

(3)e. Want to prove the commutation relations between components of kinematic ($\hat{H} = \hat{p} + e\vec{A}(x)/c$) and magnetic ($\hat{k} = \hat{p} - e\vec{A}(x)/c$) momenta given that $A(x) = \left\{-\frac{By}{2}, \frac{Bx}{2}, 0\right\}$

$$(i) [\hat{\pi}_x, \hat{\pi}_y] = -\frac{i\hbar^2}{c}. \text{ Claim:}$$

$$\begin{aligned} [\hat{\pi}_x, \hat{\pi}_y] &= [p_x - eA_x(x)/c, p_y - eA_y(y)/c] \\ &= [p_x, p_y] + [-p_x, eA_y(y)/c] + [eA_x(x)/c, p_y] + [eA_x/c, eA_y/c] \\ &= 0 - \left(i\hbar \frac{\partial A_y}{\partial x}\right) \frac{e}{c} + \frac{e}{c} \left(i\hbar \frac{\partial A_x}{\partial y}\right) + \frac{e}{c} (0) \\ &= \left(-i\hbar \frac{\partial A_y}{\partial x}\right) \frac{e}{c} + \left(i\hbar \frac{\partial A_x}{\partial y}\right) \frac{e}{c} \\ &= -\frac{i\hbar e}{c} \left(-\frac{\partial A_y}{\partial x} + \frac{\partial A_x}{\partial y}\right) \\ &= -\frac{i\hbar e}{c} \left(\frac{\partial A_x}{\partial y} - \frac{\partial A_y}{\partial x}\right) = -i\frac{\hbar e}{c} B \quad \vec{B} = \underline{\underline{B}} \Rightarrow B_x = \frac{\partial A_y}{\partial z} - \frac{\partial A_z}{\partial y} \quad \text{Since } B \text{ is uniform, } B_x = \underline{\underline{B}} \\ &\text{Let } l_B = \sqrt{\frac{\hbar c}{eB}} \Rightarrow \hat{l}_B = \frac{\hbar c}{eB} \Rightarrow \frac{\hbar^2}{l_B^2} = \frac{\hbar^2}{\hbar c} \cdot \frac{eB}{\hbar c} = \frac{e\hbar B}{c} \\ &\Rightarrow [\hat{\pi}_x, \hat{\pi}_y] = -\frac{i\hbar^2}{\underline{\underline{l}_B^2}} \end{aligned}$$

$$(ii) [\hat{k}_x, \hat{k}_y] = [p_x - eA_x/c, p_y - eA_y/c]$$

$$\begin{aligned} &= [p_x, p_y] - [p_x, eA_y/c] - [eA_x/c, p_y] + [eA_x/c, eA_y/c] \\ &= [p_x, p_y] - \frac{e}{c} [p_x, A_y] - \frac{e}{c} [A_x, p_y] + \frac{e}{c} [A_x, A_y] \\ &= 0 - \frac{e}{c} [p_x, A_y] - \frac{e}{c} [A_x, p_y] + \frac{e}{c} (0) \\ &= 0 - \frac{e}{c} \left[\frac{i\hbar \partial A_y}{\partial x} + i\hbar \frac{\partial A_x}{\partial y} \right] = 0 \end{aligned}$$

$$= -\frac{i\hbar e}{c} \left[\frac{\partial A_x}{\partial y} - \frac{\partial A_y}{\partial x} \right] = \frac{i\hbar e}{c} (-B_z) \Rightarrow [\hat{k}_x, \hat{k}_y] = \frac{i\hbar}{\underline{\underline{l}_B^2}} \quad \text{where } l_B = \sqrt{\frac{\hbar c}{eB}}$$

$$[\hat{k}_x, \hat{\pi}_x] = [p_x - e/c A_x, p_x - e/c A_x]$$

$$= [p_x, p_x] + [p_x, e/c A_x] - [e/c A_x, p_x] - [e/c A_x, e/c A_x]$$

$$= [p_x, \frac{e}{c} \hat{A}_x] + [\hat{p}_x, \frac{e}{c} \hat{A}_x]$$

$$= \frac{e}{c} ([\hat{p}_x, \hat{A}_x] + [\hat{A}_x, \hat{p}_x]) = 0$$

∴ By symmetry, $[\hat{k}_y, \hat{\pi}_y] = 0$

$$(i) [\hat{u}_x, \hat{\pi}_y] = [p_x - e/c A_x, p_y - e/c A_y]$$

$$= [p_x, p_y] - [p_x, \frac{e}{c} A_y] - [\frac{e}{c} A_x, p_y] + [\frac{e}{c} A_x, \frac{e}{c} A_y]$$

$$= -[p_x, \frac{e}{c} A_y] - [\frac{e}{c} A_x, p_y] + [\frac{e}{c} A_x, A_y]$$

$$= -\frac{e}{c} [p_x, A_y] - \frac{e}{c} [A_x, p_y] + \frac{e}{c} \left(\frac{i\hbar \partial A_y}{\partial x} + \frac{i\hbar \partial A_x}{\partial y} \right)$$

$$= \frac{e}{c} [p_y, A_x] - \frac{e}{c} [p_x, A_y] = \frac{i\hbar e}{c} \left(\frac{\partial A_x}{\partial y} - \frac{\partial A_y}{\partial x} \right) =$$

Given that $A = \left\{-\frac{By}{2}, \frac{Bx}{2}, 0\right\}$, it is evident that $\frac{\partial A_y}{\partial x} + \frac{\partial A_x}{\partial y} = \frac{B_x}{2} - \frac{B_y}{2} = 0$

$$\therefore [\hat{u}_x, \hat{\pi}_y] = 0 \Rightarrow \underline{\underline{\text{QED}}}$$

$[\hat{u}_y, \hat{\pi}_x] = 0$ by similar arguments

b. Introduce ladder operators: \hat{a} , \hat{a}^\dagger , \hat{b} and \hat{b}^\dagger as follows:

$$\hat{a} = \frac{l_B}{\sqrt{2\hbar}} (\hat{\pi}_x - i\hat{\pi}_y) \quad \hat{a}^\dagger = \frac{l_B}{\sqrt{2\hbar}} (\hat{\pi}_x + i\hat{\pi}_y)$$

$$\hat{b} = \frac{l_B}{\sqrt{2\hbar}} (\hat{k}_y - i\hat{u}_x) \quad \hat{b}^\dagger = \frac{l_B}{\sqrt{2\hbar}} (\hat{k}_y + i\hat{u}_x)$$

Want to show that: $[\hat{a}, \hat{a}^\dagger] = 1$ and $[\hat{b}, \hat{b}^\dagger] = 1$ and $[\hat{a}, \hat{b}] = [\hat{a}, \hat{b}^\dagger] = [\hat{a}^\dagger, \hat{b}] = [\hat{a}^\dagger, \hat{b}^\dagger] = 0$

$$(i) [\hat{a}, \hat{a}^\dagger] = \left[\frac{l_B}{\sqrt{2\hbar}} (\hat{\pi}_x - i\hat{\pi}_y), \frac{l_B}{\sqrt{2\hbar}} (\hat{\pi}_x + i\hat{\pi}_y) \right]$$

$$= \frac{l_B^2}{2\hbar^2} [\hat{\pi}_x - i\hat{\pi}_y, \hat{\pi}_x + i\hat{\pi}_y]$$

$$= \frac{l_B^2}{2\hbar^2} ([\hat{\pi}_x, \hat{\pi}_x] - i[\hat{\pi}_x, i\hat{\pi}_y] - i[\hat{\pi}_y, \hat{\pi}_x] - i[\hat{\pi}_y, i\hat{\pi}_y])$$

$$= \frac{l_B^2}{2\hbar^2} ((i[\hat{\pi}_x, \hat{\pi}_y]) - (i[\hat{\pi}_y, \hat{\pi}_x]))$$

$$= \frac{l_B^2}{2\hbar^2} \left(-i^2 \frac{\hbar^2}{l_B^2} - i^2 \frac{\hbar^2}{l_B^2} \right) = \frac{l_B^2}{2\hbar^2} \left(\frac{2\hbar^2}{l_B^2} \right) = 1 \Rightarrow \underline{\underline{\text{QED}}}$$

$$[\hat{b}, \hat{b}^\dagger] = \left[\frac{l_B}{\sqrt{2\hbar}} (\hat{k}_y - i\hat{u}_x), \frac{l_B}{\sqrt{2\hbar}} (\hat{k}_y + i\hat{u}_x) \right]$$

$$= \frac{l_B^2}{2\hbar^2} [(\hat{k}_y - i\hat{u}_x), (\hat{k}_y + i\hat{u}_x)]$$

$$= \frac{l_B^2}{2\hbar^2} ([\hat{k}_y, \hat{k}_y] - [\hat{k}_y, i\hat{u}_x] + [\hat{u}_x, \hat{k}_y] - [\hat{u}_x, i\hat{u}_x])$$

$$= \frac{l_B^2}{2\hbar^2} (i[\hat{k}_y, \hat{u}_x] + i[\hat{u}_x, \hat{k}_y])$$

$$= \frac{l_B^2}{2\hbar^2} (-i \cdot \frac{\hbar^2}{l_B^2} - i \cdot \frac{\hbar^2}{l_B^2}) = \frac{l_B^2}{2\hbar^2} \left(\frac{2\hbar^2}{l_B^2} \right) = 1 \Rightarrow \underline{\underline{\text{QED}}}$$

The remaining commutation relations $[\hat{a}, \hat{b}]$, $[\hat{a}, \hat{b}^\dagger]$, $[\hat{a}^\dagger, \hat{b}]$, $[\hat{a}^\dagger, \hat{b}^\dagger]$ all involve commutators of \hat{k} and $\hat{\pi}$, which are equal to zero. Hence, all of the above relations $\Rightarrow [\hat{a}, \hat{b}] = [\hat{a}, \hat{b}^\dagger] = [\hat{a}^\dagger, \hat{b}] = [\hat{a}^\dagger, \hat{b}^\dagger] = 0$

c. Let $\hat{N}_a = \hat{a}^\dagger \hat{a}$ and $\hat{N}_b = \hat{b}^\dagger \hat{b}$

Want to show that \hat{H} can be presented as $\hat{H} = \hbar\omega(\hat{a}^\dagger \hat{a} + \frac{1}{2})$

$$= \hbar\omega \left[\frac{l_B}{\sqrt{2\hbar}} (\hat{\pi}_x + i\hat{\pi}_y) \frac{l_B}{\sqrt{2\hbar}} (\hat{\pi}_x - i\hat{\pi}_y) + \frac{1}{2} \right] = \hat{H} \quad (\text{where the definitions of } \hat{a} \text{ and } \hat{a}^\dagger \text{ have been substituted into the expression})$$

$$= \frac{\hbar\omega l_B^2}{2\hbar^2} [(\hat{\pi}_x + i\hat{\pi}_y)(\hat{\pi}_x - i\hat{\pi}_y) + \frac{1}{2}]$$

$$= \frac{\hbar\omega l_B^2}{2\hbar^2} [\hat{\pi}_x^2 - i\hat{\pi}_x i\hat{\pi}_y + i\hat{\pi}_y \hat{\pi}_x + \hat{\pi}_y^2 + \frac{1}{2}]$$

$$= \frac{\hbar\omega l_B^2}{2\hbar^2} [\hat{\pi}_x^2 - i \left(-\frac{i\hbar^2}{l_B^2} \right) + \hat{\pi}_y^2 + \frac{1}{2}]$$

$$= \frac{\hbar\omega l_B^2}{2\hbar^2} [\hat{\pi}_x^2 - \frac{\hbar^2}{l_B^2} + \hat{\pi}_y^2 + \frac{1}{2}]$$

$$= \omega [\hat{\pi}_x^2 - \frac{1}{2} + \frac{1}{2} + \hat{\pi}_y^2] \cdot \frac{l_B^2}{2\hbar^2} = \frac{\hbar c}{2m} \cdot \frac{1}{2} = \frac{c}{2m}$$

$$\therefore \frac{c}{2m} \cdot \frac{c}{2m} = \frac{c}{2m}$$

Set $c = 1$

$$= \hat{H} = \frac{1}{2m} [\hat{\pi}_x^2 + \hat{\pi}_y^2] = \frac{1}{2m} ((p_x + \frac{e}{c} A_x)^2 + (p_y + \frac{e}{c} A_y)^2)$$

$$= \frac{1}{2m} (p_x^2 + p_y^2 + (eA_x)^2 + (eA_y)^2) = \hat{H} \Rightarrow \underline{\underline{\text{QED}}}$$

d. Since $[\hat{H}, \hat{N}_a] = 0$ and $[\hat{H}, \hat{N}_b] = 0$, this implies that $[\hat{N}_a, \hat{N}_b] = 0$ and thus, the eigenvalues of both of these operators constitute "good quantum numbers" and can therefore represent simultaneously measurable observables (i.e. n and l represent observables that can be measured simultaneously). One can write the state $|n, l\rangle$ as follows:

$$|n, l\rangle = |n\rangle \otimes |l\rangle$$

$$\hat{N}_a |n, l\rangle = \hat{N}_a |n\rangle \otimes \hat{N}_b |l\rangle = n |n, l\rangle$$

$$\hat{N}_b |n, l\rangle = |n\rangle \otimes \hat{N}_b |l\rangle = l |n, l\rangle$$

One can note that acting on the ground state with \hat{a} or \hat{a}^\dagger = 0

$$\hat{a}|0, 0\rangle = \hat{a}^\dagger |0, 0\rangle = 0$$

Expressing $|n\rangle$ as:

$$|n\rangle = \frac{(\hat{a}^\dagger)^n}{\sqrt{n!}} |0\rangle$$

and similarly for $|l\rangle$:

$$|l\rangle = \frac{(\hat{b}^\dagger)^l}{\sqrt{l!}} |0\rangle$$

One can rewrite the state $|n, l\rangle$ as:

$$|n, l\rangle = \frac{(\hat{a}^\dagger)^n}{\sqrt{n!}} |0\rangle \otimes \frac{(\hat{b}^\dagger)^l}{\sqrt{l!}} |0\rangle = \frac{(\hat{a}^\dagger)^n (\hat{b}^\dagger)^l}{\sqrt{n! l!}} |0, 0\rangle$$

Hence, the spectrum can be constructed with the operators \hat{a}^\dagger and \hat{b}^\dagger