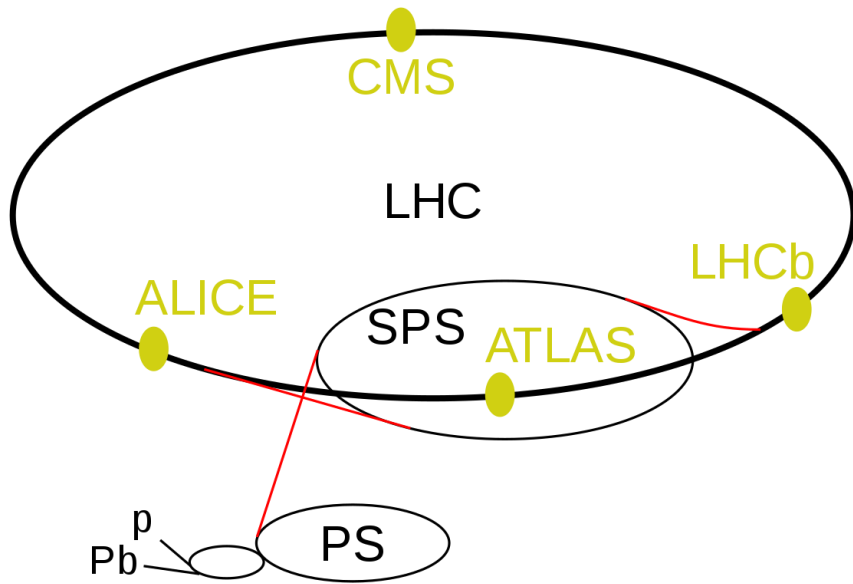


The Search for Axion Like Particles (ALPs) in B Meson Decays at the LHCb

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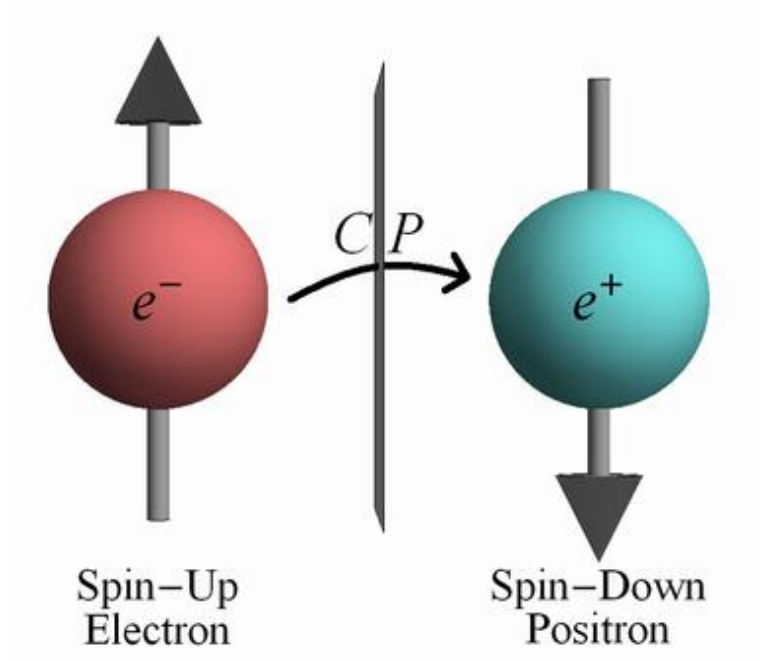


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Background and Motivation

- CP symmetry is preserved in EM interactions but violated by weak interactions (Cronin & Fitch, 1964)
- CP violation is not observed experimentally in the strong force, despite being theoretically allowed (**Strong CP Problem**). Significant limitation of the Standard Model



| | | | | | |
|------------------------------|---|---|--|--|---|
| mass → charge → spin → | <div>≈2.3 MeV/c² 2/3 1/2 u up</div> | <div>≈1.275 GeV/c² 2/3 1/2 c charm</div> | <div>≈173.07 GeV/c² 2/3 1/2 t top</div> | <div>0 0 1 g gluon</div> | <div>≈126 GeV/c² 0 0 0 H Higgs boson</div> |
| QUARKS | <div>≈4.8 MeV/c² -1/3 1/2 d down</div> | <div>≈95 MeV/c² -1/3 1/2 s strange</div> | <div>≈4.18 GeV/c² -1/3 1/2 b bottom</div> | <div>0 0 1 γ photon</div> | |
| | <div>0.511 MeV/c² -1 1/2 e electron</div> | <div>105.7 MeV/c² -1 1/2 μ muon</div> | <div>1.777 GeV/c² -1 1/2 τ tau</div> | <div>91.2 GeV/c² 0 1 Z Z boson</div> | |
| LEPTONS | <div><2.2 eV/c² 0 1/2 ν_e electron neutrino</div> | <div><0.17 MeV/c² 0 1/2 ν_μ muon neutrino</div> | <div><15.5 MeV/c² 0 1/2 ν_τ tau neutrino</div> | <div>80.4 GeV/c² ±1 1 W W boson</div> | GAUGE BOSONS |

The Strong CP Problem

- CP violation is **theoretically permitted** in the strong force (QCD) but there is **no experimental evidence of this**
- τ QED Lagrangian (electromagnetism):

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\psi}(i\gamma^\mu\mathcal{D}_\mu - m_e)\psi$$

- Consider the QCD (strong) Lagrangian written in the following form:

$$\mathcal{L}_{QCD} = -\frac{1}{4}G_{\mu\nu}G^{\mu\nu} - \boxed{\frac{g_s^2\theta}{32\pi^2}G_{\mu\nu}\tilde{G}^{\mu\nu}} + \bar{\psi}(i\gamma^\mu D_\mu - me^{i\theta'\gamma_5})\psi$$

- The effects of the θ -dependent term are not observed experimentally. **Hence, θ must be very small**

The Strong CP Problem and its Resolution

- Experimental measurements of neutron EDM $\Rightarrow |\theta| < 10^{-10}$

Solution: Promote θ to a **dynamic field** by adding a new symmetry that is spontaneously broken (Peccei & Quinn, 1977)*

- Spontaneous breaking of this **PQ symmetry** introduces a new pseudoscalar (**spin 0 and odd parity**) particle known as the (QCD) axion. (No experimental evidence of this)

Axion Like Particles (ALPs)

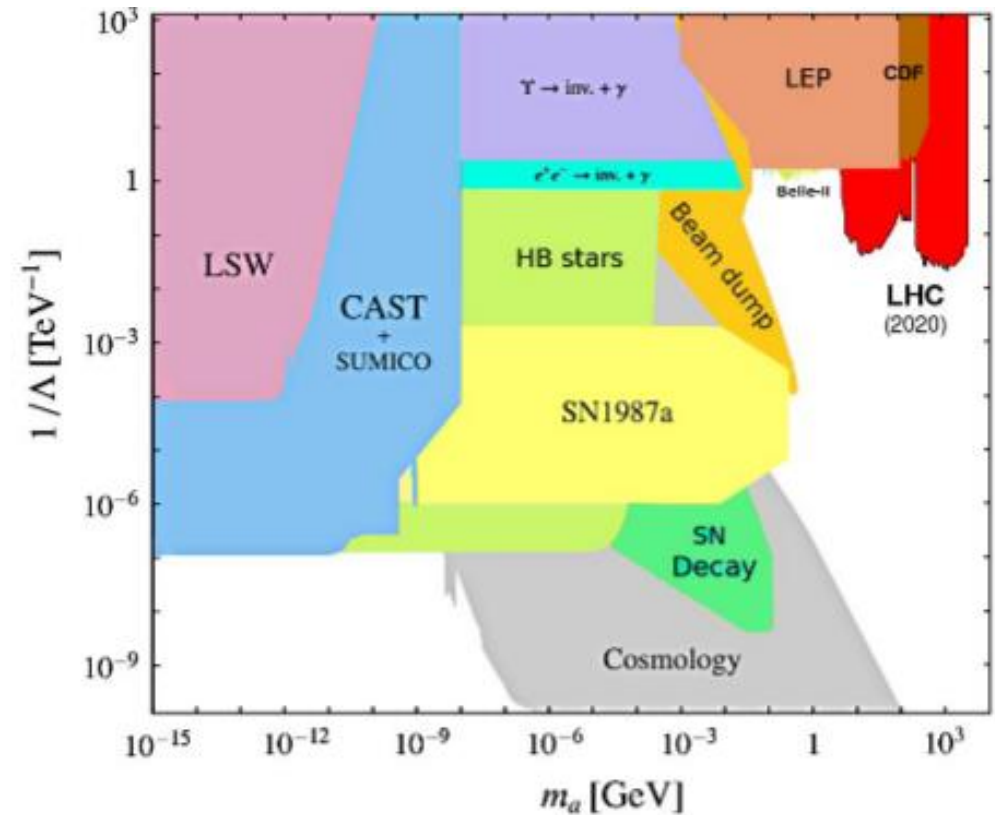
- Spontaneous breaking of an approximate symmetry (**not PQ**) can also generate other **axion-like particles (ALPs)**.
- Masses and couplings to photons are independent for ALPs and are therefore far less constrained
- Couple predominantly to pairs of gauge bosons (e.g. $gg, \gamma\gamma, ZZ, \gamma Z, W^\pm$ etc.) depending on the model being considered*

*Ringwald (2014) *Axions and Axion-Like Particles*: <https://arxiv.org/pdf/1407.0546.pdf>

** Isern et al. (2018) *Axions and the Cooling of White Dwarf Stars* <https://arxiv.org/pdf/0806.2807.pdf>

Experimental Searches for Axions and ALPs

- Spin-selection rules => light pseudoscalars naturally couple to photons
- Search strategies generally exploit the (inverse) Primakoff effect
- Notable search strategies (excluding collider searches):
 - **LSW (Light Shining Through Walls) Experiments**
 - Any Light Particles Search (ALPS I)
 - ALPS II
 - **Helioscope Searches**
 - International Axion Observatory (IAXO)
 - CERN Axion Space Telescope (CAST)
 - **Haloscope Searches**
 - Axion Dark Matter Experiment (ADMX)
 - PIXIE
 - PRISM CMB

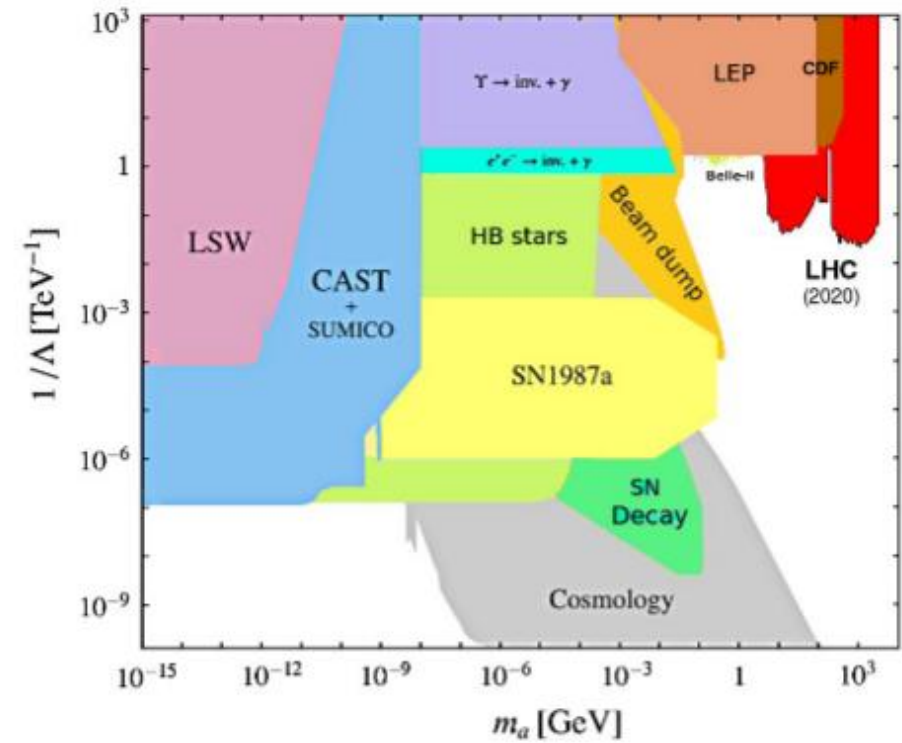
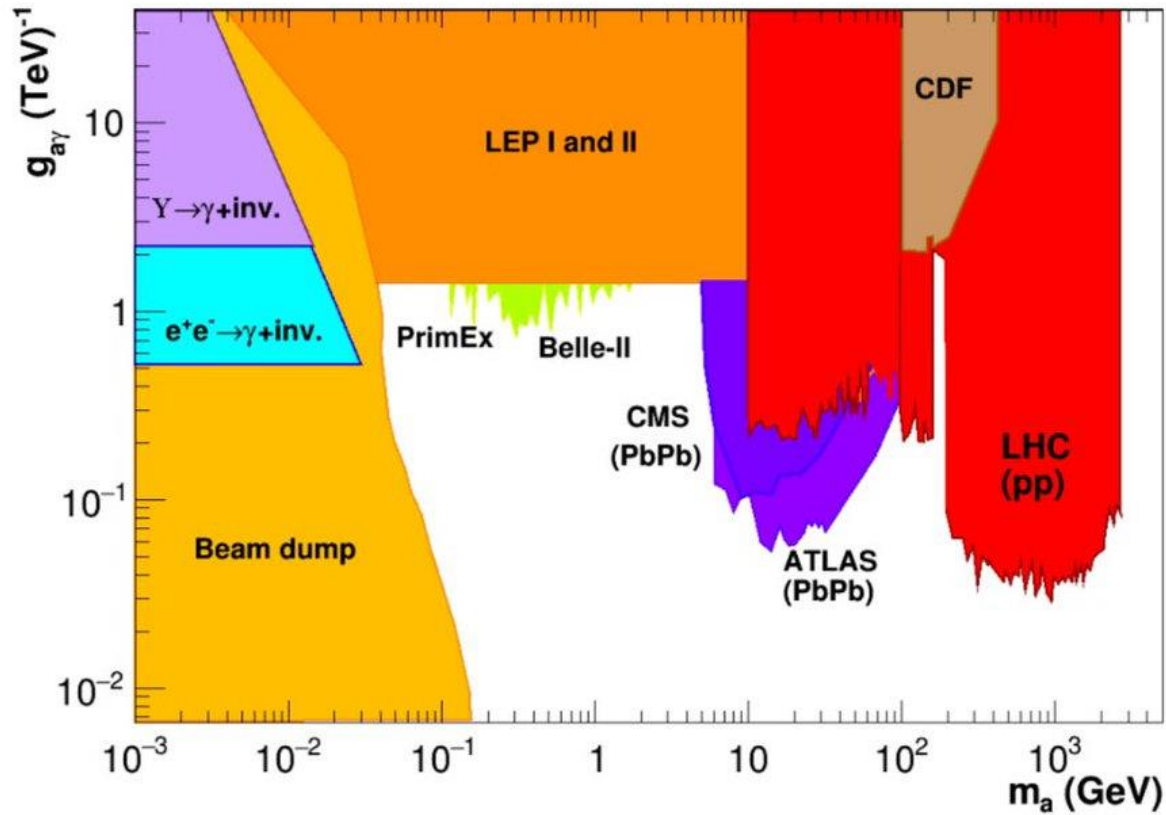


Source: A. Ringwald. *Axions and axion-like particles*, 2014.

Strategy: Search for ALPs at Colliders

0. Set limit on branching fraction of decay of interest using Monte Carlo (MC) simulated data to determine if analysis is viable/worth pursuing
1. Event selection (i.e. impose constraints on kinematic and shape variables to distinguish signal from background within MC simulated data)
2. Check optimised event selection against a real data sample to verify that MC simulation models the data
3. Perform a fit to extract the signal yield
4. Estimate systematic errors

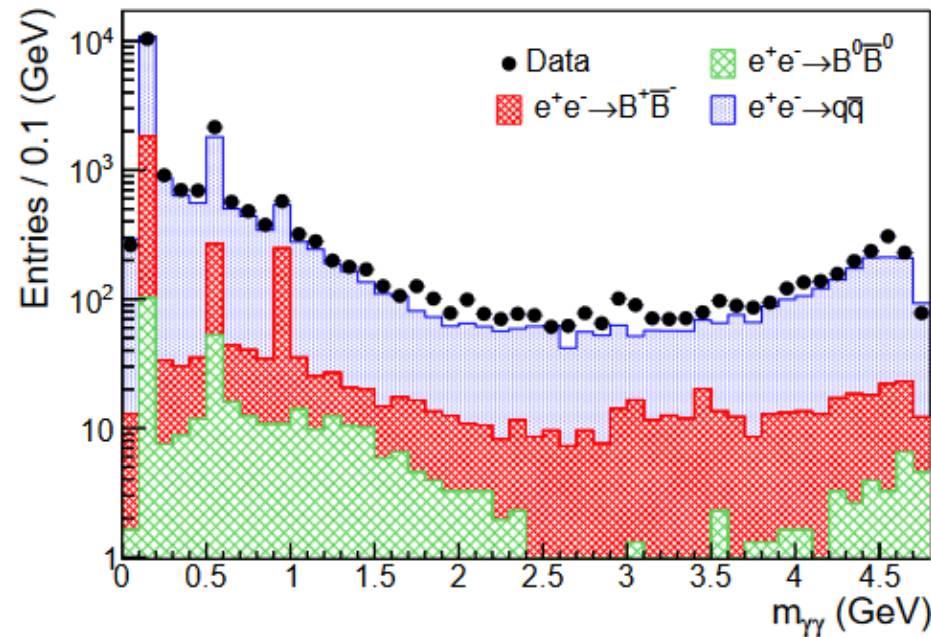
Summary of Mass and Coupling Constraints



Strategy: Search for ALPs at LHCb

- Seek diphoton resonance structures
- Promising decay channel for search: $B^0 \rightarrow K^{*0} a_0, a_0 \rightarrow \gamma\gamma$
- ALPs produced in B -meson decays have a maximal mass of

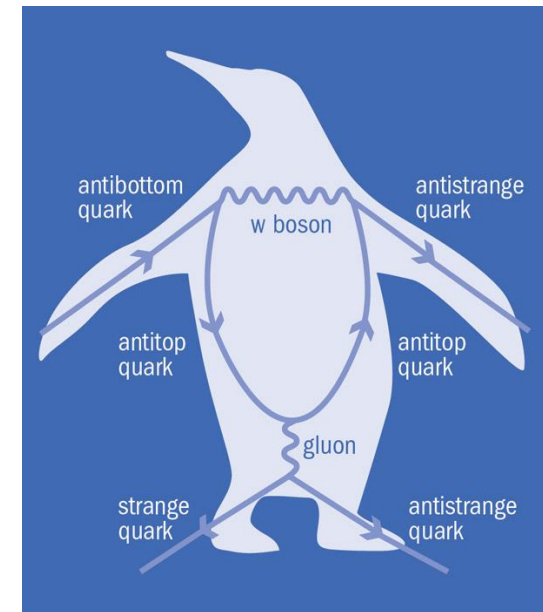
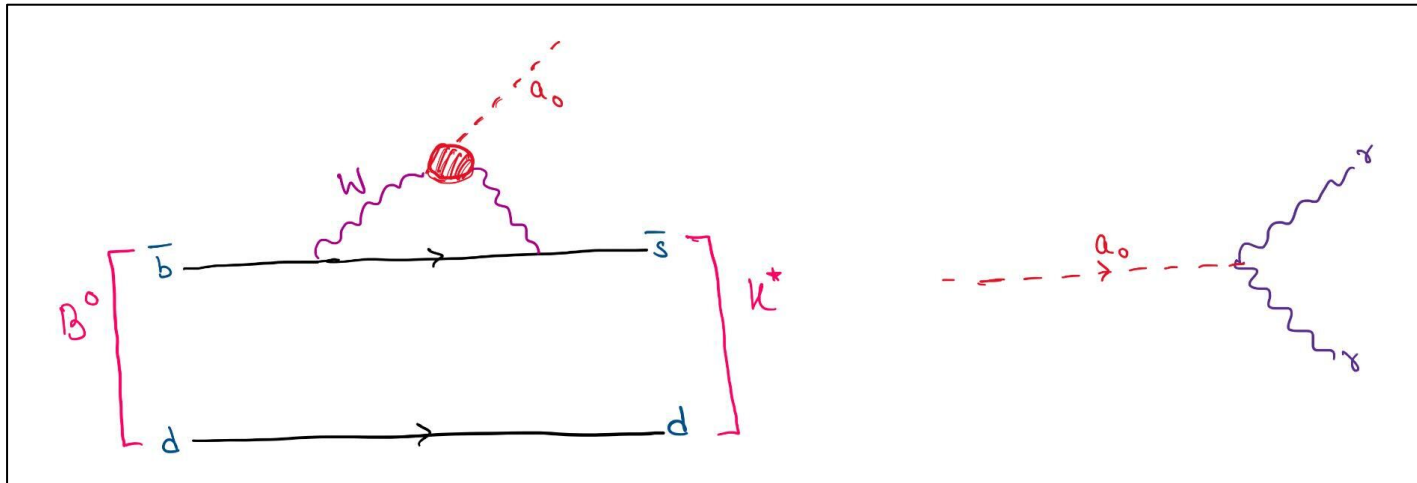
$$m_{a_0} = m_{B^0} - m_{K^*} = 5279.26 - 493.68 \text{ MeV} = 4785.58 \text{ MeV}$$



The BABAR Collaboration: *Search for Axion Like Particles in B Meson Decays*: <https://arxiv.org/pdf/2111.01800.pdf>

The $B^0 \rightarrow K^{*0} a_0, a_0 \rightarrow \gamma\gamma$ Decay

- Consider model where ALP couples to weak gauge bosons W^\pm , and gives rise to observable signatures (zero coupling with gluons)

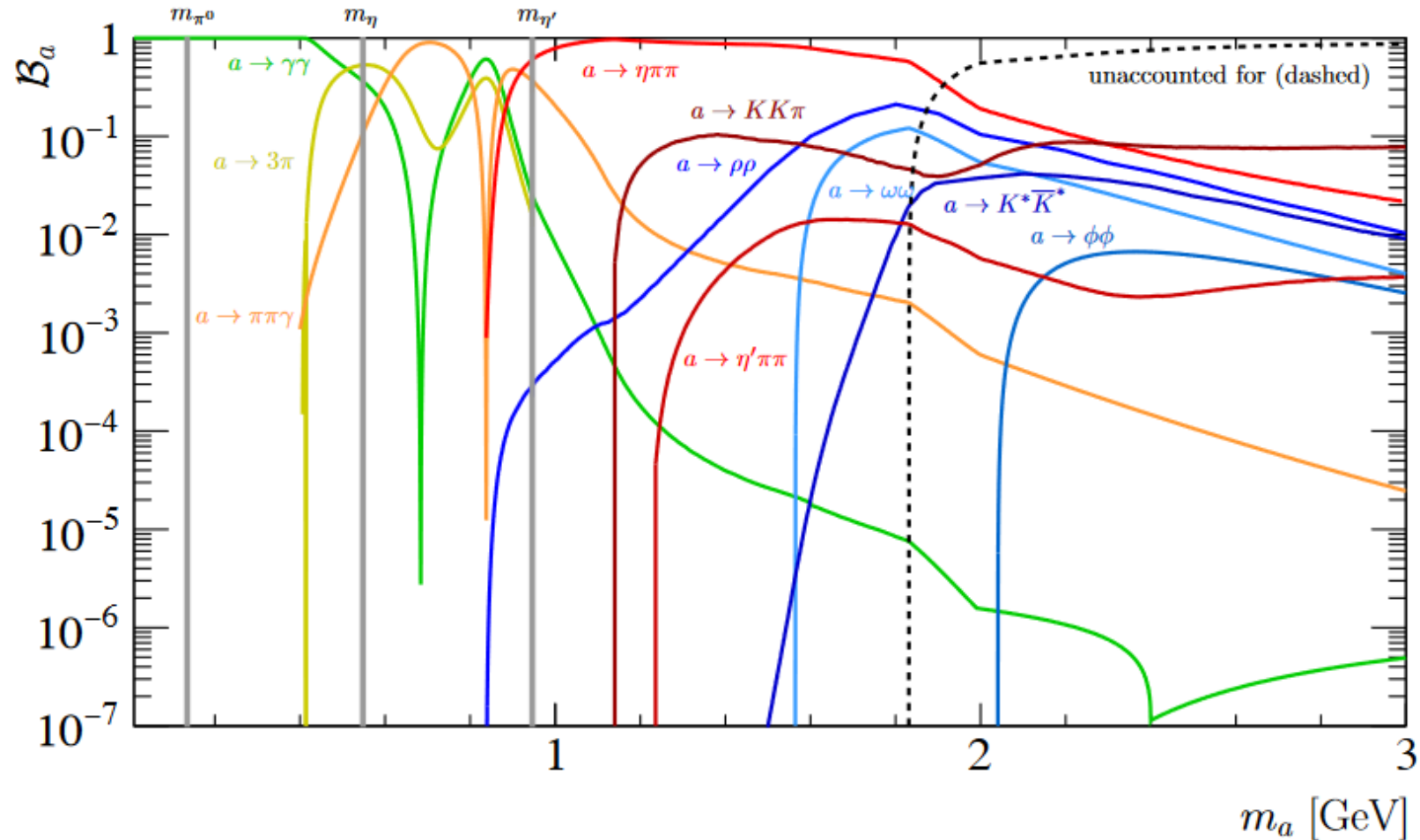


- Flavour Changing Neutral Current (FCNC) process ($\bar{b} \rightarrow \bar{s}$ quark transition)
- Electroweak penguin decay that proceeds at one-loop level

(R) Image Source: <https://cerncourier.com/a/chasing-new-physics-with-electroweak-penguins/>

Source: <https://arxiv.org/abs/1611.09355>

Plot of ALP Branching Fraction vs ALP Mass



<https://journals.aps.org/prl/supplemental/10.1103/PhysRevLett.123.031803/supplemental.pdf>