Hyperbolic plane \mathbb{H}^2

This worksheet illustrates some features of <u>SageManifolds</u> (v1.0, as included in SageMath 7.5) on computations regarding the hyperbolic plane.

Click <u>here</u> to download the worksheet file (ipynb format). To run it, you must start SageMath with the Jupyter notebook, via the command sage -n jupyter

NB: a version of SageMath at least equal to 7.5 is required to run this worksheet:

```
In [1]: version()
Out[1]: 'SageMath version 7.5, Release Date: 2017-01-11'
```

First we set up the notebook to display mathematical objects using LaTeX formatting:

```
In [2]: %display latex
```

We also define a viewer for 3D plots (use 'threejs' or 'jmol' for interactive 3D graphics):

```
In [3]: viewer3D = 'jmol' # must be 'threejs', jmol', 'tachyon' or None (defaul
t)
```

We declare \mathbb{H}^2 as a 2-dimensional differentiable manifold:

```
In [4]: H2 = Manifold(2, 'H2', latex_name=r'\mathbb{H}^2', start_index=1)
    print(H2)
    H2
```

2-dimensional differentiable manifold H2

```
\mathsf{Out}[4]:\ \mathbb{H}^2
```

We shall introduce charts on \mathbb{H}^2 that are related to various models of the hyperbolic plane as submanifolds of \mathbb{R}^3 . Therefore, we start by declaring \mathbb{R}^3 as a 3-dimensional manifold equiped with a global chart: the chart of Cartesian coordinates (X,Y,Z):

```
In [5]: R3 = Manifold(3, 'R3', latex_name=r'\mathbb{R}^3', start_index=1)

X3.<X,Y,Z>=R3.chart()

Out[5]: (\mathbb{R}^3,(X,Y,Z))
```

Hyperboloid model

The first chart we introduce is related to the **hyperboloid model of** \mathbb{H}^2 , namely to the representation of \mathbb{H}^2 as the upper sheet (Z>0) of the hyperboloid of two sheets defined in \mathbb{R}^3 by the equation $X^2+Y^2-Z^2=-1$:

```
In [6]: X_{hyp}.<X,Y>= H2.chart()

X_{hyp}

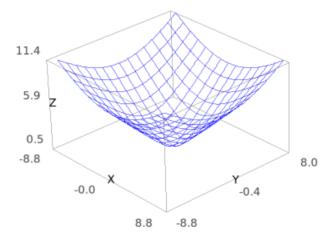
Out[6]: (\mathbb{H}^2,(X,Y))
```

The corresponding embedding of \mathbb{H}^2 in \mathbb{R}^3 is

In [7]: Phi1 = H2.diff_map(R3, [X, Y,
$$sqrt(1+X^2+Y^2)$$
], name='Phi_1', latex_nam e=r'\Phi_1') Phi1.display()

Out[7]:
$$\Phi_1$$
: \mathbb{H}^2 \longrightarrow \mathbb{R}^3 (X,Y) \longmapsto $(X,Y,Z) = (X,Y,\sqrt{X^2+Y^2+1})$

By plotting the chart $(\mathbb{H}^2, (X, Y))$ in terms of the Cartesian coordinates of \mathbb{R}^3 , we get a graphical view of $\Phi_1(\mathbb{H}^2)$:

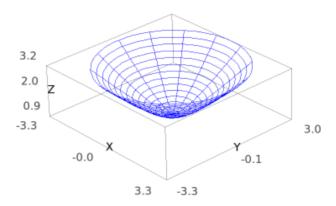


A second chart is obtained from the polar coordinates (r,φ) associated with (X,Y). Contrary to (X,Y), the polar chart is not defined on the whole \mathbb{H}^2 , but on the complement U of the segment $\{Y=0,x\geq 0\}$:

Open subset U of the 2-dimensional differentiable manifold H2

Note that (y!=0, x<0) stands for $y \neq 0$ OR x < 0; the condition $y \neq 0$ AND x < 0 would have been written [y!=0, x<0] instead.

```
In [10]: X \text{ pol.} < r, ph> = U.chart(r'r:(0,+00) ph:(0,2*pi):(varphi')
            X_pol
Out[10]: (U, (r, \varphi))
In [11]: X_pol.coord_range()
Out[11]: r: (0, +\infty); \varphi: (0, 2\pi)
            We specify the transition map between the charts (U, (r, \varphi)) and (\mathbb{H}^2, (X, Y)) as X = r \cos \varphi,
            Y = r \sin \varphi:
In [12]: pol to hyp = X pol.transition map(X hyp, [r*cos(ph), r*sin(ph)])
            pol_to_hyp
Out[12]: (U, (r, \varphi)) \to (U, (X, Y))
In [13]: pol_to_hyp.display()
Out[13]: \int X = r \cos(\varphi)
            Y = r \sin(\varphi)
In [14]: pol_to_hyp.set_inverse(sqrt(X^2+Y^2), atan2(Y, X))
In [15]: pol_to_hyp.inverse().display()
Out[15]: \int r = \sqrt{X^2 + Y^2}
                  = \arctan(Y, X)
            The restriction of the embedding \Phi_1 to U has then two coordinate expressions:
In [16]: Phil.restrict(U).display()
Out[16]: \Phi_1: U \longrightarrow \mathbb{R}^3
                   (X,Y) \longmapsto (X,Y,Z) = (X,Y,\sqrt{X^2+Y^2+1})
                   (r, \varphi) \longmapsto (X, Y, Z) = (r \cos(\varphi), r \sin(\varphi), \sqrt{r^2 + 1})
```



In [18]: Phil._coord_expression
$$\underbrace{\left\{ \left(\left(\mathbb{H}^2, (X,Y) \right), \left(\mathbb{R}^3, (X,Y,Z) \right) \right) : \left(X,Y,\sqrt{X^2+Y^2+1} \right) \right\} }$$

Metric and curvature

The metric on \mathbb{H}^2 is that induced by the Minkowksy metric on \mathbb{R}^3 : \$\$ \eta = \mathrm{d}X\otimes \mathrm{d}Y\otimes\mathrm{d}Y

• \mathrm{d}Z\otimes\mathrm{d}Z \$\$

Out[20]:
$$g = \left(\frac{Y^2 + 1}{X^2 + Y^2 + 1}\right) dX \otimes dX + \left(-\frac{XY}{X^2 + Y^2 + 1}\right) dX \otimes dY + \left(-\frac{XY}{X^2 + Y^2 + 1}\right) dY \otimes dX + \left(\frac{X^2 + 1}{X^2 + Y^2 + 1}\right) dY \otimes dY$$

The expression of the metric tensor in terms of the polar coordinates is

In [21]: g.display(X_pol.frame(), X_pol)

Out[21]:
$$g = \left(\frac{1}{r^2 + 1}\right) dr \otimes dr + r^2 d\varphi \otimes d\varphi$$

The Riemann curvature tensor associated with g is

In [22]: Riem = g.riemann()
print(Riem)

Tensor field Riem(g) of type (1,3) on the 2-dimensional differentiable manifold H2

In [23]: Riem.display(X pol.frame(), X pol)

Out[23]: $\text{Riem}(g) = -r^2 \frac{\partial}{\partial r} \otimes d\varphi \otimes dr \otimes d\varphi + r^2 \frac{\partial}{\partial r} \otimes d\varphi \otimes d\varphi \otimes dr + \left(\frac{1}{r^2 + 1}\right) \frac{\partial}{\partial \varphi}$ $\otimes dr \otimes dr \otimes d\varphi + \left(-\frac{1}{r^2 + 1}\right) \frac{\partial}{\partial \varphi} \otimes dr \otimes d\varphi \otimes dr$

The Ricci tensor and the Ricci scalar:

In [24]: Ric = g.ricci()
print(Ric)

Field of symmetric bilinear forms $\operatorname{Ric}(g)$ on the 2-dimensional different iable manifold H2

In [25]: Ric.display(X_pol.frame(), X_pol)

Out[25]: $\operatorname{Ric}(g) = \left(-\frac{1}{r^2 + 1}\right) dr \otimes dr - r^2 d\varphi \otimes d\varphi$

In [26]: Rscal = g.ricci_scalar()
print(Rscal)

Scalar field r(g) on the 2-dimensional differentiable manifold H2

In [27]: Rscal.display()

Out[27]: $r(g): \mathbb{H}^2 \longrightarrow \mathbb{R}$ $(X,Y) \longmapsto -2$ on $U: (r,\varphi) \longmapsto -2$

Hence we recover the fact that (\mathbb{H}^2, g) is a space of **constant negative curvature**.

In dimension 2, the Riemann curvature tensor is entirely determined by the Ricci scalar R according to

$$R^{i}_{jlk} = \frac{R}{2} \left(\delta^{i}_{k} g_{jl} - \delta^{i}_{l} g_{jk} \right)$$

Let us check this formula here, under the form $R^i_{jlk} = -Rg_{j[k}\delta^i_{l]}$:

```
In [28]: delta = H2.tangent_identity_field()
Riem == - Rscal*(g*delta).antisymmetrize(2,3) # 2,3 = last positions o
f the type-(1,3) tensor g*delta
```

Out[28]: True

Similarly the relation Ric = (R/2) g must hold:

```
In [29]: Ric == (Rscal/2)*g
```

Out[29]: True

Poincaré disk model

The Poincaré disk model of \mathbb{H}^2 is obtained by stereographic projection from the point S=(0,0,-1) of the hyperboloid model to the plane Z=0. The radial coordinate R of the image of a point of polar coordinate (r,φ) is

$$R = \frac{r}{1 + \sqrt{1 + r^2}}.$$

Hence we define the Poincaré disk chart on \mathbb{H}^2 by

```
In [30]: X_Pdisk.<R,ph> = U.chart(r'R:(0,1) ph:(0,2*pi):\varphi')
X_Pdisk
```

Out[30]: $(U, (R, \varphi))$

Out[31]:
$$R: (0,1); \varphi: (0,2\pi)$$

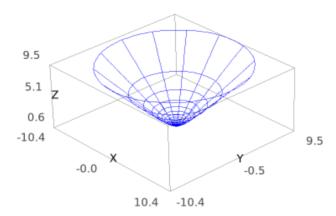
and relate it to the hyperboloid polar chart by

Out[32]: $(U,(r,\varphi)) \rightarrow (U,(R,\varphi))$

Out[33]:
$$\begin{cases} R = \frac{r}{\sqrt{r^2+1}+1} \\ \varphi = \varphi \end{cases}$$

Out[34]:
$$\begin{cases} r = -\frac{2R}{R^2 - 1} \\ \varphi = \varphi \end{cases}$$

A view of the Poincaré disk chart via the embedding $\Phi_1\colon$



The expression of the metric tensor in terms of coordinates (R, φ) :

Out[36]:
$$g = \left(\frac{4}{R^4 - 2R^2 + 1}\right) dR \otimes dR + \left(\frac{4R^2}{R^4 - 2R^2 + 1}\right) d\varphi \otimes d\varphi$$

We may factorize each metric component:

Out[37]:
$$g = \frac{4}{(R+1)^2(R-1)^2} dR \otimes dR + \frac{4R^2}{(R+1)^2(R-1)^2} d\varphi \otimes d\varphi$$

Cartesian coordinates on the Poincaré disk

Let us introduce Cartesian coordinates (u, v) on the Poincaré disk; since the latter has a unit radius, this amounts to define the following chart on \mathbb{H}^2 :

Out[38]:
$$(\mathbb{H}^2, (u, v))$$

On U, the Cartesian coordinates (u,v) are related to the polar coordinates (R,φ) by the standard formulas:

Out[39]: $(U, (R, \varphi)) \to (U, (u, v))$

Out[40]:
$$\begin{cases} u = R\cos(\varphi) \\ v = R\sin(\varphi) \end{cases}$$

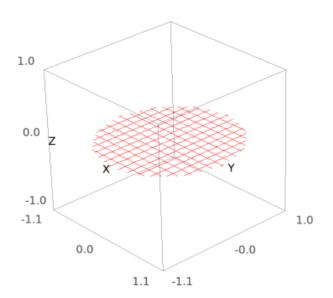
Out[41]:
$$\begin{cases} R = \sqrt{u^2 + v^2} \\ \varphi = \arctan(v, u) \end{cases}$$

The embedding of \mathbb{H}^2 in \mathbb{R}^3 associated with the Poincaré disk model is naturally defined as

Out[42]:
$$\Phi_2$$
: $\mathbb{H}^2 \longrightarrow \mathbb{R}^3$
 $(u, v) \longmapsto (X, Y, Z) = (u, v, 0)$

Let us use it to draw the Poincaré disk in \mathbb{R}^3 :

In [43]: graph_disk_uv = X_Pdisk_cart.plot(X3, mapping=Phi2, number_values=15)
show(graph_disk_uv, viewer=viewer3D, figsize=7)



On U, the change of coordinates $(r, \varphi) \to (u, v)$ is obtained by combining the changes $(r, \varphi) \to (R, \varphi)$ and $(R, \varphi) \to (u, v)$:

Out[44]: $(U, (r, \varphi)) \to (U, (u, v))$

In [45]: pol_to_Pdisk_cart.display()

Out[45]:
$$\begin{cases} u = \frac{r\cos(\varphi)}{\sqrt{r^2+1}+1} \\ v = \frac{r\sin(\varphi)}{\sqrt{r^2+1}+1} \end{cases}$$

Still on U, the change of coordinates $(X,Y) \to (u,v)$ is obtained by combining the changes $(X,Y) \to (r,\varphi)$ with $(r,\varphi) \to (u,v)$:

Out[46]: $(U, (X, Y)) \to (U, (u, v))$

In [47]: hyp_to_Pdisk_cart_U.display()

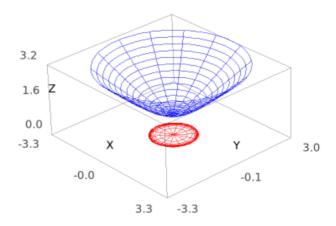
Out[47]:
$$\begin{cases} u = \frac{X}{\sqrt{X^2 + Y^2 + 1} + 1} \\ v = \frac{Y}{\sqrt{X^2 + Y^2 + 1} + 1} \end{cases}$$

We use the above expression to extend the change of coordinates $(X,Y) \to (u,v)$ from U to the whole manifold \mathbb{H}^2 :

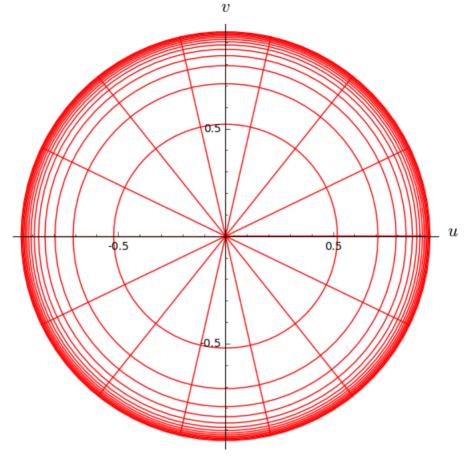
Out[48]:
$$(\mathbb{H}^2, (X, Y)) \rightarrow (\mathbb{H}^2, (u, v))$$

Out[49]:
$$\begin{cases} u = \frac{X}{\sqrt{X^2 + Y^2 + 1} + 1} \\ v = \frac{Y}{\sqrt{X^2 + Y^2 + 1} + 1} \end{cases}$$

Out[50]:
$$\begin{cases} X = -\frac{2u}{u^2 + v^2 - 1} \\ Y = -\frac{2v}{u^2 + v^2 - 1} \end{cases}$$



In [52]: X_pol.plot(X_Pdisk_cart, ranges={r: (0, 20)}, number_values=15)
Out[52]:



Metric tensor in Poincaré disk coordinates (u, v)

From now on, we are using the Poincaré disk chart $(\mathbb{H}^2, (u, v))$ as the default one on \mathbb{H}^2 :

In [54]: g.display(X_hyp.frame())

Out[54]:

$$g = \left(\frac{u^4 + v^4 + 2(u^2 + 1)v^2 - 2u^2 + 1}{u^4 + v^4 + 2(u^2 + 1)v^2 + 2u^2 + 1}\right) dX \otimes dX$$

$$+ \left(-\frac{4uv}{u^4 + v^4 + 2(u^2 + 1)v^2 + 2u^2 + 1}\right) dX \otimes dY$$

$$+ \left(-\frac{4uv}{u^4 + v^4 + 2(u^2 + 1)v^2 + 2u^2 + 1}\right) dY \otimes dX$$

$$+ \left(\frac{u^4 + v^4 + 2(u^2 + 1)v^2 + 2u^2 + 1}{u^4 + v^4 + 2(u^2 + 1)v^2 + 2u^2 + 1}\right) dY \otimes dY$$

Out[55]:

$$g = \left(\frac{4}{u^4 + v^4 + 2(u^2 - 1)v^2 - 2u^2 + 1}\right) du \otimes du$$

$$+ \left(\frac{4}{u^4 + v^4 + 2(u^2 - 1)v^2 - 2u^2 + 1}\right) dv \otimes dv$$

Out[56]:
$$g = \frac{4}{(u^2 + v^2 - 1)^2} du \otimes du + \frac{4}{(u^2 + v^2 - 1)^2} dv \otimes dv$$

Hemispherical model

The **hemispherical model of** \mathbb{H}^2 is obtained by the inverse stereographic projection from the point S=(0,0,-1) of the Poincaré disk to the unit sphere $X^2+Y^2+Z^2=1$. This induces a spherical coordinate chart on U:

In [57]:
$$X_{spher.} = U.chart(r'th:(0,pi/2):\theta ph:(0,2*pi):\varphi') X_{spher}$$

Out [57]: $(U, (\theta, \varphi))$

From the stereographic projection from S, we obtain that

$$\sin\theta = \frac{2R}{1 + R^2}$$

Hence the transition map:

Out[58]: $(U,(R,\varphi)) \rightarrow (U,(\theta,\varphi))$

Out[59]:
$$\begin{cases} \theta = \arcsin\left(\frac{2R}{R^2+1}\right) \\ \varphi = \varphi \end{cases}$$

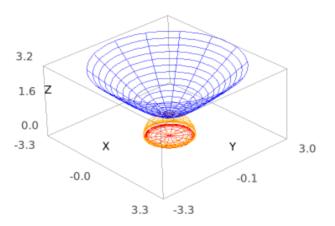
Out[60]:
$$\begin{cases} R = \frac{\sin(\theta)}{\cos(\theta)+1} \\ \varphi = \varphi \end{cases}$$

In the spherical coordinates (θ, φ) , the metric takes the following form:

Out[61]:
$$g = \frac{1}{\cos(\theta)^2} d\theta \otimes d\theta + \frac{\sin(\theta)^2}{\cos(\theta)^2} d\varphi \otimes d\varphi$$

The embedding of \mathbb{H}^2 in \mathbb{R}^3 associated with the hemispherical model is naturally:

Out[62]:
$$\Phi_3$$
: \mathbb{H}^2 \longrightarrow \mathbb{R}^3 on U : (R, φ) \longmapsto $(X, Y, Z) = \left(\frac{2R\cos(\varphi)}{R^2+1}, \frac{2R\sin(\varphi)}{R^2+1}, -\frac{R^2-1}{R^2+1}\right)$ on U : (θ, φ) \longmapsto $(X, Y, Z) = (\cos(\varphi)\sin(\theta), \sin(\varphi)\sin(\theta), \cos(\theta))$



Poincaré half-plane model

The **Poincaré half-plane model of** \mathbb{H}^2 is obtained by stereographic projection from the point W=(-1,0,0) of the hemispherical model to the plane X=1. This induces a new coordinate chart on \mathbb{H}^2 by setting (x,y)=(Y,Z) in the plane X=1:

In [64]:
$$X_{hplane}. = H2.chart('x y:(0,+oo)')$$

 X_{hplane}
Out[64]: $(\mathbb{H}^2,(x,y))$

The coordinate transformation $(\theta, \varphi) \to (x, y)$ is easily deduced from the stereographic projection from the point W:

```
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In [65]:
                                 1+sin(th)*cos(ph)),
                                                                                                                                                                                                                    2*cos(th)/(1+sin(th
                                 )*cos(ph))])
                                spher_to_hplane
Out[65]: (U, (\theta, \varphi)) \rightarrow (U, (x, y))
In [66]: spher to hplane.display()
                                                                2 \sin(\varphi) \sin(\theta)
Out[66]:
                                                                      2 \cos(\theta)
In [67]: Pdisk_to_hplane = spher_to_hplane * Pdisk_to_spher
                                Pdisk to hplane
Out[67]: (U, (R, \varphi)) \to (U, (x, y))
In [68]: Pdisk_to_hplane.display()
Out[68]:
In [69]: Pdisk_cart_to_hplane_U = Pdisk_to_hplane * Pdisk_to_Pdisk_cart.inverse(
                                Pdisk cart to hplane U
Out[69]: (U, (u, v)) \to (U, (x, y))
In [70]: Pdisk_cart_to_hplane_U.display()
Out[70]:
                                                               u^2+v^2+2 u+1
                                Let us use the above formula to define the transition map (u, v) \to (x, y) on the whole manifold \mathbb{H}^2
                                (and not only on U):
In [71]: Pdisk_cart_to_hplane = X_Pdisk_cart.transition_map(X_hplane, Pdisk_cart
                                  _to_hplane_U(u,v))
                                Pdisk_cart_to_hplane
Out[71]: (\mathbb{H}^2, (u, v)) \to (\mathbb{H}^2, (x, y))
In [72]: Pdisk_cart_to_hplane.display()
Out[72]:
```

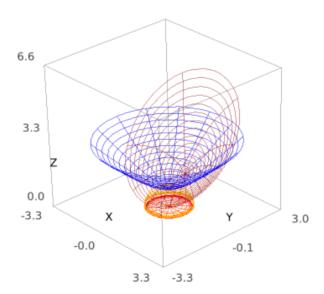
Out[73]:
$$\begin{cases} u = -\frac{x^2 + y^2 - 4}{x^2 + (y+2)^2} \\ v = \frac{4x}{x^2 + (y+2)^2} \end{cases}$$

Since the coordinates (x, y) correspond to (Y, Z) in the plane X = 1, the embedding of \mathbb{H}^2 in \mathbb{R}^3 naturally associated with the Poincaré half-plane model is

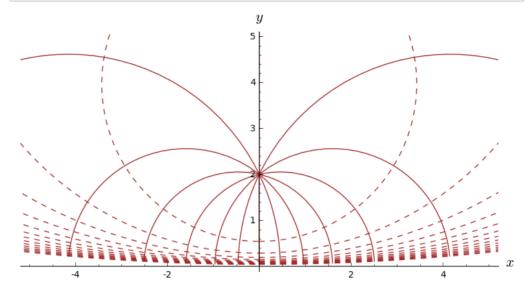
Out[74]:
$$\Phi_4$$
: $\mathbb{H}^2 \longrightarrow \mathbb{R}^3$

$$(u,v) \longmapsto (X,Y,Z) = \left(1, \frac{4v}{u^2+v^2+2u+1}, -\frac{2(u^2+v^2-1)}{u^2+v^2+2u+1}\right)$$

$$(x,y) \longmapsto (X,Y,Z) = (1,x,y)$$



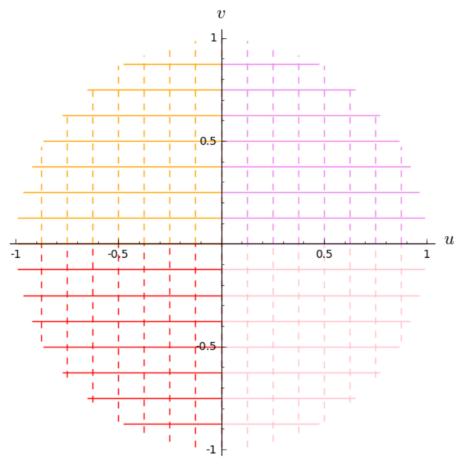
Let us draw the grid of the hyperboloidal coordinates (r, φ) in terms of the half-plane coordinates (x, y):

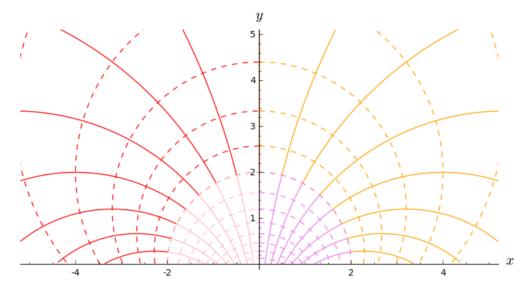


The solid curves are those along which r varies while φ is kept constant. Conversely, the dashed curves are those along which φ varies, while r is kept constant. We notice that the former curves are arcs of circles orthogonal to the half-plane boundary y=0, hence they are geodesics of (\mathbb{H}^2,g) . This is not surprising since they correspond to the intersections of the hyperboloid with planes through the origin (namely the plane $\varphi=\mathrm{const}$). The point (x,y)=(0,2) corresponds to r=0.

We may also depict the Poincaré disk coordinates (u, v) in terms of the half-plane coordinates (x, y):







The expression of the metric tensor in the half-plane coordinates (x, y) is

In [80]: g.display(X_hplane.frame(), X_hplane)

Out[80]:
$$g = \frac{1}{v^2} dx \otimes dx + \frac{1}{v^2} dy \otimes dy$$

Summary

9 charts have been defined on \mathbb{H}^2 :

```
In [81]: H2.atlas()  \begin{bmatrix} \left(\mathbb{H}^2, (X,Y)\right), \left(U, (X,Y)\right), \left(U, (r,\varphi)\right), \left(U, (R,\varphi)\right), \left(\mathbb{H}^2, (u,v)\right), \left(U, (u,v)\right), \\ \left(U, (\theta,\varphi)\right), \left(\mathbb{H}^2, (x,y)\right), \left(U, (x,y)\right) \end{bmatrix}
```

There are actually 6 main charts, the other ones being subcharts:

```
In [82]: H2.top\_charts()
Out[82]: \left[\left(\mathbb{H}^2,(X,Y)\right),\left(U,(r,\varphi)\right),\left(U,(R,\varphi)\right),\left(\mathbb{H}^2,(u,v)\right),\left(U,(\theta,\varphi)\right),\left(\mathbb{H}^2,(x,y)\right)\right]
```

The expression of the metric tensor in each of these charts is

In [83]: $\begin{aligned} & \text{for chart in H2.top_charts():} \\ & \text{show}(\texttt{g.display(chart.frame(), chart)}) \end{aligned} \\ & g = \left(\frac{Y^2 + 1}{X^2 + Y^2 + 1}\right) \mathrm{d}X \otimes \mathrm{d}X + \left(-\frac{XY}{X^2 + Y^2 + 1}\right) \mathrm{d}X \otimes \mathrm{d}Y \\ & + \left(-\frac{XY}{X^2 + Y^2 + 1}\right) \mathrm{d}Y \otimes \mathrm{d}X + \left(\frac{X^2 + 1}{X^2 + Y^2 + 1}\right) \mathrm{d}Y \otimes \mathrm{d}Y \end{aligned} \\ & g = \left(\frac{1}{r^2 + 1}\right) \mathrm{d}r \otimes \mathrm{d}r + r^2 \mathrm{d}\varphi \otimes \mathrm{d}\varphi \\ & g = \frac{4}{(R+1)^2(R-1)^2} \mathrm{d}R \otimes \mathrm{d}R + \frac{4R^2}{(R+1)^2(R-1)^2} \mathrm{d}\varphi \otimes \mathrm{d}\varphi \end{aligned} \\ & g = \frac{4}{\left(u^2 + v^2 - 1\right)^2} \mathrm{d}u \otimes \mathrm{d}u + \frac{4}{\left(u^2 + v^2 - 1\right)^2} \mathrm{d}v \otimes \mathrm{d}v \end{aligned} \\ & g = \frac{1}{\cos\left(\theta\right)^2} \mathrm{d}\theta \otimes \mathrm{d}\theta + \frac{\sin\left(\theta\right)^2}{\cos\left(\theta\right)^2} \mathrm{d}\varphi \otimes \mathrm{d}\varphi \end{aligned} \\ & g = \frac{1}{v^2} \mathrm{d}x \otimes \mathrm{d}x + \frac{1}{v^2} \mathrm{d}y \otimes \mathrm{d}y \end{aligned}$

For each of these charts, the non-vanishing (and non-redundant w.r.t. the symmetry on the last 2 indices) **Christoffel symbols of** g are