

Classical Mechanics And Relativity

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Plan of the Course

28 Sep 2020

- Lectures with notes
- Notes sent earlier and
- Recorded lectures posted earlier
- Regular Tutorials (Most essential) and help sessions from your tutors

Online interactions in B slot: 9:30 AM to 11 AM
Exceptions and additional sessions will be notified

- Assignments
 - Formal assignment sheets
 - Problems posed during lectures and discussions
- Technical skill
- conceptual clarity
 - Partly discussed in tutorials/discussions
 - You may form groups with about 5 students in each group and discuss and submit joint solutions

Resource Material

Books

- No single prescribed text book
 - Landau and Lifschitz Vol 1 (for NRCM)
 - Goldstein (for NRCM)
 - Landau and Lifschitz Vol 2 (for relativity) not in great depth

- Attendance automatic
- Submission of assignments not mandatory
- Grade based on Minor and Major performances
 - Minor of 2 hour duration
 - Major of three hour duration
- Both the exams will be open book open internet (tentative plan)

Huge syllabus

- Space Time Symmetries
- Kinematic symmetries
- Newton's Laws of Motion
- Galilean Relativity
 - Contact forces
 - Action at a distance
 - Galilean transformations (with limitations)
 - **The Galilean group***

Topics covered II

More formal aspects

- Constraints
- The action principle item generalised coordinates
- Nöther symmetries and conservation laws
- Scaling symmetries and Virial theorem

Topics covered III

Physical example 1

- Central forces
- The $\frac{1}{r}$ potential
- Keplerian Orbits

Topics covered III

Physical example 2

- Simple harmonic Oscillator
- Coupled Oscillators and normal modes
- Transition to continuum mechanics

Topics Covered IV

Hamiltonian formalism

- Legendre transformations
- Hamilton's equations and Phase space
- Canonical transformations
- Poisson brackets: Dynamics and Symmetries

Topics Covered V

Rigid Body dynamics

- Moment of inertia tensor
- Body fixed and space fixed axes
- The Top and gyroscopic motion
- Rotating coordinate systems

Topics Covered VII

Hamilton Jacobi equations

- Separation of variables
- Glimpses of perturbation theory

Topics Covered VIII

Special Relativity

- Incompatibility between Electrodynamics and Galilean transformations
- Michelson Morley and other experiments
- The Lorentz transformations
- Length contraction, time dilatation and relativity of simultaneity
- Four vectors and tensors
- Relativistic dynamics

Additional Topics

Wishful thinking

- The three body problem
- deterministic chaos
- and what not

Scope of Classical Mechanics

- Classical: *not quantum*

- Newtonian or Nonrelativistic classical mechanics
- Relativistic Classical Mechanics

Lecture 2

Relevant Natural Scales

Quantum Domain: $\hbar \sim 10^{-34}$

j-s

Einsteinian domain:

$c = 2.99792458 \times 10^8$ m/s

Lecture 2

Boundaries and Constraints

Non Quantum \implies

Angular momentum $L \gg \hbar$

Nonrelativistic \implies

Speeds $v \ll c$

Lecture 2

Practice Examples

- Gravitational interaction between two neighbouring galaxies
- A star of one solar mass at a distance of 10^6 m from a black hole of 10^6 solar mass
- An electron, at room temperature, in a magnetic field of 1T
- An electron initially at rest, in an electrostatic field of 10^8 V/m
- A system of 10^6 hydrogen atoms in a container at a temperature $10^{-11} K$. Any constraint on container size?
- The interior of a neutron star

Non Relativistic CM

- Large mass
- Weak fields
- High temperatures
- Low pressures

Considering Every Thing

- No preferred location or origin (homogeneity)
- No preferred orientation (isotropy)
- No preferred event to mark time (homogeneity)

$$D = 2+1$$

- The Euclidean plane \mathbb{R}^2
- A sphere of "radius" R
- A torus?
- Cyclic time (Circle)
- Time coordinate given by \mathbb{R}^1

Geometric Aspects

Sphere

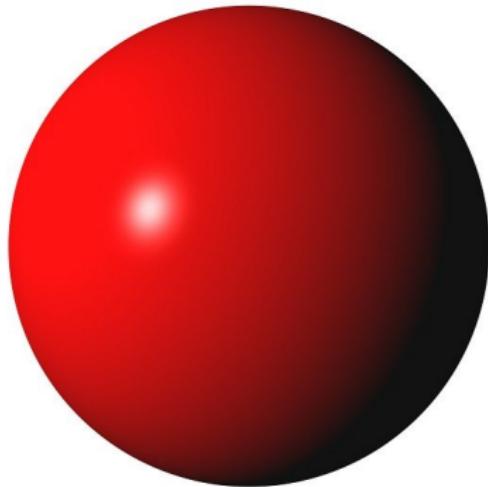


Figure: homogeneous and isotropic

Geometric aspects

Torus????

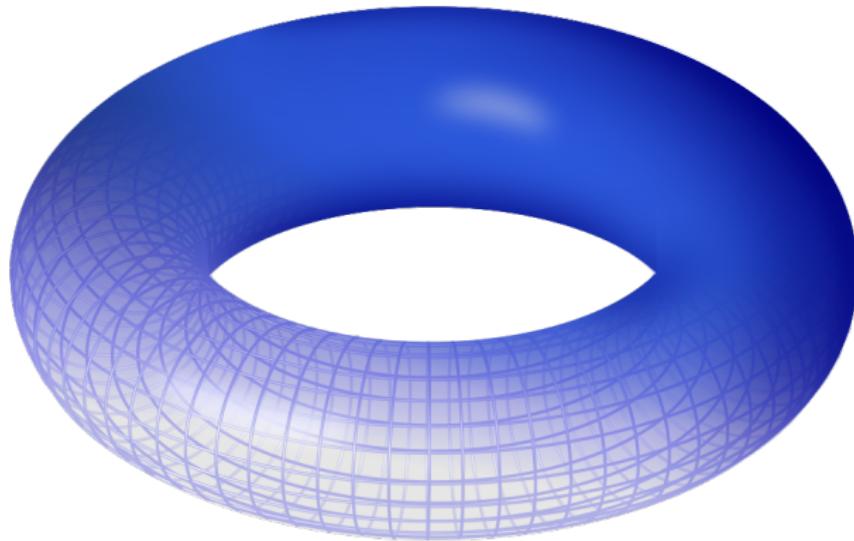


Figure: homogeneous and isotropic?

Geometric Aspects

Observations

- Space is three dimensional
- Space is Euclidean. Global rectangular Cartesian coordinate system
- Time is homogeneous

Geometric Aspects

Distance and angle

Let $P_1(x_1, y_1, z_1)$ and $P_2(x_2, y_2, z_2)$. Then,

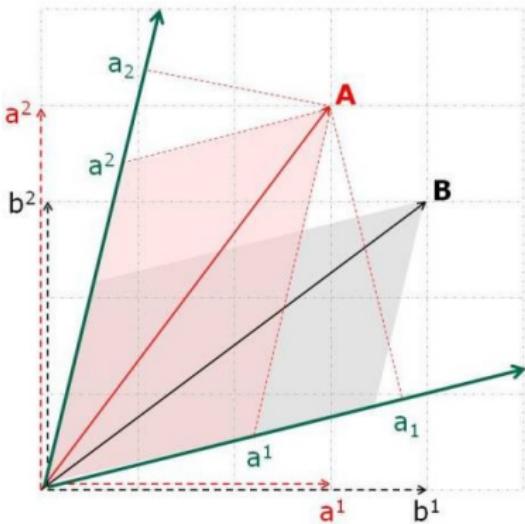
$$d(P_1, P_2) = \left\{ (x_2 - x_1)^2 + (y_2 - y_1)^2 + (z_2 - z_1)^2 \right\}^{\frac{1}{2}}$$
$$\equiv \|\vec{r}_1 - \vec{r}_2\|$$

$$\angle(P_1, P_2) = \frac{x_1 x_2 + y_1 y_2 + z_1 z_2}{\|\vec{r}_1\| \|\vec{r}_2\|} \equiv \frac{\vec{r}_1 \cdot \vec{r}_2}{\|\vec{r}_1\| \|\vec{r}_2\|}$$

Geometric Aspects

Other coordinate systems

Oblique coordinate system

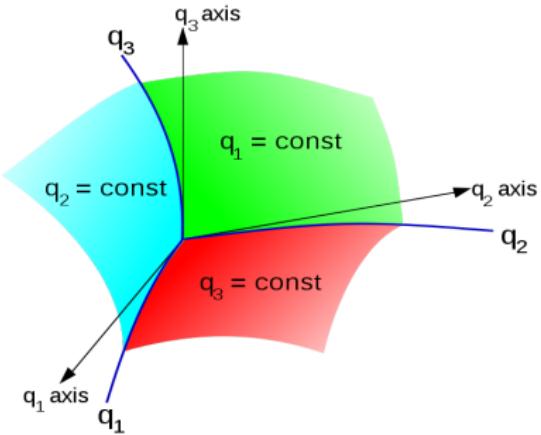


Credit:

<http://www.mysearch.org.uk/website1/html/297.Components.html>

Geometric Aspects

Curvilinear systems



Credit:

<https://commons.wikimedia.org/w/index.php?curid=2074623>

Kinematic Constraints

Galilean Transformations

- Lengths of Rods (Distances between points)
- Rates of clocks
- Masses of bodies

Laws of Dynamics

- Forces
- Frames
- Newton's Laws

Basic hypothesis

- Forces are of physical origin
- Prior knowledge
 - Contact forces: Push and Pull
 - Muscular response
 - Intrinsic existence
- Forces are additive (vectors)

$$\vec{F}_t = \vec{F}_1 + \vec{F}_2$$

Forces

Towards Free Bodies

- Forces decrease with separations
- Remove the sources or
- Neutralise

Definition: A body is free if the net force acting on it vanishes

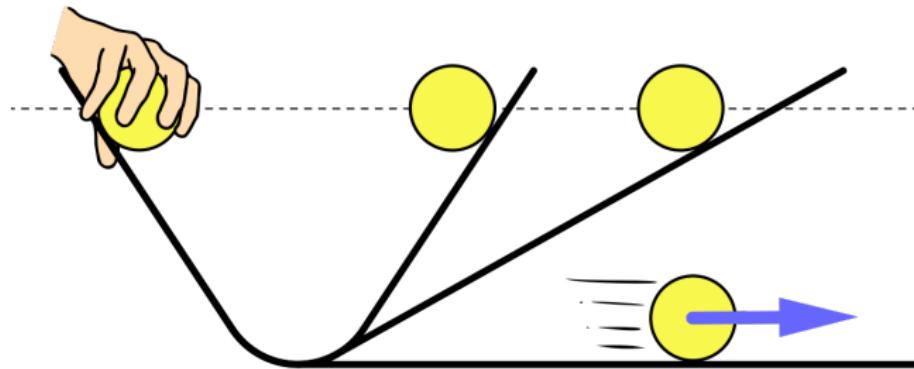
Inertial Frames (IF)

Special coordinate systems in which a free body does not accelerate

Corollary: Free bodies have only uniform velocity in IF

A thought experiment

Galileo



Credit: <https://commons.wikimedia.org/wiki/File:Thought-experiment-on-inertia.svg>

Static Transformations T

- Let \mathcal{C} be an *IF*. Then $\bar{\mathcal{C}} = T\mathcal{C}$ is an *IF*.
- Special Case: The Affine Group
 - \mathcal{A} , the set of all linear transformations (translations + rotations + oblique) forms a group, preserving Galilean law of inertia.

The second Law

Momentum

\vec{F} : an applied force; \vec{P} the momentum of the body.
Then,

$$\frac{d\vec{P}}{dt} = \vec{F}$$

- LHS is purely kinematic (effect)
- RHS is dynamics, known independently (Cause)

Examples

- Lorentz
- Hooke
- Gravitation
- Van der Waals
- Stokes

Definition of Momentum

entwined fundamental question

Inertia or Mass

$$\vec{P} = M(V) \vec{V}$$

The Fundamental Question

$$M(V) = ?$$

- Rate of moving clocks
- Length of moving rods
- Magnitude of uppermost speed (if any)

Newtonian Mechanics

Summary of the framework

- Space is Euclidean
- Space is homogeneous and isotropic
- Time is isotropic
- Mass length and clock rates do not depend on the state of motion of bodies

Galilean Transformations

Connecting inertial Frames

Canonical coordinate systems: Rectangular Cartesian

Affine transformations + transformations to moving frames: $S \rightarrow S'$

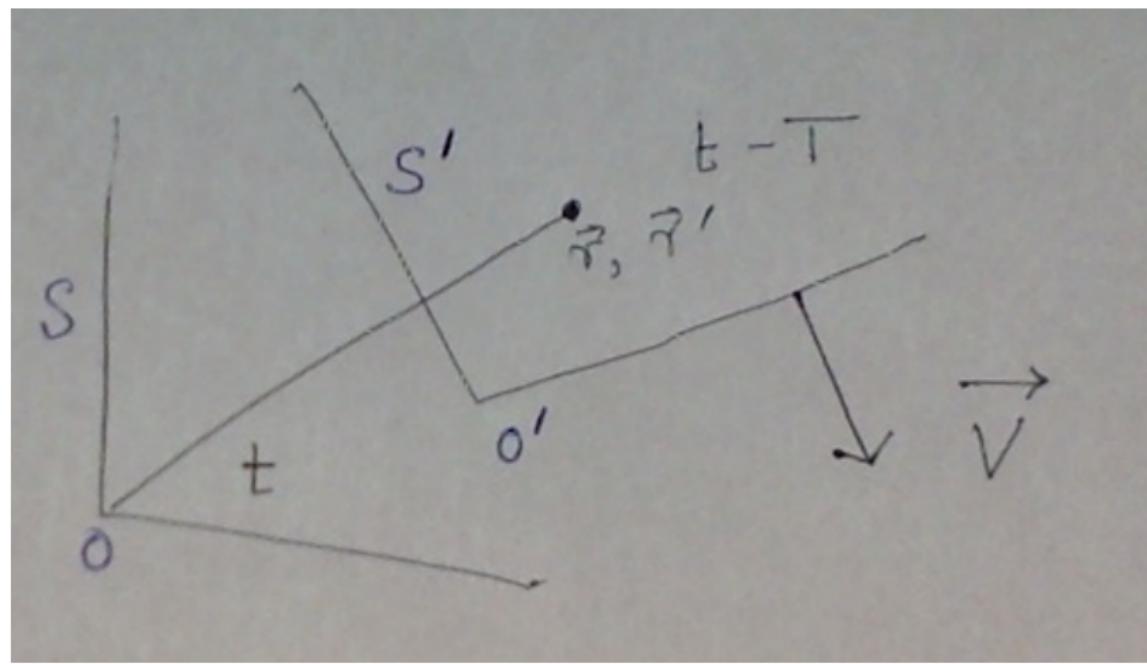
$$t' = t - T$$

$$x'_i = \mathcal{R}_{ij}x_j - R_j - V_i t$$

Note the choice of the order

Illustration

Galilean transformation



Force Laws

Interactions are mutual

- Oscillator: $\vec{F}(\vec{r}) = -k(\vec{r} - \vec{R});$
- Gravitation: $\vec{F}(\vec{r}) = -\frac{GMm}{|(\vec{r}-\vec{R})|^3}(\vec{r} - \vec{R})$

Facts and assumptions are consistent with Galilean Transformations between inertial frames

Principle of Galilean Relativity

Newton v/s Leibniz

<https://plato.stanford.edu/entries/spacetime-theories/>

Systems and Constraints

- Unconstrained systems
- Constrained systems
 - Classification
 - Implementation
 - generalised coordinates

N particles labelled $a = 1, \dots, N$

Respective masses m_a

Mutual interactions

External Forces

Description

Initial conditions and Forces

Positions : \vec{r}_a ; $a = 1, \dots, N$

Velocities : \vec{v}_a ; $a = 1, \dots, N$

Forces : $\vec{F}_{a \rightarrow b}$ Force on b due to a

No self interaction.

$$\vec{F}_T^{(a)} = \sum_{b \neq a} \vec{F}_{b \rightarrow a} + \vec{F}_{ext}^a$$

$$\vec{F}_T^{(a)} = \vec{F}_T^{(a)}(\{\vec{r}_b, \vec{v}_b\}); b = 1, \dots N$$

$$\vec{F}_{b \rightarrow a} = \vec{F}_{b \rightarrow a}(\vec{r}_{ba}; \vec{v}_{ba}); b \neq a$$

$$\vec{F}_{ext}^a = \vec{F}_{ext}^a(\vec{r}_a, \vec{v}_a)$$

A set of coupled second order ODE.

Highly Nonlinear

$3N + 3N$ initial conditions

Configuration Space (\mathcal{C})

degrees of freedom, dof = $3N = \dim(\mathcal{C})$

$$\mathcal{C} = \mathbb{R}^{3N}$$

Extend to Momentum Space and Phase space

Examples

- Particle constrained to move within a sphere
- A bead constrained to move on a rigid wire of some shape
- An insect on the surface of an expanding balloon
- Particle moving between two sticky walls
- Two bodies at the ends of a rigid rod
- Two particles stuck at the two ends of a thread of length L
- A rigid body
- A person on a unicycle (disc or coin on a rough surface)

Preliminaries on Constraints

Definitions

Definition 1: A constraint is bilateral if it can be expressed as

$$f(t, \vec{r}_a, \vec{v}_a) = 0$$

Definition 2: A constraint is unilateral if it can be expressed as

$$f(t, \vec{r}_a, \vec{v}_a) \leq 0$$

Definition 3: A bilateral constraint is geometric, or finite If f is velocity independent

Definition 4: Else, kinematic

Important Class; constraints linear in velocities:

$$\sum_1^N \vec{c}_a \cdot \vec{v}_a + D = 0$$

Let $f(\vec{r}_a, t) = 0$. Then,

$$\sum_a \nabla_{\vec{r}_a} f \cdot d\vec{r}_a + \frac{\partial f}{\partial t} dt = 0$$

But converse operation yields an arbitrary constant.
A differential constraint need not have finite form.
Integrability conditions (Discussed later)

Classification

The eight fold way

(Scleronic, Rheonomic)

×

(Holonomic, Non-holonomic)

×

(Conservative, Non-conservative)

Note: Italics in the last line

Scleronic constraints

Definition: The constraints are time independent

Rheonomic constraints

Definition: The constraints are time dependent

Holonomic constraints

Definition: The coordinates can be given arbitrary infinitesimal displacements without violating constraints

The constraints finite, or integrable

They can be written in closed form

Their general form: $f(\{\vec{r}_a\}, t) = 0$. Remember this.

Nonholonomic constraints

Definition: The constraints are not integrable. Not holonomic Not dealt with in detail; Only some examples

Examples

- two masses connected to a deformable rod
- Two masses connected to a rod expanding at a constant rate
- A disc rolling on a rough surface
- two molecules at the tips of a protein chain

$$\sum_a \vec{I}_a^\alpha \cdot \vec{v}_a = 0; \alpha = 1, \dots, k$$

Also First and Second class constraints

- $f(t, \vec{r}_a, \vec{v}_a) = 0 \implies$ redundant coordinates may be eliminated (not uniquely).
- *Generalised* coordinates. $\# \text{ GC} = \#\text{dof}$

Lecture 5

Free systems: Generalized coordinates

- Free systems ($D = 2 + 1$)
- Rectangular cartesian coordinates (x, y)
- Polar coordinates (r, θ)
- parabolic corrdinates ($x = \xi_1 \xi_2; y = \frac{1}{2}(\xi_1^2 - \xi_2^2)$)
What is parabolic here?

Note: GC need not have dimension L^1

Example: Time independent Constraints

Geometric constraints ($D = 3 + 1$)

- Particle constrained to move along a line
$$ax + by + cz + d = 0$$
- Particle constrained to move on a sphere
$$x^2 + y^2 + z^2 = R^2$$
- Three particles at the vertices of an equilateral triangle, with one vertex held fixed.
- Same as above but with the centroid held fixed.

Recall Definition:

$$f(t, \{\vec{r}_a\}_1^N) = 0$$

- A particle on an expanding surface
- Two masses at the end of a rod with length of fixed time dependence
- Three bodies at the vertices of a dilating equilateral triangle

Possible displacements

m finite constraints:

$$f_\alpha(t, \{\vec{r}_a\}) = 0; \quad \alpha = 1, \dots, m \implies$$

m differential constraints

$$\sum_a \frac{\partial f_\alpha}{\partial \vec{r}_a} \cdot \vec{dr}_a + \frac{\partial f_\alpha}{\partial t} dt = 0$$

Definition: $d\vec{r}_a$ are possible displacements. $\vec{v}_a = \frac{d\vec{r}_a}{dt}$ are possible velocities

Virtual displacements

Differences in possible displacements

Consider $\delta \vec{r}_a = d_1 \vec{r}_i - d_2 \vec{r}_a; a = 1, \dots, N$. Then, with only positions involved,

$$\sum_a \frac{\partial f_\alpha}{\partial \vec{r}_a} \cdot \vec{\delta r}_a = 0$$

Definition: $\delta \vec{r}_a$ are called virtual displacements

Question: When is Virtual = Possible?

Examples

Sphere

A particle at some latitude and longitude

- $d_1 \vec{r}$ along the latitude
- $d_2 \vec{r}$ along the longitude.

$\delta \vec{r} = d_1 \vec{r} - d_2 \vec{r}$ is virtual and possible

Counter example

Surface in motion

Sphere translating with velocity \vec{u} .

Family of possible velocities:

$$\vec{v} = \vec{v}_t - \vec{u}$$

Virtual \neq Possible.

Repeat for expanding sphere

Constraints and Forces

Constraints as forces

- System: N Particles. # coordinates: $3N$
- Constraints: $\# = m$ (linearly independent)
- # dof = $3N - m$

The Fundamental problem

Forces and Reactions

- \vec{F}_b : Force acting on the b^{th} particle
- Newton: $\vec{a}_b \neq \frac{\vec{F}_b}{m_b}$
- Question: Constraints on \vec{a}_b ???

Ans: Differentiate $f_\alpha(t, \{\vec{r}_a\}) = 0 \rightsquigarrow$

$$\sum_a \frac{\partial f_\alpha}{\partial \vec{r}_a} \cdot \vec{v}_a + \frac{\partial f_\alpha}{\partial t} = 0$$

The fundamental Problem

Forces and Reactions

Condition on Accelerations

$$\sum_a \left\{ \frac{\partial f_\alpha}{\partial \vec{r}_a} \cdot \vec{a}_a + \sum_b \frac{\partial^2 f_\alpha}{\partial \vec{r}_a \partial \vec{r}_b} \cdot \vec{v}_a \vec{v}_b + 2 \frac{\partial^2 f_\alpha}{\partial \vec{r}_i \partial t} \cdot \vec{v}_i \right\} + \frac{\partial^2 f_\alpha}{\partial t^2} = 0$$

Must be combined with external (Free) forces. Write

$$m_a \vec{a}_a = \vec{F}_a + \vec{R}_a$$

Question: How to determine the reactions?

Principle of Virtual Work

Additional solvability conditions

The Principle: Reaction forces do not do any work under virtual displacements:

$$\sum_a \vec{R}_a \cdot \delta \vec{r}_a = 0$$

$$\sum_a (m_i \vec{a}_a - \vec{F}_a) \cdot \delta \vec{r}_a = 0$$

Form of \vec{R}_a

Lagrange multipliers

Recall:

$$\sum_a \frac{\partial f_\alpha}{\partial \vec{r}_a} \cdot \vec{\delta r}_a = 0; \quad \alpha = 1, \dots k$$

Trick: Multiply each constraint by a Lagrange parameter

$$\sum_a \sum_\alpha \lambda_\alpha \frac{\partial f_\alpha}{\partial \vec{r}_a} \cdot \vec{\delta r}_a = 0$$

Now identify

$$\vec{R}_a = \sum_\alpha \lambda_\alpha \frac{\partial f_\alpha}{\partial \vec{r}_a}$$

The "full" equation

Accommodative Nature

The full expression is

$$m_b \vec{a}_b = \vec{F}_b + \sum_{\alpha} \lambda_{\alpha} \frac{\partial f_{\alpha}}{\partial \vec{r}_b}$$

Physical significance

Lagrange multipliers

- Normal reaction
- Tension

An Example

A particle on a vertical circle of Radius R in G field

Equations:

$$x^2 + y^2 - R^2 = 0$$

$$\frac{d^2x}{dt^2} = \lambda x$$

$$\frac{d^2y}{dt^2} = \lambda y - g$$

Solution worked out on board, and also in the relevant chapter that is posted.