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# General Relativity: Examination

(1)

(1). This is a reasonable equation.

- Reasonable equation

- Nonsense since operating with the covariant derivative twice would not lead to a contravariant  $(3,0)$  tensor.

- Reasonable equation.

- Reasonable equation.

(2)  $ds^2 = xy dx^2 + x^2 dy^2$

$$\therefore g_{\mu\nu} = \begin{pmatrix} xy & 0 \\ 0 & x^2 \end{pmatrix}$$

∴ 5 Possible Christoffel symbols:

$$\Gamma_{xx}^x, \Gamma_{xy}^x, \Gamma_{xx}^y, \Gamma_{xy}^y$$

$$\Gamma_{yy}^y, \Gamma_{yx}^y, \Gamma_{yy}^x$$

$$\Gamma_{xx}^{xx} = \frac{1}{2} \partial_x g^{xx} (\partial_x g_{xx} + \partial_y g_{xx} \partial_x g_{yy})$$

Diagonal metric  $\Rightarrow \Gamma_{xx}^y = -\frac{1}{2g_{yy}} \partial_y g_{xx}$

$$= -\frac{1}{2x^2} \partial_y (xy) = -\frac{1}{2x}$$

$$\Gamma_{xx}^x = \partial_x \ln \sqrt{|g_{xx}|}$$

$$= \partial_x \left( \ln \sqrt{xy} \right) = \frac{\left(\frac{1}{2}y\right)}{\sqrt{xy}} = \frac{1}{2y}$$

(3)

$$(2) \text{ and } ds^2 = g_{xy} dx^2 + x^2 dy^2 \text{ mit } g_{xy} \text{ ist diagonal}$$

(2)

$$\therefore g_{\mu\nu} = \begin{pmatrix} xy & 0 \\ 0 & x^2 \end{pmatrix} \quad \text{ist diagonal}$$

$\therefore$  Metric is diagonal. Christoffel symbols are:

$$\Gamma_{xx}^x, \Gamma_{xy}^x, \Gamma_{yx}^x, \Gamma_{yy}^x, \Gamma_{xx}^y, \Gamma_{xy}^y, \Gamma_{yx}^y, \Gamma_{yy}^y$$

$$\therefore \Gamma_{xx}^x = \partial_x \ln \sqrt{|g_{xx}|} = \partial_x \ln \sqrt{xy} = \frac{1}{2} - \frac{1}{2} x^{-1/2} y^{1/2}$$

$$\Gamma_{yy}^y = \partial_y \ln \sqrt{|g_{yy}|} = \partial_y \ln \sqrt{xy} = 0$$

$$\Gamma_{xy}^x = \partial_y \ln \sqrt{|g_{xx}|} = \partial_y \ln \sqrt{xy} = \frac{1}{2} y^{-1/2} \frac{1}{(xy)^{1/2}} = \frac{1}{2y}$$

$$\Gamma_{yx}^x = \frac{1}{2y} \quad (\text{symmetry}), \quad \Gamma_{yy}^x = \frac{\partial_x (g_{yy})}{2g_{xx}} = \frac{-1}{2xy} = -\frac{1}{2x}$$

$$= -\frac{1}{2}$$

$$\Gamma_{xy}^y = \frac{-1}{2g_{yy}} \partial_y (g_{xx}) = -\frac{1}{2x^2}, \quad \Gamma_{xy}^y = \Gamma_{yx}^y = \partial_x \ln \sqrt{|g_{yy}|} = \partial_x (\ln \sqrt{x^2}) = \frac{1}{2}$$

$$\therefore \Gamma_{yx}^y = \partial_x (\ln (x)) = \frac{1}{x}$$

$\therefore$  To compute  $R_{\mu\nu}^{\alpha\beta}$ , we need  $R_{\nu\mu\alpha}^{\beta}$  first:  
 (Riemann tensor):  $R_{\mu\nu}^{\alpha\beta} = R_{\mu\nu}^{\alpha\beta} - \frac{1}{2} g_{\mu\nu} g^{\alpha\beta} R$

(2) Compute  $R^x_{yxy}$  and then  $R^y_{yyx}$ . We can then obtain other components by symmetry

$$\begin{aligned} R^x_{yxy} &= \partial_x \Gamma^x_{yy} - \partial_y \Gamma^x_{yx} - \Gamma^x_{xx} \Gamma^x_{yy} + \Gamma^x_{xy} \Gamma^y_{yy} - \Gamma^x_{yx} \Gamma^x_{yy} \\ &= \partial_x \left( \frac{1}{2y} \right) = 0 - \partial_y \Gamma^x_{yx} \left( \frac{1}{2y} \right) - \\ &= -\partial_y \left( \frac{1}{2y} \right) - \left( \frac{1}{2x} \right) \left( -\frac{1}{y} \right) = \frac{1}{2y} - \frac{1}{2x} \\ &= \frac{-1}{2y^2} - \frac{1}{2xy} - \frac{1}{2y} = -\frac{2}{2xy} \\ &= \frac{1}{2y^2} - \frac{1}{2y} = \frac{x+2y}{4xy} \end{aligned}$$

Since Riemann tensor is antisymmetric in last two indices we have

$$R^x_{yyx} = -R^x_{yxy} = -\frac{x+2y}{4xy}$$

$$\begin{aligned} \text{Compute } R^y_{xxy} &= -R^y_{xyx} \\ &= R^y_{yxy} = -\frac{x+2y}{4xy} \end{aligned}$$

is identical to the above, up to  $-\partial_y \rightarrow -\partial_x \left( \frac{1}{2x} \right)$

$$R^y_{yxy} = -R^y_{xyx} = -\frac{x+2y}{4x^2y} \quad (\text{it is symmetric})$$

with  $R^x_{yyx} \times$  (SEE END OF PAGE 7)

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$R^x_{xxx}$  and  $R^y_{yyy}$  vanish by definition of Riemann tensor.

$$R_{xx} \text{ s.t. } R^x_{yyxx} = 0 \Rightarrow R^x_{xyxy} = 0$$

$$\therefore R^x_{xx} = g^{xx} (R^x_{xxxx} + R^y_{xyxx}) = g^{yy} R^y_{yyxy} = \frac{x+2y}{4x^2y}$$

$$\text{Similarly, } R_{yy} = g^{yy} R^x_{yyxy} + R^y_{yyyy} = \frac{x+2y}{4x^2y^2}$$

$$\therefore R_{xyxy} = \begin{pmatrix} \frac{x+2y}{4x^2y} & 0 \\ 0 & \frac{x+2y}{4x^2y^2} \end{pmatrix}$$

$$\therefore R^y_{xy} = g^{xy} R_{xy} \\ = \frac{1}{xy} \cdot \frac{x+2y}{4x^2y} + \frac{1}{x^2} \cdot \frac{x+2y}{4xy^2} = \frac{x+2y}{4x^3y^2}$$

$$= \frac{x+2y}{4x^3y^2} + \frac{x+2y}{4x^3y^2} = \frac{2x+4y}{4x^3y^2} = \frac{x+2y}{2x^3y^2}$$

(3) Claim:

$$\nabla_u T^\mu_v = \frac{1}{\sqrt{-g}} \partial_u (\sqrt{-g} T^\mu_v) - \Gamma^\lambda_{uv} T^\mu_\lambda$$

$$\nabla_u T^\mu_v = \partial_u T^\mu_v + \Gamma^\mu_{u\sigma} T^\sigma_v$$

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$$\begin{aligned}
 &= \partial_u T^u_v + (\sqrt{-g})^{-1} \partial_\sigma (\sqrt{-g}) T^{\sigma u} - \Gamma_{\mu\nu}^\sigma T^{\mu u} \\
 \nabla_u T^u_v &= \partial_u \nabla_v T^u + \Gamma_{\mu\nu}^\sigma \nabla_v T^{\mu u} - \Gamma_{\mu\nu}^\sigma T^{\mu u} \\
 &= \partial_u \nabla_v T^u + (\sqrt{-g} \partial_\sigma) \sqrt{-g} \nabla_v T^{\mu u} - \Gamma_{\mu\nu}^\sigma T^{\mu u} \\
 &= \partial_u T^u_v + \frac{1}{\sqrt{-g}} \partial_\sigma (\sqrt{-g} T^u_v) - \Gamma_{\mu\nu}^\sigma T^{\mu u} \Rightarrow QED
 \end{aligned}$$

(where  $\lambda = \sigma$ )

Want to now show that:

$$\nabla_u T^u_v = \frac{1}{\sqrt{-g}} \partial_u (\sqrt{-g} T^u_v) - \frac{1}{2} (\partial_v g_{\mu\nu}) T^{\mu u}$$

Case where  $\sigma = \mu \neq v$

$$= \frac{1}{\sqrt{-g}} \partial_u T^u_v + \frac{1}{\sqrt{-g}} \partial_u (\sqrt{-g} T^u_v) + \Gamma_{\mu\nu}^\mu T^{\mu u}$$

$$\therefore \Gamma_{\mu\nu}^\mu = \partial_\mu \ln(\sqrt{|g_{\mu\nu}|}) = \frac{1}{\sqrt{-g}} \frac{1}{2} g^{\mu\lambda} \frac{\partial g_{\mu\lambda}}{\partial \mu} = -\frac{1}{2} g^{\mu\lambda} \frac{\partial g_{\mu\lambda}}{\partial \mu}$$

$$= \frac{1}{2\sqrt{-g}} \frac{1}{2} g^{\mu\lambda} \frac{\partial g_{\mu\lambda}}{\partial \mu} T^{\mu u}$$

$$\therefore \nabla_u T^u_v = \frac{1}{\sqrt{-g}} \partial_u (\sqrt{-g} T^u_v) + \partial_v (\ln(\sqrt{|g_{\mu\nu}|})) T^{\mu u}$$

$$= \frac{1}{\sqrt{-g}} \partial_u (\sqrt{-g} T^u_v) + \frac{1}{2} \partial_v (g_{\mu\nu}) T^{\mu u}$$

$$T^u_\mu g^{\mu\lambda} = \frac{1}{\sqrt{-g}} \partial_u (\sqrt{-g} T^u_v) - \frac{1}{2} \partial_v (g_{\mu\nu}) T^{\mu u}$$

(6)

### (4)a Schwarzschild Metric (General Form)

$$ds^2 = \left(1 - \frac{R_s}{r}\right) dt^2 + \left(1 - \frac{R_s}{r}\right)^{-1} dr^2 + r^2 d\Omega^2$$

$$= -\left(1 - \frac{2GM}{r}\right) dt^2 + \left(1 - \frac{2GM}{r}\right)^{-1} dr^2 + r^2 d\Omega^2$$

$$\text{Geodesic eqn general form: } \ddot{x}^\mu + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} = 0$$

$$\frac{d^2x^\mu}{d\tau^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} = 0$$

$$\therefore \frac{d^2r}{d\tau^2} + \Gamma_{tt}^r \frac{dr}{d\tau} \frac{dt}{d\tau} + \Gamma_{rr}^r \frac{dr}{d\tau} \frac{dt}{d\tau} = 0 \quad (\text{set } \theta = \phi \text{ and } \Omega = 0)$$

$$= \frac{d^2r}{d\tau^2} + \Gamma_{tt}^r \frac{dr}{d\tau} \frac{dt}{d\tau} + \Gamma_{rr}^r \frac{dr}{d\tau} \frac{dt}{d\tau} = 0 \quad (\text{set } \theta = \phi \text{ and } \Omega = 0)$$

$$\therefore \frac{d^2t}{d\tau^2} + \frac{-2GM}{(2GM-r)r} \frac{dr}{d\tau} \frac{dt}{d\tau} = 0 \quad (1)$$

$$\frac{d^2r}{d\tau^2} + \Gamma_{tt}^r \left(\frac{dt}{d\tau}\right)^2 + \Gamma_{rr}^r \left(\frac{dr}{d\tau}\right)^2 = 0 + \Gamma_{00}^r \left(\frac{dt}{d\tau}\right)^2 + \Gamma_{\varphi\varphi}^r \left(\frac{d\varphi}{d\tau}\right)^2$$

$$\Rightarrow \frac{d^2r}{d\tau^2} + \frac{-GM(2GM-r)}{(2GM-r)r} \left(\frac{dt}{d\tau}\right)^2 + \frac{GM}{(2GM-r)r} \left(\frac{dr}{d\tau}\right)^2 + 2GM-r \left(\frac{d\varphi}{d\tau}\right)^2$$

$$+ (2GM-r)\sin^2(\theta) \left(\frac{d\varphi}{d\tau}\right)^2 = 0 \quad (2)$$

$$\frac{d^2\varphi}{d\tau^2} + 2\Gamma_{r\varphi}^{\varphi} \frac{dr}{d\tau} \frac{d\varphi}{d\tau} + \Gamma_{\varphi\varphi}^{\varphi} \left(\frac{d\varphi}{d\tau}\right)^2 = 0$$

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$$= \frac{d^2\theta}{dt^2} + \frac{2}{r} \frac{dr}{dt} \frac{d\theta}{dt} - \cos(\theta) \sin(\theta) \left( \frac{d\varphi}{dt} \right)^2 = 0$$

$$\frac{d^2\varphi}{dt^2} - 2r^2 \frac{d\varphi}{dt} \frac{dr}{dt} + 2r^2 \frac{d\theta}{dt} \frac{d\varphi}{dt} = 0$$

$$= \frac{d^2\varphi}{dt^2} + \frac{2}{r} \frac{d\varphi}{dt} \frac{dr}{dt} + 2 \cot(\theta) \frac{d\theta}{dt} \frac{d\varphi}{dt} = 0 \quad (4)$$

b.  $R = R_s$

(2) Recalculating Riemann and Ricci tensors for Question 2.

By definition, the Riemann tensor is given by:

$$R^{\mu}_{\nu\alpha\beta} = \partial_\alpha \Gamma^{\mu}_{\nu\beta} - \partial_\beta \Gamma^{\mu}_{\nu\alpha} + \Gamma^{\mu}_{\alpha\sigma} \Gamma^{\sigma}_{\nu\beta} - \Gamma^{\mu}_{\beta\sigma} \Gamma^{\sigma}_{\nu\alpha}$$

$$R^x_{xxx} = \partial_x \Gamma^x_{xx} - \partial_x \Gamma^x_{xk} + \Gamma^x_{x\lambda} \Gamma^{\lambda}_{xx} - \Gamma^x_{x\lambda} \Gamma^{\lambda}_{xx}$$

$$= \partial_x \left( \frac{1}{2x} \right) - \partial_x \left( \frac{1}{2x} \right) + \Gamma^x_{xx} \Gamma^x_{xx} - \Gamma^x_{xx} \Gamma^x_{xx}$$

$$R^x_{xyx} = \cancel{\partial_y \left( \frac{1}{2x} \right)} - \cancel{\partial_y \left( \frac{1}{2x} \right)} + \Gamma^x_{xy} \Gamma^y_{xx} - \Gamma^x_{xy} \Gamma^y_{xx}$$

$\lambda = x \Rightarrow$  the above trivially vanishes:

$$\lambda = y \Rightarrow \Gamma^x_{xy} \Gamma^y_{xx} - \Gamma^x_{xy} \Gamma^y_{xx} = 0$$

Similarly for  $R^y_{yyy} = 0$

$$R^x_{xxy} = \cancel{\partial_x \Gamma^x_{xy}} - \cancel{\partial_y \Gamma^x_{xy}} + \Gamma^x_{x\lambda} \Gamma^{\lambda}_{xy} - \Gamma^x_{x\lambda} \Gamma^{\lambda}_{xy}$$

$x = y = \lambda = x \Rightarrow = 0$ ,  $\lambda = y \Rightarrow = 0$  (given Christoffel symbols).

By symmetry  $R^x_{xyx} = 0$  (Similarly, any entry with more than two  $y$ 's is zero)

(8)

The non-vanishing components are  $R^x_{yxy}$ ,  $R^x_{yyx}$ ,  $R^y_{xxy}$ ,  $R^y_{xyx}$ ,  $R^y_{yyy}$ .

$R^x_{yxy}$  and  $R^x_{yyx}$

Hence  $R^x_{yxy} = R^x_{yyx}$

$$= \partial_x \left( \frac{\Gamma^x_{yy}}{2y} - \frac{\Gamma^x_{yy}}{2y} \right) + \Gamma^x_{xx} \Gamma^x_{yy} - \Gamma^x_{yy} \Gamma^x_{xx}$$

$$= \partial_y \left( \frac{1}{2y} \right) + \left( \frac{1}{2y} \right) \left( -\frac{1}{2y} \right) - \frac{1}{2y^2}$$

$$= \frac{1}{2y^2} - \frac{1}{4x^2y^2} = \frac{x+2y}{4x^2y^2}$$

$R^y_{xxy}$  is computed with additional eqd

$$\text{By symmetry, } R^x_{yyx} = -\frac{x+2y}{4x^2y^2}$$

$$\text{Hence } R^y_{xxy} = -R^y_{yyx} = \frac{x+2y}{4x^2y^2}$$

is identical, up to  $(\partial_y \leftrightarrow \partial_x)$ , and due to the symmetry in the form of the Christoffel symbols

Ricci tensor = ~~computation follows under eqd~~

$$R_{\mu\nu} = R^{\sigma}_{\mu\sigma\nu} = R^x_{yxy} + R^y_{xyx}$$

$$\therefore R_{xx} = R^{\sigma}_{x\sigma} = R^x_{xxy} + R^y_{xyx} = 0 + \frac{x+2y}{4x^2y^2}$$

$$R_{xy} = R^{\sigma}_{x\sigma y} = R^x_{xyx} + R^y_{xxy} = 0$$

$$R_{yy} = -R^{\sigma}_{y\sigma y} = 0$$

$$R_{yy} = R^{\sigma}_{y\sigma y} = R^x_{yyx} + R^y_{yyy}$$

(9)

$$\frac{\partial xy}{\partial x} = 0 \quad \frac{\partial xy}{\partial y} = 1$$

Value  $\begin{pmatrix} \frac{\partial xy}{\partial x} & 0 \\ 0 & \frac{\partial xy}{\partial y} \end{pmatrix}$

Ricci scalar was computed earlier (trace of  $R_{\mu\nu}$ )

$$\frac{\partial xy}{\partial x} + \frac{\partial xy}{\partial x} = \frac{\partial^2 xy}{\partial x^2}$$

Want to write:

$\partial_\mu A_\nu - \partial_\nu A_\mu$  in covariant form

$$= \nabla_\mu A_\nu - \nabla_\nu A_\mu \quad (\text{Carroll Ch. 4 p. 153})$$

2. b.  $\gamma^\mu$  = parallel transported along  $x^\mu(\tau)$ ,  $u^\mu = \frac{dx^\mu}{d\tau}$

It must obey the eqs of parallel transport

$$(\text{lie } \nabla_u \gamma = u^\mu \nabla_\mu \gamma^\mu = 0)$$

Let  $x^\mu(\tau)$  be an affine geodesic:

$$\Rightarrow x^\mu(\tau) = \alpha \tau + \beta, \quad \alpha, \beta \in \mathbb{R}$$

$$\therefore u^\mu = \dot{x}^\mu = g^{\mu\nu} g_{\nu\lambda} u^\lambda = g_{\mu\nu} u^\nu = g_{\mu\nu}(\alpha) =$$

(6)

$$= \frac{x+2y}{4xy^2} + 0 = \frac{x+2y}{4xy^2}$$

$$\therefore R_{\mu\nu} = \begin{pmatrix} \frac{x+2y}{4xy^2} & 0 \\ 0 & \frac{x+2y}{4y^2x} \end{pmatrix}$$

(Ricci scalar was computed earlier = Trace of  $R_{\mu\nu}$ )

$$= \frac{x+2y}{4y^2x} + \frac{x+2y}{4y^2x} = \frac{x+2y}{2x^3y^2}$$

(5) Q. Want to write:

$\partial_\mu A_\nu - \partial_\nu A_\mu$  in covariant form

$$= \nabla_\mu A_\nu - \nabla_\nu A_\mu \quad (\text{Carroll Ch. 4 p. 153})$$

b.  $y^\mu$  = parallel transported along  $x^\nu(\tau)$ ,  $u^\mu = \frac{dx^\mu}{d\tau}$

It must obey the eqn of parallel transport

$$(\text{i.e. } \nabla_u^\mu y^\nu = u^\lambda \nabla_\lambda y^\mu = 0)$$

Let  $x^\mu(\tau)$  be an affine geodesic:

$$\Rightarrow x^\mu(\tau) = \alpha \tau + \beta, \quad \alpha, \beta \in \mathbb{C}$$

$$u^\mu = \dot{x}^\mu = \frac{dx^\mu}{d\tau} = g_{\mu\nu} g^{\nu\lambda} \dot{x}^\lambda = g_{\mu\nu} \delta^\lambda_\nu (\alpha) =$$

(9)

(P)

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$$\therefore u^\mu u_\mu = \alpha^2 = \text{constant.}$$

$$ds^2 = 0 \quad ds^2 + 0 =$$

prop

prop

$$Y^\mu u_\mu = Y^0 \alpha_0 + \alpha Y^1 + \alpha Y^2 + \alpha Y^3 \dots$$

$\approx \alpha Y^\mu$  (dot product is conserved since  $u_\mu = \text{constant}$ ).

$$Y^\mu u_\mu =$$

$$\text{Know that } \nabla_\mu Y^\nu = 0 \Rightarrow Y^\nu \partial_\mu Y^\mu + \Gamma^\nu_{\mu\nu} =$$

(~~So we want to express the covariant derivative in terms of partial derivatives~~)

$$\partial_\mu Y^\nu + \Gamma^\nu_{\mu\nu} = 0$$

$$\therefore Y^\mu u_\mu = Y^\mu \frac{\partial x_\mu}{\partial t}$$

stays at hand.

modified formulas are A, B - A, B

(Euler-Lagrange Equations) A, B - A, B

$x^\mu = x^\mu(t)$  paths between following = X - d

constant between large set nodes from E

$$(0 = P^\mu \partial_\mu x^\nu - L^\mu \delta x^\nu)$$

boundary conditions d (J) are fed

$$A \oplus d, B \oplus d \in \mathbb{R}^{d \times n} = (J)^m$$

Hamiltonian and Legendre transformation

## 2 RG&TC-Code

```
In[54]:= xCoord = {t, x, θ, φ};  
g = {  
  {-x y, 0, 0, 0},  
  {0, x y t, 0, 0},  
  {0, 0, z, 0},  
  {0, 0, 0, x t}  
};  
RGtensors[g, xCoord]  
  
gdd = 
$$\begin{pmatrix} -x y & 0 & 0 & 0 \\ 0 & t x y & 0 & 0 \\ 0 & 0 & z & 0 \\ 0 & 0 & 0 & t x \end{pmatrix}$$
  
LineElement = -x y d[t]^2 + z d[θ]^2 + t x d[φ]^2 + t x y d[x]^2  
  
gUU = 
$$\begin{pmatrix} -\frac{1}{x y} & 0 & 0 & 0 \\ 0 & \frac{1}{t x y} & 0 & 0 \\ 0 & 0 & \frac{1}{z} & 0 \\ 0 & 0 & 0 & \frac{1}{t x} \end{pmatrix}$$
  
gUU computed in 0. sec  
Gamma computed in 0. sec  
Riemann(dddd) computed in 0. sec  
Riemann(Uddd) computed in 0. sec  
Ricci computed in 0. sec  
Weyl computed in 0. sec  
Einstein computed in 0. sec  
  
Out[56]=  
All tasks completed in 0.  
  
In[57]:= (* Ricci Scalar *)  
  
In[58]:= R  
Out[58]=  
- 
$$\frac{1}{2 t^2 x y}$$
  
  
In[59]:= (* Einstein Tensor *)  
  
In[60]:= EUd  
Out[60]=  

$$\left\{ \left\{ -\frac{1}{4 t^2 x y}, 0, 0, 0 \right\}, \left\{ 0, \frac{1}{4 t^2 x y}, 0, 0 \right\}, \left\{ 0, 0, \frac{1}{4 t^2 x y}, 0 \right\}, \left\{ 0, 0, 0, \frac{1}{4 t^2 x y} \right\} \right\}$$
  
  
In[61]:= (* Christoffel Symbol *)
```

In[62]:= **GUdd // MatrixForm**

Out[62]//MatrixForm=

$$\left( \begin{array}{c|c|c|c} \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ \frac{1}{2} \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 0 \\ \frac{1}{2y} \end{pmatrix} \\ \begin{pmatrix} 0 \\ \frac{1}{2t} \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} \frac{1}{2t} \\ 0 \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} \\ \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} \\ \begin{pmatrix} 0 \\ 0 \\ 0 \\ \frac{1}{2t} \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \end{pmatrix} & \begin{pmatrix} \frac{1}{2t} \\ 0 \\ 0 \\ 0 \end{pmatrix} \end{array} \right)$$

In[63]:= **Part[GUdd, 1, 2, 2]****Part[GUdd, 2, 2, 1]**

Out[63]=

$$\frac{1}{2}$$

Out[64]=

$$\frac{1}{2t}$$

In[65]:= **(\* Riemann tensor \*)**

In[66]:= **RUddd**

Out[66]=

$$\left\{ \left\{ \left\{ \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, -\frac{1}{4t}, 0, 0 \right\}, \left\{ \frac{1}{4t}, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\} \right\}, \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0, -\frac{1}{4ty} \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ \frac{1}{4ty}, 0, 0, 0 \right\} \right\} \right\}, \left\{ \left\{ \left\{ 0, -\frac{1}{4t^2}, 0, 0 \right\}, \left\{ \frac{1}{4t^2}, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, \frac{1}{4ty} \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, -\frac{1}{4ty}, 0, 0 \right\} \right\} \right\}, \left\{ \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0, -\frac{1}{4t^2} \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ \frac{1}{4t^2}, 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, -\frac{1}{4t} \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, \frac{1}{4t}, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\} \right\}$$

In[67]:= (\* Ricci Tensor \*)

In[68]:= **Rdd**

Out[68]=

$$\left\{ \left\{ \frac{1}{2t^2}, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\}, \left\{ 0, 0, 0, 0 \right\} \right\}$$

In[69]:= **Part[Rdd, 1, 1]**

Out[69]=

$$\frac{1}{2t^2}$$

```
In[70]:= xCoord = {t, r, φ};
g = {{-((1 + 2 G M)/r), 0, 0},
      {0, ((1 + 2 G M)/r), 0},
      {0, 0, r^2}}
```

Out[71]=

$$\left\{ \left\{ -1 - \frac{2GM}{r}, 0, 0 \right\}, \left\{ 0, 1 + \frac{2GM}{r}, 0 \right\}, \left\{ 0, 0, r^2 \right\} \right\}$$

In[72]:= **RGtensors[g, xCoord]**

$$g_{dd} = \begin{pmatrix} -1 - \frac{2GM}{r} & 0 & 0 \\ 0 & 1 + \frac{2GM}{r} & 0 \\ 0 & 0 & r^2 \end{pmatrix}$$

$$\text{LineElement} = \frac{(2GM+r) d[r]^2}{r} - \frac{(2GM+r) d[t]^2}{r} + r^2 d[\varphi]^2$$

$$g_{UU} = \begin{pmatrix} -\frac{r}{2GM+r} & 0 & 0 \\ 0 & \frac{r}{2GM+r} & 0 \\ 0 & 0 & \frac{1}{r^2} \end{pmatrix}$$

`gUU computed in 0. sec`

`Gamma computed in 0. sec`

`Riemann(dddd) computed in 0. sec`

`Riemann(Uddd) computed in 0. sec`

`Ricci computed in 0. sec`

`Weyl computed in 0. sec`

*Testing for 3-dim conformal flatness...*

••• Outer: Heads Times and List at positions 3 and 2 are expected to be the same. [i](#)

`Einstein computed in 0. sec`

`Out[72]=`

All tasks completed in 0.

`In[73]:= GUdd`

`Out[73]=`

$$\left\{ \left\{ \left\{ 0, -\frac{GM}{r(2GM+r)}, 0 \right\}, \left\{ -\frac{GM}{r(2GM+r)}, 0, 0 \right\}, \{0, 0, 0\} \right\}, \right. \\ \left. \left\{ \left\{ -\frac{GM}{r(2GM+r)}, 0, 0 \right\}, \left\{ 0, -\frac{GM}{r(2GM+r)}, 0 \right\}, \left\{ 0, 0, -\frac{r^2}{2GM+r} \right\} \right\}, \right. \\ \left. \left\{ \{0, 0, 0\}, \left\{ 0, 0, \frac{1}{r} \right\}, \left\{ 0, \frac{1}{r}, 0 \right\} \right\} \right\}$$

`In[74]:= xCoord = {\psi, \theta, \varphi};`

```
g = {{1, 0, 0},
     {0, Sin[\psi]^2, 0},
     {0, 0, Sin[\psi]^2 * Sin[\varphi]^2}}
```

`Out[75]=`

$$\{\{1, 0, 0\}, \{0, \text{Sin}[\psi]^2, 0\}, \{0, 0, \text{Sin}[\varphi]^2 \text{Sin}[\psi]^2\}\}$$

`In[76]:= RGtensors[g, xCoord]`

$$g_{\theta\theta} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \sin[\psi]^2 & 0 \\ 0 & 0 & \sin[\varphi]^2 \sin[\psi]^2 \end{pmatrix}$$

$$\text{LineElement} = d[\psi]^2 + d[\theta]^2 \sin[\psi]^2 + d[\varphi]^2 \sin[\varphi]^2 \sin[\psi]^2$$

$$g_{UU} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \csc[\psi]^2 & 0 \\ 0 & 0 & \csc[\varphi]^2 \csc[\psi]^2 \end{pmatrix}$$

`gUU` computed in 0. sec

`Gamma` computed in 0. sec

`Riemann(dddd)` computed in 0. sec

`Riemann(Uddd)` computed in 0. sec

`Ricci` computed in 0. sec

`Weyl` computed in 0. sec

*Testing for 3-dim conformal flatness...*

Outer: Heads Times and List at positions 3 and 2 are expected to be the same. [i](#)

`Einstein` computed in 0. sec

`Out[76]=`

All tasks completed in 0.

`In[77]:= RUdd`

`Out[77]=`

$$\left\{ \left\{ \left\{ \left\{ 0, 0, 0 \right\}, \left\{ 0, 0, 0 \right\}, \left\{ 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, \sin[\psi]^2, 0 \right\}, \left\{ -\sin[\psi]^2, 0, 0 \right\}, \left\{ 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, \sin[\varphi]^2 \sin[\psi]^2 \right\}, \left\{ 0, 0, 0 \right\}, \left\{ -\sin[\varphi]^2 \sin[\psi]^2, 0, 0 \right\} \right\} \right\}, \left\{ \left\{ \left\{ 0, -1, 0 \right\}, \left\{ 1, 0, 0 \right\}, \left\{ 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0 \right\}, \left\{ 0, 0, 0 \right\}, \left\{ 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0 \right\}, \left\{ 0, 0, -\cos[\psi]^2 \sin[\varphi]^2 \right\}, \left\{ 0, \cos[\psi]^2 \sin[\varphi]^2, 0 \right\} \right\} \right\}, \left\{ \left\{ \left\{ 0, 0, -1 \right\}, \left\{ 0, 0, 0 \right\}, \left\{ 1, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0 \right\}, \left\{ 0, 0, \cos[\psi]^2 \right\}, \left\{ 0, -\cos[\psi]^2, 0 \right\} \right\}, \left\{ \left\{ 0, 0, 0 \right\}, \left\{ 0, 0, 0 \right\}, \left\{ 0, 0, 0 \right\} \right\} \right\} \right\}$$

`In[78]:= GUdd`

`Out[78]=`

$$\left\{ \left\{ \left\{ 0, 0, 0 \right\}, \left\{ 0, -\cos[\psi] \sin[\psi], 0 \right\}, \left\{ 0, 0, -\cos[\psi] \sin[\varphi]^2 \sin[\psi] \right\} \right\}, \left\{ \left\{ 0, \cot[\psi], 0 \right\}, \left\{ \cot[\psi], 0, 0 \right\}, \left\{ 0, 0, 0 \right\} \right\}, \left\{ \left\{ 0, 0, \cot[\psi] \right\}, \left\{ 0, 0, 0 \right\}, \left\{ \cot[\psi], 0, \cot[\varphi] \right\} \right\} \right\}$$

`In[79]:=`

`In[80]:=`

`In[81]:=`

`In[97]:= xCoord = {x, y};  
g = {{x*y, 0}, {0, x^2}}`

`Out[98]=`

$$\left\{ \left\{ x y, 0 \right\}, \left\{ 0, x^2 \right\} \right\}$$

`In[99]:= RGtensors[g, xCoord]`

$$g_{dd} = \begin{pmatrix} xy & 0 \\ 0 & x^2 \end{pmatrix}$$

$$\text{LineElement} = xy d[x]^2 + x^2 d[y]^2$$

$$g_{UU} = \begin{pmatrix} \frac{1}{xy} & 0 \\ 0 & \frac{1}{x^2} \end{pmatrix}$$

$g_{UU}$  computed in 0. sec

Gamma computed in 0.015 sec

Riemann(dddd) computed in 0. sec

Riemann(Uddd) computed in 0. sec

Ricci computed in 0. sec

Weyl computed in 0. sec

## Conformally Flat

Einstein computed in 0. sec

## Einstein Space

Out[99]=

All tasks completed in 0.015625

In[85]:= **GUdd**

Out[85]=

$$\left\{ \left\{ \left\{ \frac{1}{2x}, \frac{1}{2y} \right\}, \left\{ \frac{1}{2y}, -\frac{1}{y} \right\} \right\}, \left\{ \left\{ -\frac{1}{2x}, \frac{1}{x} \right\}, \left\{ \frac{1}{x}, 0 \right\} \right\} \right\}$$

In[86]:= **RUdd**

Out[86]=

RUdd

In[100]=

### **RUddd // MatrixForm**

Out[100]//MatrixForm=

$$\begin{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} & \begin{pmatrix} 0 & \frac{x+2y}{4xy^2} \\ -\frac{x+2y}{4xy^2} & 0 \end{pmatrix} \\ \begin{pmatrix} 0 & -\frac{x+2y}{4x^2y} \\ \frac{x+2y}{4x^2y} & 0 \end{pmatrix} & \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} \end{pmatrix}$$

In[88]:= **Rdd**

Out[88]=

$$\left\{ \left\{ \frac{x+2y}{4x^2y}, 0 \right\}, \left\{ 0, \frac{x+2y}{4xy^2} \right\} \right\}$$

In[89]:= **R**

Out[89]=

$$\frac{x+2y}{2x^3y^2}$$

```
In[90]:= xCoord = {t, r, θ, φ}
Out[90]= {t, r, θ, φ}

In[91]:= g = {{-(1 - 2 * G * M / r), 0, 0, 0}, {0, (1 - 2 * G * M / r)^(-1), 0, 0}, {0, 0, r^2, 0}, {0, 0, 0, r^2 * Sin[θ]^2}}
Out[91]= {{-1 + 2 G M / r, 0, 0, 0}, {0, 1 / (1 - 2 G M / r), 0, 0}, {0, 0, r^2, 0}, {0, 0, 0, r^2 Sin[θ]^2} }

In[92]:= RGtensors[g, xCoord]
gdd = {{-1 + 2 G M / r, 0, 0, 0}, {0, 1 / (1 - 2 G M / r), 0, 0}, {0, 0, r^2, 0}, {0, 0, 0, r^2 Sin[θ]^2}}
LineElement = -r d[r]^2 / (2 G M - r) + (2 G M - r) d[t]^2 / r + r^2 d[θ]^2 + r^2 d[φ]^2 Sin[θ]^2
gUU = {{r / (2 G M - r), 0, 0, 0}, {0, -2 G M / r, 0, 0}, {0, 0, 1 / r^2, 0}, {0, 0, 0, Csc[θ]^2}}
gUU computed in 0. sec
Gamma computed in 0. sec
Riemann(dddd) computed in 0. sec
Riemann(Uddd) computed in 0. sec
Ricci computed in 0. sec
Weyl computed in 0. sec
```

## Ricci Flat

```
Out[92]= All tasks completed in 0.
```

```
In[93]:= GUdd
Out[93]= {{{{0, -(G M) / ((2 G M - r) r), 0, 0}, {-G M / ((2 G M - r) r), 0, 0, 0}, {0, 0, 0, 0}, {0, 0, 0, 0}}, {{{-G M (2 G M - r) / r^3, 0, 0, 0}, {0, G M / ((2 G M - r) r), 0, 0}, {0, 0, 2 G M - r, 0}, {0, 0, 0, (2 G M - r) Sin[θ]^2}}}, {{{0, 0, 0, 0}, {0, 0, 1 / r, 0}, {0, 1 / r, 0, 0}, {0, 0, 0, -Cos[θ] Sin[θ]}}, {{0, 0, 0, 0}, {0, 0, 0, 1 / r}, {0, 0, Cot[θ], 0}, {0, 1 / r, Cot[θ], 0}}}}
```