
Two different eigenstates:

$$|0\rangle , \quad |1\rangle .$$

$$\sigma_3 = 1 \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} - 1 \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} = E_1 - E_2 .$$

Probability for each eigenstate is the square of the projection:

$$\mathbb{P}(1) = ||E_1 |\varphi\rangle||^2 = |\alpha|^2 \qquad \mathbb{P}(-1) = ||E_2 |\varphi\rangle||^2 = |\beta|^2 .$$

Bloch sphere representation of Qubits

Theorem 6

$$|\varphi\rangle = \cos\left(\frac{\theta}{2}\right) + e^{i\phi}\sin\left(\frac{\theta}{2}\right)|1\rangle \ .$$