

6) Electric and magnetic field operators in Heisenberg picture:

$$\hat{E}(\underline{r}, t) = i \sum_{\underline{k}s} \sqrt{\frac{2\pi\hbar\omega_{\underline{k}}}{V}} \underline{e}_{\underline{k}s} (\hat{a}_{\underline{k}s} e^{i(\underline{k}\cdot\underline{r}-\omega t)} - \hat{a}_{\underline{k}s}^\dagger e^{-i(\underline{k}\cdot\underline{r}-\omega t)})$$

$$\hat{B}(\underline{r}, t) = i \sum_{\underline{k}s} \sqrt{\frac{2\pi\hbar c^2}{\omega_{\underline{k}} V}} [\underline{k} \times \underline{e}_{\underline{k}s}] (\hat{a}_{\underline{k}s} e^{i(\underline{k}\cdot\underline{r}-\omega t)} - \hat{a}_{\underline{k}s}^\dagger e^{-i(\underline{k}\cdot\underline{r}-\omega t)})$$

- Required to compute average values for $\hat{E}(\underline{r}, t)$ and $\hat{B}(\underline{r}, t)$ in the classical single mode state (occupations of all photonic modes = 0 except for the single mode)
- For $\hat{B}(\underline{r}, t)$ $\underline{k} = k_0 \underline{e}_x$, polarization $\underline{e}_0 = \underline{e}_y$ and the occupation is as the coherent state α .

Solution:

$$\langle \hat{E}(\underline{r}, t) \rangle = \langle \alpha | \hat{E} | \alpha \rangle$$

$$\begin{aligned} &= \langle \alpha | \hat{E} | \alpha \rangle = i \sqrt{\frac{2\pi\hbar\omega_{\underline{k}}}{V}} \langle \alpha | \underline{e}_{\underline{k}s} (\hat{a}_{\underline{k}s} e^{i(\underline{k}\cdot\underline{r}-\omega t)} - \hat{a}_{\underline{k}s}^\dagger e^{-i(\underline{k}\cdot\underline{r}-\omega t)}) | \alpha \rangle \\ &= i \sqrt{\frac{2\pi\hbar\omega_{\underline{k}}}{V}} \underline{e}_{\underline{k}s} \langle \alpha | \hat{a}_{\underline{k}s} e^{i(\underline{k}\cdot\underline{r}-\omega t)} - \hat{a}_{\underline{k}s}^\dagger e^{-i(\underline{k}\cdot\underline{r}-\omega t)} | \alpha \rangle \\ &= i \sqrt{\frac{2\pi\hbar\omega_{\underline{k}}}{V}} \underline{e}_{\underline{k}s} \left(\alpha e^{i(\underline{k}\cdot\underline{r}-\omega t)} - \alpha^* e^{-i(\underline{k}\cdot\underline{r}-\omega t)} \right) \\ &= 2i \cdot i \sqrt{\frac{2\pi\hbar\omega_{\underline{k}}}{V}} \underline{e}_{\underline{k}s} \cdot \frac{1}{2i} \left(\alpha e^{i(\underline{k}\cdot\underline{r}-\omega t)} - \alpha^* e^{-i(\underline{k}\cdot\underline{r}-\omega t)} \right) \\ &= -2 \sqrt{\frac{2\pi\hbar\omega_{\underline{k}}}{V}} \underline{e}_{\underline{k}s} \frac{1}{2i} \left(|\alpha| e^{i\delta} e^{i(\underline{k}\cdot\underline{r}-\omega t)} - |\alpha| e^{-i\delta} e^{-i(\underline{k}\cdot\underline{r}-\omega t)} \right) \\ &= -2 \sqrt{\frac{2\pi\hbar\omega_{\underline{k}}}{V}} \underline{e}_{\underline{k}s} \frac{1}{2i} \left(|\alpha| e^{i(\underline{k}\cdot\underline{r}-\omega t+\delta)} - |\alpha| e^{-i(\underline{k}\cdot\underline{r}-\omega t+\delta)} \right) \\ &= -2 \sqrt{\frac{2\pi\hbar\omega_{\underline{k}}}{V}} \underline{e}_{\underline{k}s} \cdot |\alpha| \sin(\underline{k}\cdot\underline{r}-\omega t+\delta) \end{aligned}$$

Similarly for $\langle \hat{B}(\underline{r}, t) \rangle$

$$\begin{aligned} &= \langle \alpha | \hat{B}(\underline{r}, t) | \alpha \rangle \\ &= i \sqrt{\frac{2\pi\hbar c^2}{\omega_{\underline{k}} V}} \langle \alpha | [\underline{k} \times \underline{e}_{\underline{k}s}] (\hat{a}_{\underline{k}s} e^{i(\underline{k}\cdot\underline{r}-\omega t)} - \hat{a}_{\underline{k}s}^\dagger e^{-i(\underline{k}\cdot\underline{r}-\omega t)}) | \alpha \rangle \\ &= i \sqrt{\frac{2\pi\hbar c^2}{\omega_{\underline{k}} V}} [\underline{k} \times \underline{e}_{\underline{k}s}] \langle \alpha | \hat{a}_{\underline{k}s} e^{i(\underline{k}\cdot\underline{r}-\omega t)} - \hat{a}_{\underline{k}s}^\dagger e^{-i(\underline{k}\cdot\underline{r}-\omega t)} | \alpha \rangle \\ &= i \sqrt{\frac{2\pi\hbar c^2}{\omega_{\underline{k}} V}} [\underline{k} \times \underline{e}_{\underline{k}s}] \left[\alpha e^{i(\underline{k}\cdot\underline{r}-\omega t)} - \alpha^* e^{-i(\underline{k}\cdot\underline{r}-\omega t)} \right] \\ &\quad \cdot \text{Let } \alpha = |\alpha| e^{i\delta} \\ &\Rightarrow i \sqrt{\frac{2\pi\hbar c^2}{\omega_{\underline{k}} V}} [\underline{k} \times \underline{e}_{\underline{k}s}] \left[|\alpha| e^{i(\underline{k}\cdot\underline{r}-\omega t+\delta)} - |\alpha| e^{-i(\underline{k}\cdot\underline{r}-\omega t+\delta)} \right] \\ &= 2i \cdot i \sqrt{\frac{2\pi\hbar c^2}{\omega_{\underline{k}} V}} [\underline{k} \times \underline{e}_{\underline{k}s}] \frac{1}{2i} \left[|\alpha| e^{i(\underline{k}\cdot\underline{r}-\omega t+\delta)} - |\alpha| e^{-i(\underline{k}\cdot\underline{r}-\omega t+\delta)} \right] \\ &= -2 \sqrt{\frac{2\pi\hbar c^2}{\omega_{\underline{k}} V}} [\underline{k} \times \underline{e}_{\underline{k}s}] \left[|\alpha| \sin(\underline{k}\cdot\underline{r}-\omega t+\delta) \right] \\ &= -2 |\alpha| \sqrt{\frac{2\pi\hbar c^2}{\omega_{\underline{k}} V}} [\underline{k} \times \underline{e}_{\underline{k}s}] \sin(\underline{k}\cdot\underline{r}-\omega t+\delta) \end{aligned}$$

- In both of the above scenarios, the average values of \hat{E} and \hat{B} for the coherent state $|\alpha\rangle$ are of the form:

$$\vec{E} = E_0 \sin(\underline{k}\cdot\underline{r}-\omega t+\delta) \underbrace{\hat{e}}_{\text{electric field vector}} \quad \text{and} \quad E_0 = -2|\alpha| \sqrt{\frac{2\pi\hbar\omega_{\underline{k}}}{V}}$$

and

$$\vec{B} = B_0 \sin(\underline{k}\cdot\underline{r}-\omega t+\delta) \underbrace{[\hat{k} \times \hat{e}]}_{\text{direction of magnetic field vector}} \quad \text{and} \quad B_0 = -2|\alpha| \sqrt{\frac{2\pi\hbar c^2}{V\omega_{\underline{k}}}}$$

- Take $\underline{k} = k_0 \underline{e}_x$ and take $\underline{e} = \underline{e}_y = \underline{e}_y$. Hence:

$$\vec{E}(\underline{r}, t) = E_0 \sin(k_0 \underline{e}_x \cdot (\hat{x}, \hat{y}, \hat{z}) - \omega t + \delta) \underline{e}_y$$

$$\Rightarrow \vec{E}(x, t) = E_0 \sin(k_0 x - \omega t + \delta) \underline{e}_y$$

$$\begin{aligned} \vec{B}(\underline{r}, t) &= B_0 [k_0 \underline{e}_x \times \underline{e}_y] \sin(k_0 x - \omega t + \delta) \\ &= B_0 k_0 \underline{e}_z \sin(k_0 x - \omega t + \delta) \end{aligned}$$

(for \vec{E} and \vec{B})

- Note that both of the above expressions are in phase with one another, and are also orthogonal to one another, thereby mimicking the spatial and the temporal behaviour of the classical EM wave.