

An Extension Of The Quadric Model

Spencer T. Parkin

Abstract. An extension of the quadric model set forth in [3] is found in which the rigid body motions are represented by versors applicable to any quadric surface. Extending the model yet further to include a specific form of quartic surface, we find that the versors of the conformal model may be used to transform any such surface. Results of a computer program implementing this model are presented.

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1. Introduction

In the original paper [3], a model for quadric surfaces was presented based upon the ideas of projective geometry. What was unfortunate about this model, however, was its lack of support for the rigid body transformations. It was predicted in the conclusion of that paper that a better model for quadric surfaces may exist that is more like the conformal model of geometric algebra. The present paper details what may be such a model. We'll find that the rigid body transformations can be incorporated into the model by using an alternative method of encoding the quadric form. An extension to the quadric form will then allow us to support the conformal transformations at the expense of expanding our model to necessarily include a specific form of quartic surfaces. Both extensions of the original quadric model will use the same geometric algebra.

2. The Geometric Algebra

We begin here with a description of the structure of the geometric algebra upon which our model will be imposed. This geometric algebra will contain

the following vector spaces.

Notation	Basis	
\mathbb{V}^e	$\{e_i\}_{i=1}^n$	(2.1)
\mathbb{V}^o	$\{o\} \cup \{e_i\}_{i=1}^n$	
\mathbb{V}^∞	$\{e_i\}_{i=1}^n \cup \{\infty\}$	
\mathbb{V}	$\{o\} \cup \{e_i\}_{i=1}^n \cup \{\infty\}$	

The set of vectors $\{e_i\}_{i=1}^n$ forms an orthonormal set of basis vectors for the n -dimensional Euclidean vector space \mathbb{V}^e , which we'll use to represent n -dimensional Euclidean space. The vectors o and ∞ are the familiar null-vectors representing the points at origin and infinity taken from the conformal model of geometric algebra. An inner-product table for these basis vectors is given as follows, where $1 \leq i < j \leq n$.

\cdot	o	e_i	e_j	∞	
o	0	0	0	-1	(2.2)
e_i	0	1	0	0	
e_j	0	0	1	0	
∞	-1	0	0	0	

We will now let $\mathbb{G}(\mathbb{V})$ denote the Minkowski geometric algebra generated by \mathbb{V} . For each vector space in table (2.1), we will let an over-bar above this vector space denote an identical copy of that vector space. The vector space \mathbb{W} will denote the smallest vector space containing each of \mathbb{V} and $\bar{\mathbb{V}}$ as vector subspaces. In symbols, one may write

$$\mathbb{G}(\mathbb{W}) = \mathbb{G}(\mathbb{V} \oplus \bar{\mathbb{V}}) \quad (2.3)$$

to illustrate the structure of $\mathbb{G}(\mathbb{W})$ in terms of its two isomorphic Minkowski geometric sub-algebras $\mathbb{G}(\mathbb{V})$ and $\mathbb{G}(\bar{\mathbb{V}})$.

We will use over-bar notation to distinguish between vectors taken from \mathbb{V} with vectors taken from $\bar{\mathbb{V}}$. Though not necessary, we can work exclusively in $\mathbb{G}(\mathbb{V})$ by defining the over-bar notation as an outermorphic isomorphism between $\mathbb{G}(\mathbb{V})$ and $\mathbb{G}(\bar{\mathbb{V}})$. Doing so, we see that for any element $E \in \mathbb{G}(\mathbb{V})$, we may define $\bar{E} \in \mathbb{G}(\bar{\mathbb{V}})$ as

$$\bar{E} = SE\tilde{S}, \quad (2.4)$$

where S is the versor given by

$$S = \left(\frac{1}{\sqrt{2}} \right)^{n+2} (1 + e_- \bar{e}_-)(1 - e_+ \bar{e}_+) \prod_{i=1}^n (1 - e_i \bar{e}_i). \quad (2.5)$$

This definition is non-circular if we let the over-bars in equation (2.5) be purely notation. The vectors e_- and e_+ , taken from [2], are defined as

$$e_- = \frac{1}{2}\infty + o, \quad (2.6)$$

$$e_+ = \frac{1}{2}\infty - o. \quad (2.7)$$

The vectors \bar{e}_- and \bar{e}_+ are defined similarly in terms of \bar{o} and $\bar{\infty}$. Defined this way, it is important to realize that, unlike the over-bar function defined

in [3], here we do not have the property that for any vector $w \in \mathbb{W}$, we have $\overline{\overline{w}} = w$. This is because $\overline{\overline{o}} = -o$ and $\overline{\overline{\infty}} = -\infty$.

3. The Form Of Quadric Surfaces In $\mathbb{G}(\mathbb{W})$

We now give a formal definition under which elements $E \in \mathbb{G}(\mathbb{W})$ are representative of n -dimensional quadric surfaces in the extended model.

Definition 3.1. Referring to an element $E \in \mathbb{G}(\mathbb{W})$ as a quadric surface, it is representative of such an n -dimensional surface as the set of all points $p \in \mathbb{V}^o$ such that

$$0 = p \wedge \overline{p} \cdot E. \quad (3.1)$$

From this definition it can be seen that the general form of a quadric $E \in \mathbb{G}(\mathbb{W})$ is given by

$$E = \sum_{i=1}^k a_i \wedge \overline{b}_i, \quad (3.2)$$

where each of $\{a_i\}_{i=1}^k$ and $\{b_i\}_{i=1}^k$ is a sequence of k vectors taken from \mathbb{V}^∞ . To see why, realize that the form (3.2) can always be reduced to the form

$$E = \sum_{i=1}^n \sum_{j=i}^n \lambda_{ij} e_i \overline{e}_j + \sum_{i=1}^n \lambda_i e_i \overline{\infty} + \lambda \infty \overline{\infty}, \quad (3.3)$$

where each of λ_{ij} , λ_i , and λ are scalars, in the sense that this reduced form represents the same surface as that in equation (3.2) under Definition 3.1. We then see that this form (3.3), when it is substituted into equation (3.1), produces a polynomial equation of degree 2 in the vector components of $p - o$. Doing so with $p = o + x$, where $x \in \mathbb{V}^e$, we get the equation

$$0 = - \sum_{i=1}^n \sum_{j=i}^n \lambda_{ij} (x \cdot e_i)(x \cdot e_j) + \sum_{i=1}^n \lambda_i (x \cdot e_i) - \lambda, \quad (3.4)$$

which we may recognize as the equation of an n -dimensional quadric surface.

In practice, a computer program might take such a bivector of the form (3.2) and extract from it the coefficients of the quadric polynomial (3.4) it represents. It could then render the surface using traditional methods, such as those used to render the traced surfaces in Figure 1 far below.

Of course, using geometric algebra on paper, it is undesirable and unnecessary to think of quadrics in terms of polynomial equations. A better way to think of quadrics is in terms of an element of the algebra whose decomposition produces the parameters characterizing the quadric surface. For example, many common quadrics are the solution set in \mathbb{V}^e of the equation

$$0 = -r^2 + (x - c)^2 + \lambda((x - c) \cdot v)^2, \quad (3.5)$$

in the variable x . (An explanation of the parameters r , c , v and λ was given in [3].) Then, factoring out $-p \wedge \overline{p}$, we see that the element $E \in \mathbb{G}(\mathbb{W})$, given by

$$\Omega + \lambda v \overline{v} + 2(c + \lambda(c \cdot v)v)\overline{\infty} + (c^2 + \lambda(c \cdot v)^2 - r^2)\infty \overline{\infty} \quad (3.6)$$

is representative of this very same quadric by Definition 3.1, where Ω is defined as

$$\Omega = \sum_{i=1}^n e_i \bar{e}_i. \quad (3.7)$$

Canonical forms similar to (3.6) can be found in [3] that are easily translated into the extended model.

4. Transformations Supported By The Model

The main result of this section will depend upon the following lemma.

Lemma 4.1. *For any versor $V \in \mathbb{G}(\mathbb{W})$, and any four vectors $a, b, c, d \in \mathbb{V}$, we have*

$$V^{-1}aV \wedge \overline{V^{-1}bV} \cdot c \wedge \bar{d} = a \wedge \bar{b} \cdot V\bar{V}(c \wedge \bar{d})(V\bar{V})^{-1}. \quad (4.1)$$

Proof. We begin by first establishing that

$$V^{-1}aV \wedge \overline{V^{-1}bV} \cdot c \wedge \bar{d} \quad (4.2)$$

$$= -(V^{-1}aV \cdot c)(V^{-1}bV \cdot d) \quad (4.3)$$

$$= -(a \cdot VcV^{-1})(b \cdot VdV^{-1}) \quad (4.4)$$

$$= a \wedge \bar{b} \cdot VcV^{-1} \wedge \overline{VdV^{-1}}. \quad (4.5)$$

We now notice that

$$VcV^{-1} \quad (4.6)$$

$$= V\bar{V}V^{-1}cV^{-1} \quad (4.7)$$

$$= (-1)^m V\bar{V}c\bar{V}^{-1}V^{-1} \quad (4.8)$$

$$= (-1)^m V\bar{V}c(V\bar{V})^{-1}, \quad (4.9)$$

where m is the number of vectors taken together in a geometric product to form V . We then notice that

$$\overline{VdV^{-1}} \quad (4.10)$$

$$= VV^{-1}\bar{V}dV^{-1} \quad (4.11)$$

$$= (-1)^{m^2} V\bar{V}V^{-1}\bar{d}V^{-1} \quad (4.12)$$

$$= (-1)^{m^2+m} V\bar{V}dV^{-1}\bar{V}^{-1} \quad (4.13)$$

$$= (-1)^{2m^2+m} V\bar{V}dV^{-1}V^{-1} \quad (4.14)$$

$$= (-1)^m V\bar{V}d(V\bar{V})^{-1}. \quad (4.15)$$

It now follows that

$$a \wedge \bar{b} \cdot VcV^{-1} \wedge \overline{VdV^{-1}} = a \wedge \bar{b} \cdot V\bar{V}(c \wedge \bar{d})(V\bar{V})^{-1}. \quad (4.16)$$

□

We're now ready to prove the main result as follows.

Theorem 4.2. *Let $E \in \mathbb{G}(\mathbb{W})$ be a bivector of the form (3.2). Let $p, p' \in \mathbb{V}^o$ be a pair of points related by a versor $V \in \mathbb{G}(\mathbb{V})$ by the equation*

$$p' = o \cdot V^{-1} p V \wedge \infty. \quad (4.17)$$

Now let $E' \in \mathbb{G}(\mathbb{W})$ be a bivector given by

$$E' = V \bar{V} E (V \bar{V})^{-1}. \quad (4.18)$$

Then, if E' is of the form (3.2), then the set of all points $p \in \mathbb{V}^o$ such that

$$0 = p' \wedge p' \cdot E \quad (4.19)$$

is exactly the set of all points $p \in \mathbb{V}^o$ such that

$$0 = p \wedge p \cdot E'. \quad (4.20)$$

Proof. The theorem goes through by the following chain of equalities.

$$(o \cdot V^{-1} p V \wedge \infty) \wedge (\overline{o \cdot V^{-1} p V \wedge \infty}) \cdot E \quad (4.21)$$

$$= V^{-1} p V \wedge \overline{V^{-1} p V} \cdot E \quad (4.22)$$

$$= p \wedge \bar{p} \cdot (V \bar{V}) E (V \bar{V})^{-1}. \quad (4.23)$$

The first equality holds by the fact that E is of the form (3.2), while the second equality holds by Lemma 4.1. \square

The key motivation behind Theorem 4.2 is the observation that the desired transformation of E by V is given by the algebraic variety of equation (4.19), because an understanding of how V^{-1} transforms p gives us an understanding of what type of geometry we get from equation (4.19) in terms of E and V . The theorem then shows that this is also the algebraic variety of equation (4.20), thereby giving us a means of performing desired transformations on elements representative of quadric surfaces in $\mathbb{G}(\mathbb{W})$.

We can now apply Theorem 4.2 to show that the rigid body transformations are supported by our extended model. Letting $\pi \in \mathbb{V}$ be a dual plane of the conformal model, given by

$$\pi = v + (c \cdot v) \infty, \quad (4.24)$$

where $v \in \mathbb{V}^e$ is a unit-length vector indicating the norm of the plane, and where $c \in \mathbb{V}^e$ is a vector representing a point on the plane, we see that for any homogenized point $p \in \mathbb{V}^o$, we have

$$-\pi p \pi^{-1} = o + x - 2((x - c) \cdot v)v + \lambda \infty, \quad (4.25)$$

where $p = o + x$ with $x \in \mathbb{V}^e$, and where the scalar $\lambda \in \mathbb{R}$ is of no consequence. Letting $V = \pi$, the point $p' \in \mathbb{V}^o$ of consequence here is given by equation (4.17), from which we can recognize an orthogonal reflection about the plane π . It now follows by Theorem 4.2 that $\pi \bar{\pi}$ is a versor capable of reflecting any quadric surface about the plane π . Being able to perform planar reflections of any quadric in any plane, it now follows that we can always find a versor $V \in \mathbb{G}(\mathbb{W})$ capable of performing any rigid body motion applicable to any quadric surface. The development of the rigid body motions, (combinations

of translations and rotations), by planar reflections, is well known, and can be found in section 2.7 of [2].

Notice that not all versors of the conformal model are applicable in our extended model of quadric surfaces. This is because they fail to satisfy the condition of Theorem 4.2 that E' be of the form (3.2).

5. Extending The Model Yet Further

Interestingly, if we were not content with the rigid by motions of quadrics, then we really could find what is, for example, the spherical reflection of, say, an infinitely long cylinder in a sphere. To do this, we change Definition 3.1 into the following definition.

Definition 5.1. For any element $E \in \mathbb{G}(\mathbb{W})$, we may refer to it as an n -dimensional quartic surface as the set of all points $p \in \mathbb{V}^e$ such that

$$0 = P(p) \wedge \overline{P}(p) \cdot E, \quad (5.1)$$

where $P : \mathbb{V}^e \rightarrow \mathbb{V}$ is the conformal mapping, defined as

$$P(p) = o + p + \frac{1}{2}p^2\infty. \quad (5.2)$$

A version of Theorem 4.2 is then easily found such that if $V \in \mathbb{G}(\mathbb{W})$ is any versor of the conformal model, and if E is a surface under Definition 5.1, then the element $E' \in \mathbb{G}(\mathbb{W})$, given by equation (4.18), must, by Definition 5.1, be representative of the desired transformation of E by the versor V . The general polynomial equation arising from the form of such elements E in Definition 5.1 is much more involved than what we have in equation (3.4). Nevertheless, it is possible to extract a specific form of a quartic polynomial equation in the vector components of p from equation (5.1). The result being unsightly, it will not be presented here. Suffice it to say, a computerized algebraic system was used to find the polynomial form. In any case, it is easy to see from equation (5.1) that the degree of the resulting polynomial will be 4.

Now notice that under Definition 5.1, canonical forms such as (3.6) are still valid. This is because

$$P(p) \wedge \overline{P}(p) \cdot E = (o + p) \wedge \overline{(o + p)} \cdot E \quad (5.3)$$

in the case that E is of the form (3.2). This allows us to use what we already know about quadrics in the old model with the new model.

Putting theory into practice, the author wrote a piece of computer software that implements this conformal-like model for the special class of quartic surfaces of equation (5.1). Giving the program the following script as input, the output of the program is given in Figure 1. The script is easy for anyone to read, even if they are not familiar with its language. It is given here to illustrate how one might use the model with the aide a computer system.

```
/*
* Calculate the surface that is the
* reflection of a cylinder in a sphere.
```

```

*/
do
(
  /* Make the cylinder. */
  v = e2,
  c = -7*e1,
  r = 2,
  cylinder = Omega - v~bar(v) + 2*c*nib + (c.c - r*r)*ni~nib,
  bind_quadric(cylinder),
  geo_color(cylinder,0,1,0),

  /* Make the sphere. */
  c = 0,
  r = 6,
  sphere = no + c + 0.5*(c.c - r*r)*ni,
  bind_dual_sphere(sphere),
  geo_color(sphere,1,0,0,0.2),

  /* Make the reflection of the cylinder in the sphere. */
  V = sphere*bar(sphere),
  reflection = V*cylinder*V~,
  bind_conformal_quartic(reflection),
  geo_color(reflection,0,0,1),
)

```

The functions beginning with the word “bind” create and bind an entity to the given element of the geometric algebra that is responsible for interpreting that element as a surface under Definition 5.1 or as a surface under the definition given by the conformal model. The computer program can then use traditional methods to render the surface from the extracted polynomial equation. For example, the polynomial equation in x , y and z for the reflected surface presented in Figure 1 is given by

$$\begin{aligned}
 0 = & 28.8x^2 + 11.2x^3 + x^4 + 11.2xy^2 + 2x^2y^2 + \\
 & 11.2xz^2 + 2x^2z^2 + y^4 + 2y^2z^2 + 28.8z^2 + z^4.
 \end{aligned} \tag{5.4}$$

It is interesting how a bit of reasoning in geometric algebra has given us a means to finding this polynomial equation. Of course, while such equations lend themselves to computer algorithms, they are not practical on paper. This is where the canonical forms of elements become useful. A parameterization of the reflected surface of Figure 1 in terms of lengths, vectors and positions, however, does not seem obvious, but is plausible and surely exists.

References

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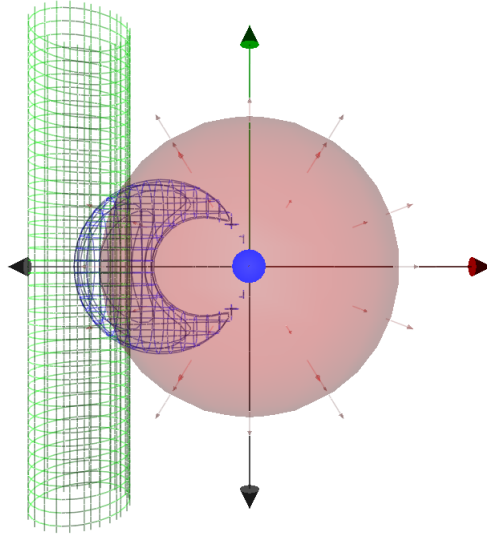


FIGURE 1. The blue reflection of a green cylinder in a red sphere. Traces in various planes were used to render the cylinder and reflection surfaces. Numerical instability in the tracing algorithm leads to problems in drawing the part of the reflected surface near the origin.

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Spencer T. Parkin
 2113 S. Claremont Dr.
 Bountiful, Utah 84010
 USA
 e-mail: spencer.parkin@gmail.com