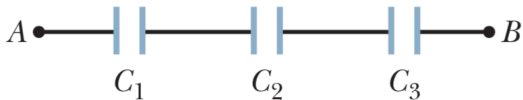


General Physics II

Solution #2

2020/10/16

P2-1. In a three-capacitor, $C_1 = 10.0 \mu\text{F}$, $C_2 = 20.0 \mu\text{F}$, and $C_3 = 25.0 \mu\text{F}$. If no capacitor can withstand a potential difference of more than 100 V without failure, what are (a) the magnitude of the maximum potential difference that can exist between points A and B and (b) the maximum energy that can be stored in the three-capacitor arrangement?

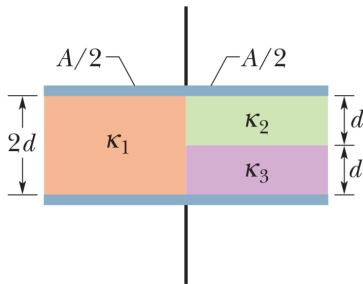


Solution:

(a) They each store the same charge, so the maximum voltage is across the smallest capacitor. With 100 V across $10\ \mu\text{F}$, then the voltage across the $20\ \mu\text{F}$ capacitor is 50 V and the voltage across the $25\ \mu\text{F}$ capacitor is 40 V. Therefore, the voltage across the arrangement is 190 V.

(b) Using $U = \frac{q^2}{2C}$, we sum the energies on the capacitors and obtain $U_{\text{total}} = 0.095\ \text{J}$.

P2-2. A parallel-plate capacitor of plate area $A = 10.5 \text{ cm}^2$ and plate separation $2d = 7.12 \text{ mm}$. The left half of the gap is filled with material of dielectric constant $\kappa_1 = 21.0$; the top of the



right half is filled with material of dielectric constant $\kappa_2 = 42.0$; the bottom of the right half is filled with material of dielectric constant $\kappa_3 = 58.0$. What is the capacitance?

Solution: Let

$$C_1 = \varepsilon_0(A/2)\kappa_1/2d = \varepsilon_0 A\kappa_1/4d,$$

$$C_2 = \varepsilon_0(A/2)\kappa_2/d = \varepsilon_0 A\kappa_2/2d,$$

$$C_3 = \varepsilon_0 A\kappa_3/2d.$$

Note that C_2 and C_3 are effectively connected in series, while C_1 is effectively connected in parallel with the C_2 - C_3 combination. Thus,

$$\begin{aligned} C &= C_1 + \frac{C_2 C_3}{C_2 + C_3} = \frac{\varepsilon_0 A\kappa_1}{4d} + \frac{(\varepsilon_0 A/d)(\kappa_2/2)(\kappa_3/2)}{\kappa_2/2 + \kappa_3/2} \\ &= \frac{\varepsilon_0 A}{4d} \left(\kappa_1 + \frac{2\kappa_2\kappa_3}{\kappa_2 + \kappa_3} \right) \end{aligned}$$

With $A = 1.05 \times 10^{-3} \text{ m}^2$, $d = 3.56 \times 10^{-3} \text{ m}$, $\kappa_1 = 21.0$, $\kappa_2 = 42.0$ and $\kappa_3 = 58.0$, we find the capacitance to be

$$\begin{aligned} C &= \frac{(8.85 \times 10^{-12})(1.05 \times 10^{-3})}{4(3.56 \times 10^{-3})} \\ &\cdot \left(21.0 + \frac{2 \times 42.0 \times 58.0}{42.0 + 58.0} \right) \\ &= 4.55 \times 10^{-11} \text{ F}. \end{aligned}$$

P2-3. The rain-soaked shoes of a person may explode if ground current from nearby lightning vaporizes the water. The sudden conversion of water to water vapor causes a dramatic expansion that can rip apart shoes. Water has density 1000 kg/m^3 and requires 2256 kJ/kg to be vaporized. If horizontal current lasts 2.00 ms and encounters water with resistivity $150 \text{ } \Omega \cdot \text{m}$, length 12.0 cm , and vertical cross-sectional area $15 \times 10^{-5} \text{ m}^2$, what average current is required to vaporize the water?

Solution: The mass of the water over the length is

$$m = \rho AL = 1000 \cdot 15 \times 10^{-5} \cdot 0.12 = 0.018 \text{ kg}$$

and the energy required to vaporize the water is

$$Q = Lm = 2256 \cdot 0.018 = 4.06 \times 10^4 \text{ J.}$$

The thermal energy is supplied by joule heating of the resistor:

$$Q = P\Delta t = I^2 R \Delta t$$

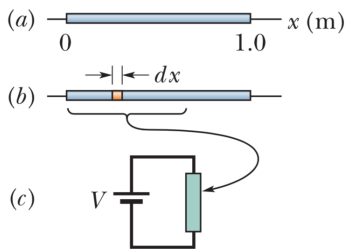
Since the resistance over the length of water is

$$R = \frac{\rho_w L}{A} = \frac{150 \cdot 0.120}{15 \times 10^{-5}} = 1.2 \times 10^5 \Omega,$$

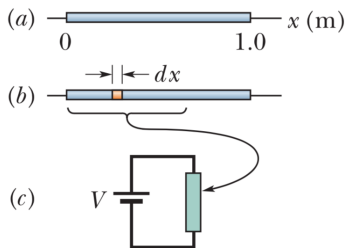
the average current required to vaporize water is

$$I = \sqrt{\frac{Q}{R\Delta t}} = \sqrt{\frac{4.06 \times 10^4}{1.2 \times 10^5 \cdot 2.0 \times 10^{-3}}} = 13.0 \text{ A}.$$

P2-4. There is a rod of resistive material (Figure a). The resistance per unit length of the rod increases in the positive direction of the x axis. At any position x along the rod, the resistance dR of a narrow (differential) section of width dx is given by $dR = 5.00x \, dx$, where dR is in ohms and x is in meters. Figure b shows such a narrow section.



You are to slice off a length of the rod between $x = 0$ and some position $x = L$ and then connect that length to a battery with potential difference $V = 5.0 \text{ V}$ (Figure c). You want the current in the length to transfer energy to thermal energy at the rate of 200 W . At what position $x = L$ should you cut the rod?



Solution: From $P = V^2/R$, we have

$$R = 5.0^2/200 = 0.125 \, \Omega$$

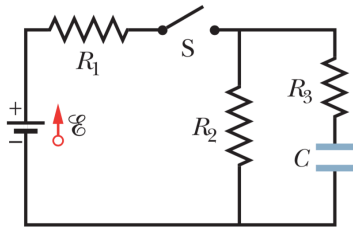
To meet the conditions of the problem statement, we must therefore set

$$\int_0^L 5.00x \, dx = 0.125 \, \Omega$$

Thus,

$$\frac{5}{2}L^2 = 0.125 \quad \Rightarrow \quad L = 0.224 \, \text{m}.$$

P2-5. In a circuit, $\mathcal{E} = 1.2 \text{ kV}$, $C = 6.5 \text{ } \mu\text{F}$, $R_1 = R_2 = R_3 = 0.73 \text{ M}\Omega$. With C completely uncharged, switch S is suddenly closed (at $t = 0$). At $t = 0$, what are (a)



current i_1 in resistor 1, (b) current i_2 in resistor 2, and (c) current i_3 in resistor 3? At $t = \infty$ (that is, after many time constants), what are (d) i_1 , (e) i_2 , and (f) i_3 ? What is the potential difference V_2 across resistor 2 at (g) $t = 0$ and (h) $t = \infty$? (i) Sketch V_2 versus t between these two extreme times.

Solution: Let i_1 be the current in R_1 and take it to be positive if it is to the right. Let i_2 be the current in R_2 and take it to be positive if it is downward. Let i_3 be the current in R_3 and take it to be positive if it is downward. The junction rule produces $i_1 = i_2 + i_3$, the loop rule applied to the left-hand loop produces

$$\varepsilon - i_1 R_1 - i_2 R_2 = 0,$$

and the loop rule applied to the right-hand loop produces

$$i_2 R_2 - i_3 R_3 = 0.$$

Since the resistances are all the same we can simplify the mathematics by replacing R_1 , R_2 , and R_3 with R .

(a) Solving the three simultaneous equations, we find

$$i_1 = \frac{2\varepsilon}{3R} = \frac{2 \cdot 1.2 \times 10^3}{3 \cdot 0.73 \times 10^6} = 1.1 \times 10^{-3} \text{ A}$$

(b) $i_2 = \frac{\varepsilon}{3R} = \frac{1.2 \times 10^3}{3 \cdot 0.73 \times 10^6} = 5.5 \times 10^{-4} \text{ A},$

(c) and $i_3 = i_2 = 5.5 \times 10^{-4} \text{ A}.$

At $t = \infty$ the capacitor is fully charged and the current in the capacitor branch is 0. Thus, $i_1 = i_2$, and the loop rule yields $\varepsilon - i_1 R_1 - i_1 R_2 = 0.$

(d) The solution is $i_1 = \frac{\varepsilon}{2R} = \frac{1.2 \times 10^3}{2 \cdot 0.73 \times 10^6} = 8.2 \times 10^{-4} \text{ A}$

(e) and $i_2 = i_1 = 8.2 \times 10^{-4}$ A.

(f) As stated before, the current in the capacitor branch is $i_3 = 0$.

We take the upper plate of the capacitor to be positive. This is consistent with current flowing into that plate. The junction equation is $i_1 = i_2 + i_3$, and the loop equations are

$$\varepsilon - i_1 R - i_2 R = 0$$

$$-\frac{q}{C} - i_3 R + i_2 R = 0.$$

We use the first equation to substitute for i_1 in the second and obtain

$$\varepsilon - 2i_2 R - i_3 R = 0$$

Thus $i_2 = (\varepsilon - i_3 R)/2R$.

We substitute this expression into the third equation above to obtain

$$-\frac{q}{C} - i_3 R + \frac{\varepsilon}{2} - \frac{i_3 R}{2} = 0.$$

Now we replace i_3 with dq/dt to obtain

$$\frac{3R}{2} \frac{dq}{dt} + \frac{q}{C} = \frac{\varepsilon}{2}.$$

This is just like the equation for an RC series circuit, except that the time constant is $\tau = 3RC/2$ and the impressed potential difference is $\varepsilon/2$. The solution is

$$q = \frac{C\varepsilon}{2}(1 - e^{-2t/3RC}).$$

The current in the capacitor branch is

$$i_3(t) = \frac{dq}{dt} = \frac{\varepsilon}{3R} e^{-2t/3RC}.$$

The current in the center branch is

$$i_2(t) = \frac{\varepsilon}{2R} - \frac{i_3}{2} = \frac{\varepsilon}{2R} - \frac{\varepsilon}{6R} e^{-2t/3RC} = \frac{\varepsilon}{6R} (3 - e^{-2t/3RC})$$

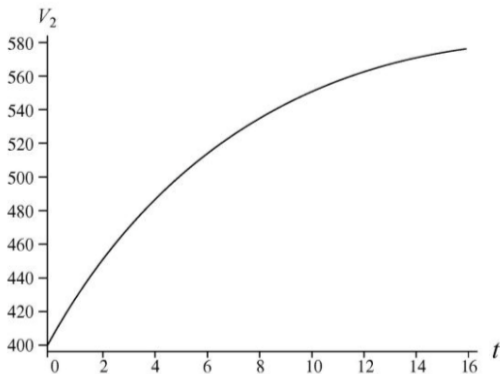
and the potential difference across R_2 is

$$V_2(t) = i_2 R = \frac{\varepsilon}{6} (3 - e^{-2t/3RC}).$$

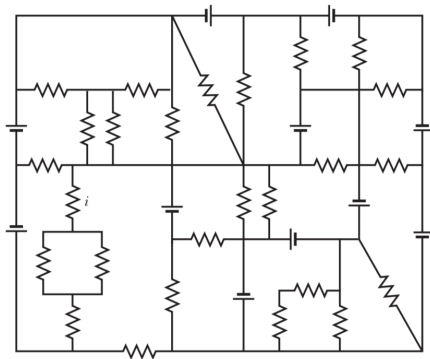
(g) For $t = 0$, $e^{-2t/3RC} = 1$ and
 $V_2 = \varepsilon/3 = (1.2 \times 10^3)/3 = 4.0 \times 10^2 \text{ V}.$

(h) For $t = \infty$, $e^{-2t/3RC} \rightarrow 0$ and
 $V_2 = \varepsilon/2 = (1.2 \times 20^3)/2 = 6.0 \times 10^2 \text{ V}.$

- (i) A plot of V_2 as a function of time is shown in the following graph.



P2-6. What are the (a) size and (b) direction (up or down) of current i , where all resistances are $4.0\ \Omega$ and all batteries are ideal and have an emf of $10\ \text{V}$? (*Hint:* This can be answered using only mental calculation.)



Solution: The resistor by the letter i is above three other resistors; together, these four resistors are equivalent to a resistor $R = 10\ \Omega$ (with current i). As if we were presented with a maze, we find a path through R that passes through any number of batteries (10, it turns out) but no other resistors, which — as in any good maze — winds “all over the place.” Some of the ten batteries are opposing each other (particularly the ones along the outside), so that their net emf is only $\varepsilon = 40\text{ V}$.

(a) The current through R is then $i = \varepsilon/R = 4.0\text{ A}$.

(b) The direction is upward in the figure.

P2-7. A metal sphere of radius 15 cm has a net charge of 3.0×10^{-8} C.

- (a) What is the electric field at the sphere's surface?
- (b) If $V = 0$ at infinity, what is the electric potential at the sphere's surface?
- (c) At what distance from the sphere's surface has the electric potential decreased by 500 V?

Solution:

(a) The magnitude of the electric field is

$$\begin{aligned} E &= \frac{\sigma}{\epsilon_0} = \frac{q}{4\pi\epsilon_0 R^2} = \frac{(3.0 \times 10^{-8})(8.99 \times 10^9)}{(0.15)^2} \\ &= 1.2 \times 10^4 \text{ N/C} \end{aligned}$$

(b) $V = RE = (0.15)(1.2 \times 10^4) = 1.8 \times 10^3 \text{ V}$

(c) Let the distance be x . Then

$$\Delta V = V(x + R) - V = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{R + x} - \frac{1}{R} \right) = -500 \text{ V}$$

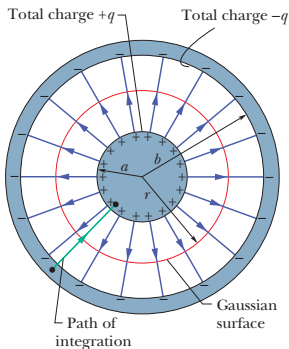
which gives

$$x = \frac{R\Delta V}{-V - \Delta V} = \frac{(0.15)(-500)}{-1800 + 500} = 5.8 \times 10^{-2} \text{ m}$$

P2-8. Consider two concentric spherical shells, of radii a and b . Show that the capacitance of the shells is

$$C = 4\pi\epsilon_0 \frac{ab}{b - a}.$$

What is the capacitance to a single isolated spherical conductor of radius R , then?



Solution: As a Gaussian surface we draw a sphere of radius r concentric with the two shells. We have

$$q = \epsilon_0 EA = \epsilon_0 E(4\pi r^2).$$

Solving for E , we obtain

$$E = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2}.$$

Hence,

$$V = - \int_{-}^{+} \vec{E} \cdot d\vec{s} = - \frac{q}{4\pi\epsilon_0} \int_b^a \frac{dr}{r^2} = \frac{q}{4\pi\epsilon_0} \frac{b-a}{ab}.$$

Solving for C , we find

$$C = 4\pi\epsilon_0 \frac{ab}{b-a}.$$

We can assign a capacitance to a single isolated spherical conductor of radius R by assuming that the “missing plate” is a conducting sphere of infinite radius. To find the capacitance of the conductor, we take the limit $b \rightarrow \infty$ and replace a by R . Thus we find

$$C = 4\pi\epsilon_0 R.$$

P2-9. Show that the curl of a central force $\vec{F}(\vec{r}) = f(r)\hat{r}$ is zero, i.e.,

$$\nabla \times \vec{F}(\vec{r}) = 0.$$

Hence, central forces are conservative.

Solution: There are a number of ways of showing that.

(1) If you only know what we taught in class, you can write

$$\vec{F}(\vec{r}) = g(r)\vec{r} = g(r)(x\hat{i} + y\hat{j} + z\hat{k}),$$

where $g(r) = f(r)/r$. The x component of $\nabla \times \vec{F}$ is

$$\begin{aligned}(\nabla \times \vec{F})_x &= \frac{\partial F_z}{\partial y} - \frac{\partial F_y}{\partial z} \\&= \frac{\partial g(r)}{\partial r} \frac{y}{r} z - \frac{\partial g(r)}{\partial r} \frac{z}{r} y \\&= 0.\end{aligned}$$

So, the curl of a central force is always zero.

(2) One may as well show this by

$$\nabla \times \vec{F} = \nabla \times (g\vec{r}) = (\nabla g) \times \vec{r} + g\nabla \times \vec{r}.$$

Note that ∇g is radial, hence $(\nabla g) \times \vec{r} = 0$. Obviously, $\nabla \times \vec{r} = 0$.

(3) You may prove this straightforwardly in spherical coordinate system (you can find the formula in an advanced book on E&M). In fact, we can write down the potential energy explicitly as

$$\vec{F} = -\nabla V,$$

where

$$V = - \int f(r) dr.$$

Curl is a well named mathematical term—it denotes the degree of rotation in the vector field (circulation). For this reason, if you go all the way around in a vector field, you'll find that the total integral along that path will depend on the curl of the field in question. If a force had a curl, you could go all the way around and have some net work done, and so it would be nonconservative. A conservative force, on the other hand, cancels itself back out as you go on a closed loop.

The central force has a spherical symmetry and is pointing radially inward/outward. Naturally, its curl is zero.

P2-10. Consider a two-dimensional electric field

$$\vec{E}(x, y) = \frac{-y\hat{i} + x\hat{j}}{x^2 + y^2}.$$

- (a) Calculate the curl of the field $\nabla \times \vec{E}$.
- (b) Show that the circulation of the field

$$\Gamma = \oint_C \vec{E} \cdot d\vec{s} = 2\pi$$

around a unit circle centered at origin.

Therefore, a vanishing curl does not implies, in general, that the force is conservative. They are equivalent only when the space is simply connected.

Solution:

(a) Explicitly, we have

$$E_x = -\frac{y}{x^2 + y^2},$$

$$E_y = \frac{x}{x^2 + y^2}.$$

Taking partial derivatives, we have

$$\frac{\partial E_x}{\partial y} = -\frac{1}{x^2 + y^2} + y \frac{2y}{x^2 + y^2},$$

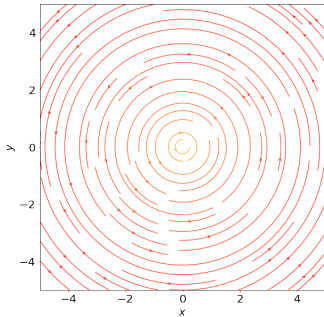
$$\frac{\partial E_y}{\partial x} = \frac{1}{x^2 + y^2} - x \frac{2x}{x^2 + y^2}.$$

Therefore,

$$\nabla \times \vec{E} = \frac{\partial E_x}{\partial y} - \frac{\partial E_y}{\partial x} = 0.$$

- (b) The field goes in a circumferential direction with no radial component.

$$\begin{aligned}\Gamma &= \oint_C \vec{E} \cdot d\vec{s} \\ &= \int_0^{2\pi} \left(\frac{\hat{\theta}}{r} \right) \cdot (rd\theta \hat{\theta}) \\ &= \int_0^{2\pi} d\theta \\ &= 2\pi\end{aligned}$$



Such a configuration is known as a free vortex. Notice the mathematical singularity at the origin. The space is the punctured plane $(x, y) \neq (0, 0)$. It is not simply connected.