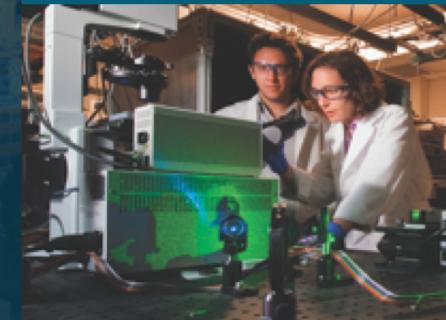


Guest Lecture Stanford ME469: Common low-Mach Discretization Approaches



Sandia
National
Laboratories



PRESENTED BY

Stefan P. Domino

Computational Thermal and Fluid Mechanics

Sandia National Laboratories SAND2018-4536 PE



Sandia National Laboratories is a multimission laboratory managed and operated by National Technology & Engineering Solutions of Sandia, LLC, a wholly owned subsidiary of Honeywell International Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA0003525.

Common low-Mach Discretization Approaches: Outline

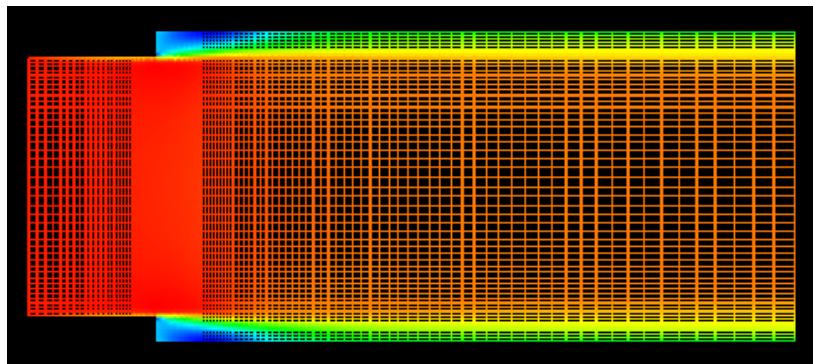


- Why Unstructured?
- Unstructured Element Typoes
- Model PDE for Discretization
- Finite Element Method (FEM)
- Control-Volume Finite Element Method (CVFEM)
- Edge-based Vertex-Centered (EBVC)
- Cell-centered Finite Volume (FV)
- Staggered arrangement
- Conclusions

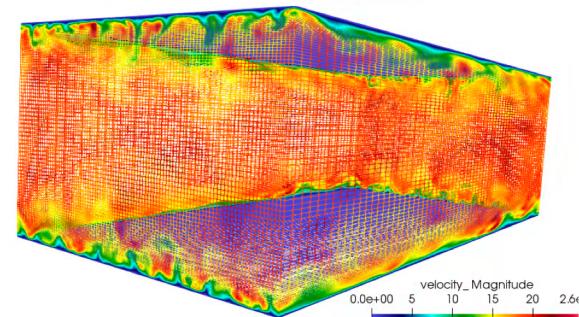
Structured vs Unstructured



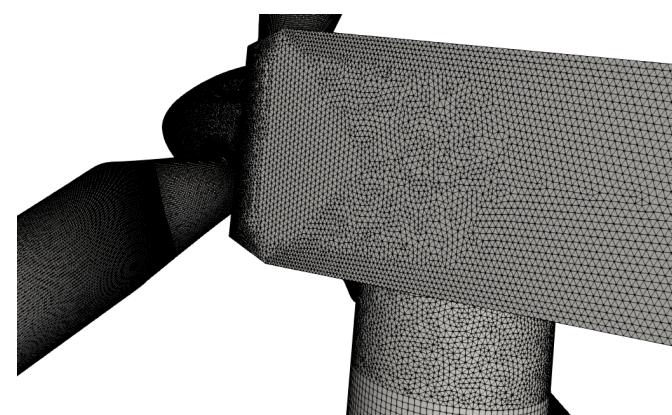
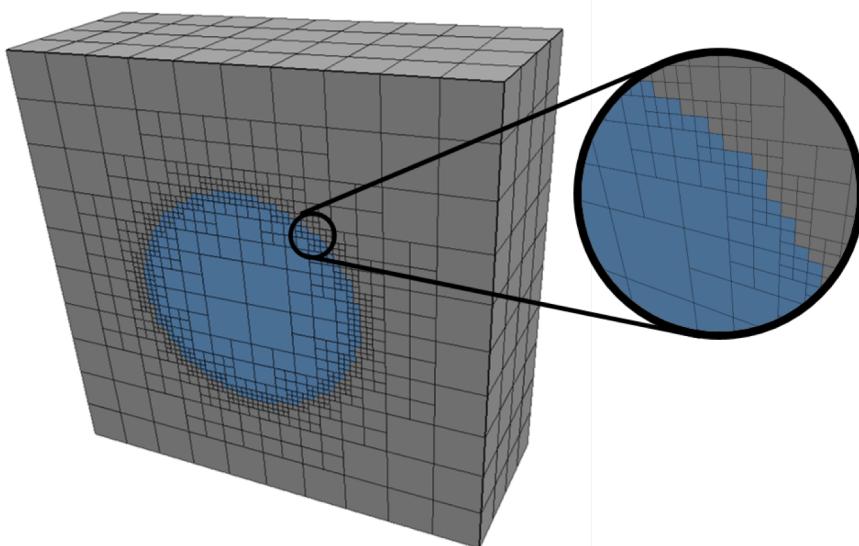
- Many times, the canonical flows of interest are in simplified geometries that allow for cartesian meshes – with “stair-stepping”



RANS-based backward facing step (Domino, 2012)



Re τ 395 plane-channel (Jofre, Domino, Iaccarino, 2018)



Often times, not!

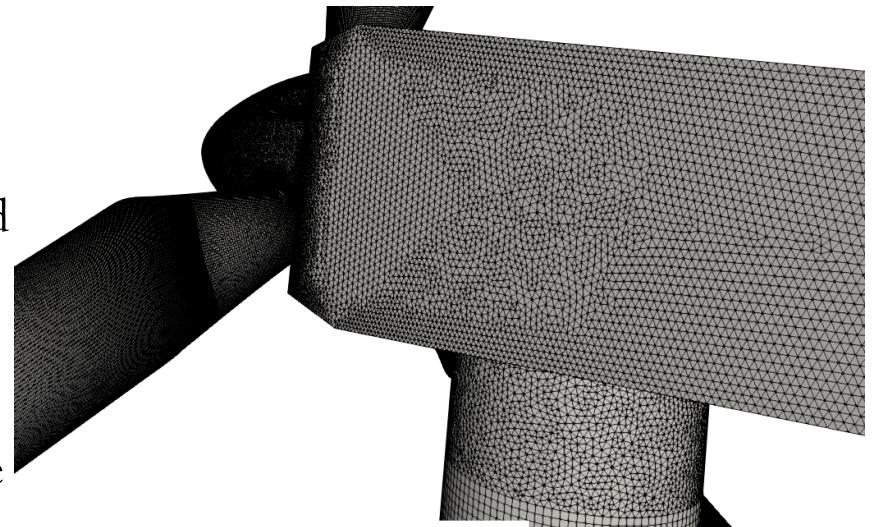
Reality: Meshing time for complex applications remains a significant bottleneck!



- Many applications of interest to SNL contain complex geometries
- low-Mach fluids users interested in high-quality simulation results tend towards hexahedral-based topologies (if possible)
- However, if a scheme is “design-order” accurate, any topology may suffice as it is simply a matter of mesh size and efficiency – not unlike the active discussion on low- vs higher-order
- Sometimes, the penetration of a low-Mach fluids physics addition in common analysis is high as the meshing can be prohibitively complex



Very complex world - stair-stepped!

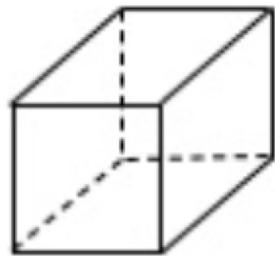


UUR
Example:
Vestas V27
225 kw
hybrid low-
order
hex/tet/pyr
/wedge

Examples of Various Topologies



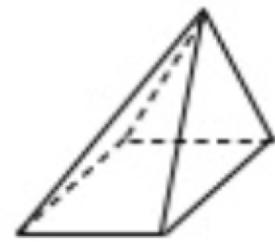
Hex8



Tet4



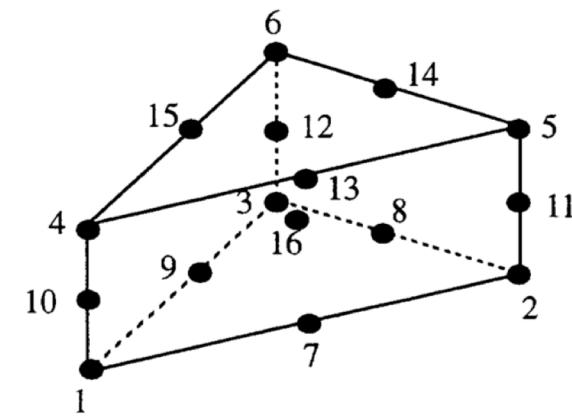
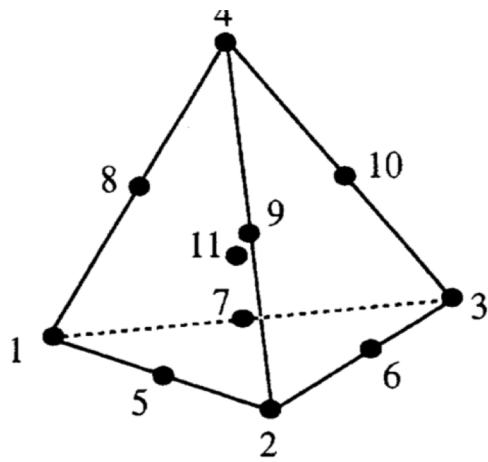
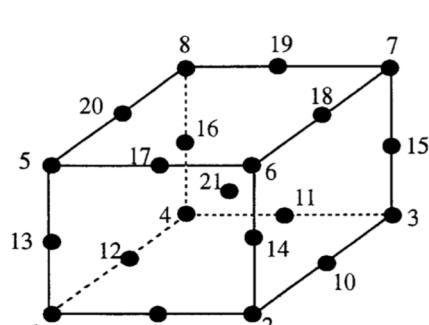
Pyramid5



Wedge6



Arbitrary

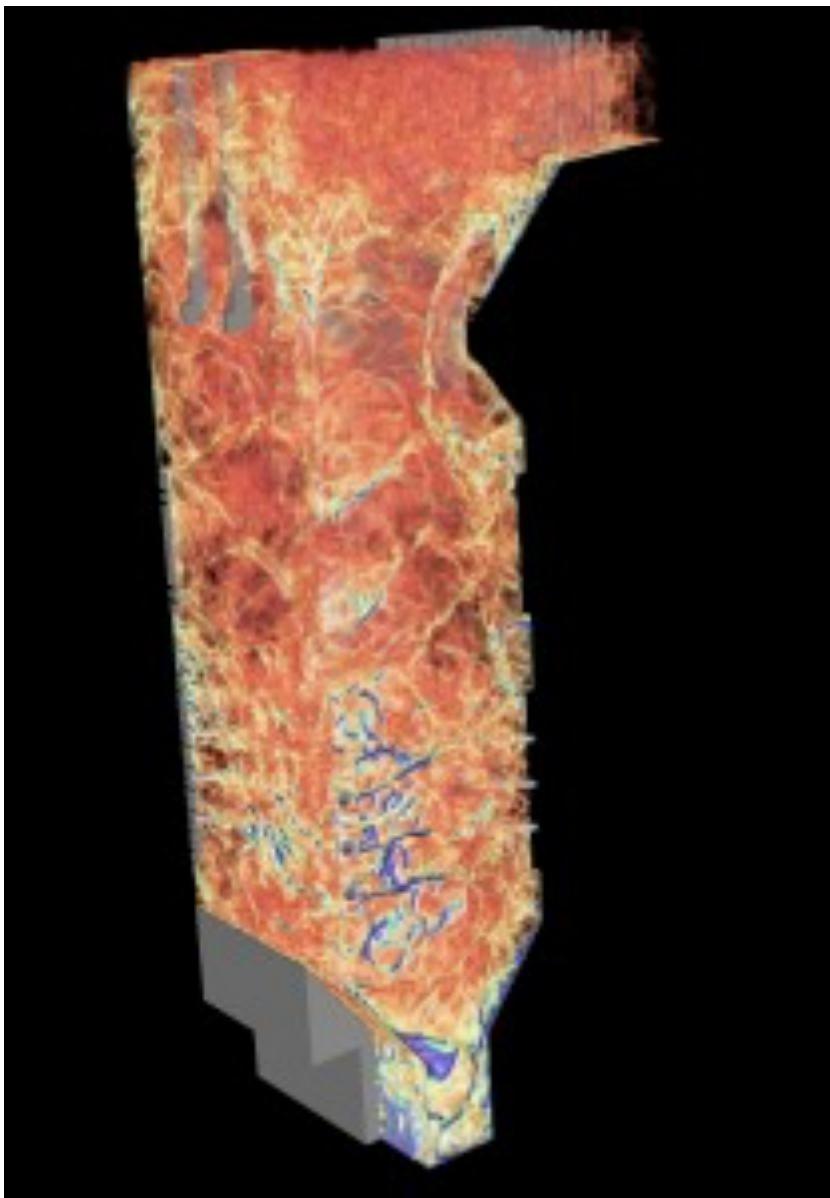


Higher-order promoted elements (Hex27, Tet10, Wedge16, Hex64, etc.)

In Fairness....

6

The Carbon-Capture Multidisciplinary Simulation Center



15MW coal-fired boiler volume rendered image of large ($90 \mu\text{m}$) particles

Staggered schemes have been demonstrated to support complex applications

Cut-cells and embedded approaches help

Fundamentals of Discretization



- Given a partial differential equation (PDE) and associated volumetric form:

$$\rho C_p \frac{\partial T}{\partial t} - \frac{\partial}{\partial x_j} \lambda \frac{\partial T}{\partial x_j} = 0 \quad \int \rho C_p \frac{\partial T}{\partial t} dV - \int \frac{\partial}{\partial x_j} \lambda \frac{\partial T}{\partial x_j} dV = 0$$

- Applying Gauss Divergence provides the standard finite volume form:

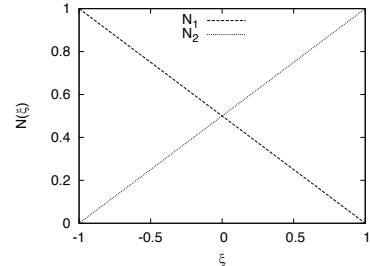
$$\int \frac{\partial q_j}{\partial x_j} dV = \int q_j n_j dS \quad \int \rho C_p \frac{\partial T}{\partial t} dV - \int \lambda \frac{\partial T}{\partial x_j} n_j dS = 0$$

- We can also multiple PDE by an arbitrary test function, w, and integrate over a volume,

$$\int w \rho C_p \frac{\partial T}{\partial t} dV - \int w \frac{\partial}{\partial x_j} \lambda \frac{\partial T}{\partial x_j} n_j dS = 0 \quad \text{Next, integrate by parts and apply Gauss-Divergence:}$$

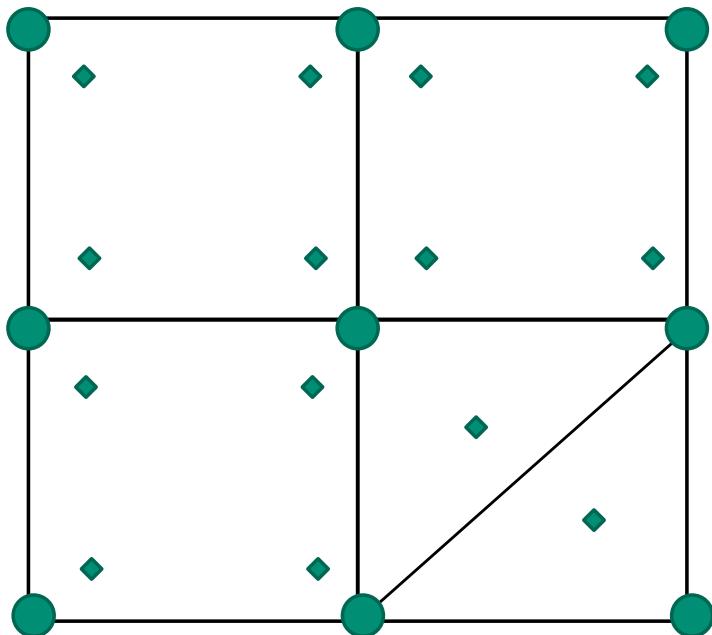
$$\int w \rho C_p \frac{\partial T}{\partial t} dV + \int \frac{\partial w}{\partial x_j} \lambda \frac{\partial T}{\partial x_j} dV - \int w \lambda \frac{\partial T}{\partial x_j} dS = 0 \quad \text{Sometimes, you see T as } T^H - \text{the trial space}$$

The Finite Element Method (FEM)



- Consider an alternative approach in which a finite element method (FEM) is employed
- Define an underlying nodal basis with the element:

$$T(x_k) \approx \sum_{i=1}^{npe} N_i(x_k) T_i \quad \frac{\partial T(x_k)}{\partial x_j} \approx \sum_{i=1}^{npe} \frac{\partial N_i(x_k)}{\partial x_j} T_i$$



- On a -1:1 range, $\pm \sqrt{3}/3$

- Consider a simple heat conduction model PDE,

$$\rho C_p \frac{\partial T}{\partial t} - \frac{\partial}{\partial x_j} \lambda \frac{\partial T}{\partial x_j} = 0$$

- Integrate using a test function,

$$\int w \rho C_p \frac{\partial T}{\partial t} dV - \int w \frac{\partial}{\partial x_j} \lambda \frac{\partial T}{\partial x_j} dV = 0$$

- Integration-by-parts (with G-D) provides:

$$\int w \rho C_p \frac{\partial T}{\partial t} dV + \int \frac{\partial w}{\partial x_j} \lambda \frac{\partial T}{\partial x_j} dV - \int w \lambda \frac{\partial T}{\partial x_j} n_j dS = 0$$

- Iterate element quadrature points
- Note that N can be arbitrary in order (shown here for a linear)

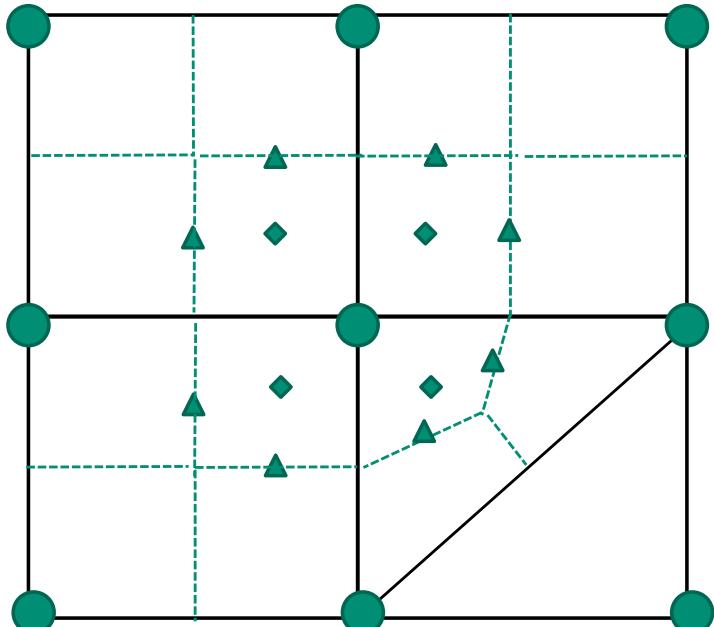
The Hybrid Control-Volume Finite Element Method (CVFEM)



- A combination between the edge-based vertex-centered and FEM is the method known as Control Volume Finite Element ()
- A dual mesh is constructed to obtain flux and volume quadrature locations

- As with FEM, a basis is defined:

$$T(x_k) \approx \sum_{i=1}^{npe} N_i(x_k) T_i \quad \frac{\partial T(x_k)}{\partial x_j} \approx \sum_{i=1}^{npe} \frac{\partial N_i(x_k)}{\partial x_j} T_i$$



Dual-volume definition

- Integration-by-parts over test function w:

$$\int w \rho C_p \frac{\partial T}{\partial t} dV + \int \frac{\partial w}{\partial x_j} \lambda \frac{\partial T}{\partial x_j} dV - \int w \lambda \frac{\partial T}{\partial x_j} n_j dS = 0$$

- However, define a test function, w, as a piece-wise constant function (Heavyside) to be 1 inside the dual volume and 0 outside. Gradient is a Dirac-delta function:

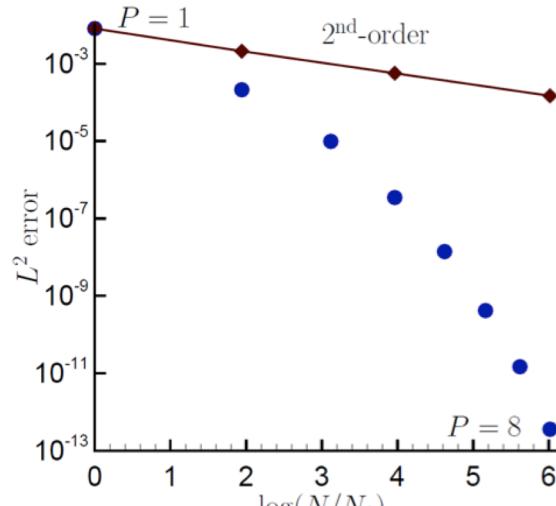
$$\frac{\partial w}{\partial x_j} = -n_j \delta(x_j - xIP_j)$$

- Leading to: $\int \rho C_p \frac{\partial T}{\partial t} dV - \int \lambda \frac{\partial T}{\partial x_j} n_j dS = 0$

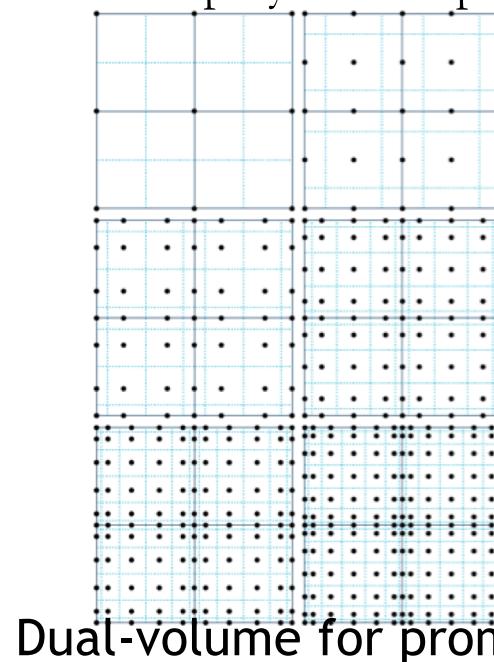
Control-Volume Finite Element Method Attributes



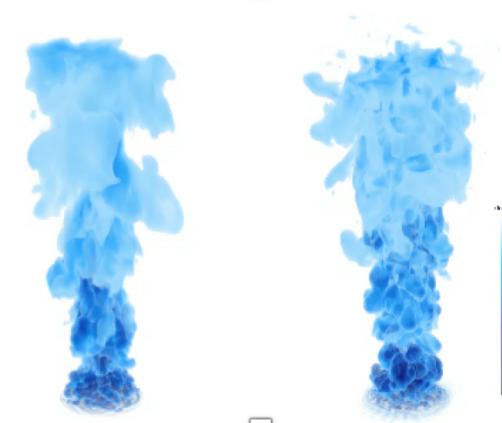
- The CVFEM method, therefore, is a finite volume scheme that is locally conservative, i.e., momentum leaving one dual volume face enters the adjacent dual volume
- However, the gradient operator, like its FEM counterpart is absent of any error due to non-orthogonality
- Since the test-function, w , is different from the underlying basis representation, this method can be considered a Petrov-Galerkin method
- The method can also be promoted in polynomial space



Spectral convergence



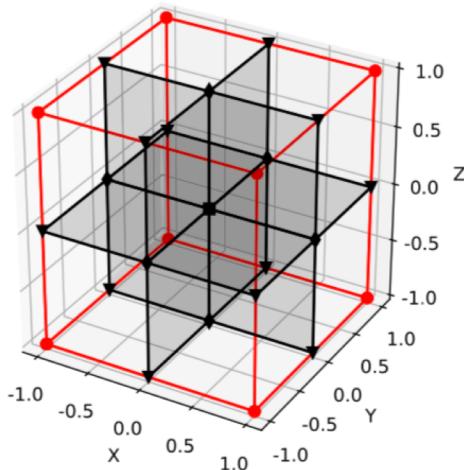
Dual-volume for promoted quad4



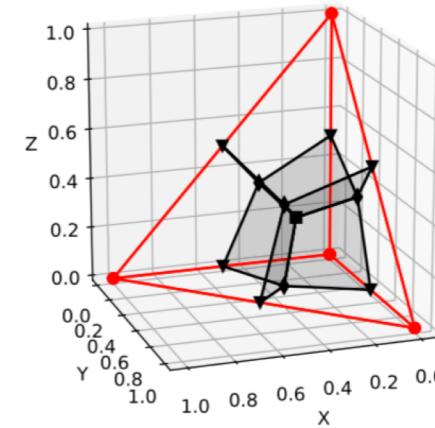
$P=1$ (left) and $P=4$ (right)
Helium plume (VR-density)



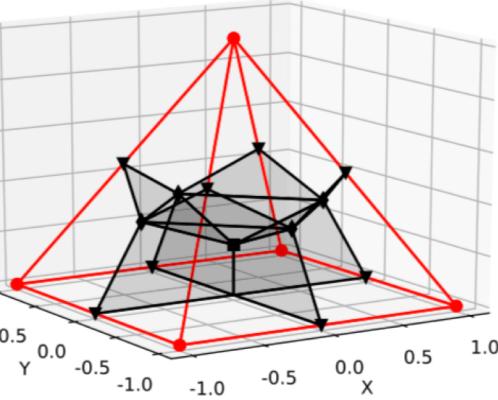
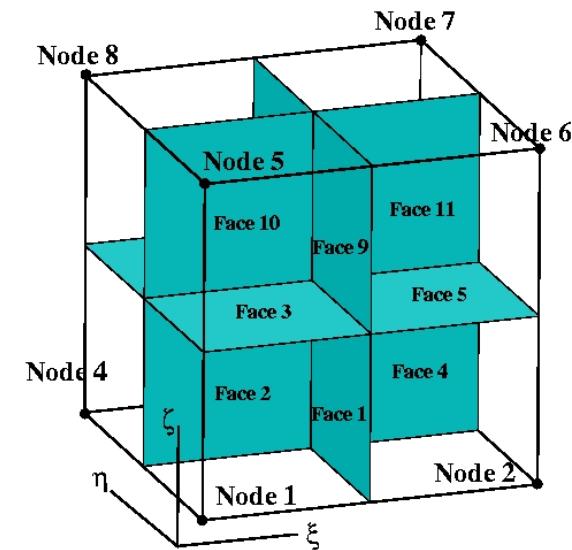
Dual Volume Definitions for Hybrid Meshes



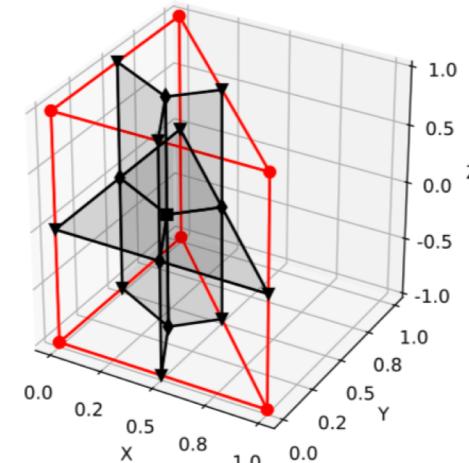
(a) Hexahedral topology (Hex8).



(b) Tetrahedral topology (Tet4).



(c) Pyramid topology (Pyramid5).



(d) Wedge topology (Wedge6).

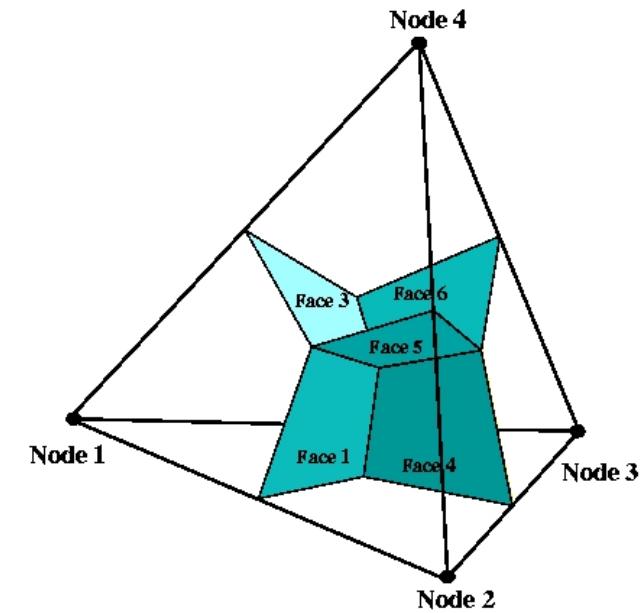


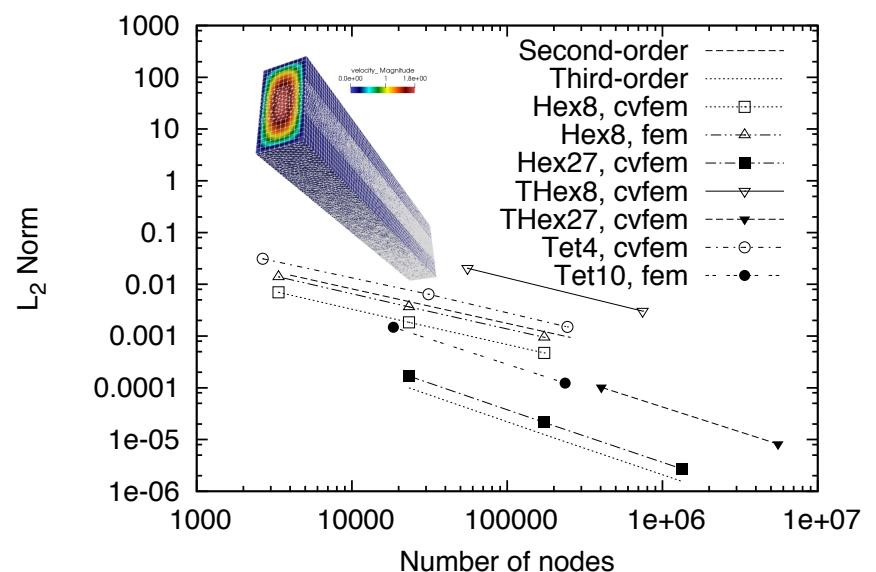
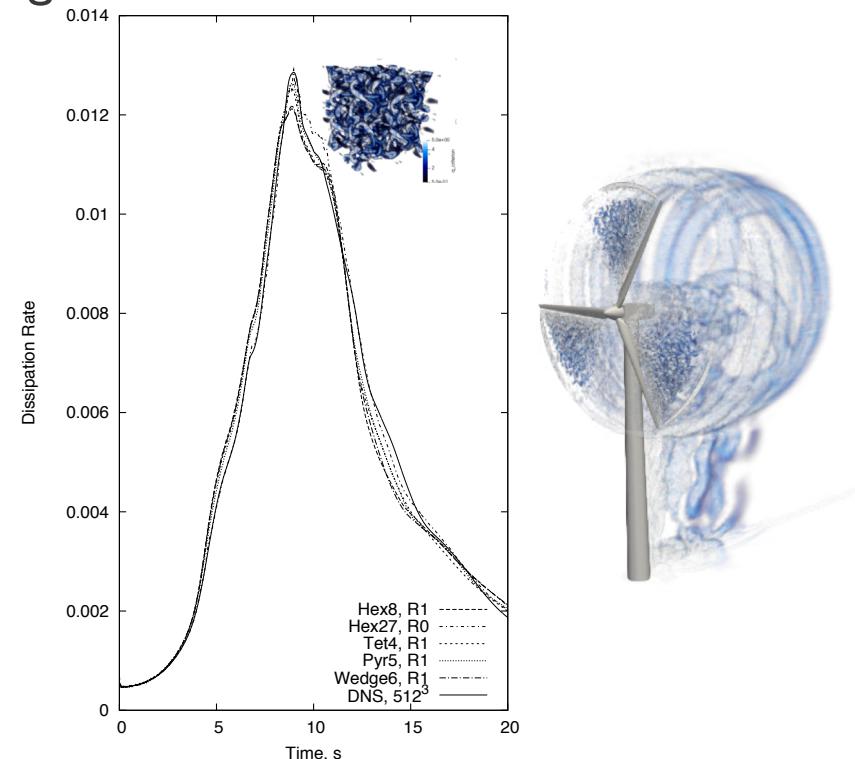
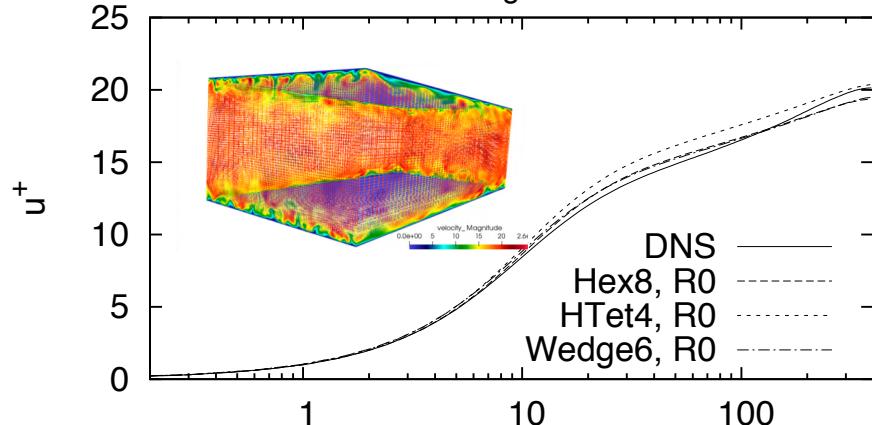
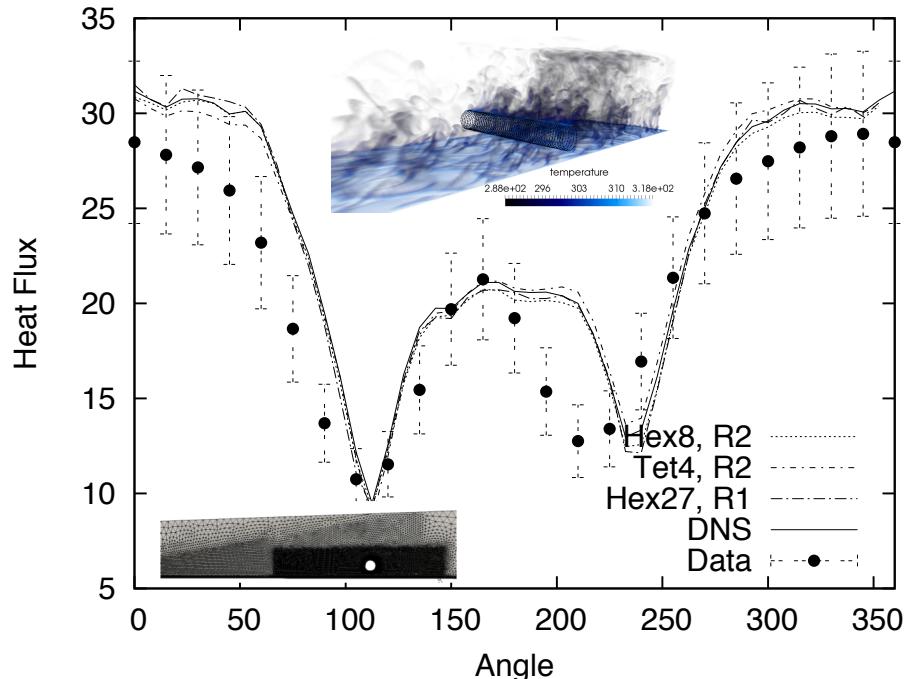
Fig. 1. CVFEM element and dual-volume definition for the low-order topologies.

Domino, et. al, “An assessment of atypical mesh topologies for low-Mach large-eddy simulation” 2019

Recent Generalized Unstructured Findings



- Domino, et. al, “An assessment of atypical mesh topologies for low-Mach large-eddy simulation”
2019

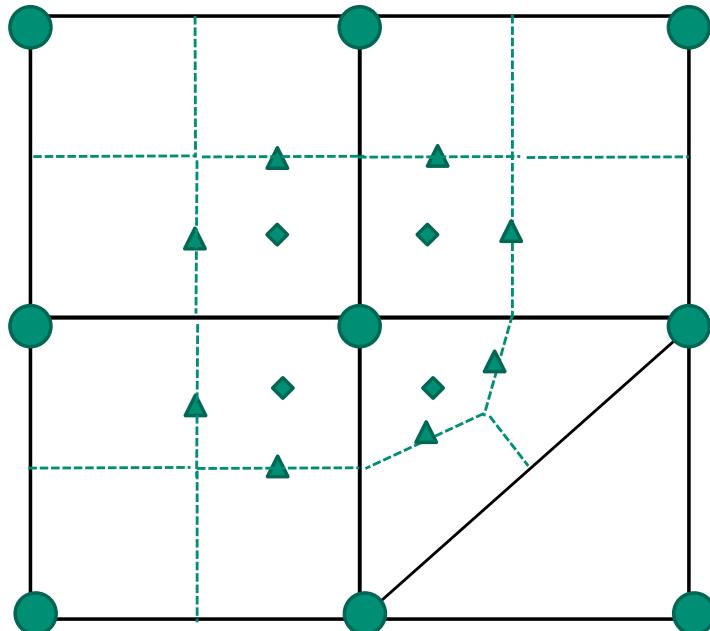


Equal Order Interpolation Edge-Based Vertex-Centered (EBVC) Finite Volume

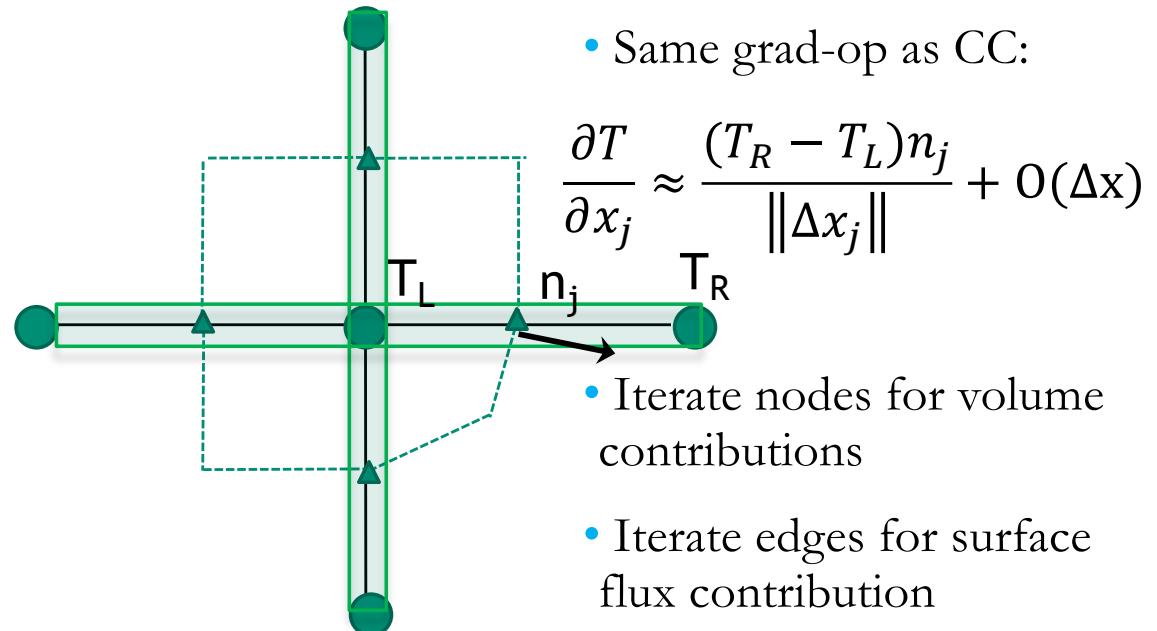
13



- All primitives are collocated at the vertices of the elements with equal-order interpolation
- A dual mesh is constructed to obtain flux and volume quadrature locations
- Classic two-state, “L” and “R” approach provides spatially second-order accuracy
 - ▲ Surface quadrature point (area summed to edge)
 - ◆ Volume quadrature point (sub-vol summed to node)



Dual-volume definition



Edge-based stencil

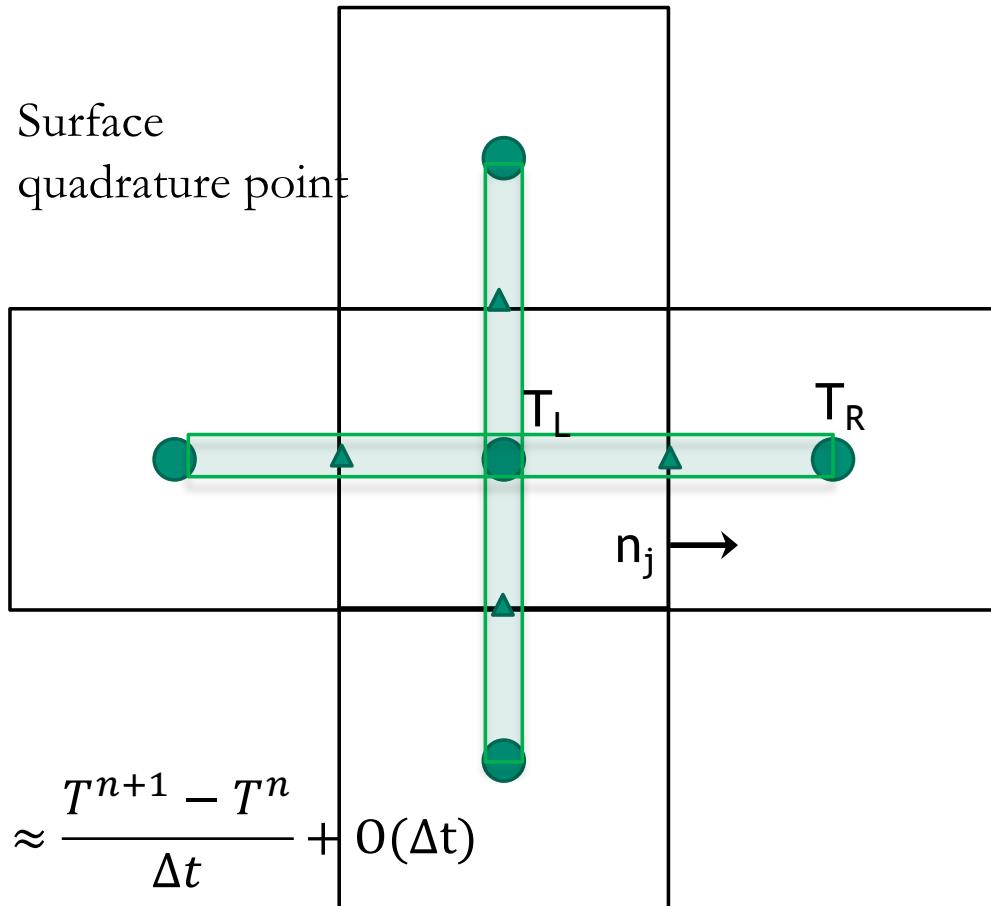
Equal-Order Interpolation Cell-Centered (CC) Finite Volume



- All primitives are collocated at the cell-center of the element with equal-order interpolation
- Classic two-state, “L” and “R” approach provides spatially second-order accuracy



Surface quadrature point



- Consider a simple heat conduction model PDE,

$$\rho C_p \frac{\partial T}{\partial t} - \frac{\partial}{\partial x_j} \lambda \frac{\partial T}{\partial x_j} = 0$$

- Integrate over a control volume and use Gauss-Divergence,

$$\int \rho C_p \frac{\partial T}{\partial t} dV - \int \lambda \frac{\partial T}{\partial x_j} n_j dS = 0$$

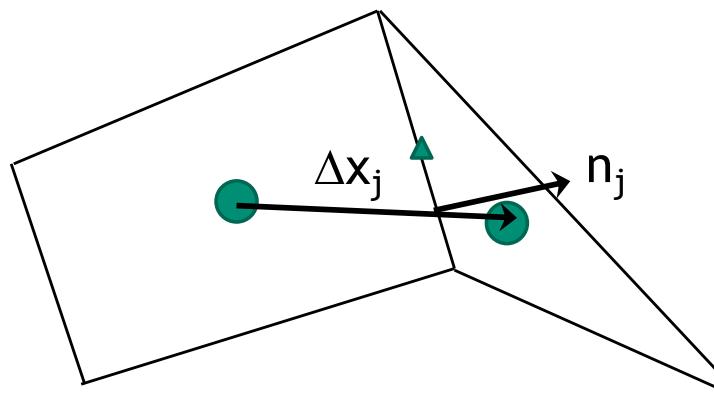
$$\text{with: } \frac{\partial T}{\partial x_j} \approx \frac{(T_R - T_L)n_j}{\|\Delta x_j\|} + O(\Delta x)$$

- Iterate element cell-centers for volume contributions (time/source)
- Iterate element faces for surface flux contribution

Typical Failings for Two-State Discretization Methods



- With two points, only a linear basis can be used.
- Therefore, unstructured CC and EBVC are limited to second-order spatial accuracy
- Non-orthogonality is problematic for gradient-operator



$$\frac{\partial T}{\partial x_j} = G_j T + [(T_R - T_L) - G_k T \Delta x_k] \frac{A_j}{A_l \Delta x_l}$$

With area vector defined by: $A_j = n_j dS$

- Above, $G_j T$ is a projected nodal gradient at the cell-center, or vertex center:
- Non-orthogonality is simply defined as the mis-alignment of the distance vector $G_j T = \frac{\int T A_j}{\int dV}$ between the two "L" and "R" states and the surface normal
- Both edge- and cell centered-based schemes show degraded accuracy on typical production meshes
- Several non-orthogonality approaches are available, for the best source, see Jasak
 - Jasak, "Error analysis and error estimation for the finite volume method with applications to fluid flow", Imperial College Dissertation, 1996

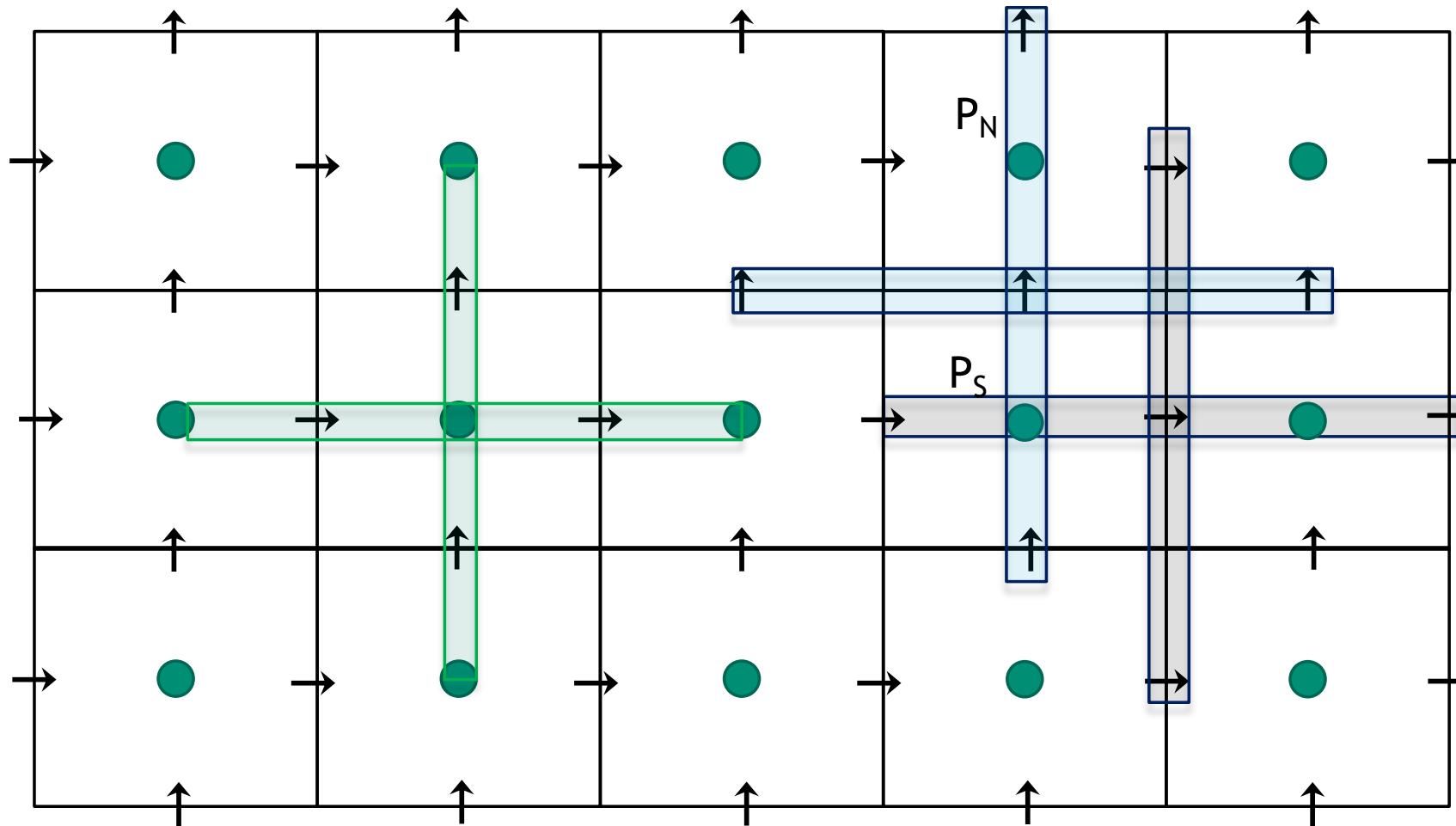
Classic Staggered Finite Volume

Stencil for CC-quantities ●

Stencil for x-velocity →

Stencil for y-velocity ↑

- Velocity degree-of-freedom is staggered relative to pressure and other primitives, e.g., enthalpy, mixture fraction, etc.



Attributes of a Staggered Scheme



- By design, non-orthogonality is absent, however, complex geometry will be stair-stepped
- From a fluids perspective, the operators are ideal, i.e., pressure gradient for momentum is compact, e.g., $(P_E - P_W)\Delta x^{-1}$
 - As will be seen in future lecture topics, the skew-adjoint nature of the Divergence operator, **D**, and Gradient operator, **G**, allows for a Laplace operator, **L** = **DG**
- Can be extended to higher-order
- Frequently, meshing complex geometries can be extremely difficult (consider our V27 example)

Common low-Mach Discretization Approaches: Conclusions



- Two-state methods, e.g., cell-centered and EBVC are attractive due to simplicity, however, suffer from non-orthogonality issues in the diffusion operator
- FEM provides a machinery to provide accurate discretizations on non-ideal meshes, however, the same diffusion operator suffers on high-aspect ratio meshes
- CVFEM is a hybrid method that contains the likeable attributes of both FV and FEM (same high-aspect ratio diffusion operator finding)
- Staggered arrangement is well suited for a class of fluid mechanics applications where low-order or simple geometries are found