

$$D_s^{*+} \longrightarrow D_s^+ e^+ e^-$$

Souvik Das

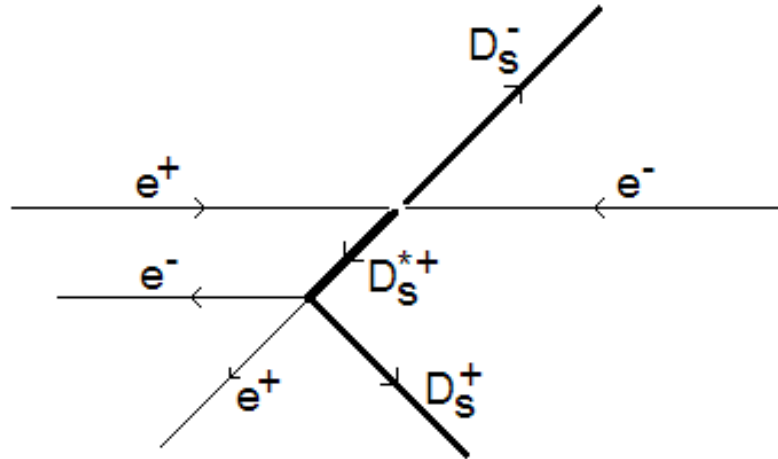
Cornell University
For the CLEO Collaboration

Contents

- What Are We Looking For?
- Predicted $D_s^{*+} \rightarrow D_s^+ e^+ e^-$ Rate
- Backgrounds
- Does Vertex Constraining Help?
- Predictions for the Background
- Low Energy Electron Tracking Efficiency

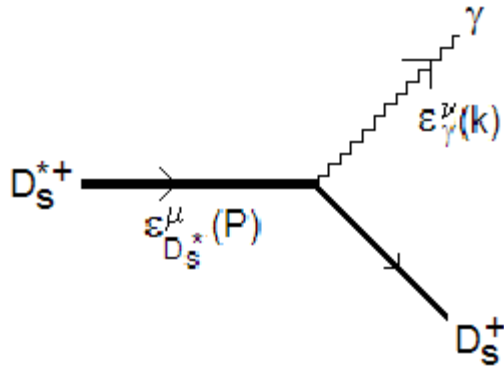
16 July 2010

What Are We Looking For?



- Searching for $D_s^{*+} \rightarrow D_s^+ e^+ e^-$ with a **blind analysis**.
- Known decay channels are:
 - $D_s^{*+} \rightarrow D_s^+ \gamma$; Branching Fraction = 94.2%
 - $D_s^{*+} \rightarrow D_s^+ \pi^0$; Branching Fraction = 5.8% [[1](#)]
- We are using e^+e^- collision data collected by the CLEO-c detector at the Cornell Electron Storage Ring (CESR) operating at $\sqrt{s} = 4170$ MeV. We have [586 ± 6 pb⁻¹](#) of data at this energy.
- $D_s^{*+} D_s^-$ Production cross section at this energy is 948 ± 36 pb (combining results from [[2](#)] and [[3](#)]). This will give us $\sim 600,000$ events to work with.

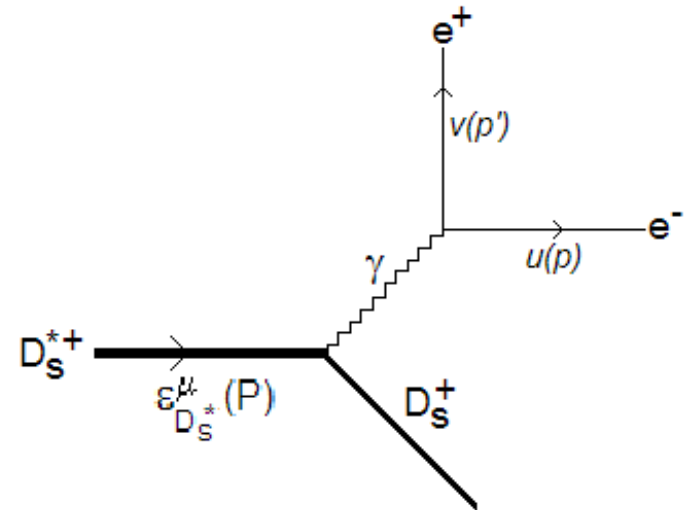
Predicted $D_s^{*\pm} \rightarrow D_s^\pm e^+ e^-$ Rate



If we write the matrix element of the D_s^* decay to a real photon in the form:

$$M = \varepsilon_{D_s^*}^\mu \varepsilon_\gamma^{*\nu} T_{\mu\nu}(P, k)$$

Where $T_{\mu\nu}(P, k)$ is a generic form factor coupling the D_s^* with a photon.



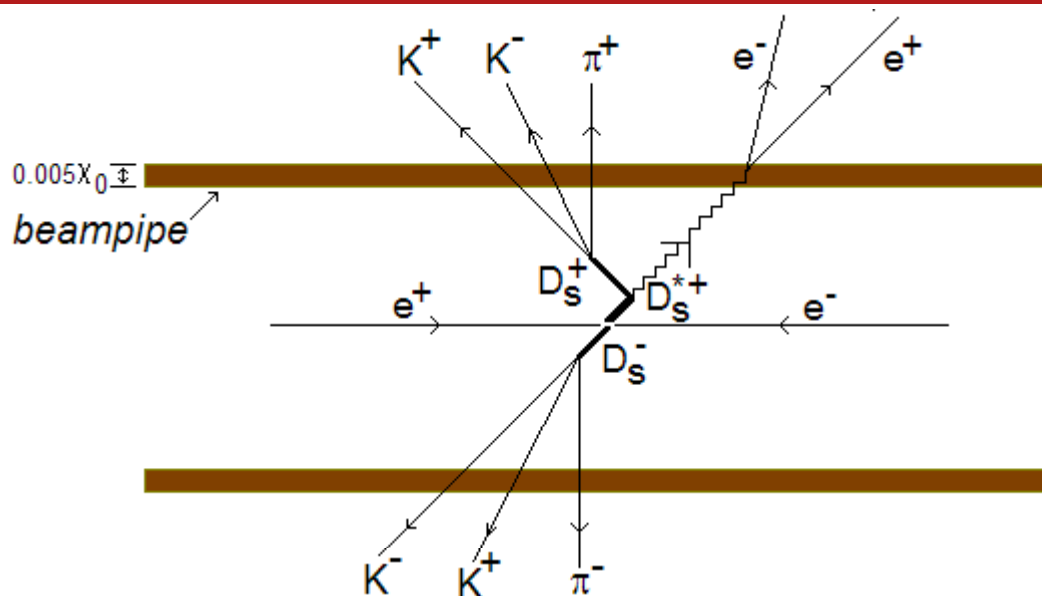
Then we can write the matrix element of the decay to e^+e^- in the form:

$$M = \varepsilon_{D_s^*}^\mu T_{\mu\nu}(P, k) \left(\frac{-ig^{\nu\sigma}}{k^2} \right) \bar{u}(p) i e \gamma_\sigma v(p')$$

Evaluating the spin-average over the initial states and spin-sum over the final states of the invariant amplitudes and integrating over the phase space of daughters, we predict the ratio of decay rates:

$$\frac{\Gamma(D_s^{*+} \rightarrow D_s^+ e^+ e^-)}{\Gamma(D_s^{*+} \rightarrow D_s^+ \gamma)} = 0.65\% = 0.90\alpha$$

Backgrounds



Photon Conversion Background

- A background that resembles the signal is expected from D_s^{*+} decaying to $D_s^+ \gamma$ and the γ converting to e^+e^- in the beam-pipe and other material.
- Given that the beam-pipe is $\sim 0.5\%$ of a radiation length, we can estimate this conversion background to occur at roughly the same rate as the signal

Combinatorial Backgrounds

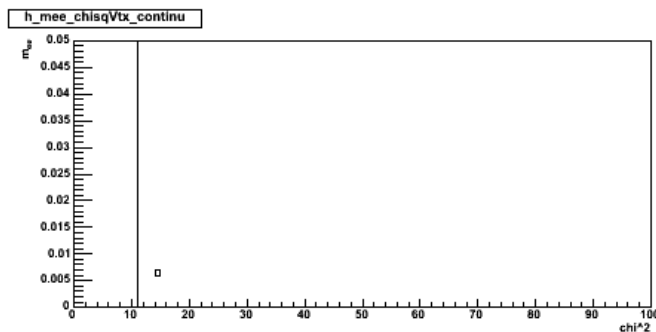
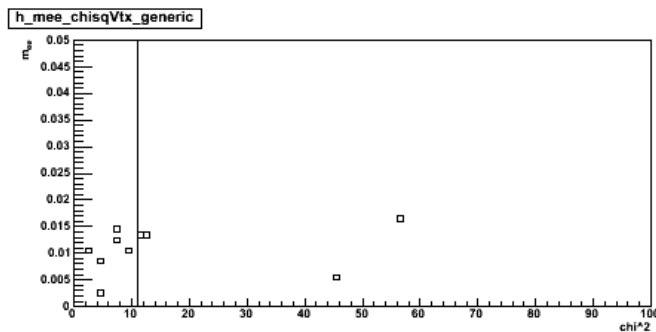
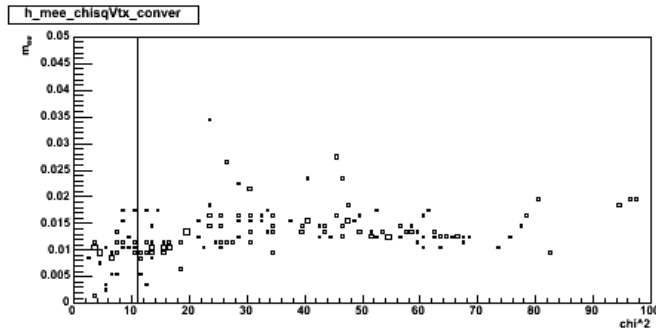
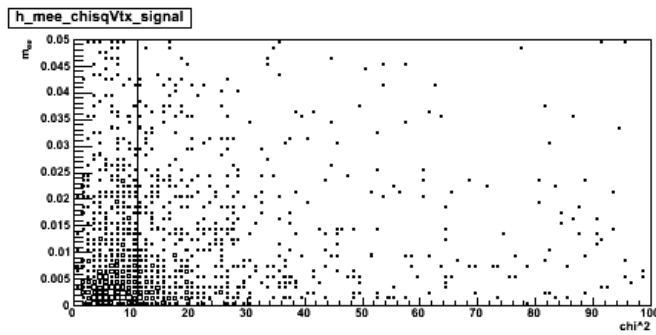
- Dalitz decay of any $\pi^0 \rightarrow \gamma e^+ e^-$ also give equally soft electrons that appear to come from interaction point
- Fake D_s tags

Can Vertex-Constraints & m_{ee} Cut Replace Δd_0 and $\Delta\varphi_0$ Cuts?

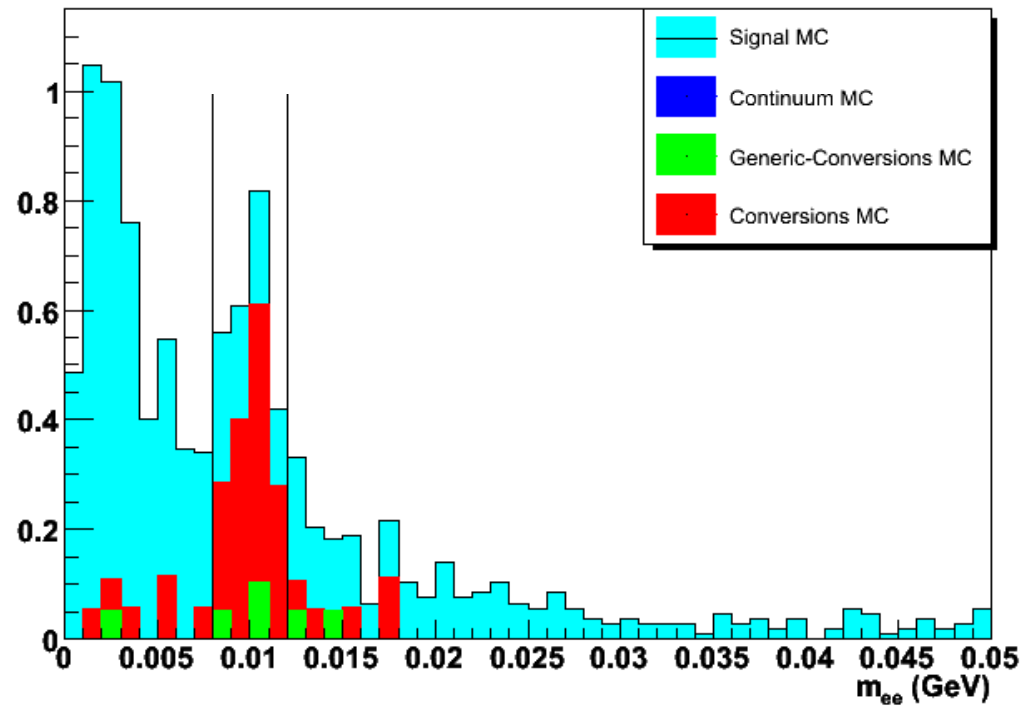
1. Vertex constrain all tracks from the D_s^* to a point,
2. Vertex constraint the e^+e^- to a point,
3. Vertex constrain all tracks from the D_s^* and the beamspot, and vertex constrain all tracks from the D_s and the beamspot.

Vertex-Constraints all tracks of D_s^* to a point.

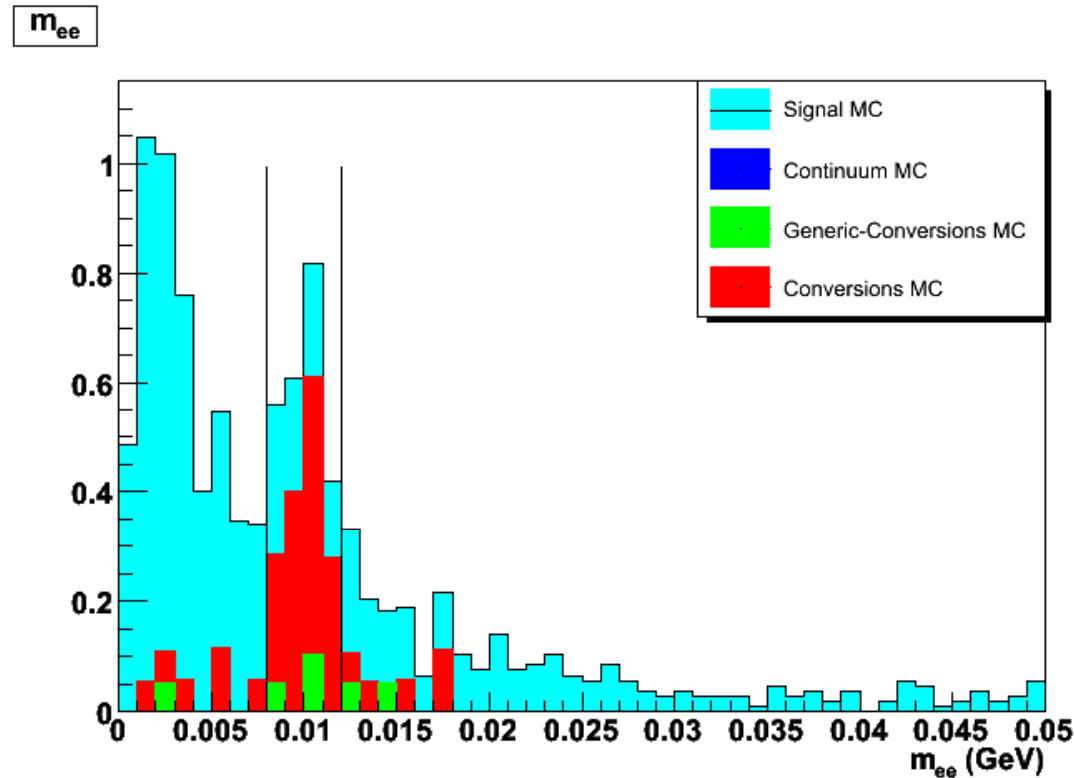
1. Cut on the χ^2 between 0 and 11.
2. Plot m_{ee} of the e^+e^- pair. Reject conversion events by rejecting peak.



m_{ee}



Vertex-Constraints all tracks of D_s^* to a point.



Expected Number of Events in 586 pb^{-1}	Electron-Fitted Samples and Criteria
Signal ($N_{\text{signalEvents}}$)	13.36
Conversion Background	1.04
Generic Background (without Conversions in e-fit)	0.42
Continuum Background	0.00
Total Background ($N_{\text{backgroundEvents}}$)	1.45
$\frac{N_{\text{signalEvents}}}{\sqrt{N_{\text{backgroundEvents}}}}$	11.1

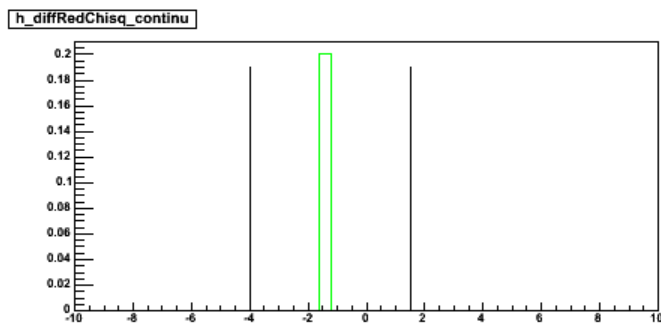
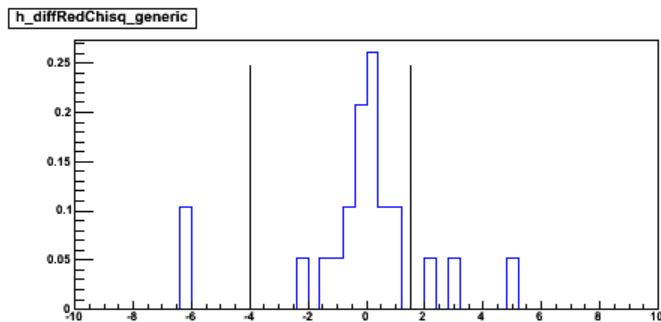
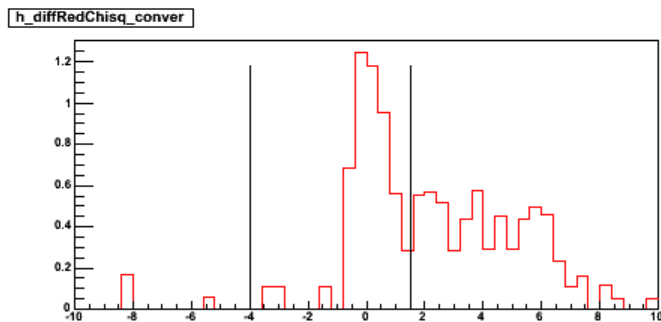
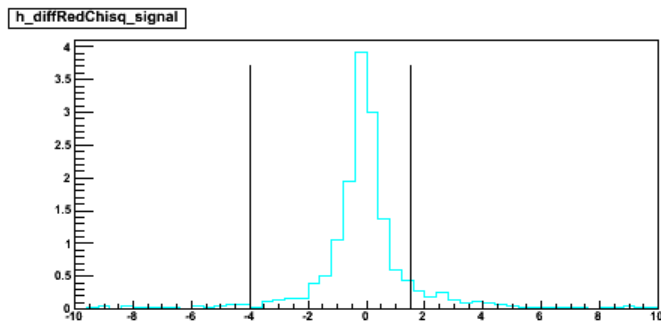
Vertex-fitted, e-fitted samples

Signal: 7.48
 Conversion bg: 0.90
 Generic - conv bg: 0.47
 Continu bg: 0
 Total bg: 1.37
 s/sqrt(b): 6.39

Vertex-Constraints e^+e^- to a point.

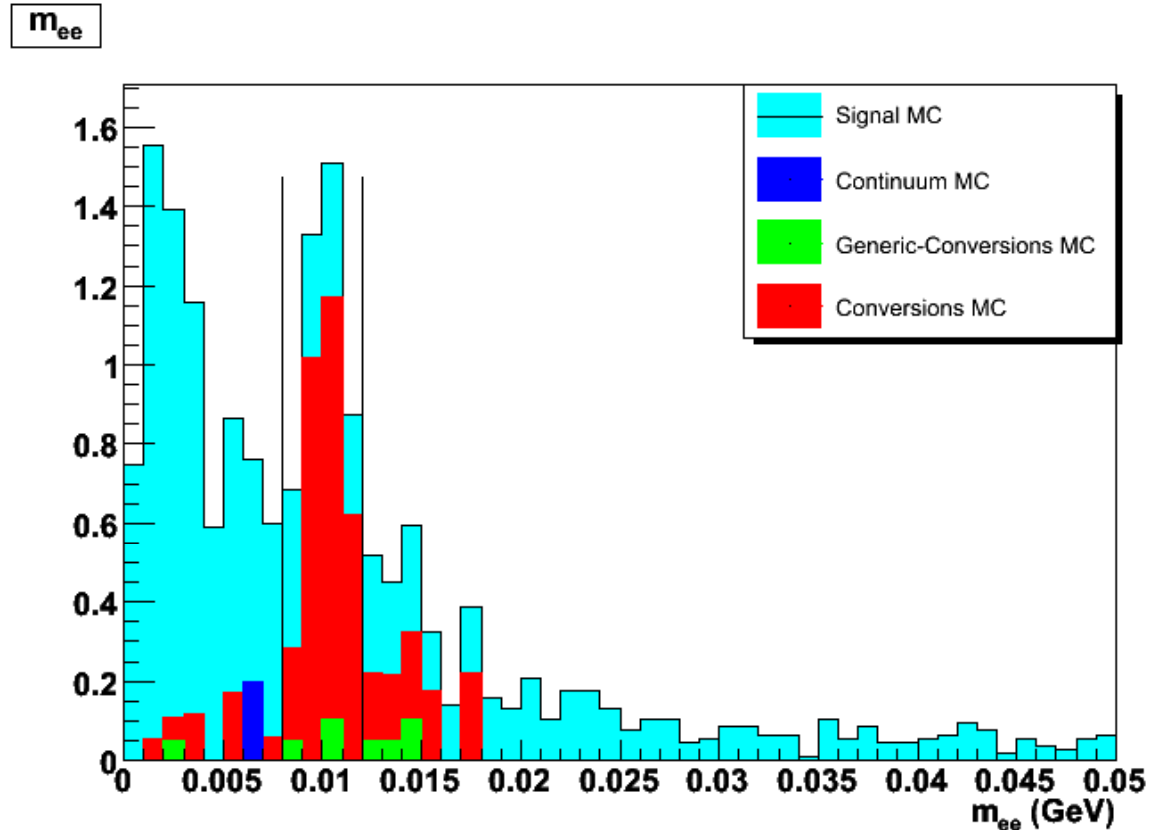
Expected Number of Events in 586 pb^{-1}	Plain Electron-Fitted Samples with All Criteria	e^+e^- Vertex-Fitted Samples with All Criteria
Signal ($N_{\text{signalEvents}}$)	13.36	8.31
Conversion Background	1.04	0.75
Generic Background (without Conversions in e-fit)	0.42	0.26
Continuum Background	0.0	0.0
Total Background ($N_{\text{backgroundEvents}}$)	1.45	1.01
$\frac{N_{\text{signalEvents}}}{\sqrt{N_{\text{backgroundEvents}}}}$	11.1	8.29

Vertex-Constraint D_s^* and beamspot, constrain D_s and beamspot



1. Plot the difference in red- χ^2 between the vertex fit of the D_s^* daughters and the beamspot and the vertex fit of the D_s daughters and the beamspot.
2. Cut between -4 and +1.5
3. Plot $m_{\{ee\}}$ between the e^+e^- and reject conversion peak.

Vertex-Constraint Ds* and beamspot, constrain Ds and beamspot

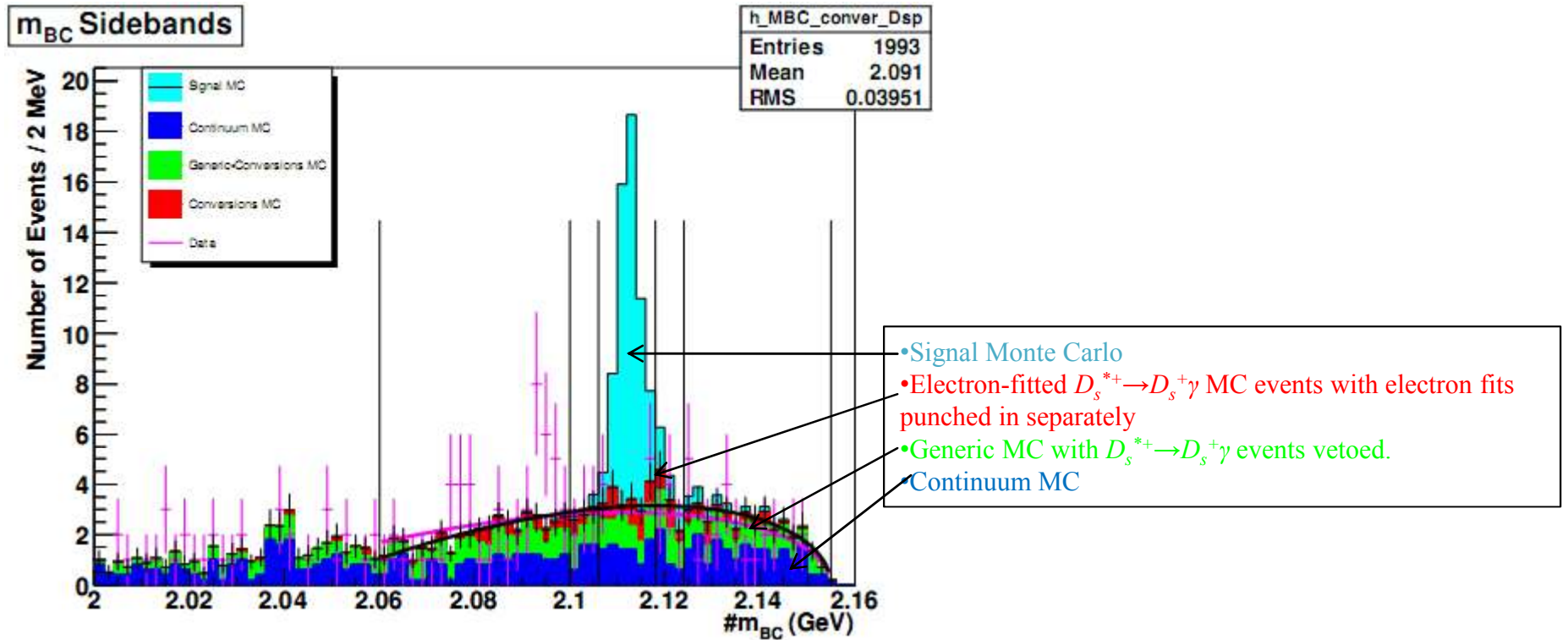


Expected Number of Events in 580 pb ⁻¹	Plain Electron-Fitted Samples with All Criteria	Vertex-Fitted Samples with All Criteria
Signal ($N_{\text{signalEvents}}$)	13.38	12.84
Conversion Background	1.04	0.92
Generic Background (without Conversions in e-fit)	0.42	0.16
Continuum Background	0.0	0.0
Total Background ($N_{\text{backgroundEvents}}$)	1.45	1.07
$\frac{N_{\text{signalEvents}}}{\sqrt{N_{\text{backgroundEvents}}}}$	11.1	12.4

Beamspot vertex-fitted samples

Signal: 11.9
 Conversion bg: 2.07
 Generic - conv bg: 0.78
 Continu bg: 0.2
 Total bg: 3.05
 s/sqrt(b): 6.82

Estimation of Background Shape from m_{BC} Sidebands

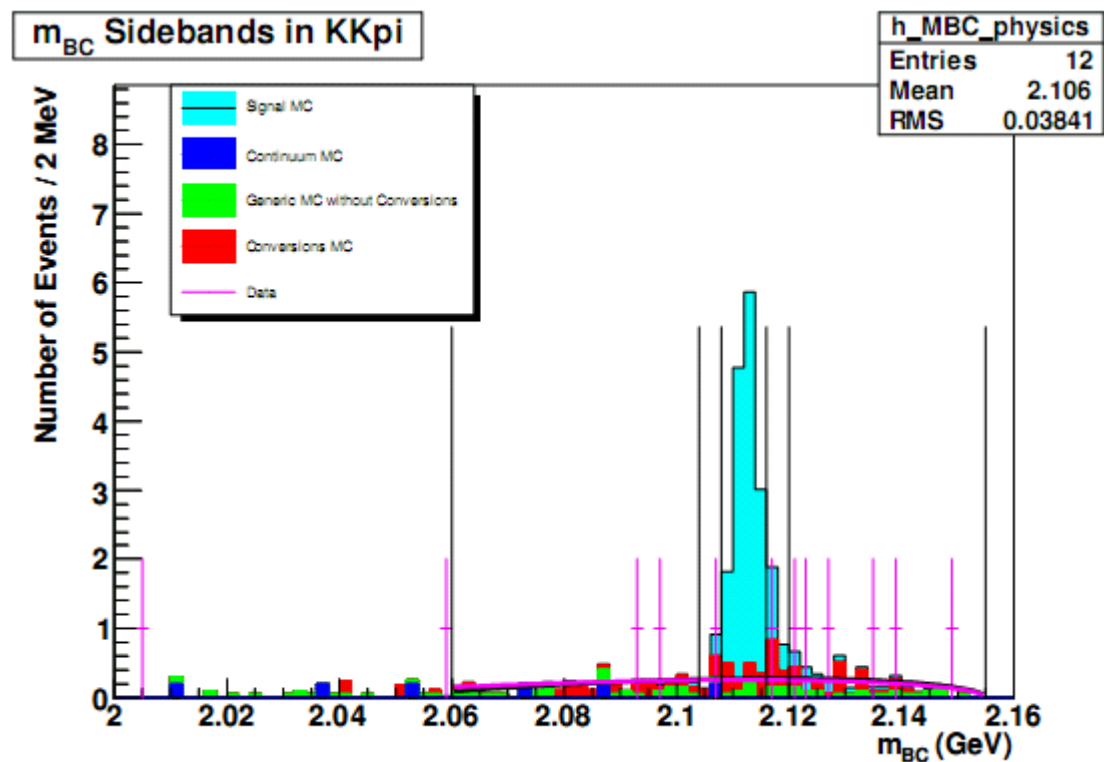


- To estimate the background in the signal region for each channel, we consider the m_{BC} distribution with all other criteria applied.
- Individual modes have low statistics. We add up the m_{BC} distributions in MC and data in all modes.
- We fit a curve to the MC background between 2.060 and 2.155 GeV. We call this the *MC shape*.
- We fit a curve to the data in the sidebands 2.060 - 2.100 & 2.124 - 2.155 GeV. This is the *data shape*.

$$N = (p_0 + p_1 m_{MBC}) \sqrt{2.155 - m_{BC}}$$

- Fit these shapes back in the individual channels to estimate background in the signal region.

Estimation of Background from m_{BC} Sidebands in the $K^+K^-\pi^+$ Mode

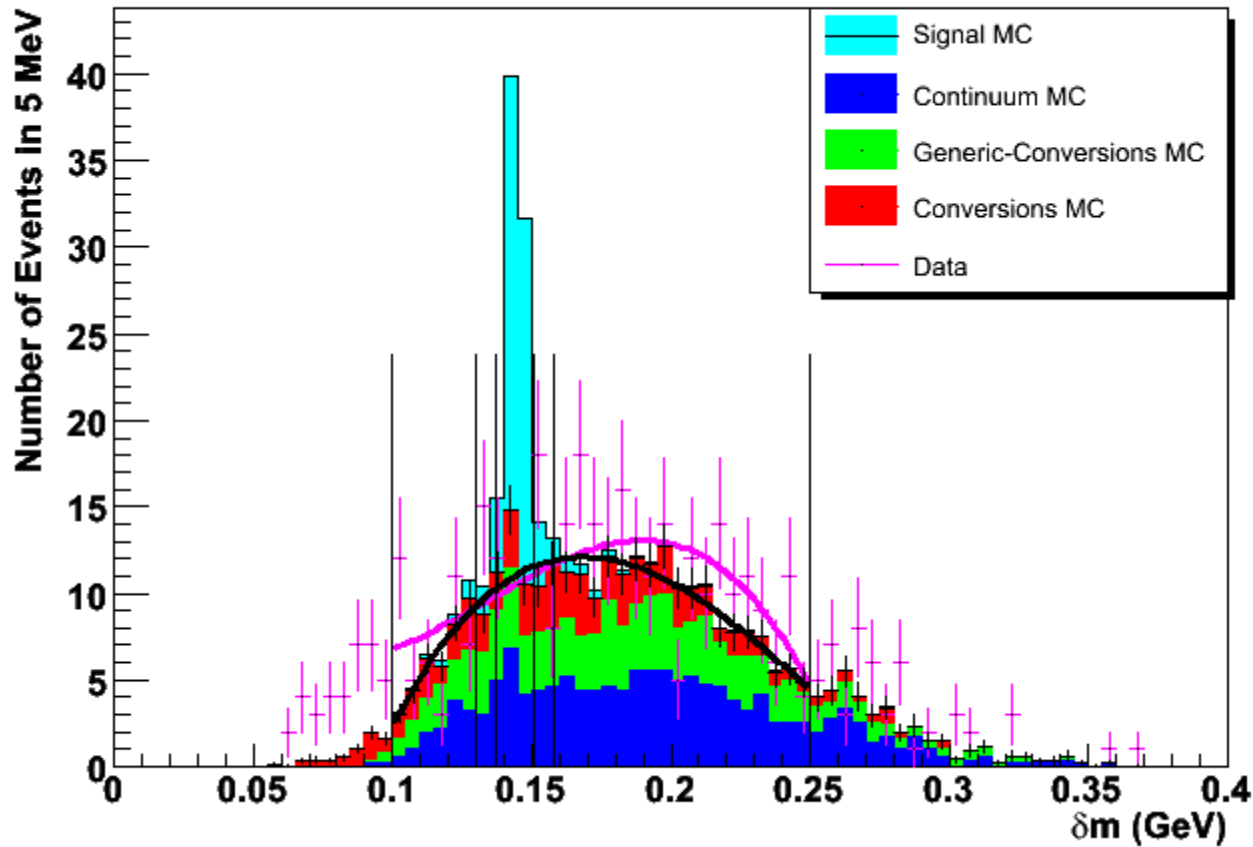


- We fit (scale) the *MC shape* and *Data shape* to the data in the sidebands of the $K^+K^-\pi^+$ (and other) mode.
- We estimate the expected background in the signal region from both shapes. The statistical uncertainty for each fit is:
$$\frac{N_{\text{expectedBackgroundFromFit}}}{\sqrt{N_{\text{sidebands}}}}$$

• For the $K^+K^-\pi^+$ channel, we estimate 1.1 ± 0.4 events from the *MC shape* and 1.0 ± 0.4 events from the *Data shape*.

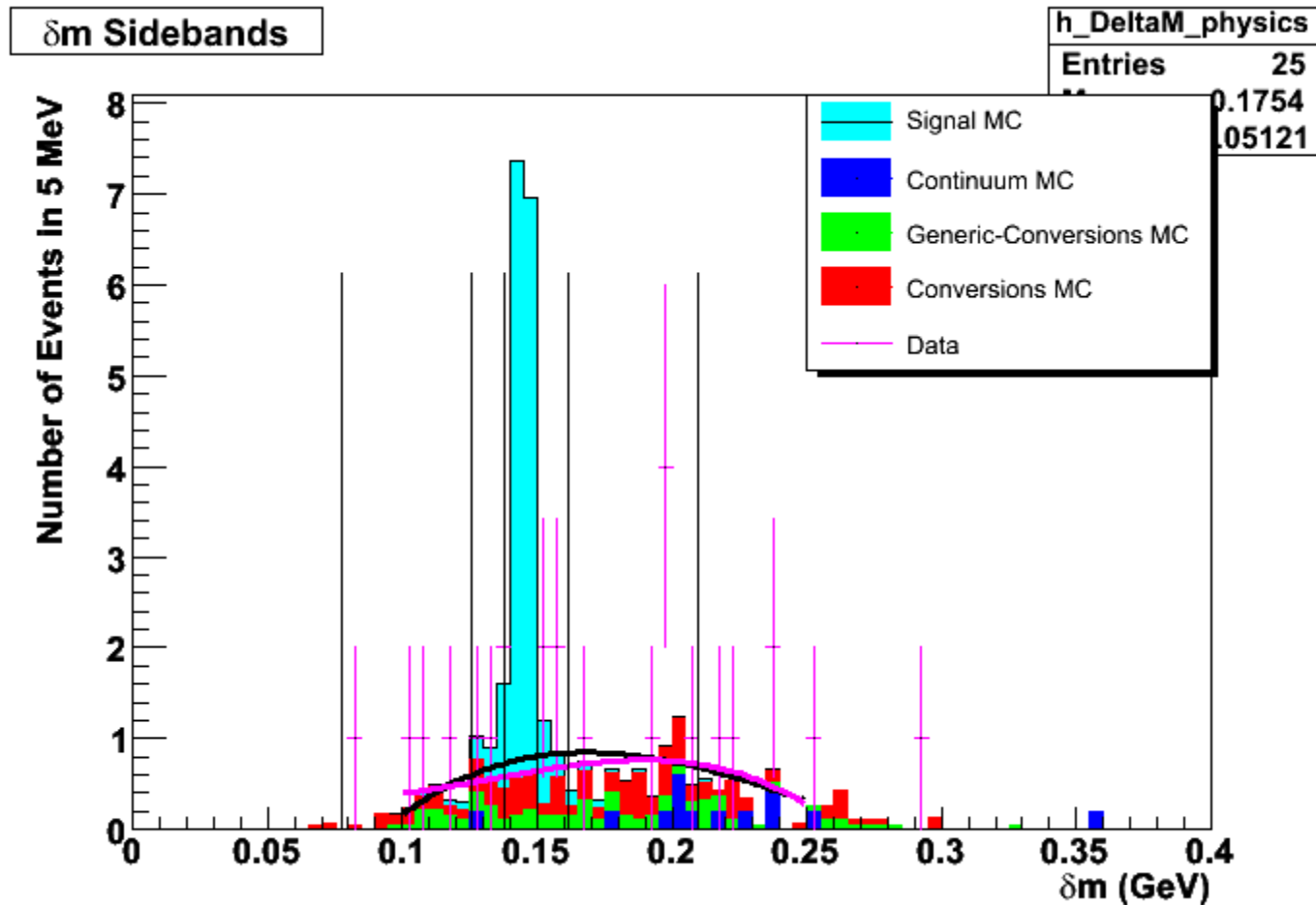
Estimation of Background Shape from δm Sidebands

δm Sidebands



- To estimate a statistical uncertainty in our estimate from fitting, we repeat our procedure with δm . Two shapes, the *MC shape* and the *Data shape* are extracted.
- We doubled the m_{BC} cut width for each mode to increase statistics.

Estimation of Background from δm Sidebands in the $K^+K^-\pi^+$ Mode



- Repeating the procedure of fitting the shapes to the individual channels, we estimate 2.1 ± 0.5 events from the *MC shape* and 1.6 ± 0.4 events from the *Data shape* in the $K^+K^-\pi^+$ channel.

Statistical Uncertainties in the Estimated Background

- Consider the mean of the *MC shape* and *Data shape* estimates from the m_{BC} distributions as the **primary estimate** because
 - the m_{BC} distribution is less peaked,
 - the difference between two estimates smaller, and
 - we did not loosen other cuts.
- The statistical error of the primary estimate is the mean of the statistical errors from the MC and Data shapes.
- Also calculate the mean of the MC shape and Data shape estimates from the δm distributions as the **secondary estimate**.
- The absolute difference between the primary and secondary estimate is recorded as the systematic uncertainty.

For the $\pi^+\eta$; $\eta \rightarrow \pi^+\pi^-\eta$; $\eta \rightarrow \gamma\gamma$ mode, we do not have any data points in the sidebands of m_{BC} or δm . To estimate the statistical error, we place a data point “by hand” at the center of the largest sideband in each case, and the statistical error is computed from that.

Estimated Background in each of the Modes

Mode	m_{BC}		δm		Background \pm (Stat) \pm (Syst)
	<i>MC Shape</i>	<i>Data Shape</i>	<i>MC Shape</i>	<i>Data Shape</i>	
$K^+K^-\pi^+$	1.1 ± 0.4	1.0 ± 0.4	2.1 ± 0.5	1.6 ± 0.4	$1.1 \pm 0.4 \pm 0.8$
$K_S K^+$	1.2 ± 0.5	1.1 ± 0.5	0.4 ± 0.2	0.3 ± 0.2	$1.1 \pm 0.5 \pm 0.8$
$\eta\pi^+$	1.5 ± 0.7	1.4 ± 0.7	1.5 ± 0.5	1.2 ± 0.4	$1.41 \pm 0.71 \pm 0.04$
$\eta'\pi^+; \eta' \rightarrow \pi^+\pi^-\eta$	0.0 ± 0.6	0.0 ± 0.7	0.0 ± 0.3	0.0 ± 0.3	$0.0 \pm 0.6 \pm 0.0$
$K^+K^-\pi^+\pi^0$	1.8 ± 0.5	1.7 ± 0.5	2.8 ± 0.6	2.2 ± 0.5	$1.7 \pm 0.5 \pm 0.8$
$\pi^+\pi^-\pi^+$	1.6 ± 0.5	1.5 ± 0.4	2.7 ± 0.6	2.1 ± 0.4	$1.6 \pm 0.5 \pm 0.8$
$K^{*+}K^{*0}$	1.8 ± 0.6	1.7 ± 0.5	2.2 ± 0.6	1.8 ± 0.5	$1.8 \pm 0.6 \pm 0.2$
$\eta\rho^+$	2.9 ± 0.6	2.7 ± 0.6	3.4 ± 0.6	2.7 ± 0.5	$2.8 \pm 0.6 \pm 0.3$
$\eta'\pi^+; \eta' \rightarrow \rho^0\gamma$	2.1 ± 0.5	1.9 ± 0.5	2.1 ± 0.6	1.7 ± 0.4	$2.0 \pm 0.5 \pm 0.1$

Projected Signal Significances assuming Signal MC Yields

Hadronic Decay Mode	Estimated Background	Projected Signal	Signal Significance
$K^+K^-\pi^+$ (1)	1.1 ± 0.9	13.9	5.32
$K_S K^+$ (2)	1.1 ± 1.0	3.1	1.73
$\eta\pi^+$ (3)	1.4 ± 0.7	1.9	1.29
$\eta'\pi^+; \eta' \rightarrow \pi^+\pi^-\eta$ (4)	0.0 ± 0.6	1.2	1.00
$K^+K^-\pi^+\pi^0$	1.7 ± 0.9	5.1	2.60
$\pi^+\pi^-\pi^+$	1.6 ± 0.9	3.8	1.93
$K^{*+}K^{*0}$ (5)	1.8 ± 0.6	2.3	1.53
$\eta\rho^+; \eta \rightarrow \gamma\gamma$	2.8 ± 0.7	5.8	2.79
$\eta'\pi^+$	2.0 ± 0.5	2.6	1.86
Sum of All Modes	13.4 ± 3.3	39.6	6.09
Sum except (1)	12.4 ± 2.8	25.7	4.63
Sum except (4)	13.4 ± 3.0	38.4	6.18
Sum except (3) & (4)	12.0 ± 2.6	36.5	6.30
Sum except (3), (4) & (5)	10.3 ± 2.3	34.2	6.42
Sum except (3), (4), (5) & (2)	9.1 ± 2.8	31.1	5.55

Unblinding Strategy

Keeping in mind that we are attempting to measure the ratio of branching fractions, K :

$$K = \frac{B(D_s^{*+} \rightarrow D_s^+ e^+ e^-)}{B(D_s^{*+} \rightarrow D_s^+ \gamma)}$$

we may write:

$$L\sigma_{D_s D_s^*} B(D_s^{*+} \rightarrow D_s^+ \gamma) B(D_s \rightarrow i) \epsilon_{D_s \gamma}^i = N_{D_s \gamma}^i \quad (61)$$

$$L\sigma_{D_s D_s^*} B(D_s^{*+} \rightarrow D_s^+ \gamma) B(D_s \rightarrow i) K \epsilon_{D_s e^+ e^-}^i = N_{D_s e^+ e^-}^i \quad (62)$$

Summing over the hadronic decay modes in consideration, we may express K as in Eq. [63](#). The value of K is what will be reported at the end of our study. The values of $\epsilon_{D_s \gamma}^i$ and $\epsilon_{D_s e^+ e^-}^i$ must be recovered from studying Monte Carlo samples.

$$K = \frac{\sum_i N_{D_s e^+ e^-}^i}{\sum_i N_{D_s \gamma}^i} \times \frac{\sum_i \epsilon_{D_s \gamma}^i B(D_s \rightarrow i)}{\sum_i \epsilon_{D_s e^+ e^-}^i B(D_s \rightarrow i)} \quad (63)$$

Unblinding Strategy – KKpi Mode

Measuring $D_s^* \rightarrow D_s \gamma$ yields and efficiency

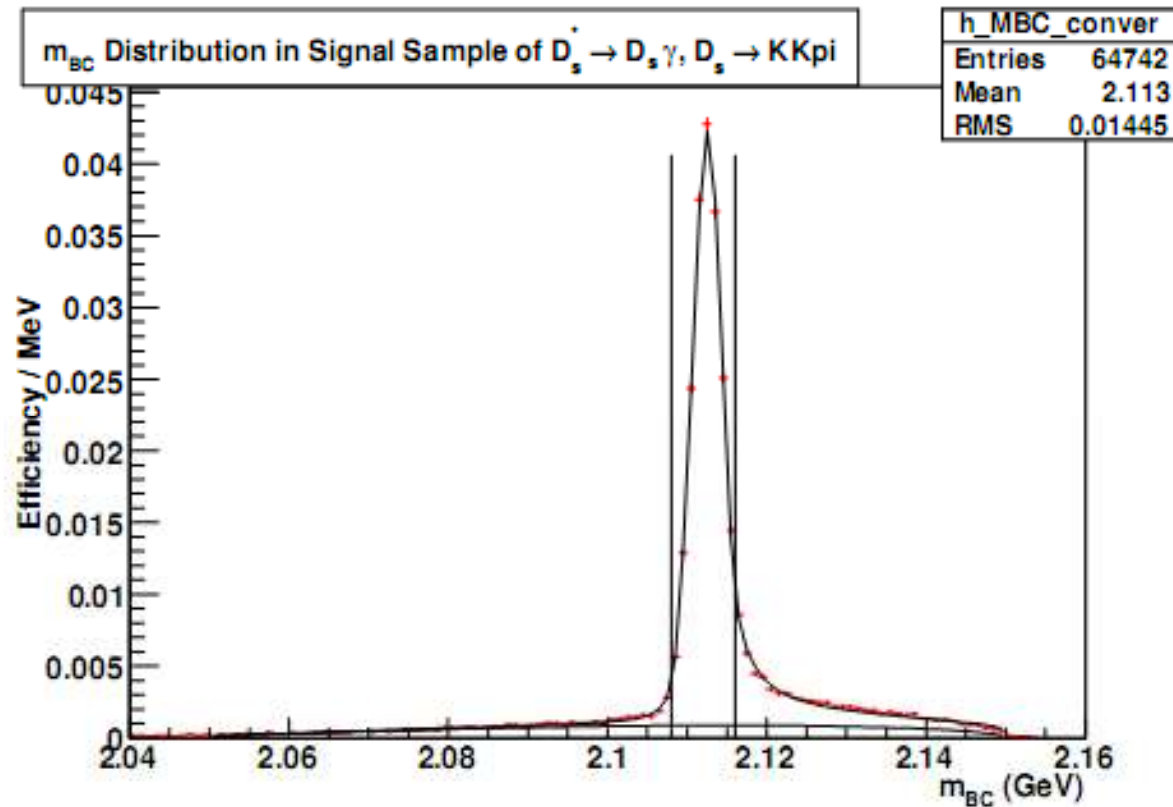


Figure 149: Distribution of m_{BC} in the signal Monte Carlo sample of $D_s^{*+} \rightarrow D_s^+ \gamma$ events where $D_s^+ \rightarrow K^+ K^- \pi^+$. The plot is normalized so as to directly read out the efficiency of the m_{BC} selection criterion.

Signal Efficiency = $19.2 \pm 0.1 \%$

Unblinding Strategy – KKpi Mode

Measuring $D_s^* \rightarrow D_s \gamma$ yields and efficiency

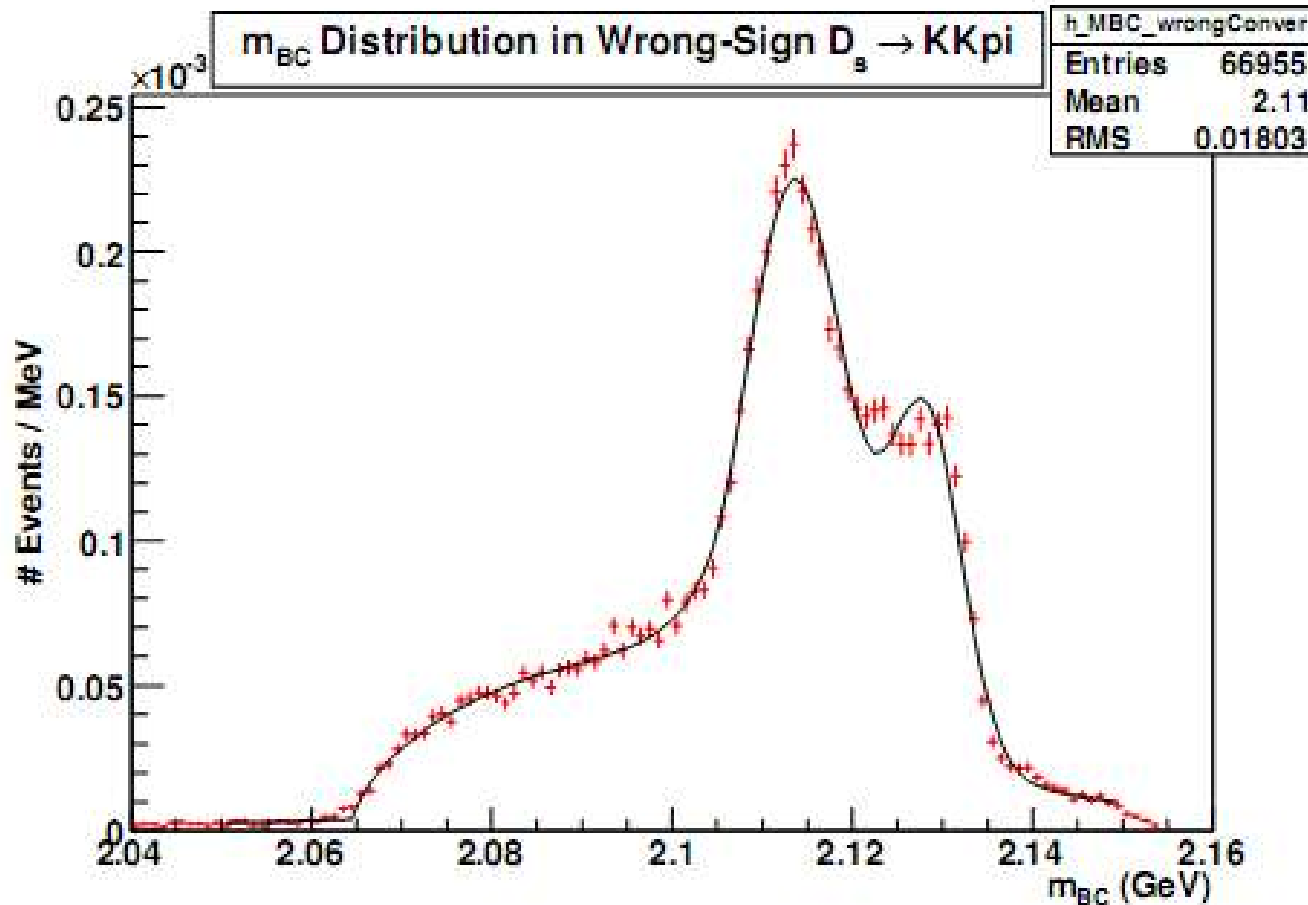


Figure 150: Peaking background in the Monte Carlo m_{BC} distribution from incorrectly reconstructing the D_s^{*+} out of the D_s^- and the γ in an event.

Unblinding Strategy – KKpi Mode

Measuring $D_s^* \rightarrow D_s \gamma$ yields and efficiency

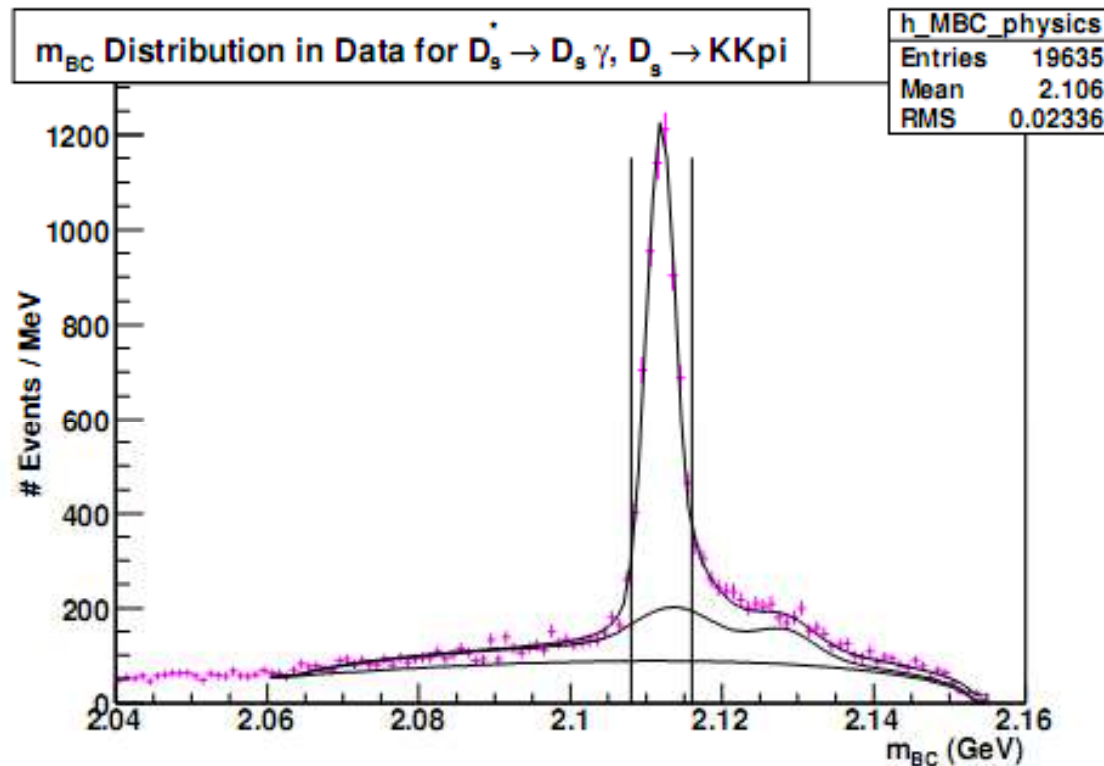


Figure 151: Distribution of m_{BC} of $D_s^{*+} \rightarrow D_s^+ \gamma$ events where $D_s^+ \rightarrow K^+ K^- \pi^+$ in 586 pb^{-1} of data. The fits are described in the text.

$N_{\{D_s \gamma\}^i}$ (signal yield above background) = 4853 events.

Unblinding Strategy – KKpi Mode

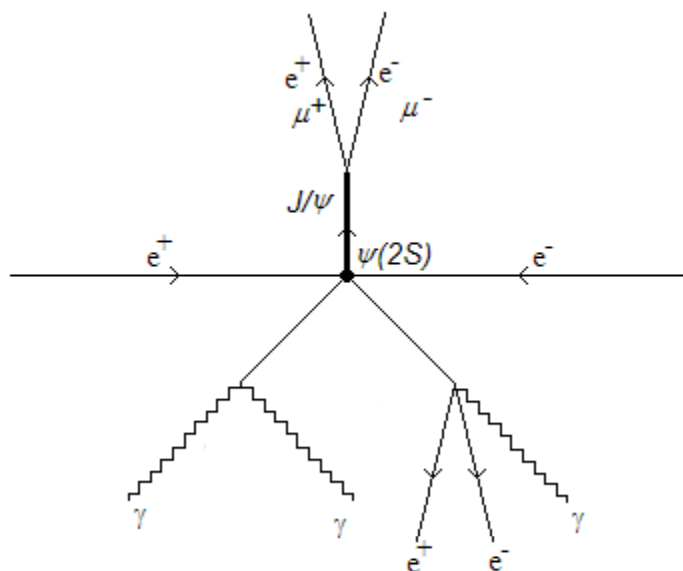
Measuring $D_s^* \rightarrow D_s$ gamma yields and efficiency

Similar procedure for delta m distribution... were some problems noticed last night.
Am yet to fix them.

Table 45: Numbers from the $K^+K^-\pi^+$ mode relevant for our measurement.

i	$N_{D_s^*e^+e^-}^i$		$N_{D_s^*\gamma}^i$	$\epsilon_{D_s^*e^+e^-}^i$	$\epsilon_{D_s^*\gamma}^i$		$B(D_s \rightarrow i)$
Mode		m_{BC}	δm		m_{BC}	δm	
$K^+K^-\pi^+$	13.9 ± 0.4 (stat) ± 0.8 (syst)	4853	4345	$0.074 \pm$	0.192 ± 0.001 (stat)	0.189 ± 0.001 (stat)	0.55 ± 0.0028

Low Energy Electron Reconstruction Efficiency



$$\begin{aligned}\psi(2S) &\rightarrow J/\psi \pi^0 \pi^0 \\ J/\psi &\rightarrow e^+ e^-; \mu^+ \mu^- \\ \pi^0 &\rightarrow \gamma \gamma \\ \pi^0 &\rightarrow \gamma e^+ e^-\end{aligned}$$

- The missing mass method hit a wall because photons radiating from the very soft electrons are often mistaken for the photon from the π^0 Dalitz decay. This leaves a photon's mass as the missing mass, which is also peaked around 0, like an electron's mass! So these end up under the inefficiency peak and we report much higher inefficiencies.
- We are now reconstructing the $\Psi(2S)$ from this channel where in one sample, the last π^0 is allowed to Dalitz decay, and in another sample the last π^0 goes to gamma gamma. The ratio of efficiencies will be very like what we have in our analysis.

Backup Slides

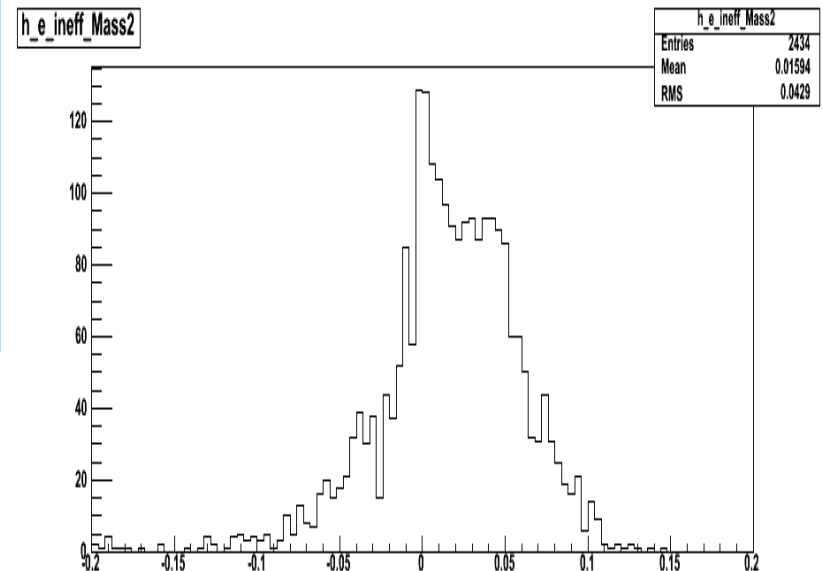
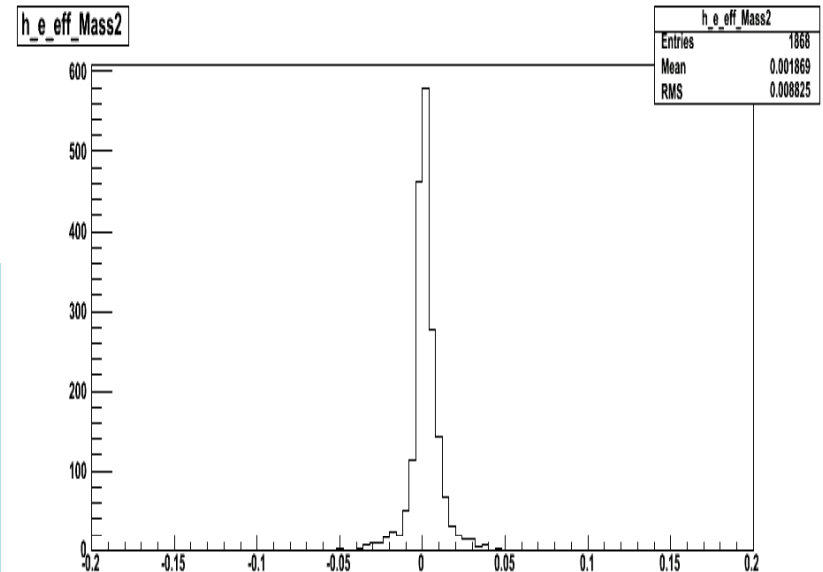
Datasets Used

Dataset	Integrated Luminosity \pm stat \pm syst
39	$55.1 \pm 0.03 \pm 0.56 \text{ pb}^{-1}$
40	$123.9 \pm 0.05 \pm 1.3 \text{ pb}^{-1}$
41	$119.1 \pm 0.05 \pm 1.3 \text{ pb}^{-1}$
47	$109.8 \pm 0.05 \pm 1.1 \text{ pb}^{-1}$
48	$178.3 \pm 0.06 \pm 1.9 \text{ pb}^{-1}$
Total	$586.2 \pm 0.11 \pm 6.1 \text{ pb}^{-1}$

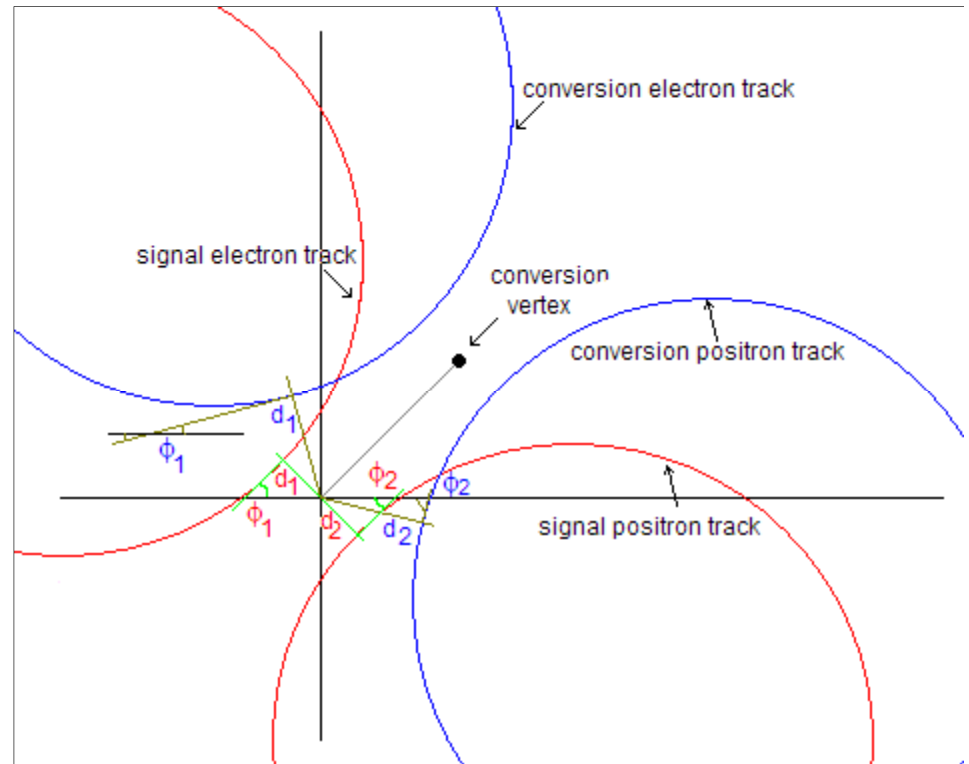
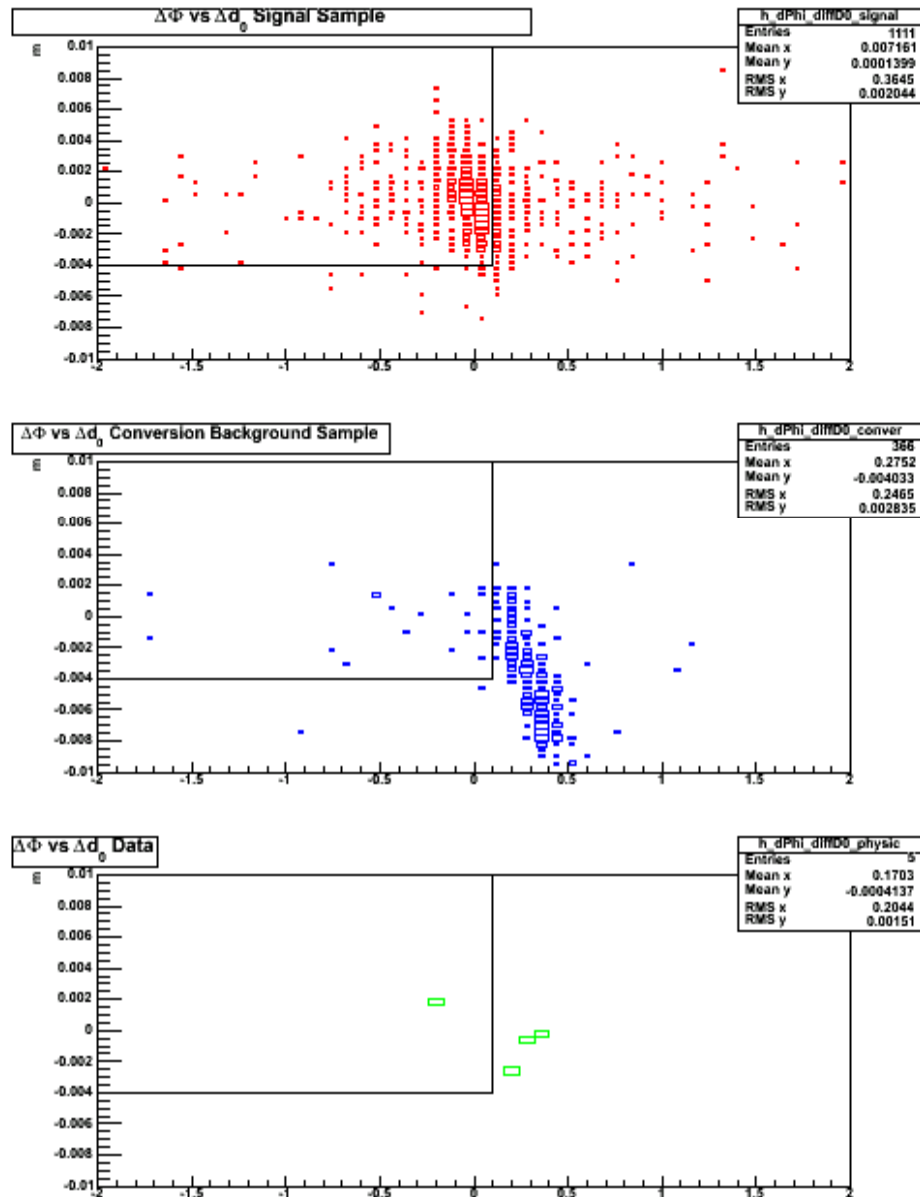
The statistical uncertainties are added in quadrature, while the systematic uncertainties are added linearly. Then these two forms of uncertainties are added in quadrature to give us $586 \pm 6 \text{ pb}^{-1}$ of integrated luminosity.

Low Energy Electron Reconstruction Efficiency

- The missing mass of this last electron is split into two plots:
 - the Efficient plot where the $\psi(2S)$ is correctly reconstructed (top plot)
 - the Inefficient plot where the $\psi(2S)$ is not correctly reconstructed (bottom plot)
- By cutting and counting, we can roughly estimate the efficiency of electron reconstruction to be $\sim 90\%$
- We will generate Monte Carlo to fit these plots for a more precise measurement.



$K^+K^-\pi^+$ Mode $\Delta\Phi$ vs Δd_0



The $\Delta\Phi$ & Δd_0 between the electron and positron in the signal (red) and conversion (blue)

Signal Monte Carlo Samples

- For signal Monte Carlo, we force the e^+e^- collision to produce a $\Psi(4160)$, and that to decay into D_s^{*+} , $D_s^{-+} + \text{c.c.}$
- We added an EVTGEN plug-in to generate vector (D_s^{*+}) to scalar (D_s^+), lepton (e^-), lepton (e^+) distributions with the invariant amplitude in consideration, apart from the invariant phase space factor.
- The D_s^+ was forced to decay through each of the previously mentioned channels. The D_s^- was allowed to decay generically.
- We fitted electrons to the electron hypothesis as well as the default pion hypothesis.
- We generated 10,000 signal MC events for each decay mode of the D_s^+ .

Background Monte Carlo

Continuum Backgrounds

- Combinatoric background from light quark (u, d, s) production. Does dominate in some channels.
- Comes with the datasets. Electrons are pion-fitted.

Generic Backgrounds

- All known physics processes at 4170 MeV involving heavy quark production.
- Comes with the datasets. Electrons are pion-fitted.
- We veto $D_s^{*+} \rightarrow D_s^+ \gamma$ events from the MC truth and replace them with privately produced and electron fitted conversion MC.

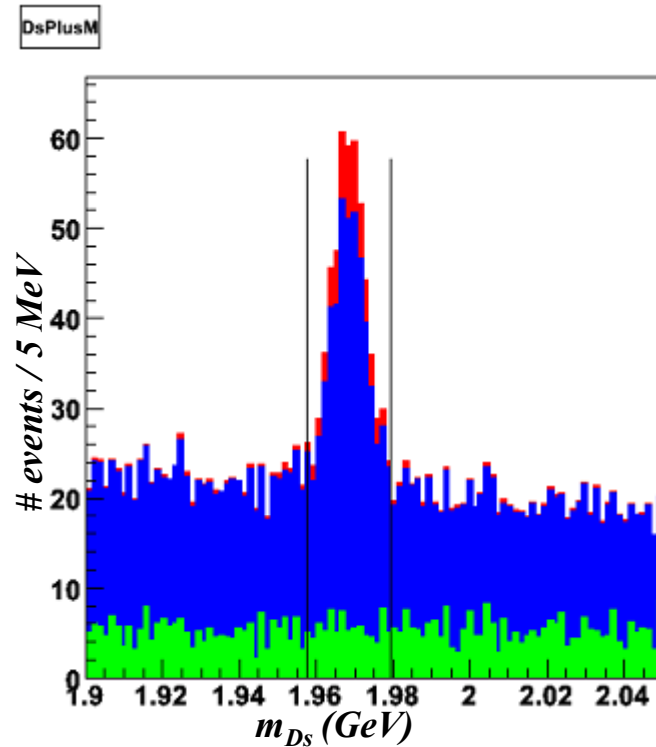
Conversion Background

- For this conversion background Monte Carlo, we force the e^+e^- collision to produce a $\Psi(4160)$, and then that to decay into the D_s^{*+}, D_s^- . The D_s^{*+} now decays via D_s^+, γ . The conversion of the photon to e^+e^- is taken care of in the detector simulation.

Selection Criteria Common to All D_s^+ Decay Modes

- Electron tracks must pass track quality cuts:
 - $10 \text{ MeV} < \text{Track Energy} < 150 \text{ MeV}$
 - $\chi^2 < 100,000$
 - $|d_0| < 5 \text{ mm}$
 - $|z_0| < 5 \text{ cm}$
 - dE/dx within 3.0σ of that expected for an electron.
- The *DTag* tools applied their default criteria for the nine investigated modes. [[5](#)]
- These cuts, and the reconstruction of a D_s^{*+} were required for filling our n-tuples on which we applied subsequent cuts.

m_{D_s} Selection Criterion for the $K^+K^-\pi^+$ Mode



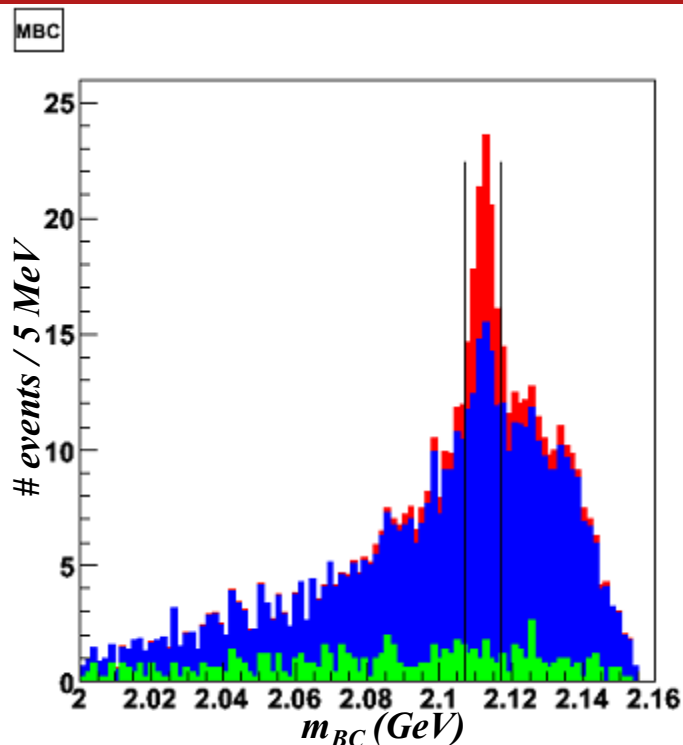
Red: Signal Monte Carlo
Blue: Generic Monte Carlo ($c\bar{c}$ production)
Green: Continuum Monte Carlo (light quarks)

Histograms normalized to 586 pb^{-1}

- We reconstruct the invariant mass m_{D_s} of a D_s from its decay products.
- Selection Criterion for this mode:

$$\left| m_{D_s} - 1.969 \text{ GeV} \right| < 0.011 \text{ GeV}$$

m_{BC} Selection Criterion for the $K^+K^-\pi^+$ Mode



Red: Signal Monte Carlo

Blue: Generic Monte Carlo ($c\bar{c}$ production)

Green: Continuum Monte Carlo (light quarks)

Histograms normalized to 586 pb^{-1}

- We know the energy of the CESR beam to high precision. Given the masses of the D_s^* and D_s , we can calculate the energy carried away by the D_s^*

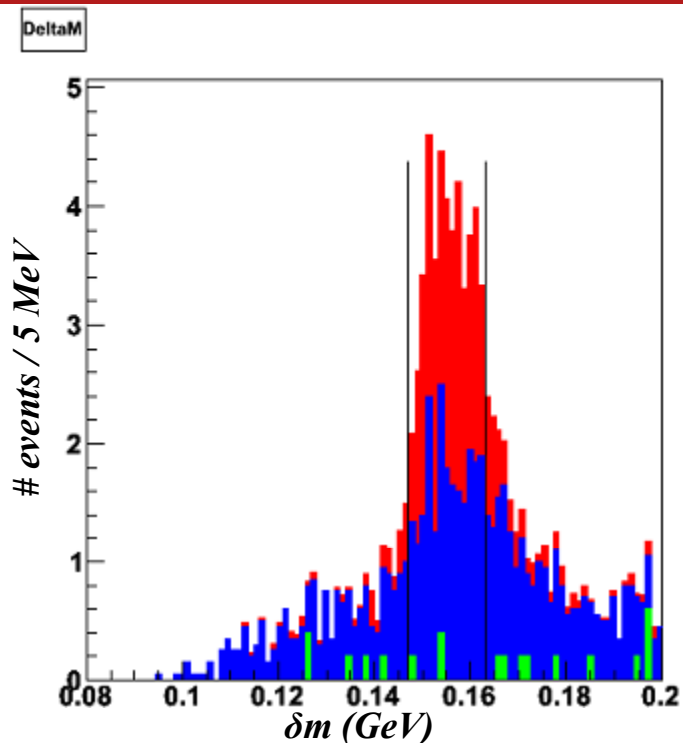
- We define the beam-constrained mass of the D_s^* as:

$$m_{BC} = \sqrt{E^2(D_s^{*+} beam) - P^2(K^+ K^- \pi^+ e^+ e^-)}$$

- Selection Criterion for this mode:

$$|m_{BC} - 2.112 \text{ GeV}| < 0.004 \text{ GeV}$$

δm Selection Criterion for the $K^+K^-\pi^+$ Mode



Red: Signal Monte Carlo

Blue: Generic Monte Carlo (cc production)

Green: Continuum Monte Carlo (light quarks)

Histograms normalized to 586 pb⁻¹

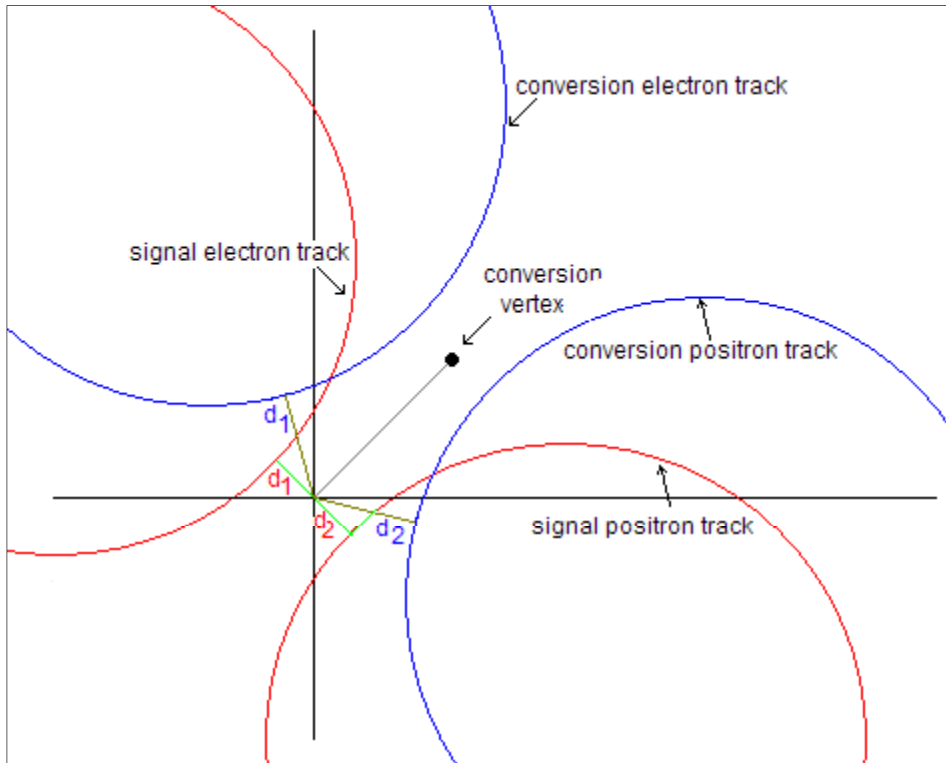
- We define δm as the mass difference between the D_s^* and the D_s where both are reconstructed from their daughters:

$$\delta m = M(K^+K^-\pi^+e^+e^-) - M(K^+K^-\pi^+)$$

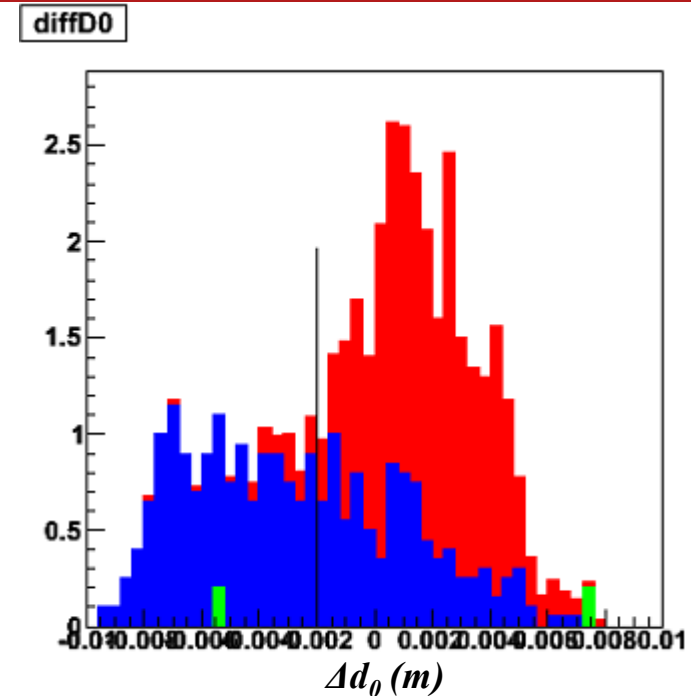
- Selection Criterion for this mode:

$$|\delta m - 0.1438 \text{ GeV}| < 0.006 \text{ GeV}$$

Δd_0 Selection Criterion for the $K^+K^-\pi^+$ Mode



Δd_0 between the electron and positron in the signal (red) and conversion (blue)



Red: Signal Monte Carlo

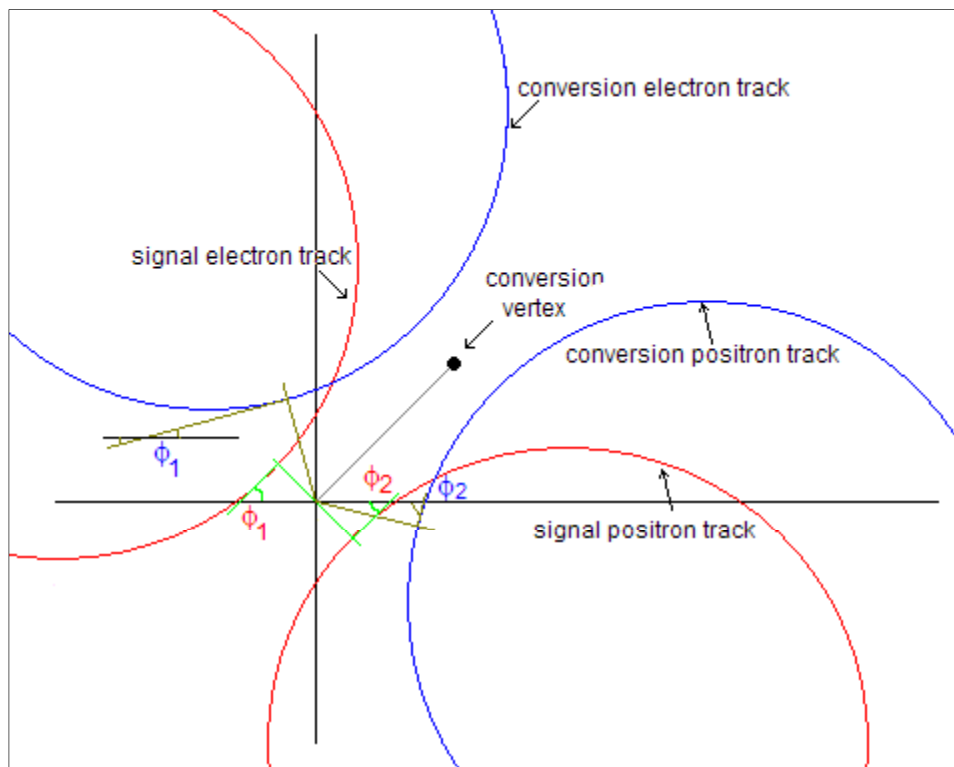
Blue: Generic Monte Carlo (heavy quarks)

Green: Continuum Monte Carlo (light quarks)

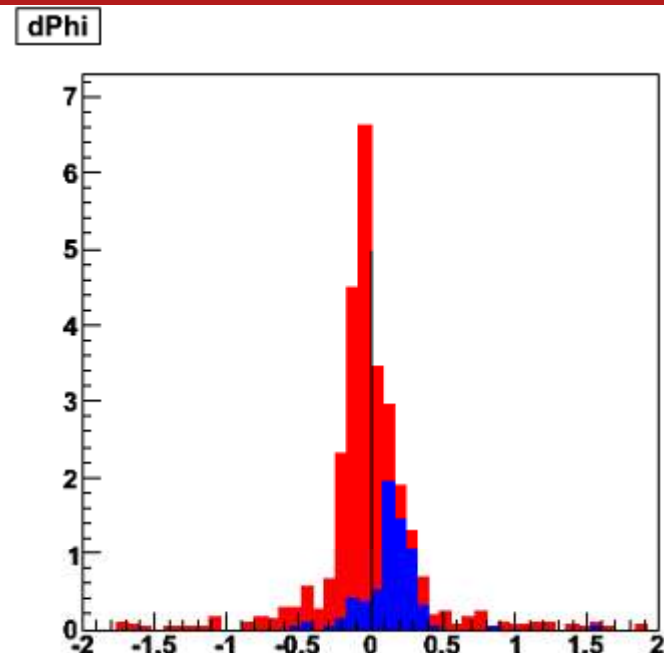
Histograms normalized to 586 pb⁻¹

- d_0 : Track's closest distance of approach to the beamline [4]
- The $\Delta d_0 = d_{0_{e^-}} - d_{0_{e^+}}$ is centered around 0 for the signal and offset from 0 for conversion backgrounds
- We require $d_1 - d_2 > -6$ mm

$\Delta\phi$ Selection Criterion for the $K^+K^-\pi^+$ Mode



$\Delta\phi$ between the electron and positron in the signal (red) and conversion (blue)



Red: Signal Monte Carlo

Blue: Generic Monte Carlo (heavy quarks)

Green: Continuum Monte Carlo (light quarks)

Histograms normalized to 586 pb⁻¹

- ϕ : Azimuth of track at origin [4]

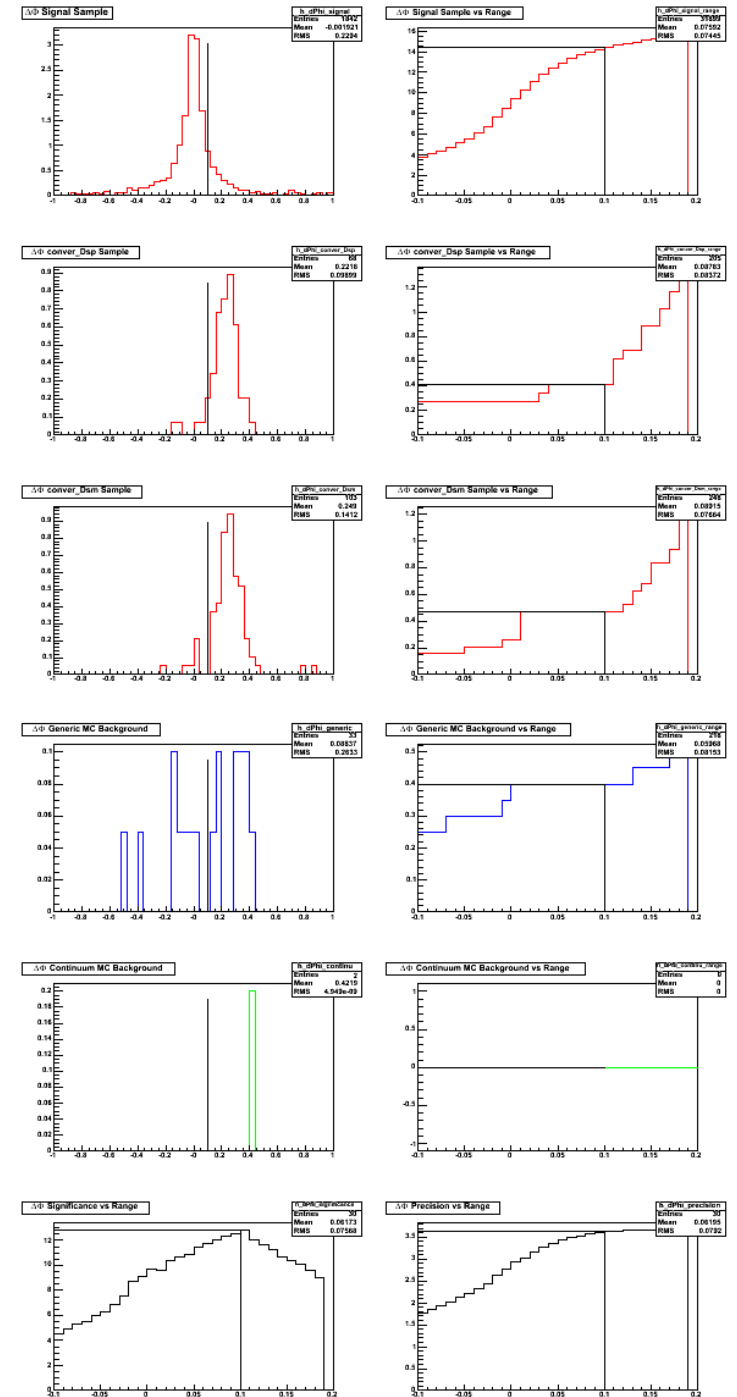
- $\Delta\phi = \phi_{e^-} - \phi_{e^+}$ is centered around 0 for the signal and offset for the conversion background.

- We require $\Delta\phi < 0.1$

- A powerful criterion against the photon conversion background.

Optimizing Selection Criteria

- We went channel by channel, criterion by criterion. (Example of $\Delta\Phi$ Selection Criterion for the $K^+K^-\pi^+$ Mode on the right)
- Plotted the signal MC, conversion MC, generic without conversion MC, and continuum MC vs variation in the cut.
- Optimized for significance $[s/\sqrt{b}]$ for low-statistics channels and precision $[s/\sqrt{(s+b)}]$ for high-statistics channels.



Prediction for Signal from Monte Carlo

Decay Mode of the D_s^+	Expected Signal Events in 586 pb ⁻¹	Expected Background Events in 586 pb ⁻¹
$K^+K^-\pi^+$	14.1	1.1
K_sK^+	3.2	0.5
$\pi^+\eta; \eta \rightarrow \gamma\gamma$	4.8	0.5
$\pi^+\eta'; \eta' \rightarrow \pi^+\pi^-\eta; \eta \rightarrow \gamma\gamma$	1.2	0.0
$K^+K^-\pi^+\pi^0$	5.1	2.2
$\pi^+\pi^-\pi^+$	3.9	2.1
$K^{*+}K^{*0}; K^{*+} \rightarrow K^0_S\pi^+; K^{*0} \rightarrow K^-\pi^+$	2.1	1.0
$\eta\rho^+; \eta \rightarrow \gamma\gamma; \rho^+ \rightarrow \pi^+\pi^0$	6.0	2.5
$\eta'\pi^+; \eta' \rightarrow \rho^0\gamma$	2.5	2.3
Total	42.9	12.2

$signal/\sqrt{background} = 12.3$, would've been **9.1** for pion-fitted data.

If $D_s^{*+} \rightarrow D_s^+e^+e^-$ exists, and our QED based estimation of its rate is correct, **we should see a clear signal over the background for it** in our data on unblinding.