



# Probing Dark Matter-Baryon Interactions using JWST Early Galaxies

### Souradeep Das

5th year BS-MS, Indian Institute of Science

Ongoing work with Ranjini Mondol, Abhijeet Singh, and Ranjan Laha

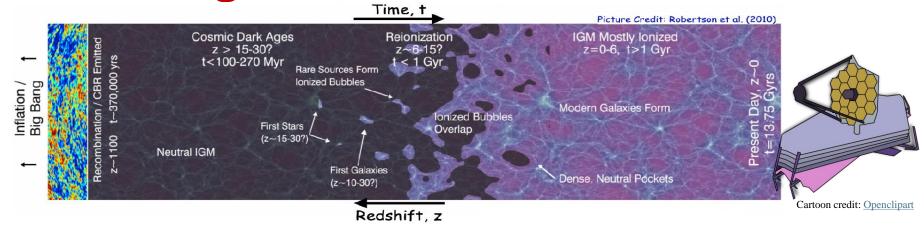
Second Neighbourhood Cosmology Meeting, IISc Bangalore

Date: December 19, 2024

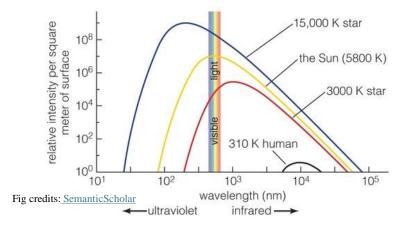
### **Contents**

- 1. What we observe using JWST (James Webb Space Telescope)?
- 2. How can we use this observation to probe DM-baryon interaction?

**Observing the Earliest Galaxies** 



Earlier galaxies are star-forming and host younger stars. Younger stars are hotter and emit mostly in UV. (A fiducial wavelength for this purpose is 150 nm)



We observe these galaxies in near-IR.

Hubble can observe  $\lambda \lesssim 1.6 \ \mu m$ JWST can observe  $\lambda \lesssim 5 \ \mu m$  in NIR

JWST can look at higher redshift galaxies!

Farthest galaxy observed by HST: GN-z11 ( $z \simeq 11.09$ )

JWST has already observed galaxies at  $z \approx 12$  and higher!

### **UV Luminosity Function**

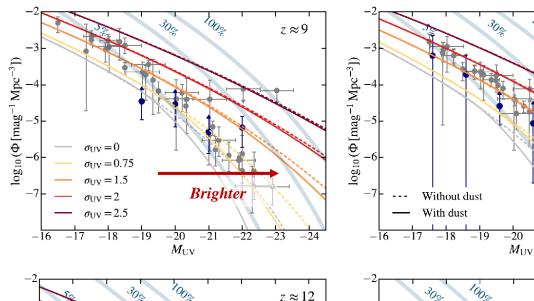
 $\Phi(M_{\mathrm{UV}})$ 

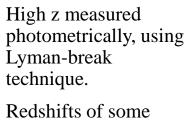
Number of galaxies per unit volume of space, with UV magnitude between  $M_{UV}$  and  $M_{UV} + dM_{UV}$ :

 $\Phi(M_{\rm UV}).{\rm d}M_{\rm UV}$ 

UV magnitude is defined in terms of the luminosity at wavelength 150 nm (in the rest frame of the galaxy)

## JWST has observed many bright galaxies at $z \sim 9 - 16$

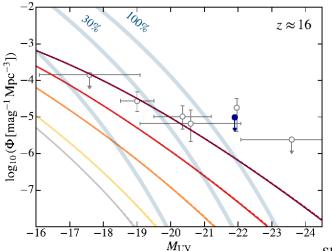




 $z \approx 10$ 

-22

Redshifts of some galaxies were later confirmed using spectroscopy.



These high-redshift galaxies are apparently much more numerous than expected from pre-JWST measurements at high-redshift.

Slide idea: Priyank Parashari, Ranjan Laha

M<sub>UV</sub> Shen et al (2023), arXiv:2305.05679

-20

JWST photo z + pre JWST JWST spec z

 $\log_{10}\left(\Phi\left[\mathrm{mag}^{-1}\,\mathrm{Mpc}^{-3}\right]\right)$ 

### **Pre-JWST expectations**

#### Low-redshift observations

Hubble, Spitzer observe galaxies at redshifts  $\lesssim 10$ 

#### **Hierarchical structure formation**

Halos form *bottom-up* : smaller halos form first, larger halos form later.

At high z, we expect less halos to have collapsed/merged

#### Reionization

First stars and galaxies produce UV radiation and ionize the neutral hydrogen. Ly $\alpha$  forest tells us about reionization history

#### **Cosmological simulations**

Predicts first stars:  $z \sim 15 - 30$ First galaxies:  $z \sim 15 - 30$ 

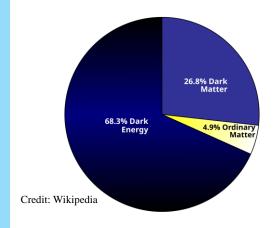
When we observe earlier galaxies with JWST, we find many more bright galaxies than we expected (extrapolated from pre-JWST models).

## Proposed solutions: Change the Astrophysical Models

Top-heavy initial stellar mass function	Yung et al., MNRAS 2024
Low dust attenuation in high-redshift galaxies	Ferrara et al., MNRAS 2023
Feedback-free starburst	Dekel et al., MNRAS 2023
Bursty star-formation and high variability	Shen et al., arXiv:2305.05679
Non-stellar sources of UV radiation	Inayoshi et al., ApJL 2022

What if, we instead changed the cosmological model?

## **Changing the Cosmological Model**



Energy budget of the ACDM model

The standard  $\Lambda$ CDM model explains almost all observations on large scales. But it is not the perfect model –

- Tensions:  $H_0$  tension,  $\sigma_8$  tension
- Not quite satisfactory at smaller scales (core-cusp problem, missing satellites problem etc.)
- The two components (dark matter and dark energy) that make up ~95% of the Universe have largely unknown properties.

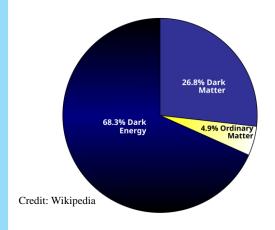
## Alterations to the ΛCDM model (some examples)

- Early Dark Energy (EDE)
- Warm Dark Matter (WDM)
- Self-Interacting Dark Matter (SIDM)
- Dark Matter Standard Model interactions
- Modifying the Primordial Power Spectrum

Souradeep Das

8

## **Changing the Cosmological Model**



Energy budget of the ACDM model

The standard  $\Lambda$ CDM model explains almost all observations on large scales. But it is not the perfect model –

- Tensions:  $H_0$  tension,  $\sigma_8$  tension
- Not quite satisfactory at smaller scales (core-cusp problem, missing satellites problem etc.)
- The two components (dark matter and dark energy) that make up ~95% of the Universe have largely unknown properties.

## Alterations to the ΛCDM model (some examples)

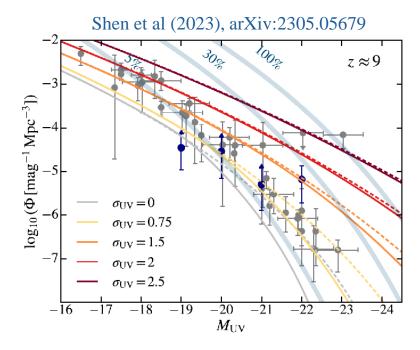
- Early Dark Energy (EDE)
- Warm Dark Matter (WDM)
- Self-Interacting Dark Matter (SIDM)
- Dark Matter Standard Model interactions
- Modifying the Primordial Power Spectrum

Some of these models predict more galaxies at high redshifts (compared to ΛCDM)!

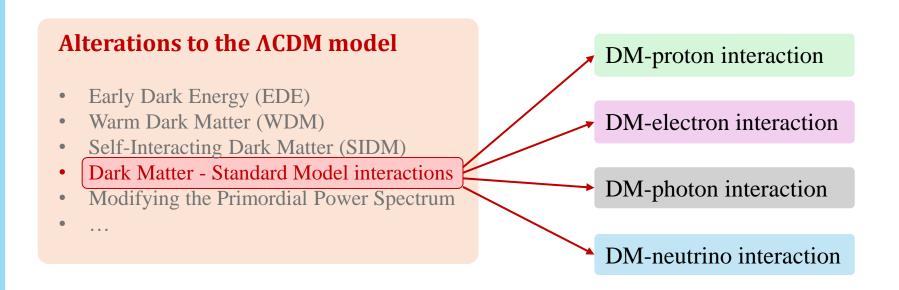
### **Do Early Galaxies contradict ΛCDM?**

### Likely not.

- Astrophysical uncertainties are large for early galaxies. Variability could be very important. (Shen et al. (2023), arXiv:2305.05679)
- Changing cosmology/matter power spectrum also changes other data at similar redshifts, e.g., HST data. (Parashari and Laha, *MNRAS* 526 (2023) 1, L63; Sabti, et al., *PRL* 132 (2024) 6, 061002)



# Can we obtain new constraints on Dark Matter using Early Galaxies?



Dark Matter-Standard Model (DM-SM) interactions generically predict lesser number of early galaxies

### Our Model: Dark Matter - Baryon scattering

DM scatters elastically with protons/electrons. (For simplicity, consider only one interaction, say with electrons)

Momentum transfer: Effect of drag – DM can be pulled out of potential wells.

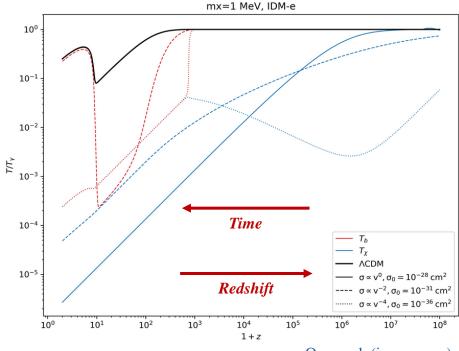
Heat transfer: DM can become warmer

Interaction cross-section:  $\sigma = \sigma_n v^n$ 

- n = 0: Heavy-mediator/contact interaction
- n = -1: Yukawa potential
- n = -2: Dipole interaction
- n = -4: Light-mediator/long-range/Coulomb interaction

For n > -3, DM decouples from baryons at some high redshift  $(z \gtrsim 10^4)$ .

For lower n, DM can recouple with baryons at low redshift (late Universe)



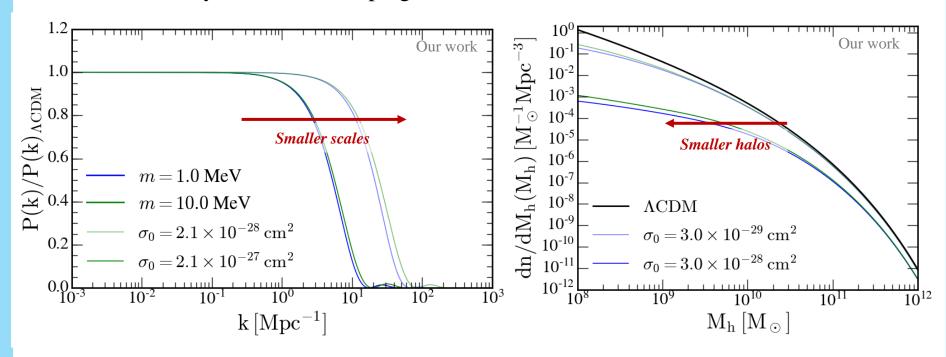
Our work (in progress)

# **Example: Effect of DM-electron interaction on Cosmology**

Assumption: Velocity-independent momentum-transfer cross-section

Matter power spectrum is suppressed at small scales by collisional damping

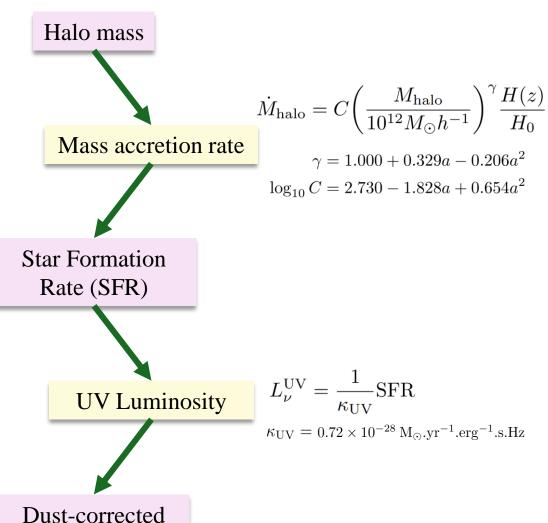
Hence, less number of small-mass halos



\*Here ACDM indicates no DM-electron interaction

## From Halo Mass to UV Luminosity

**UV** luminosity



$$\mathrm{SFR} = \epsilon_{\star} f_{\mathrm{b}} \dot{\mathrm{M}}_{\mathrm{halo}}$$

$$\epsilon_{\star}(M_{\mathrm{halo}}) = \frac{2\epsilon_{0}}{\left(M_{\mathrm{halo}}/M_{0}\right)^{-\alpha} + \left(M_{\mathrm{halo}}/M_{0}\right)^{\beta}}$$

(Star Formation Efficiency)

$$M_{\rm UV}^{\rm obs} = M_{\rm UV} + A_{\rm UV}$$
  
 $A_{\rm UV} = -0.34(21 + M_{\rm UV}^{\rm obs}) + 0.79$ 

Souradeep Das

14

## **UV Variability - major factors**

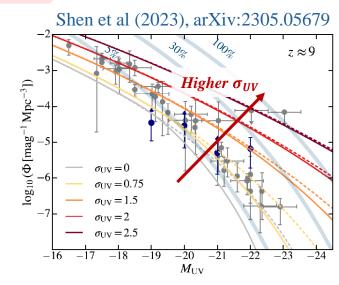
- Accretion rate of halo is not uniform
- Bursty star-formation: Star formation in short-lived 'bursts'. Irregular, clumpy morphology (observed in high-z galaxies)
- Also, variable amount of dust

Combined effect: Log-normal spread in UV luminosity for a given halo mass

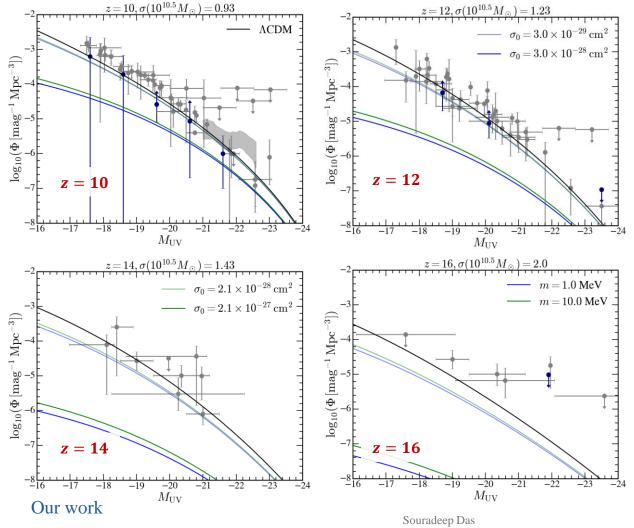
Spread in 
$$M_{\rm UV} \equiv \sigma_{\rm UV} \approx 0.5 - 2.0$$

### Variability leads to higher UV luminosity function!

This is simply because more low-mass galaxies are upscattered than high-mass galaxies down-scattered (HMF is monotonically decreasing function)



## DM-SM interaction predicts less number of **Early Galaxies!**



We assume a benchmark astrophysical model at each redshift, and change the interaction crosssection.

Benchmark model chosen so that ACDM (with no interactions) fits the JWST data best.

Larger interaction crosssections lead to more suppression of UVLF compared to  $\Lambda$ CDM.

Suppression stronger at higher redshifts.

16

## Summary

- 1. JWST has observed more numerous early, bright galaxies than we expected.
- 2. Early galaxies may be accommodated by changing astrophysical models.
- 3. DM-SM interactions can suppress the power spectrum and halo mass function at the relevant scales.
- 4. We can get observable effects on the UVLF from DM-SM interactions, which allows us to get limits of interaction strengths.
- 5. UVLF at high z can be a new probe of DM properties.
- 6. Future JWST observations should improve such constraints.

Contact: souradeepdas@iisc.ac.in

Souradeep Das

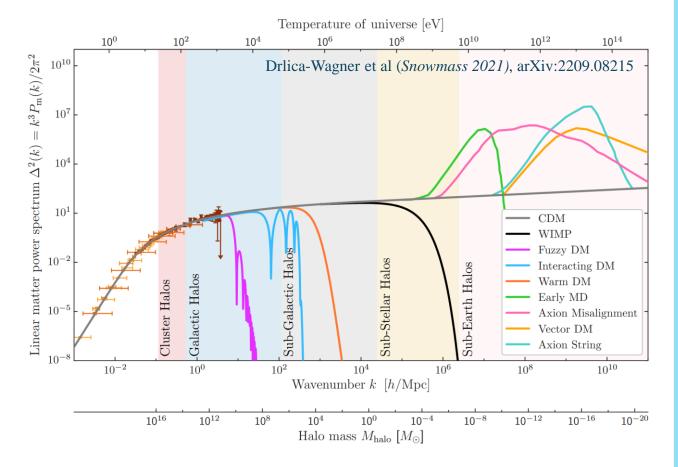
17

### Power spectrum suppression

Typical processes involved in matter power spectrum suppression:

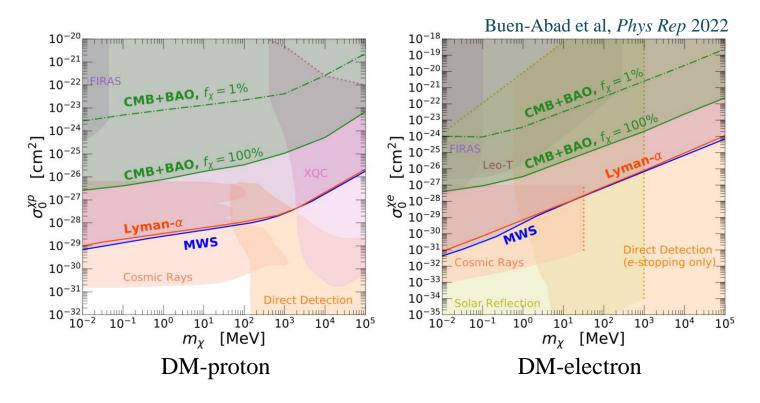
- 1. Free streaming:
  Dark matter such as
  WDM escapes from
  halos
- 2. Collisional damping
  Dark matter-baryon
  interaction prevents
  collapse

Both processes are stronger at smaller scales



### Present limits on DM-baryon interaction

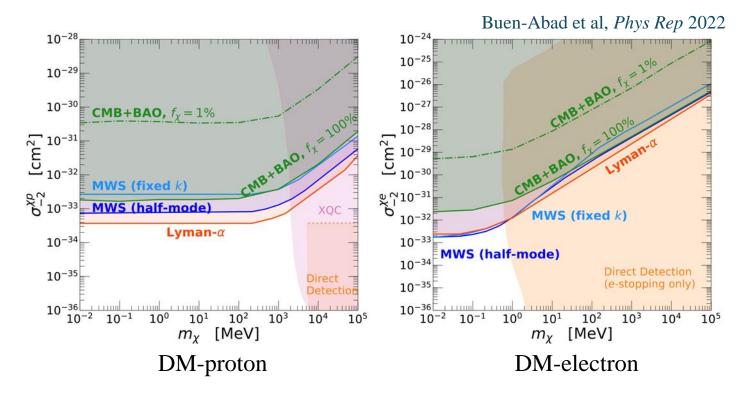
$$n = 0$$



Our limits (in progress) will be close to Lyman alpha limits

### Present limits on DM-baryon interaction

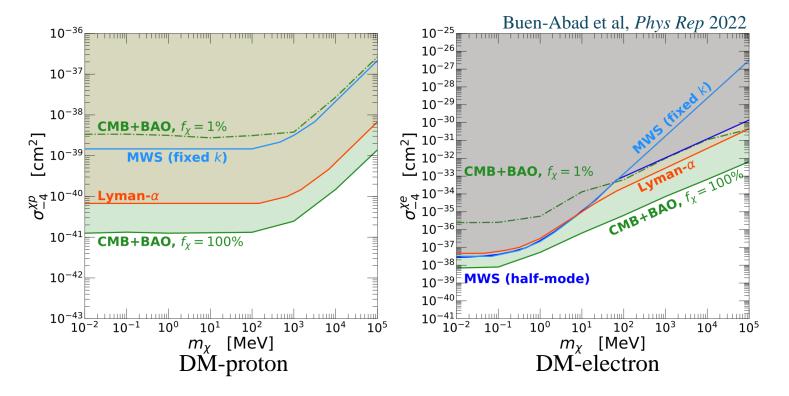
$$n = -2$$



Our limits (in progress) will be close to Lyman alpha limits

### Present limits on DM-baryon interaction

$$n = -4$$



Our limits (in progress) will be close to Lyman alpha limits

### Formalism: DM-electron Interaction

$$\begin{array}{ll} \text{Continuity} & \begin{cases} \dot{\delta_{\chi}} = -\theta_{\chi} + 3\dot{\phi} \\ \dot{\delta_{b}} = -\theta_{b} + 3\dot{\phi} \end{cases} \\ \text{Euler} & \begin{cases} \dot{\theta_{\chi}} = -aH\theta_{\chi} + k^{2}c_{\chi}^{2}\delta_{\chi} + R_{\chi-b}(\theta_{b} - \theta_{\chi}) + k^{2}\psi \\ \dot{\theta_{b}} = -aH\theta_{b} + k^{2}c_{b}^{2}\delta_{b} + R_{b-\chi}(\theta_{\chi} - \theta_{b}) \\ & + R_{b-\gamma}(\theta_{\gamma} - \theta_{b}) + k^{2}\psi \end{cases} \\ \text{Temperature} & \begin{cases} \dot{T_{\chi}} = -2aHT_{\chi} + 2R'_{\chi-b}(T_{\gamma} - T_{\chi}) \\ \dot{T_{b}} = -2aHT_{b} + 2R'_{b-\chi}(T_{\chi} - T_{b}) + 2R'_{b-\gamma}(T_{\gamma} - T_{b}) \end{cases}$$

Solved using class\_dmb

23

#### Interaction rate

$$R_{\chi-b} = a \sum_{B} \frac{Y_B \rho_b}{m_{\chi} + m_B} \sigma_n^{\chi B} c_n u_B^{n+1}$$

#### Sound speed

$$c_i^2 = \frac{k_B T_i}{\mu} \left( 1 - \frac{1}{3} \frac{\mathrm{d} \ln T_i}{\mathrm{d} \ln a} \right)$$

### **Halo Mass Function**

Number of DM halos per unit mass interval per unit volume

$$\frac{\mathrm{d}n}{\mathrm{d}\ln M} = M \frac{\rho_m}{M^2} f(\sigma(M)) \left| \frac{\mathrm{d}\ln \sigma}{\mathrm{d}\ln M} \right|$$

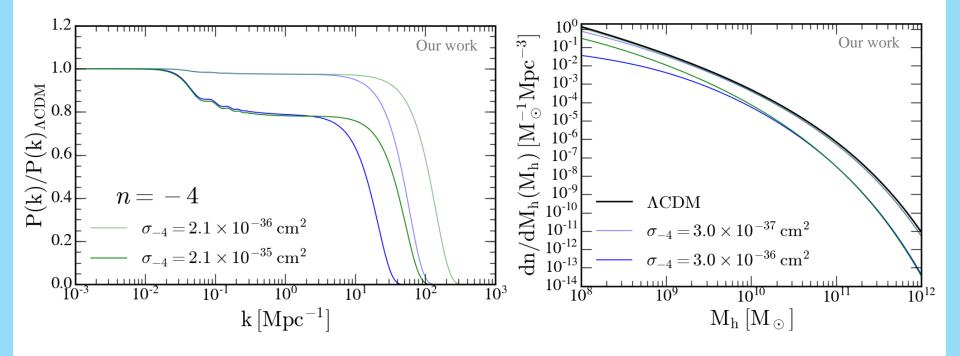
Sheth-Tormen fitting function: 
$$f(\sigma) = A\sqrt{\frac{2a}{\pi}} \left[ 1 + \left( \frac{\sigma^2}{a\delta_c^2} \right)^p \right] \frac{\delta_c}{\sigma} \exp\left( -\frac{a\delta_c^2}{2\sigma^2} \right)$$

Variance of fluctuations: 
$$\sigma^2(R) = \frac{1}{2\pi^2} \int_0^\infty \mathrm{d}k \; k^2 W^2(kR) P(k)$$

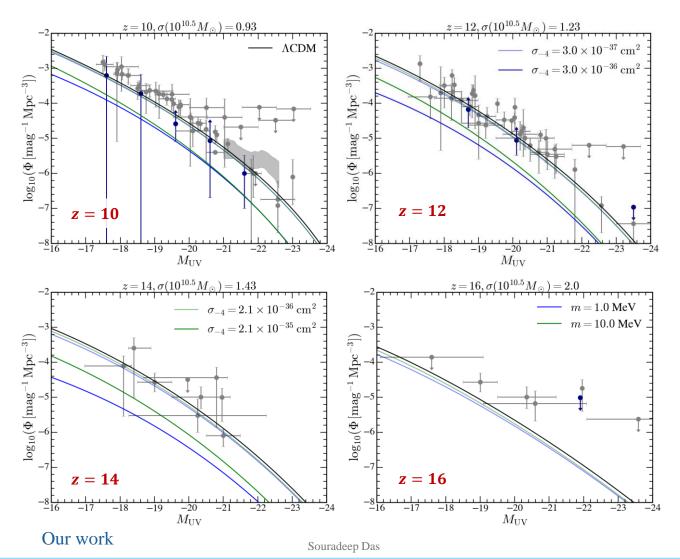
Top-hat filter function: 
$$W(kR) = \frac{3}{(kR)^3} (\sin(kR) - kR\cos(kR))$$

# Effect of DM-electron interaction on Cosmology: *Coulomb-like* Interactions

$$\sigma \propto v_{\rm rel}^{-4}$$



# **Coulomb-like** interactions and Early Galaxies



26