

A Simulation Model for Passive Attitude Control Systems Used in CubeSats

A Simulation Model for Passive Attitude Control Systems Used in CubeSats SpaceLab, Universidade Federal de Santa Catarina, Florianópolis - Brazil

A Simulation Model for Passive Attitude Control Systems Used in CubeSats February, 2021

Project Chief: Eduardo Augusto Bezerra

Authors:

Matheus Wagner

Contributing Authors:

Gabriel M. Marcelino

Revision Control:

Version	Author	Changes	Date
0.0	Gabriel M. Marcelino	Document creation	2021/02/04
0.1	Matheus Wagner	Document Writing	2021/02/14



© 2021 by SpaceLab. A Simulation Model for Passive Attitude Control Systems Used in CubeSats. This work is licensed under the Creative Commons Attribution–ShareAlike 4.0 International License. To view a copy of this license, visit http://creativecommons.org/licenses/by-sa/4.0/.

List of Figures

List of Tables

Contents

Lis	st of Figures	V
Lis	st of Tables	vii
Ν	omenclature	vii
1	Objective	1
2	Geomagnetic Field Model 2.1 Introduction	
3	Orbital Motion Model 3.1 Introduction	7 7 7
Re	eferences	9

CHAPTER 1

Objective

The objective of this document is to describe the components of a simulation model for a nanosatellite's passive attitude control system, derive the related equations and provide a software architecture for the integration of such components.

A simulation of this system must capture the phenomena related to orbital motion, attitude dynamics and the interaction between the geomagnetic field and the magnetic materials used for attitude control, hence the components of such simulation are chosen to be independent models of each phenomena containing some physically motivated interface between them, allowing for the integration of the components to provide the final model.

Geomagnetic Field Model

2.1 Introduction

The International Geomagnetic Reference Field (IGRF) is a standard mathematical description of the Earth's magnetic field widely used in studies of the magnetosphere. The model is developed and maintained by the International Association of Geomagnetism and Aeronomy (IAGA) and it's going to be the model used in this simulation.

The objective of this section is to derive a mathematical expression that defines a magnetic flux density field over a Cartesian space described in spherical coordinates, hence by the end of the derivation it must be possible to assign a vector, representing the local magnetic field, to every point of an orbit around the Earth.

2.2 Model Description

Assuming that there are no sources of magnetic field, that is, free currents, outside the Earth's surface, a scalar potential description of the magnetic field becomes feasible, since the field becomes irrotational and the existence of the potential field can be proved given the existence of a solution to Poisson's equation. Given the linearity of free-space, a well known solution to Poisson's equation can be found by the expansion of Green's function in terms of spherical harmonics, but since the source of magnetic potential, mathematically found by deriving the Poisson equation for the problem, cannot be fully known for the case of the geomagnetic field, a statistical approach is used to find the coefficients of the spherical harmonics expansion, based on measurement data, providing an approximated model for the geomagnetic field. The final equation for the model, described in [1], is:

$$V(r,\theta,\phi,t) = a \sum_{n=1}^{N} \sum_{m=0}^{n} \left(\frac{a}{r}\right)^{n+1} \left[g_{n}^{m}(t)\cos(m\phi) + h_{n}^{m}(t)\sin(m\phi)\right] P_{n}^{m}(\cos\theta)$$
 (2.1)

Where a is the Earth's radius, g_n^m , h_n^m are named Gauss Coefficients and P_n^m is the Schmidt Quasi-Normalized Associated Legendre Functions of Degree n and Order m, deeply described in [2].

Is this model, the convention used for the spherical coordinate system defines r as the radial distance from the center of the Earth, θ as the co-latitude and ϕ as the east longitude.

Based on the equation 2.1, it is possible to compute the magnetic flux density through the expression:

$$\vec{B} = -\nabla V \tag{2.2}$$

In this equation, the gradient operator is described in spherical coordinates, hence the components of the magnetic flux density vector are given by:

$$B_{r} = -\frac{\partial V}{\partial r}$$

$$B_{\theta} = -\frac{1}{r} \frac{\partial V}{\partial \theta}$$

$$B_{\phi} = -\frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi}$$

An explicit representation for the magnetic flux density vector with regard to it's basis vectors is given by:

$$\vec{B} = B_r \hat{r} + B_\theta \hat{\theta} + B_\phi \hat{\phi} \tag{2.3}$$

Where \hat{r} , $\hat{\theta}$ and $\hat{\phi}$ are unit vectors forming a basis for a curvilinear space.

Since vectors are invariant to coordinate transformations (see [3] for a detailed explanation), a direct consequence of the fact that physical quantities must not change when it's mathematical representation changes, the use of the spherical coordinates representation of the magnetic flux density field should not be a problem, as long as the basis vectors are properly defined.

A definition for the basis vectors can be constructed based on the relation between Cartesian coordinates and spherical coordinates, given by:

$$x = rsin\theta cos\phi$$
$$y = rsin\theta sin\phi$$
$$z = rcos\theta$$

To arrive at a basis all it takes is to differentiate the vector (x,y,z) with respect to (r,ϕ,θ) , arriving at a Jacobian matrix. The Jacobian matrix represents the amount of deformation caused by a transformation from Cartesian to spherical coordinates, establishing a relation between the systems. By simple differentiation rules, one gets:

$$\vec{R} = (x, y, z)$$

$$\frac{\partial R}{\partial r} = (\sin\theta\cos\phi, \sin\theta\sin\phi, \cos\theta)$$

$$\frac{\partial R}{\partial \theta} = (r\cos\theta\cos\phi, r\cos\theta\sin\phi, -r\sin\theta)$$

$$\frac{\partial R}{\partial \phi} = (-r\sin\theta\sin\phi, r\sin\theta\cos\phi, 0)$$

After normalization:

$$\hat{r} = \frac{\partial R}{\partial r} = (\sin\theta\cos\phi, \sin\theta\sin\phi, \cos\theta)$$

$$\hat{\theta} = \frac{1}{r}\frac{\partial R}{\partial \theta} = (\cos\theta\cos\phi, \cos\theta\sin\phi, -\sin\theta)$$

$$\hat{\phi} = \frac{1}{r\sin\theta}\frac{\partial R}{\partial \phi} = (-\sin\phi, \cos\phi, 0)$$

By using such basis, it is possible to verify the relation:

$$\vec{B} = B_x \hat{x} + B_y \hat{y} + B_z \hat{z} = B_r \hat{r} + B_\theta \hat{\theta} + B_\phi \hat{\phi}$$
 (2.4)

And the spherical coordinates representation of vector field can be used directly in the computation of the magnetic flux density at every point of an orbit around the Earth.

CHAPTER 3

Orbital Motion Model

3.1 Introduction

Since most LEO satellite orbits are near-circular, for this analysis a circular orbit will be assumed, based on the orbital motion model used in [4].

The objective of this model is to associate every instant in time to a position contained in a feasible orbit around the Earth.

3.2 Model Description

The mathematical description of a circular orbit is defined by five parameters: three angles that define the orbital plane with respect to the xy plane, through successive rotations; a radius, defining the distance between the Earth's center to the satellite's center of mass; and an angle that defines the angular position of the satellite's center of mass in the orbit at an instant. Note that since the orbit is approximated as a circle, both the radius and the rotation of the orbital plane are constant, implying that only the angular position in the orbit changes with time.

The construction of the proposed representation of orbital motion can be realized by successive rotations of a circular orbit in the xy plane by a set of angles α , β , γ . Hence any point P contained in the orbit can be computed using the relation:

$$P = R_{\alpha\beta\gamma}P_0 \tag{3.1}$$

Where:

$$R_{\alpha\beta\gamma} = R_{\alpha}R_{\beta}R_{\gamma} \tag{3.2}$$

$$R_{\alpha} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\alpha) & -\sin(\alpha) \\ 0 & \sin(\alpha) & \cos(\alpha) \end{bmatrix}$$

$$\begin{bmatrix} \cos(\beta) & 0 & \sin(\beta) \end{bmatrix}$$
(3.2)

$$R_{\beta} = \begin{bmatrix} \cos(\beta) & 0 & \sin(\beta) \\ 0 & 1 & 0 \\ -\sin(\beta) & \cos(\beta) \end{bmatrix}$$
(3.4)

$$R_{\gamma} = \begin{bmatrix} \cos(\gamma) & -\sin(\gamma & 0) \\ \sin(\gamma) & \cos(\gamma) & 0 \\ 0 & 0 & 1 \end{bmatrix}$$
(3.5)

(3.6)

Note that matrix multiplication is not a commutative operation, hence the order of rotations, $\alpha-\beta-\gamma$, must be respected.

Since $P_0(\theta)$ is a point contained in a circular orbit defined in the xy plane, a convenient parameterization for the curve is:

$$P_{0} = \begin{bmatrix} x \\ y \\ z \end{bmatrix} = \begin{bmatrix} r\cos\theta \\ r\sin\theta \\ 0 \end{bmatrix} \tag{3.7}$$

Where r is the constant radius, and theta is the angular position parameter.

Finally, for the link between the orbital position and time to be established, an expression for the angular position as a function of time must be defined. Assuming that the there is no energy dissipation in the orbital motion, implying that the kinetic energy of the satellite's center of mass is constant, then the angular velocity of such motion is constant and the angular position evolves linearly in time. Mathematically:

$$\theta(t) = \omega_o t = \sqrt{\frac{Gm}{(R_E + h)^3}} t \tag{3.8}$$

Where G is the gravitational constant, m is Earth's mass, R_E is Earth's surface average radius and h is the altitude, using the Earth's surface average radius as reference, as established in [4].

Bibliography

- [1] E. Thébault et al. International geomagnetic reference field: the 12th generation. *Earth, Planets and Space*, 2015.
- [2] D. Winch et al. Geomagnetism and schmidt quasi-normalization. *Geophysical Journal International*, 2005.
- [3] A. Hazel. Chapter 1 Describing the Physical World: Vectors & Tensors.
- [4] C. A. Rigo. Task scheduling for optimal power management and quality-of-service assurance in CubeSats, 2021.