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## Turbulence Assignment 1

**VERY IMPORTANT:** Submit this assignment digitally on canvas before **23:59h on December 1st, 2023**. You are allowed to work in pairs: 2 people submit 1 solution. **Show all of your steps in the derivations, state the assumptions you make, and provide source code of numerically solved tasks!** Ensure that your solutions are clear and legible. Unfortunately, if the grader cannot understand your process, you will not receive any points. Reminder: Physical quantities must have units and figures must have appropriate labels and axes (i.e. logarithmic, if applicable)..

### 1 Turbulent statistics (2.0 points)

Turbulence is characterized by the chaotic and irregular motion of flow. We have provided measurement data of a turbulent flow which you will use to calculate various statistics. Download `velocity_A1.Ex1.h5` from Canvas. The file contains data of a three-dimensional velocity field.

Reading in the data from `.h5`-files is fairly simple. We have provided some code snippets on the [online cheatsheet](#) to help you get started.

*Note: The velocity field in the `h5`-files is given as  $\mathbf{u}(x, y, z) = (u, v, w) = (v_x, v_y, v_z)$  such that  $(N_x, N_y, N_z) = (192, 96, 257)$  and  $(L_x, L_y, L_z) = (2.0, 1.0, 1.0)$ . Here, the  $z$  axis is aligned with the vertical direction whereas  $x$  and  $y$  are the horizontal axes.  $N$  indicates the number of grid points along an axis and  $L$  indicates the length of the domain along the axis.*

a) Make 2D contour plots of all velocity components at the horizontal plane at  $z = 0.5$ . Determine whether the fields satisfy periodic boundary conditions.

b) Compute the fluctuations  $u'(z)$ ,  $v'(z)$  and  $w'(z)$  at the location  $(x = 1.0, y = 0.5)$  and plot their profiles as function of height  $z$ . The fluctuations are defined relative to the horizontal ( $xy$ -)average, such that for example

$$u'(z) = u - \bar{u}(z), \quad \text{where} \quad \bar{u}(z) = \frac{1}{L_x L_y} \iint_A u(x, y, z) dx dy$$

Throughout this problem, we shall treat the horizontal average as the relevant mean flow.

c) Calculate the root-mean-square (RMS) of  $u$ ,  $v$ ,  $w$  and the turbulent kinetic energy at each of the horizontal planes in the velocity data and plot their profiles as function of height  $z$ .

d) Evaluate the PDFs of  $u'$  at the horizontal planes given by  $z = 0.5$  and  $z = 0.06$  and plot it. Compare the both PDFs with the PDF of Gaussian distribution.

e) Write down expressions for the standard deviation, skewness, and kurtosis of  $u$  solely containing averages of powers of  $u$  (so that no expression contains a fluctuation term such as  $u'$ ). Using these expressions, calculate these quantities for  $u'$  in the two horizontal planes  $z = 0.5$  and  $z = 0.06$ . Discuss how to interpret the values of these statistics in the context of your answer for (d).

## 2 Stochastic model for large scale convection roll (1.5 point)

In Rayleigh-Bénard convection, a layer of fluid heated from below and cooled from above with temperature difference  $\Delta T$ , there is a large scale circulation (LSC) of the fluid. The LSC can be described as  $T(t, \theta) = T_0 + \delta(t) \cdot \cos(\theta - \theta_0(t))$ , where  $\delta(t)$  is the amplitude/strength of the LSC and  $\theta_0(t)$  the orientation of the LSC in the azimuthal plane  $\theta$  (Brown and Ahlers, 2007). The change in  $\delta$  is determined by the balance between the buoyancy and the drag forces and is driven by stochastic fluctuations  $f_\delta(t)$  which phenomenologically represent the action of the small-scale turbulent background. Brown and Ahlers (2007) propose the simple model

$$\dot{\delta} = \frac{\delta}{\tau_\delta} \left( 1 - \sqrt{\frac{\delta}{\delta_0}} \right) + f_\delta(t), \quad (1)$$

where  $\delta_0$  is the steady-state amplitude, and  $\tau_\delta$  is a decay time scale. The azimuthal acceleration of the LSC is damped by the rotational inertia of the LSC, and again a stochastic fluctuation term  $f_\theta(t)$  is added to obtain:

$$\ddot{\theta}_0 = -\frac{\dot{\theta}_0 \delta}{\tau_\theta \delta_0} + f_\theta(t), \quad (2)$$

where  $\tau_\theta = L^2 / (2\nu R_e)$ .

Simulate the behaviour of the LSC for one day.

- Numerically solve the ODE system (1)-(2) for  $\delta(t)$  and  $\theta_0(t)$ .
- Model the stochastic terms  $f_\delta(t)$  and  $f_\theta(t)$  with Gaussian distributed white noise. Use a standard deviation of  $\sqrt{D/h}$  where  $D_\delta = 3.5 \times 10^{-5} \text{ K}^2/\text{s}$ ,  $D_\theta = 2.5 \times 10^{-5} \text{ rad}^2/\text{s}^3$ , and  $h$  is the simulation time step.
- Take  $\delta_0 = 0.25K$  (experimental value) for the simulation.
- Use the decay time scales

$$\tau_\delta = \frac{L^2}{18\nu R_e^{1/2}} \quad \tau_\theta = \frac{L^2}{2\nu R_e} \quad (3)$$

calculated from the physical parameters

$L$	$\nu$	$R_e$
0.2476 m	$6.69 \times 10^{-3} \text{ cm}^2/\text{s}$	3700

Table 1: cell height  $L$ , viscosity  $\nu$ , equilibrium Reynolds number  $R_e$

Think about the meaning of the initial condition for  $\dot{\theta}_0(t=0)$  and  $\theta_0(t=0)$ . Choose on your own and explain your choice.

*Hint(s): You should use the intrinsic standard function **randn** of NumPy/Julia/Matlab that generates normally-distributed random numbers to model turbulence  $f_\delta(t)$  and  $f_\theta(t)$  and use Euler's method to integrate the equations instead of using standard ODE solvers.*

### 3 Lagrangian description of flow (1.5 points)

Consider a two-dimensional flow given by the stream function:

$$\psi(x, y, t) = -2x^2 \sin(2t) + y^2 \cos(t) \quad (4)$$

a) Calculate the velocity field:

$$\mathbf{U}(x, y, t) = \frac{\partial \psi}{\partial y} \mathbf{i} - \frac{\partial \psi}{\partial x} \mathbf{j} \quad (5)$$

b) At  $t = 0$  we put dye in a circle, with unit radius, around the origin. The dye can be viewed as a passive scalar. Draw the time evolution of this circle. Give snapshots of what has been the circle at  $t = 0$  for  $t = 0.1$ ,  $t = 1$ , and  $t = 10$ .

c) Now we assume that the dye is no longer a passive tracer, but it has instead a density which is much larger than the fluid. The dynamics of small dye particles of such kind can be described by the following equation of motion:

$$\frac{d\mathbf{x}(t)}{dt} = \mathbf{v}(t) \quad (6)$$

$$\frac{d\mathbf{v}(t)}{dt} = \frac{1}{\tau} (\mathbf{U}(\mathbf{x}(t), t) - \mathbf{v}(t)) \quad (7)$$

where  $\tau$  represents the typical response time of the dye-particle to the fluid flow fluctuations. We have the following initial condition for the dye

$$\mathbf{v}(0) = \mathbf{U}(\mathbf{x}(0), 0) \quad (8)$$

What is the difference in the dye distribution at  $t = 10$  for values of  $\tau = 10^{-1}$ ,  $\tau = 10^0$ ,  $\tau = 10^1$ , and  $\tau = 10^2$ ? Give the corresponding snapshots. For comparison reasons you should fix the plot ranges.

*Hint: You should make sure that your  $\Delta t$  in your integration is small enough, such that the final solution is independent of the time step  $\Delta t$ .*

### 4 Linear stability of a system of ODEs (2.5 points)

Stability analysis is widely used in fluid dynamics. Roughly speaking, it involves 3 basic steps:

1. Determining stationary solution(s)  $\mathbf{s}$  of the equations,
2. Prescribing a specific form of small perturbation to  $\mathbf{s}$ ,
3. Predicting the evolution of such perturbations.

The aim of the present exercise is to apply this strategy to a problem in which a *linear* stability analysis is feasible.

To begin with, it is convenient to introduce:

$$\mathbf{f} = \begin{pmatrix} f \\ g \\ h \end{pmatrix},$$

which allows us to write a set of ordinary differential equations in a compact form as:

$$\frac{d\mathbf{f}}{dt} = \mathbf{F}(\mathbf{f}). \quad (9)$$

Then, let  $\mathbf{f}^*(t)$  equal:

$$\mathbf{f}^*(t) = \mathbf{s}(t) + \mathbf{e}(t),$$

with  $\mathbf{s}(t)$  a stationary solution (i.e., a solution for which  $d\mathbf{s}/dt = \mathbf{F}(\mathbf{s}) = 0$ ) and  $\mathbf{e}$  a small perturbation.

To describe the temporal behavior of  $\mathbf{e}$ , we differentiate  $\mathbf{f}^* = \mathbf{s} + \mathbf{e}$ :

$$\frac{d\mathbf{f}^*}{dt} = \frac{d\mathbf{s}}{dt} + \frac{d\mathbf{e}}{dt}.$$

Now, by expanding  $\mathbf{F}(\mathbf{f}^*) = \mathbf{F}(\mathbf{s} + \mathbf{e})$  around  $\mathbf{s}$  and taking only the linear term one finds:

$$\mathbf{F}(\mathbf{s} + \mathbf{e}) = \mathbf{J}\mathbf{e}, \quad (10)$$

where  $\mathbf{J}$  is the Jacobian *matrix*, a multivariate analog of  $F'(s)$  in

$$F(s + e) = F(s) + eF'(s) + O(e^2), \quad (11)$$

with stationary solution  $F(s) = 0$ . The matrix elements of  $\mathbf{J}$  are *functions* of  $f, g$ , and  $h$  and are evaluated at  $\mathbf{s}$ :

$$\mathbf{J} = \begin{bmatrix} \frac{\partial F_1}{\partial f} & \frac{\partial F_1}{\partial g} & \frac{\partial F_1}{\partial h} \\ \frac{\partial F_2}{\partial f} & \frac{\partial F_2}{\partial g} & \frac{\partial F_2}{\partial h} \\ \frac{\partial F_3}{\partial f} & \frac{\partial F_3}{\partial g} & \frac{\partial F_3}{\partial h} \end{bmatrix}_{(s_1, s_2, s_3)}$$

Here, if  $\mathbf{e}$  is an eigenvector of  $\mathbf{J}$ , then  $\mathbf{J}\mathbf{e} = \lambda \mathbf{e}$  and hence eq. (9) becomes

$$\frac{d\mathbf{e}}{dt} = \lambda \mathbf{e}.$$

Therefore, the central aspect of the linear stability analysis is the following: if  $\lambda > 0$ , then the perturbation grows exponentially fast with time.

In this exercise we study the following non-linear ordinary differential equations analytically:

$$\frac{df}{dt} = -g - h \quad (12)$$

$$\frac{dg}{dt} = f + \alpha g \quad (13)$$

$$\frac{dh}{dt} = \beta + (f - \gamma)h, \quad (14)$$

where  $f$ ,  $g$  and  $h$  depend only on  $t$ . The equations are relatively simple, consisting of a single non-linear term only, but nevertheless show interesting features. In this exercise we will study these equations in some detail. For now we will take for the parameters  $\alpha = 0.25$ ,  $\beta = 0.5$  and let  $\gamma > 0$ .

**a)** Find analytically the stationary solutions,  $\mathbf{s}$ , and show that  $\mathbf{s}$  does not exist when  $\gamma$  is lower than some critical value. How many solutions for  $\mathbf{s}$  can you find when  $\gamma$  equals the critical value? And how many when  $\gamma$  is greater than the critical value?

**b)** Construct the Jacobian matrix.

**c)** Determine the value(s) for  $\mathbf{s}$  with the critical value of  $\gamma$  you found in (a). Do the same when  $\gamma = 1$ . What are the corresponding eigenvalues of the associated Jacobian matrix?

**d)** Discuss the stability based on results you found in (c).

## 5 Numerical study (2.5 points)

In the previous exercise we studied a system of ODE's algebraically, where we have learned about some of its local behaviour. Now we will numerically integrate these ODE's and study how trajectories  $\mathbf{f}(t)$  move throughout space.

**a)** Integrate the system of equations (12) to (14), with initial condition  $\mathbf{f}(0) = (1, 0, -2)^T$  for three different cases:  $\gamma = 2.0$ ,  $\gamma = 3.2$  and  $\gamma = 3.8$ . Keep  $\alpha$ ,  $\beta$  the same, i.e.  $\alpha = 0.25$ ,  $\beta = 0.5$ . Make a parametric plot  $(f(t), g(t))$  for these three cases and describe the solutions behaviour. You might want to skip the initial transient.

**b)** Now take  $\alpha = 0.1$ ,  $\beta = 0.1$ , and  $\gamma = 18$ , and make a 3D parametric curve of  $\mathbf{f}(t)$ . Also plot the three time series  $f(t)$ ,  $g(t)$ , and  $h(t)$  separately. Which of these signals is the most intermittent? Do you observe something unusual about the trajectory  $\mathbf{f}(t)$ ? This is known as a strange attractor.

**c)** Make a Probability Density Function (PDF) of  $f(t)$ , obtained in (b), for:

- $0 < t < 100$ ,
- $30 < t < 100$ ,
- $30 < t < 500$ ,
- $30 < t < 1000$ .

What are the differences between the histograms? Elaborate on your answer in terms of statistical stationarity.

**d)** Simulate the system again with the same conditions as in (b), but now using initial conditions  $\mathbf{f}_1(0) = (1, 10^{-8}, -2)^T$ . Make a parametric plot of the new  $\mathbf{f}_1$ , and plot the difference  $\mathbf{f}_1 - \mathbf{f}$  for each component separately between the new results and the simulated values of  $\mathbf{f}$  computed in (b). Additionally repeat the plots from part (d). What do you notice about the PDFs?

**e)** Compute the power spectrum of  $f(t)$  as computed in (b). Normalize your result such that Parseval's theorem is satisfied, that is

$$\int_0^\infty |f(t)|^2 dt = \int_0^\infty |\hat{f}(k)|^2 dk \quad (15)$$

What is the spectrum at low frequencies: sharp peaks or broad noise? Does it provide an indication of chaos?