

## APPENDIX D:

### ANALYTICAL SOLUTION OF THE TWO BODY PROBLEM (KEPLERIAN MOTION)

An attractive general solution of the two-body problem, whose equation of motion is

$$\ddot{\mathbf{r}} = -\mu \mathbf{r}/r^3 \quad (\text{D.1})$$

(with  $\mu = G(m_1 + m_2)$ , the gravitational mass constant) can be obtained in the elegant form

$$\mathbf{r}(t) = f(t)\mathbf{r}_0 + g(t)\dot{\mathbf{r}}_0 \quad (\text{D.2})$$

$$\dot{\mathbf{r}}(t) = \dot{f}(t)\mathbf{r}_0 + \dot{g}(t)\dot{\mathbf{r}}_0 \quad (\text{D.3})$$

These equations reflect the following truth: since the motion is planar, coefficients  $f$  and  $g$  must exist which map the initial conditions  $(\mathbf{r}_0, \dot{\mathbf{r}}_0)$  onto the time varying position and velocity vectors  $\mathbf{r}(t)$  and  $\dot{\mathbf{r}}(t)$ . Herrick (ref. D.1), develops functional forms for  $f$  and  $g$  and several alternative solutions in detail. We summarize equations in order of solution for computation of  $\mathbf{r}(t)$  and  $\dot{\mathbf{r}}(t)$ , given  $\mathbf{r}_0, \dot{\mathbf{r}}_0$  at time  $t_0$ , for the most common case of an elliptic orbit:

#### Initial Calculations

$$\Delta t = t - t_0$$

$$r_0 = |\mathbf{r}_0| = + [\mathbf{r}_0 \cdot \mathbf{r}_0]^{1/2}$$

$$v_0 = |\dot{\mathbf{r}}_0| = + [\dot{\mathbf{r}}_0 \cdot \dot{\mathbf{r}}_0]^{1/2}$$

$$D_0 = \frac{1}{\sqrt{\mu}} \mathbf{r}_0 \cdot \dot{\mathbf{r}}_0$$

$$\frac{1}{a} = \left( \frac{2}{r_0} - \frac{v_0^2}{\mu} \right)$$

$a$  is the orbit semi major axis.

$$c_1 = e \cos E_0 = 1 - r_0/a$$

$$c_2 = e \sin E_0 = D_0/\sqrt{a}$$

$e$  is the orbit eccentricity

$$e = + [c_1^2 + c_2^2]^{1/2}$$

$E_0$  is the initial eccentric anomaly

$$\tan E_0 = c_2/c_1$$

#### Kepler's Equation

First solve the following transcendental equation (a variant of Kepler's equation) for  $\hat{E} \equiv E - E_0$  using Newton's method:

$$\sqrt{\mu} \Delta t = a^{3/2} \hat{E} - c_2 \sin \hat{E} + c_1 (1 - \cos \hat{E})$$

This usually requires three or four corrections, initiating with the estimate

$$\hat{E} \approx \sqrt{\mu} \Delta t a^{-3/2}.$$

#### Solution for $\mathbf{r}(t)$ , $\dot{\mathbf{r}}(t)$

The position and velocity at time are obtained as follows:

$$f = 1 - a(1 - \cos \hat{E})/r_0$$

$$g = \Delta t - a^{3/2}(\hat{E} - \sin \hat{E})/\sqrt{\mu}$$

$$\mathbf{r}(t) = f\mathbf{r}_0 + g\dot{\mathbf{r}}_0$$

$$r = |\mathbf{r}(t)| = [\mathbf{r} \cdot \mathbf{r}]^{1/2}$$

$$\dot{f} = -(\sqrt{\mu}a/r_0)\sin \hat{E}/r$$

$$\dot{g} = 1 - a(1 - \cos \hat{E})/r$$

$$\dot{\mathbf{r}}(t) = \dot{f}\mathbf{r}_0 + \dot{g}\dot{\mathbf{r}}_0$$

The above solution summary is a modification of a similar solution developed in detail by Herrick in reference D.1.

#### REFERENCE

- D.1 S. Herrick, **Astrodynamics**, Vol. I, Van Nostrand Reinhold, Co., London, 1971, pp. 210-212.