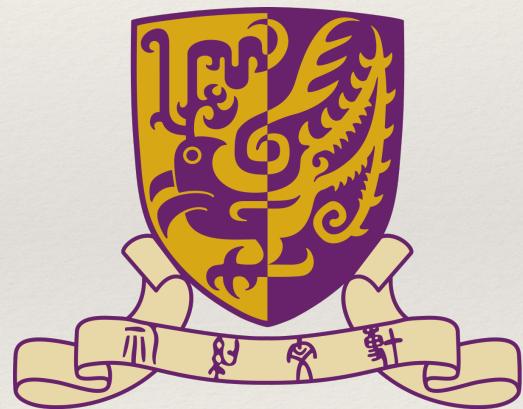
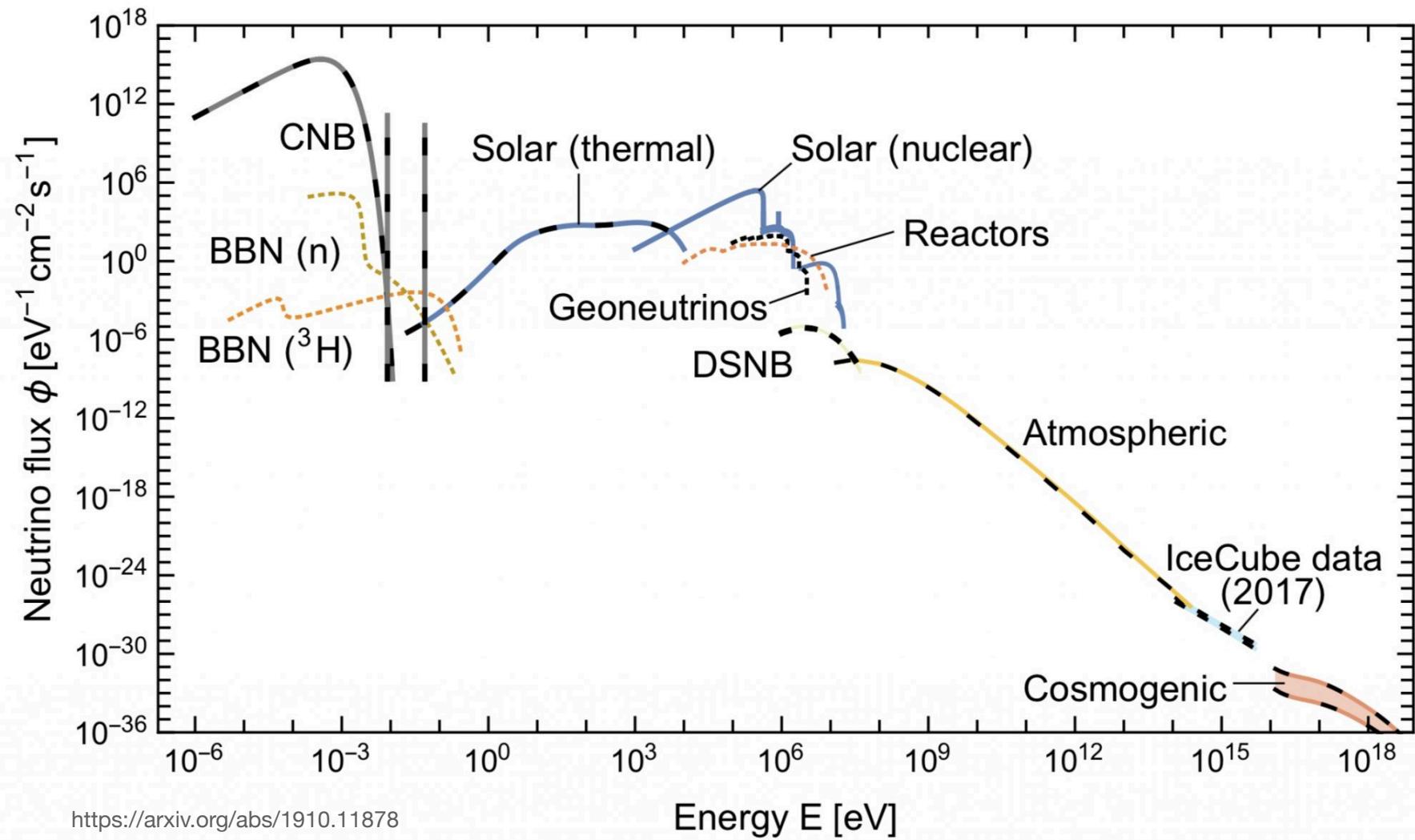


# Atmospheric Neutrinos

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- ❖ Sci Cen North Black 345
- ❖ CUHK
- ❖ Course webpage: <https://blackboard.cuhk.edu.hk>

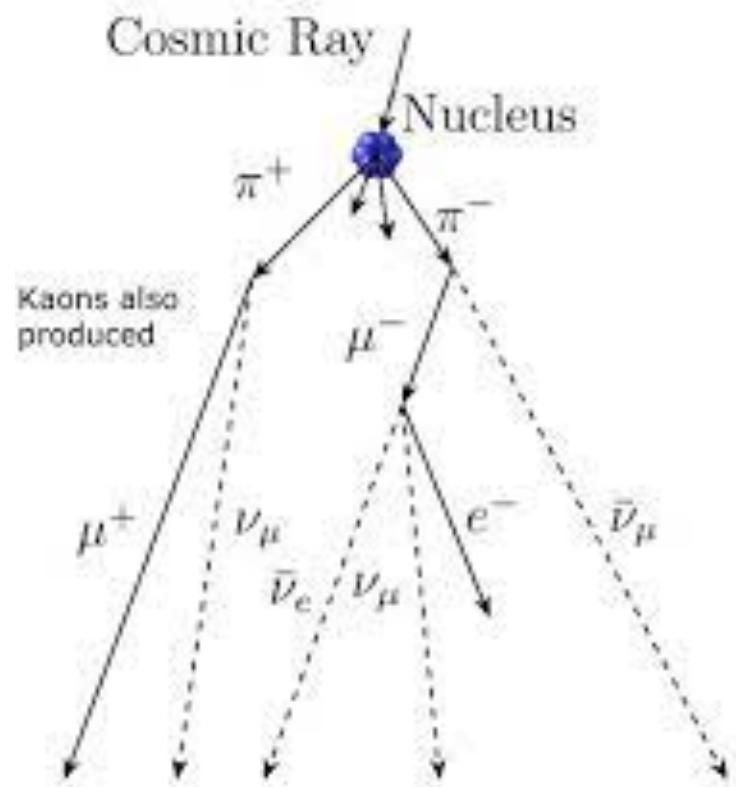


# The Grand Unified Neutrino Spectrum



# Atmospheric neutrinos

- Cosmic rays are everywhere
- They can hit the Earth atmosphere
- (CR, Proton)+ (Atmosphere, nitrogens, etc)
- Hadronic interactions
  - Pions
  - Charged pions -> Neutrinos
  - If  $\pi^+$  and  $\pi^-$  are equally produced
  - Then we should roughly expect  $\nu_\mu : \nu_e = 2 : 1$
  - Antineutrino: neutrino = 1:1



$$\frac{\nu_\mu + \bar{\nu}_\mu}{\nu_e + \bar{\nu}_e} \cong 2, \quad \frac{\nu_\mu}{\bar{\nu}_\mu} \cong \frac{\nu_e}{\bar{\nu}_e} \cong 1$$

# ATM theory

- Cosmic rays are everywhere
- They can hit the Earth atmosphere
- (CR, Proton)+ (Atmosphere, nitrogens, etc)
- Hadronic interactions
  - Pions
  - Charged pions -> Neutrinos
  - If  $\pi^+$  and  $\pi^-$  are equally produced
  - Then we should roughly expect  $\nu_\mu : \nu_e = 2 : 1$
  - Antineutrino: neutrino = 1:1

Atmospheric neutrino flux calculation using the NRLMSISE00 atmospheric model

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(Dated: May 27, 2015)

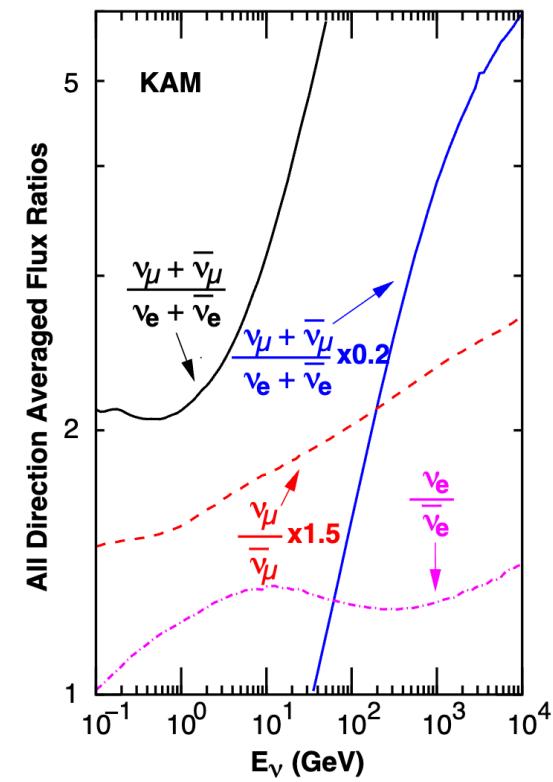
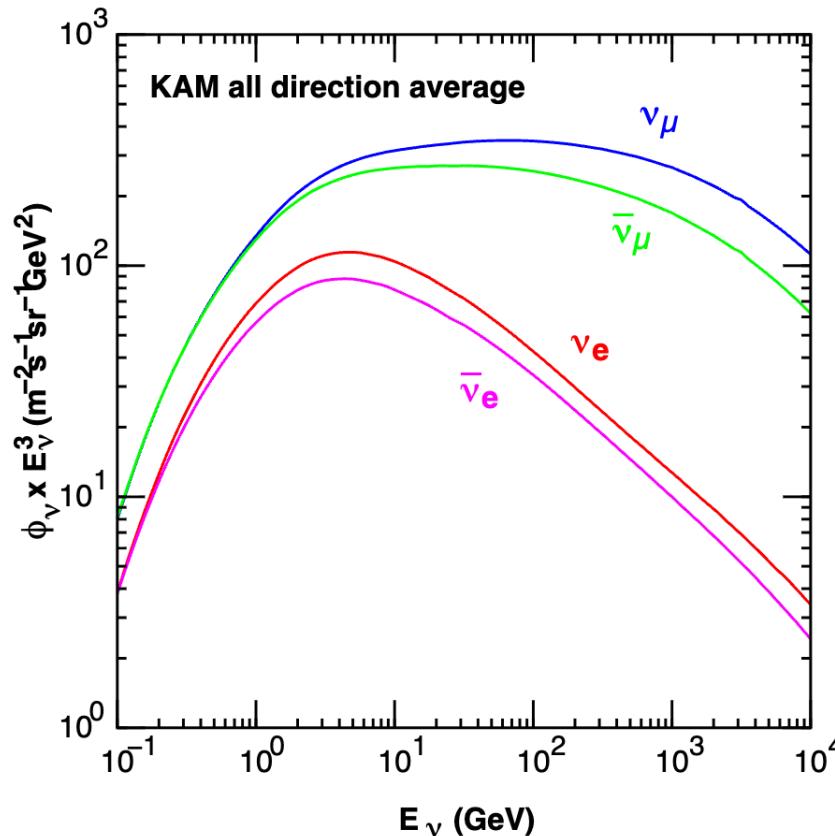
$$\pi^\pm \rightarrow \mu^\pm + \nu_\mu(\bar{\nu}_\mu)$$



$$e^\pm + \nu_e(\bar{\nu}_e) + \bar{\nu}_\mu(\nu_\mu),$$

# ATM theory

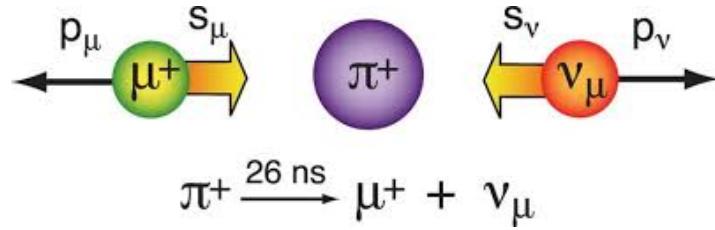
- Actual theory calculation
- $10^0 - 10^2$  GeV,  $E^{-3}$  for Muon neutrinos, then softens
- Electron neutrinos softens even faster
- (Naive expectation would be the primary cosmic-ray spectrum  $E^{-2.7}$ )
  - Then we should roughly expect  $\nu_\mu : \nu_e = 2 : 1$
  - Antineutrino: neutrino = 1:1



Clearly not what we expect!!!

# Side note on pion decay

- Charged pion decay to a muon and muon-neutrino
- It also decays into electron and electron-neutrino but with 0.01% probability.
- To understand this, first remember that only “left-handed” neutrinos participates in the weak interaction



6

## Improved Measurement of the $\pi \rightarrow e\nu$ Branching Ratio

A. Aguilar-Arevalo<sup>1</sup>, M. Aoki<sup>2</sup>, M. Blecher<sup>3</sup>, D. I. Britton<sup>4</sup>, D. A. Bryman<sup>5</sup>, D. vom Bruch<sup>5</sup>, S. Chen<sup>6</sup>, J. Comfort<sup>7</sup>, M. Ding<sup>6</sup> et al.  
(PIENU Collaboration)

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Phys. Rev. Lett. 115, 071801 – Published  
13 August, 2015

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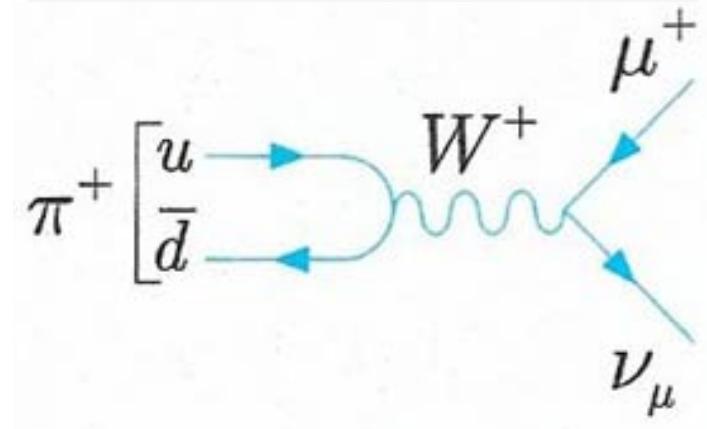
Am. score 2

Citations 71

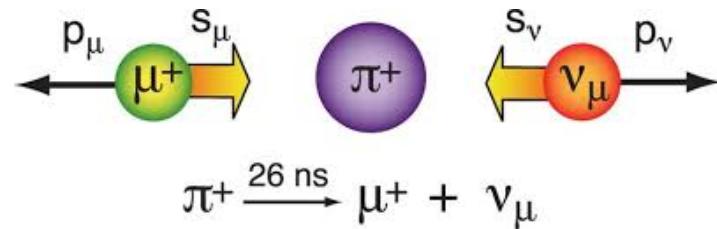
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### Abstract

A new measurement of the branching ratio  $R_{e/\mu} = \Gamma(\pi^+ \rightarrow e^+\nu + \pi^+ \rightarrow e^+\nu\gamma)/\Gamma(\pi^+ \rightarrow \mu^+\nu + \pi^+ \rightarrow \mu^+\nu\gamma)$  resulted in  $R_{e/\mu}^{\text{exp}} = [1.2344 \pm 0.0023(\text{stat}) \pm 0.0019(\text{syst})] \times 10^{-4}$ . This is in agreement with the standard model prediction and improves the test of electron-muon universality to the level of 0.1%.



# Side note on pion decay

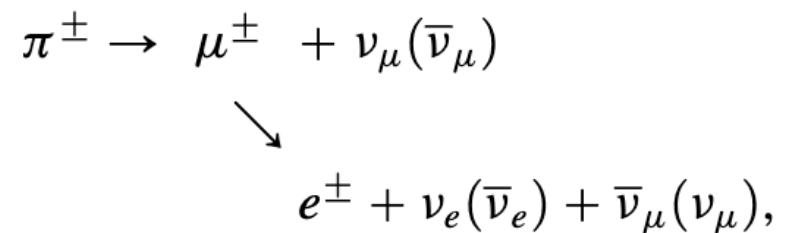


- To understand this, first remember that only “left-handed” neutrinos participates in the weak interaction
- Introduce a new concept “helicity” which is  $\vec{s} \cdot \vec{p}$ 
  - Turns out helicity = chirality when the particles are massless (Left = -1)
  - Remember this, by observing the helicity is NOT Lorentz invariant in massive case.
    - You can boost to a frame to flip the sign of momentum.
  - Chirality is Lorentz invariant (physics is Lorentz invariant)
- For Pion+ decay (Spin-0) at rest, conservation of angular momentum mean they must have opposite spin, and thus Helicity.
- So, a Negative helicity neutrino (~ left-handed), and a negative helicity anti-muon (~ left-handed)
- Luckily, the anti-muon is only (~ left-handed), so decay is still possible.
  - The larger the mass, the more deviation of Chirality from Helicity, so muon has much higher branching fraction than electrons.
  - The electron channel can also be called helicity suppressed.

$$\frac{\Gamma(\pi^+ \rightarrow e^+ \nu_e)}{\Gamma(\pi^+ \rightarrow \mu^+ \nu_\mu)} = \left( \frac{m_e^2}{m_\mu^2} \right) \left( \frac{m_\pi^2 - m_e^2}{m_\pi^2 - m_\mu^2} \right) = 1.275 \cdot 10^{-4} \quad (1.230 \pm 0.004) \cdot 10^{-4}_{PDG}$$

# ATM theory

- How to calculate ATM neutrinos
- Conceptually, the method is identical to cosmic-ray propagation. Difference are
  - Optical thick regime (cascade/shower)
  - No diffusion
  - Time evolution ( $\frac{d}{dt}$ ) -> distance evolution ( $\frac{d}{dx}$ )
  - 3D in modern calculations
  -



## A. Coupled cascade equations

The equations describe the evolution of the differential flux, defined as the differential of the particle flux  $\phi$  with respect to energy per unit area, unit solid angle, and time

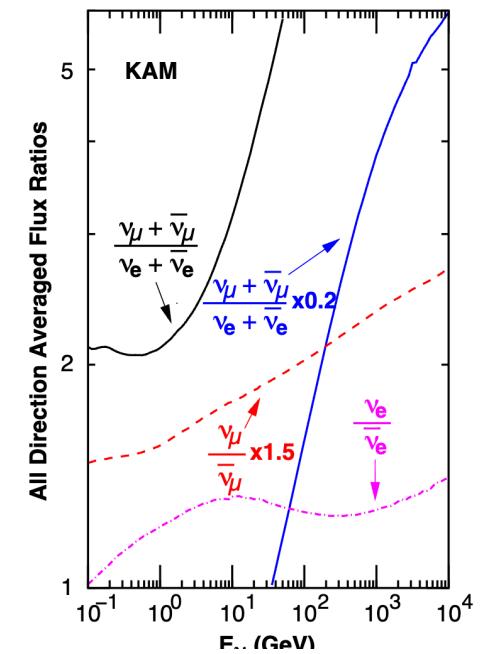
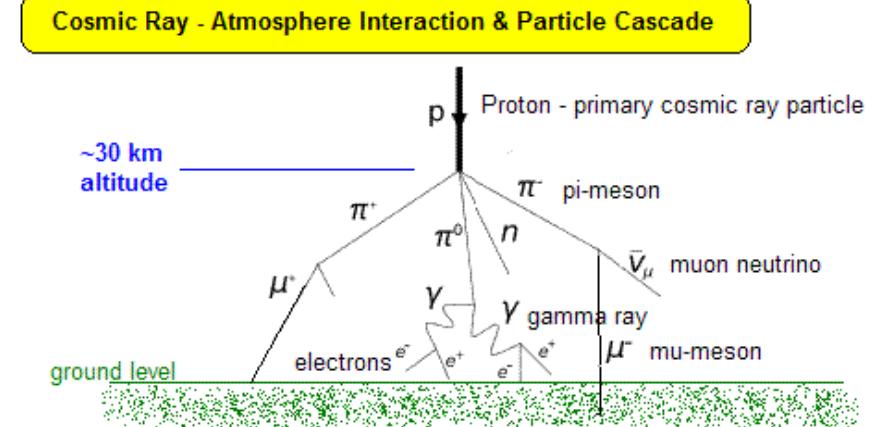
$$\Phi = \frac{d\phi}{dE} = \frac{dN}{dE \, dA \, d\Omega \, dt}. \quad (2)$$

The coupled cascade equations

$$\begin{aligned} \frac{d\Phi_h(E, X)}{dX} &= -\frac{\Phi_h(E, X)}{\lambda_{int,h}(E)} - \frac{\Phi_h(E, X)}{\lambda_{dec,h}(E, X)} \\ &\quad - \frac{\partial}{\partial E} (\mu(E)\Phi_h(E, X)) \\ &\quad + \sum_\ell \int_E^\infty dE_\ell \frac{dN_{\ell(E_\ell) \rightarrow h(E)}}{dE} \frac{\Phi_\ell(E_\ell, X)}{\lambda_{int,l}(E_\ell)} \\ &\quad + \sum_\ell \int_E^\infty dE_\ell \frac{dN_{\ell(E_\ell) \rightarrow h(E)}^{\text{dec}}}{dE} \frac{\Phi_\ell(E_\ell, X)}{\lambda_{dec,l}(E_\ell, X)}, \end{aligned} \quad (3)$$

# ATM theory

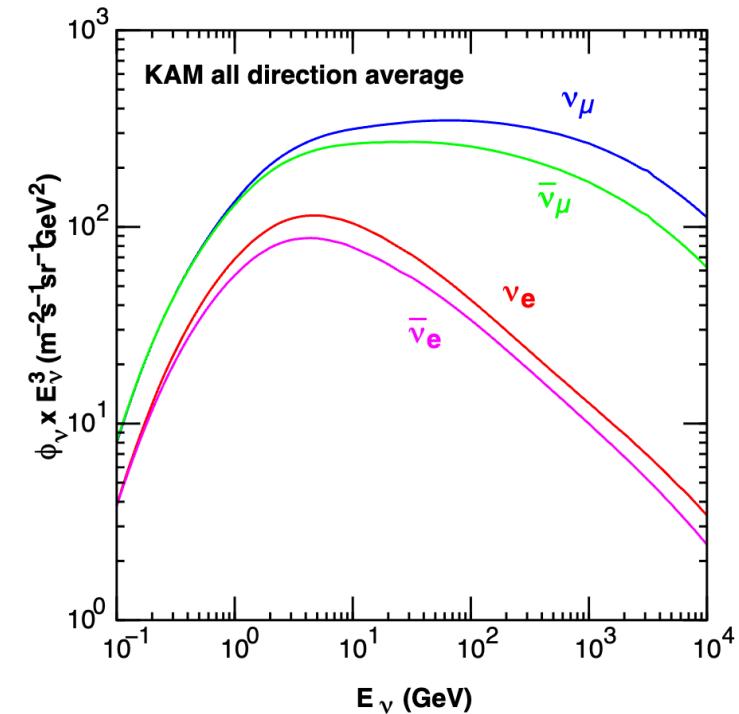
- All are known physics
- So, given the cosmic-ray flux at the top of the atmosphere, we should understand everything that goes on with atm flux.
- 1) Why electron neutrino flux are more suppressed.
- Muon decay life time = 2.2 micro second.
- $c\tau = 660m$
- If the typical interaction altitude is about 30km
- Then a muon of  $\gamma$  factor about 50, would hit the ground before decaying
- $\gamma m \sim 5 \text{ GeV}$ .



# ATM theory

- The contribution of Kaons
- Meson that contains a strange quark
- Critical energy, above that interaction > decay
  - Pions: 115 GeV
  - Kaons: 850 GeV
  - Above  $\sim 850$  GeV, mesons re-interact more than decay
  - This is why the nu-mu spectrum softens again above TeV

Particle name	Particle symbol	Antiparticle symbol	Quark content	Properties of kaons							Commonly decays to (> 5% of decays)
				Rest mass [MeV/c <sup>2</sup> ]	J <sup>G</sup>	J <sup>PC</sup>	S	C	B'	Mean lifetime [s]	
Kaon <sup>[1]</sup>	K <sup>+</sup>	K <sup>-</sup>	u s̄	493.677 ± 0.016	1/2	0 <sup>-</sup>	1	0	0	(1.2380 ± 0.0020) × 10 <sup>-8</sup>	$\mu^+ + \nu_\mu$ or $\pi^+ + \nu_\mu$ or $\pi^+ + \pi^-$ or $\pi^0 + \pi^+$ or $\pi^0 + e^- + \nu_e$
K-short <sup>[3]</sup>	K <sub>S</sub> <sup>0</sup>	self	$\frac{d\bar{s} - s\bar{d}}{\sqrt{2}}$ [t][4][5]	497.611 ± 0.013 <sup>[4]</sup>	1/2	0 <sup>-</sup>	[t]	0	0	(8.954 ± 0.004) × 10 <sup>-11</sup>	$\pi^+ + \pi^-$ or $\pi^0 + \pi^0$ or $\pi^0 + \pi^-$
K-long <sup>[6]</sup>	K <sub>L</sub> <sup>0</sup>	self	$\frac{d\bar{s} + s\bar{d}}{\sqrt{2}}$ [t][4][5]	497.611 ± 0.013 <sup>[4]</sup>	1/2	0 <sup>-</sup>	[t]	0	0	(5.116 ± 0.021) × 10 <sup>-8</sup>	$\pi^+ + e^- + \nu_e$ or $\pi^+ + \mu^- + \nu_\mu$ or $\pi^0 + \pi^0$ or $\pi^0 + \pi^-$ or $\pi^+ + \pi^-$



# ATM theory

- Angular distribution of atmospheric neutrinos

- Two extremely cases

- a) Isotropic emitter in the shell

- A collection of stars

- b) blackbody emitter in the shell  $\frac{dL}{dA} \propto \cos \theta_e$

- The case of a cavity, isotropic at the observer

Atmospheric neutrinos are in-between the two cases  
But they are always symmetric wrt horizon!

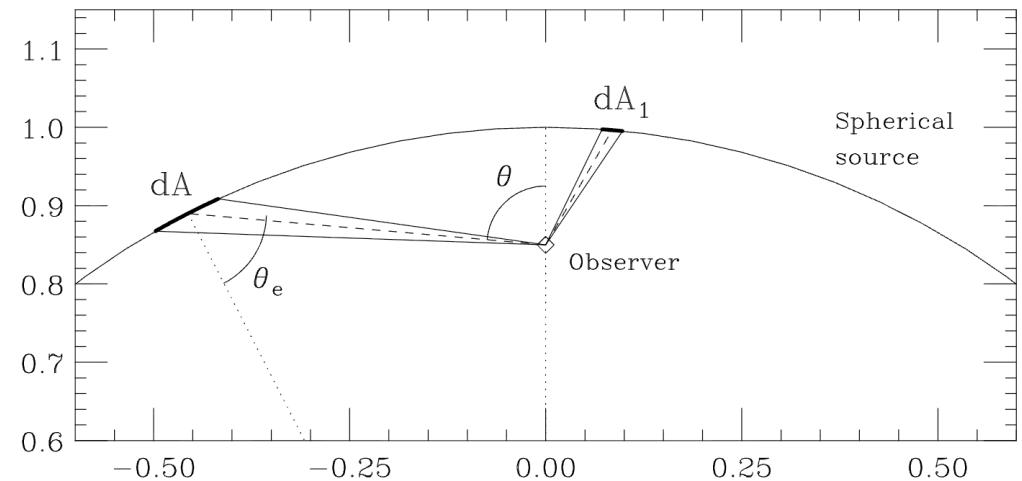
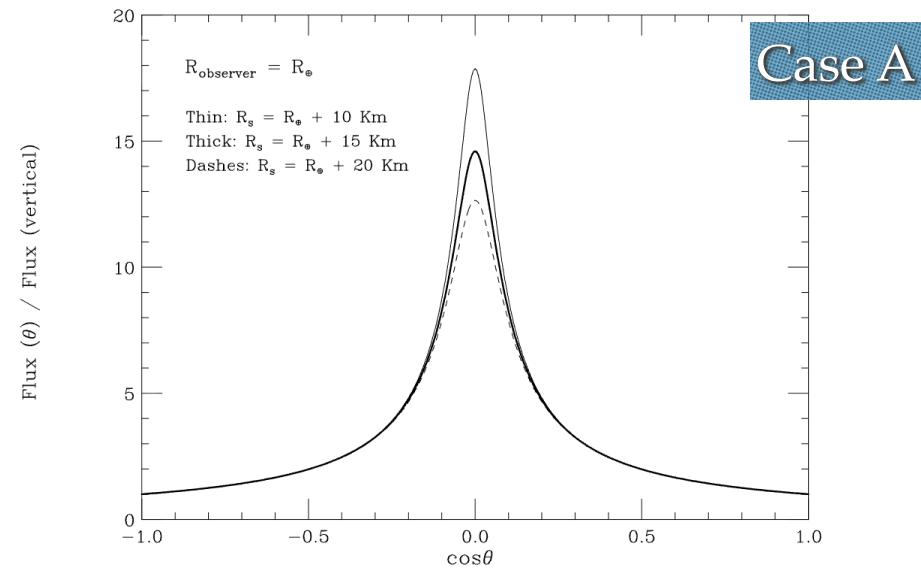
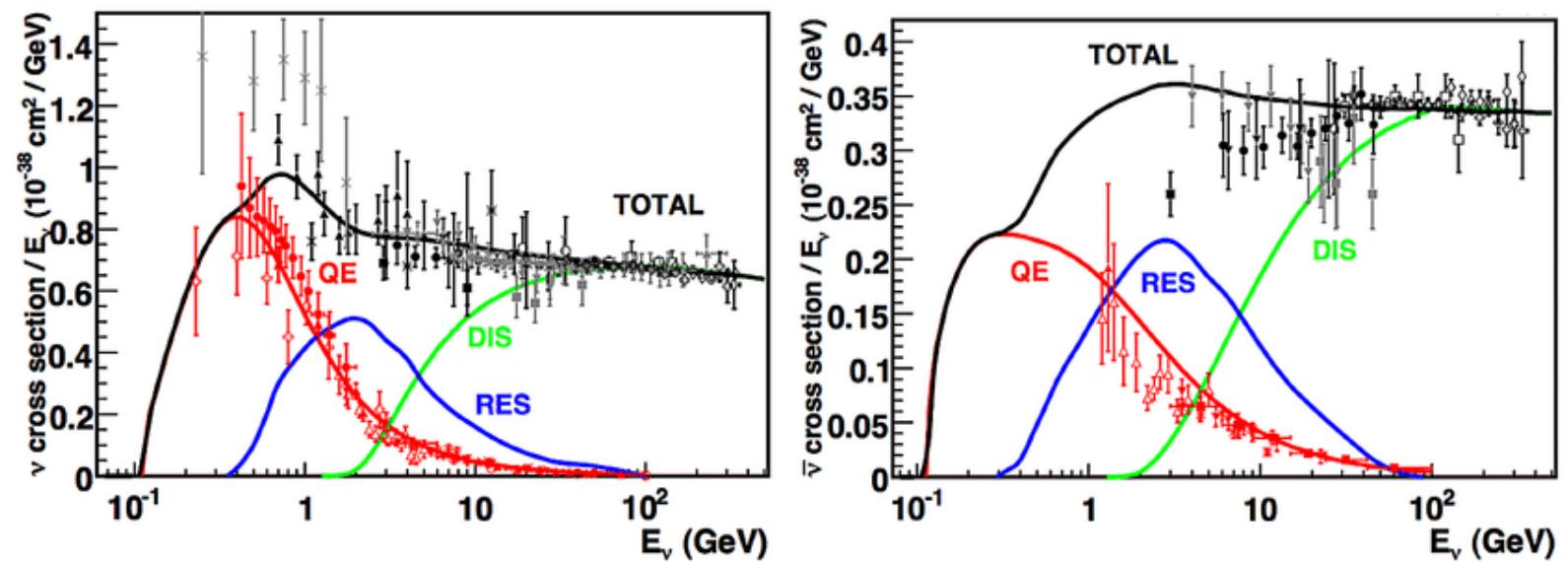
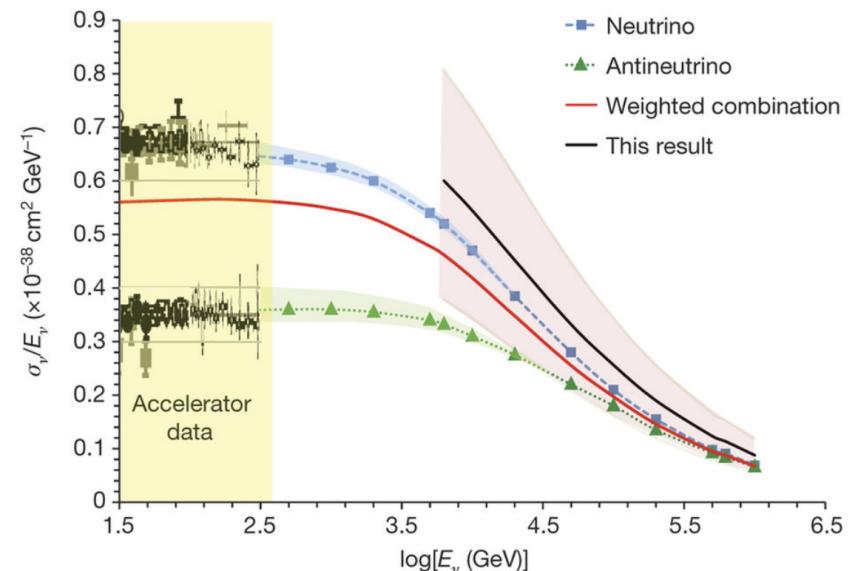


Figure 2: Geometry of the emission from a thin spherical shell with the observer located inside the shell.



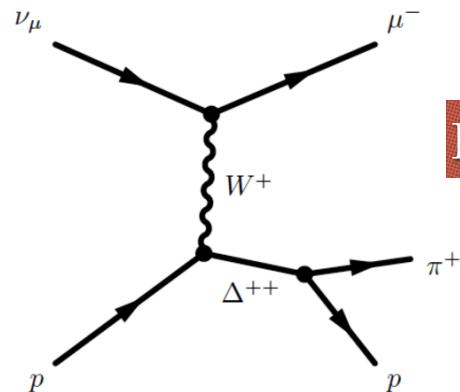
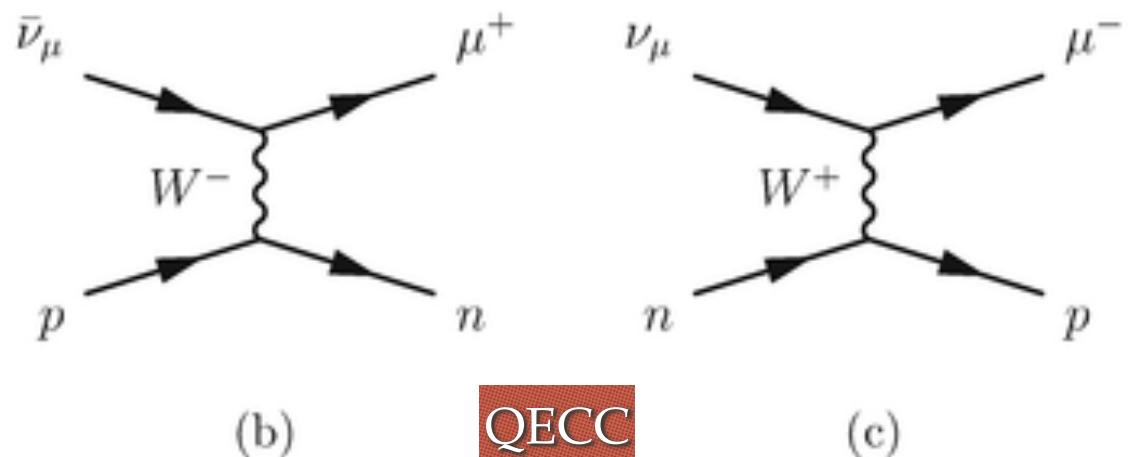
# ATM theory

- Atmospheric neutrino interactions
- Quasielastic scattering
- Resonant scattering
- Deep Inelastic scattering

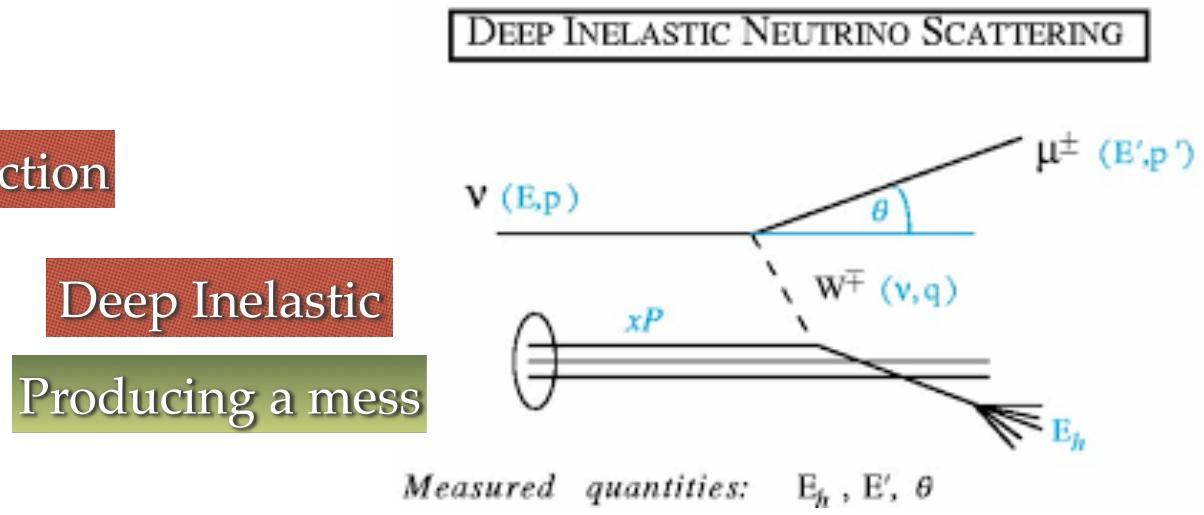


# ATM theory

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- Quasielastic scattering
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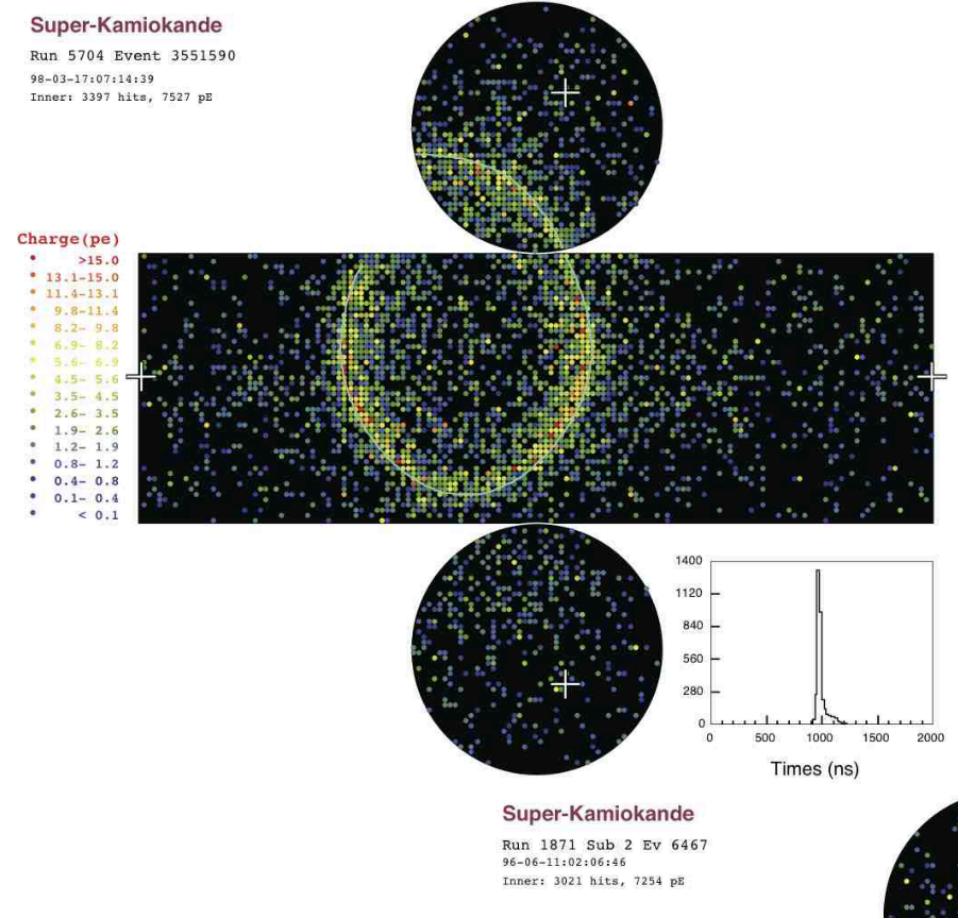
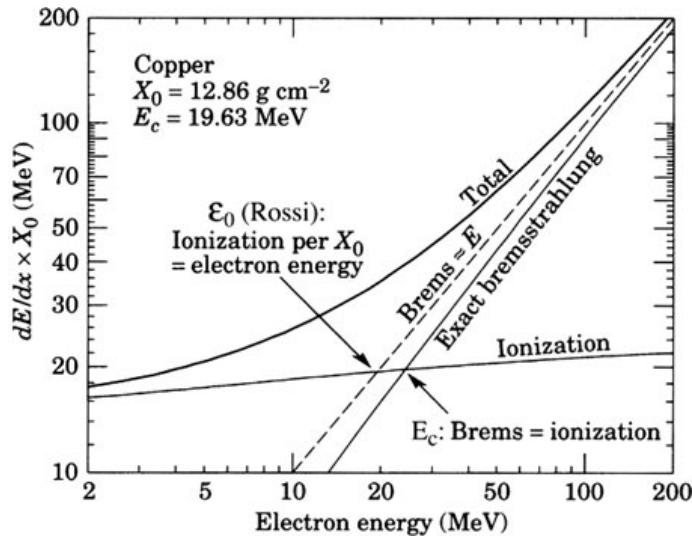


Producing pions, etc



# ATM observation Electron and muon like events

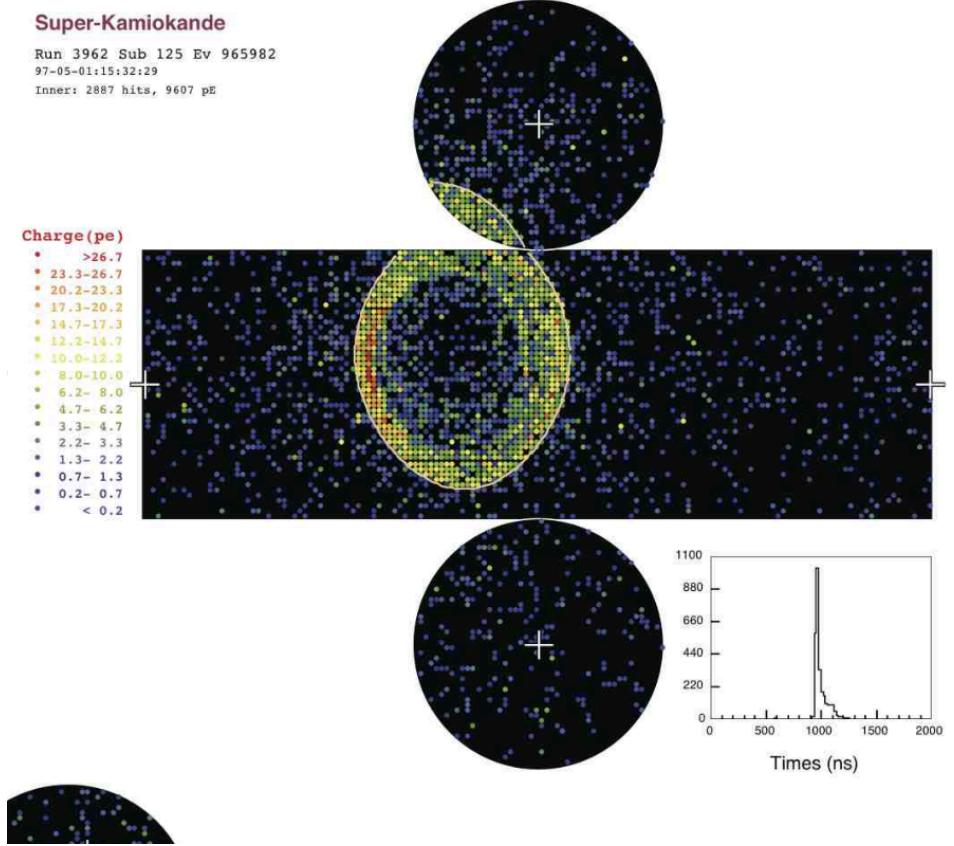
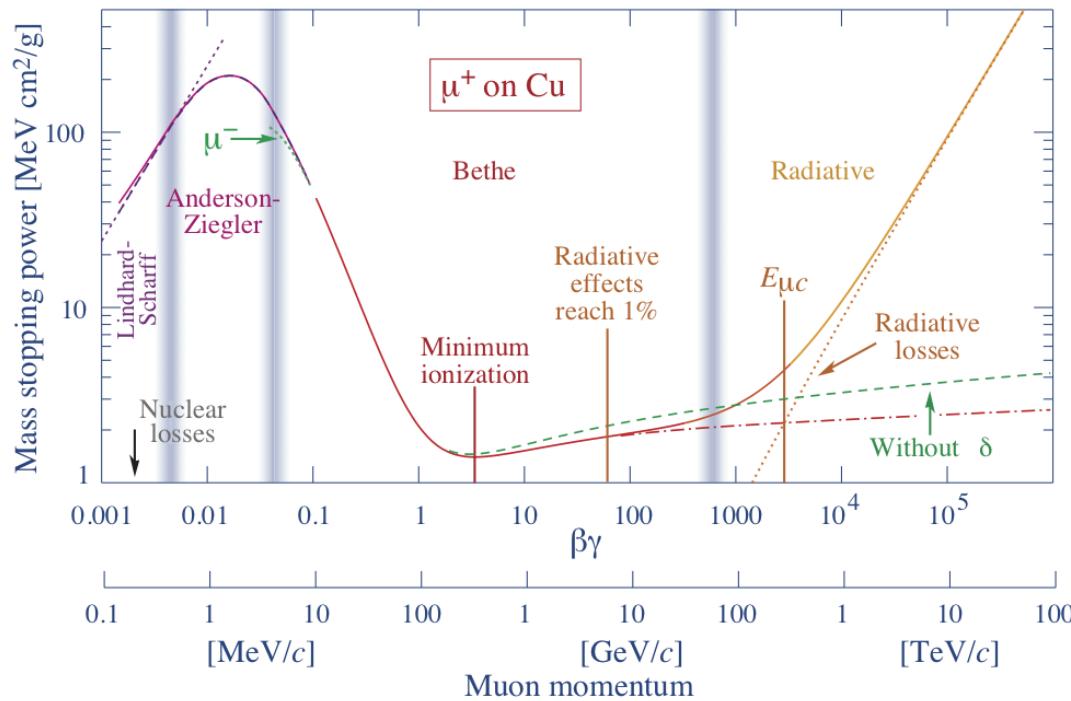
- Electron like events, **fuzzy rings**.
  - Due to electromagnetic shower initiated by the electron
  - More than 1 electrons



# ATM observation

## Electron and muon like events

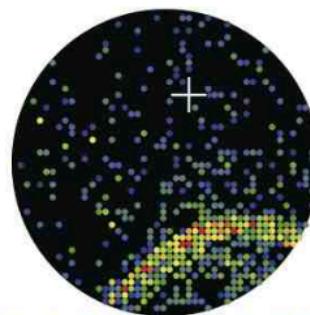
- Muon like events, **clean rings**
  - Muon energy loss is dominated by ionisation



# Multiple muon rings

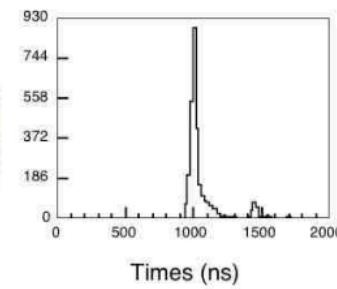
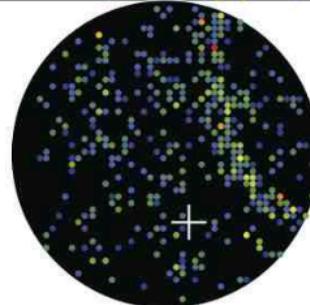
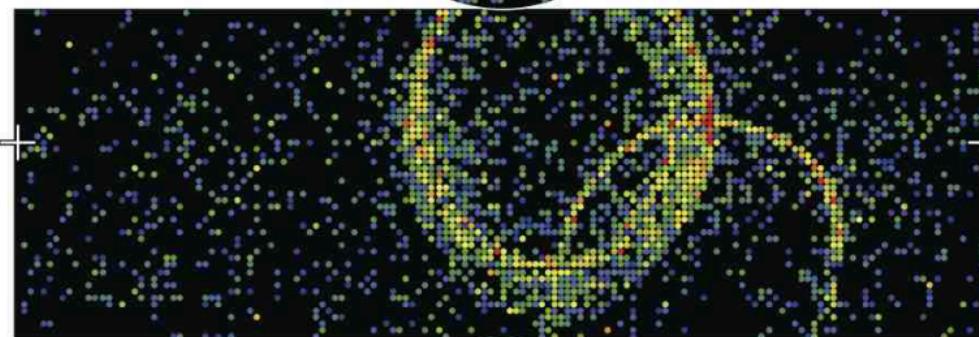
Super-Kamiokande

Run 1871 Sub 2 Ev 6467  
96-06-11:02:06:46  
Inner: 3021 hits, 7254 pE



Charge (pE)

- \* >15.0
- \* 13.1-15.0
- \* 11.4-13.1
- \* 9.8-11.4
- \* 8.2- 9.8
- \* 6.9- 8.2
- \* 5.6- 6.9
- \* 4.5- 5.6
- \* 3.5- 4.5
- \* 2.6- 3.5
- \* 1.9- 2.6
- \* 1.2- 1.9
- \* 0.8- 1.2
- \* 0.4- 0.8
- \* 0.1- 0.4
- \* < 0.1



### Super-Kamiokande I

Run 1728 Sub 4 Ev 25171

96-05-29:08:01:53

Inner: 2294 hits, 7095 pE

Outer: 4 hits, 32 pE (in-time)

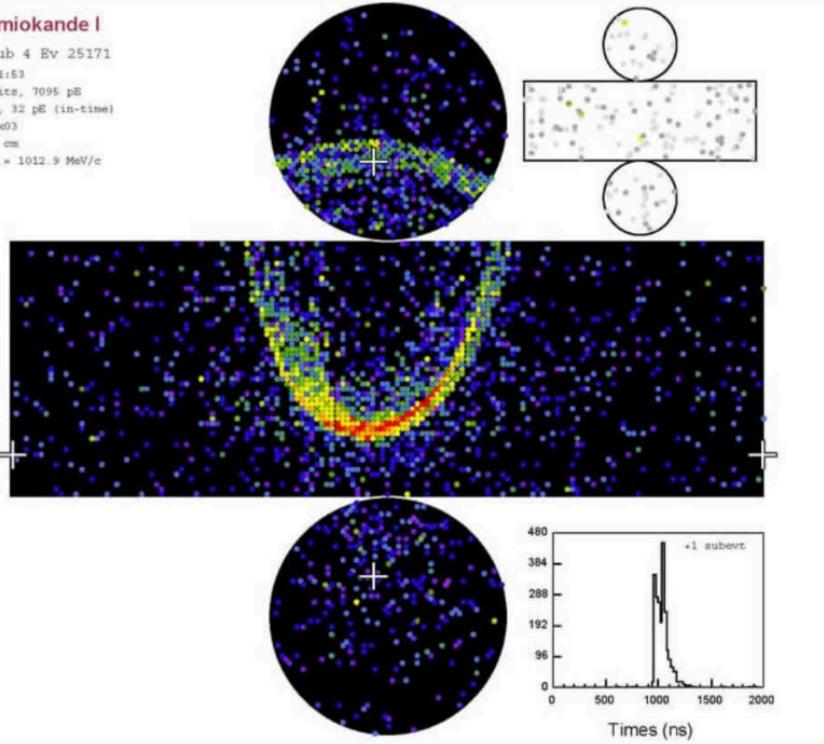
Trigger ID: 0x03

D wall: 592.8 cm

PC mu-like,  $p = 1012.9$  MeV/c

#### Charge (pe)

- >26.7
- 23.3-26.7
- 20.2-23.3
- 17.3-20.2
- 14.7-17.2
- 12.0-14.7
- 9.0-12.0
- 6.2- 9.0
- 4.7- 6.2
- 3.3- 4.7
- 2.2- 3.3
- 1.3- 2.2
- 0.7- 1.3
- 0.2- 0.7
- < 0.2



### Super-Kamiokande I

Run 1757 Sub 4 Ev 25716

96-06-03:04:51:37

Inner: 1948 hits, 5243 pE

Outer: 4 hits, 30 pE (in-time)

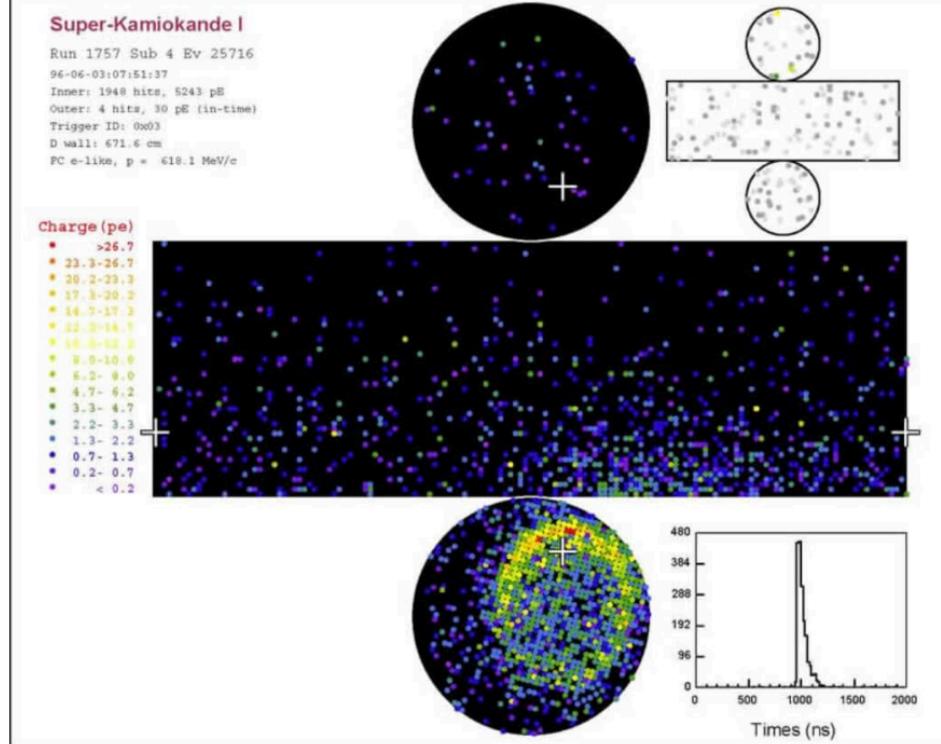
Trigger ID: 0x03

D wall: 671.6 cm

PC e-like,  $p = 618.1$  MeV/c

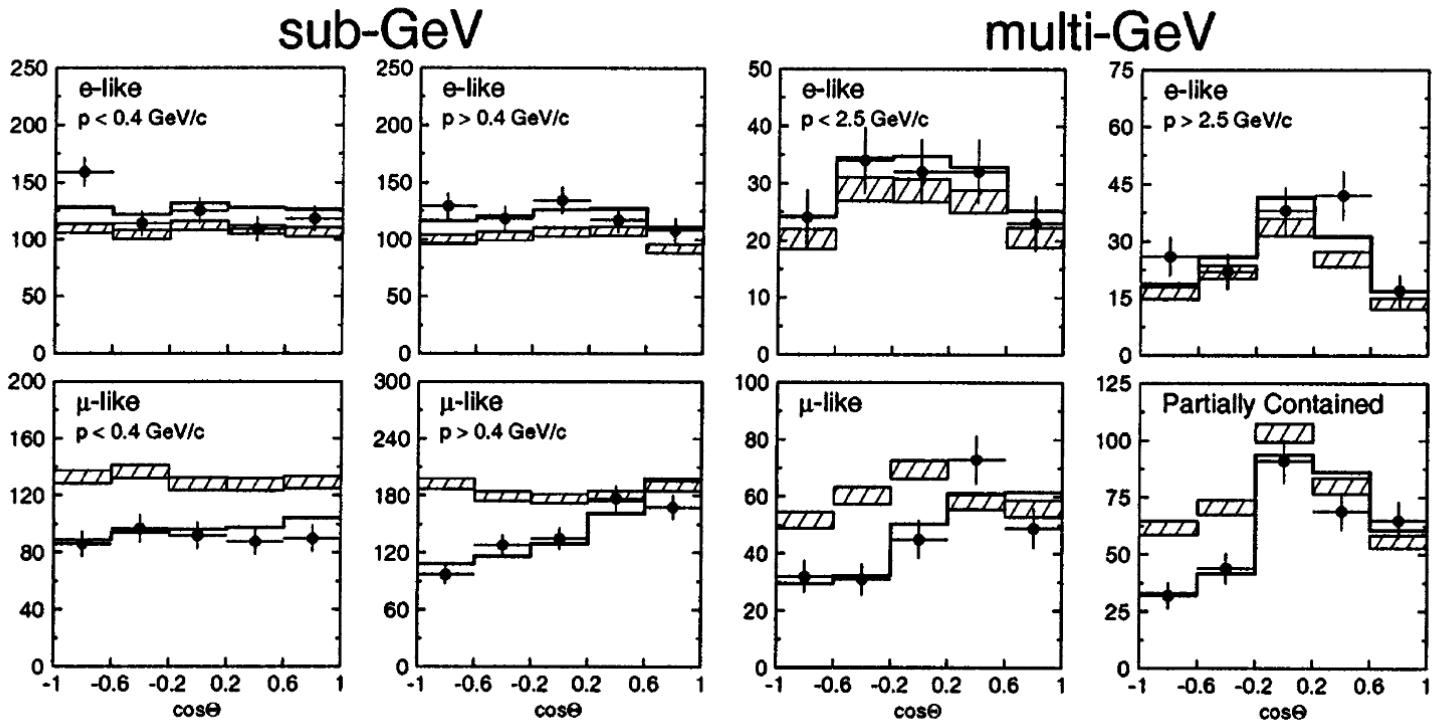
#### Charge (pe)

- >26.7
- 23.3-26.7
- 20.2-23.3
- 17.3-20.2
- 14.7-17.2
- 12.0-14.7
- 9.0-12.0
- 6.2- 9.0
- 4.7- 6.2
- 3.3- 4.7
- 2.2- 3.3
- 1.3- 2.2
- 0.7- 1.3
- 0.2- 0.7
- < 0.2



# ATM observation Angular distribution

- Nothing was going on for electron type neutrinos
- < 0.4GeV muon neutrinos
  - Overall suppression
- > 0.4 GeV muon neutrinos
  - Location dependent suppressions (disappearing)



VOLUME 81, NUMBER 8

PHYSICAL REVIEW LETTERS

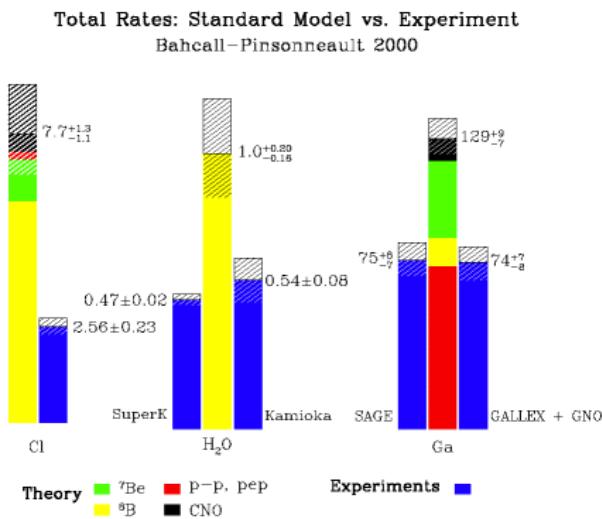
24 AUGUST 1998

## Evidence for Oscillation of Atmospheric Neutrinos

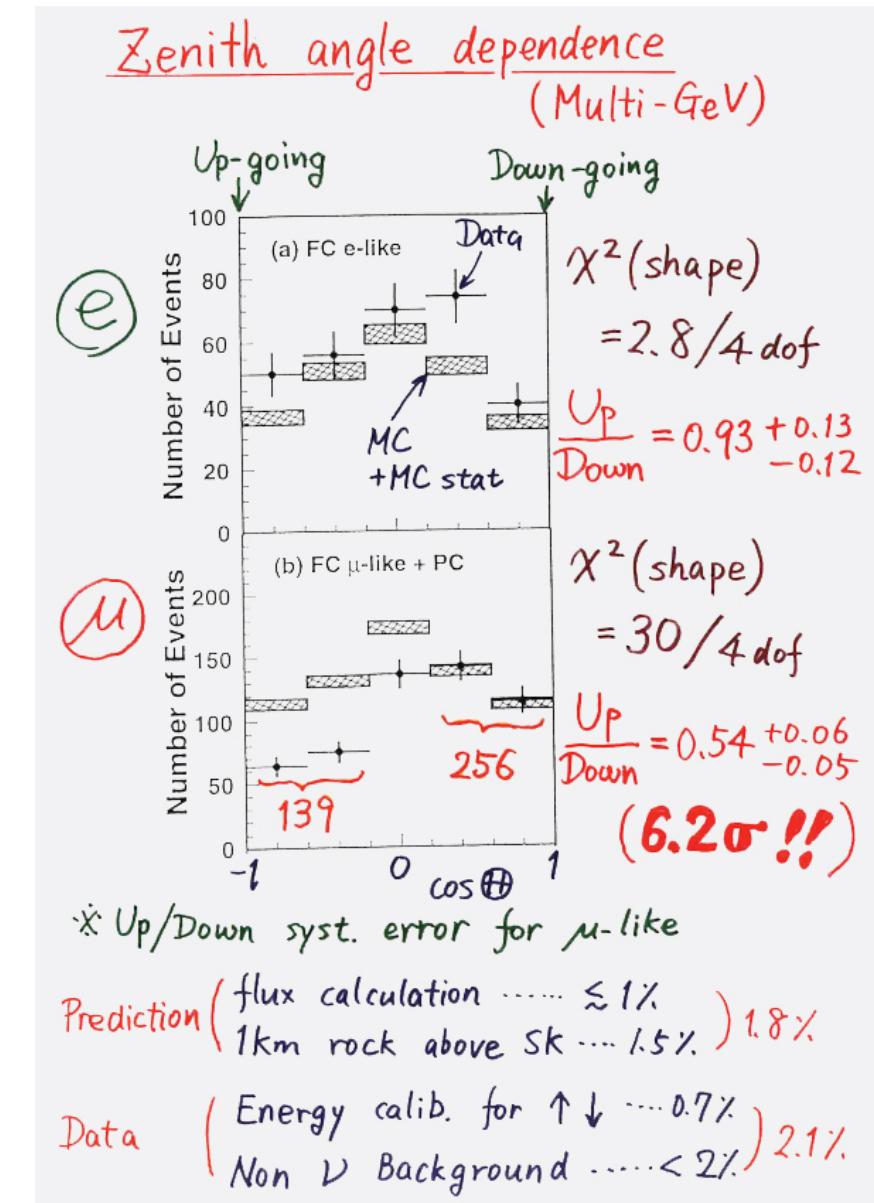
Y. Fukuda,<sup>1</sup> T. Hayakawa,<sup>1</sup> E. Ichihara,<sup>1</sup> K. Inoue,<sup>1</sup> K. Ishihara,<sup>1</sup> H. Ishino,<sup>1</sup> Y. Itow,<sup>1</sup> T. Kajita,<sup>1</sup> J. Kameda,<sup>1</sup> S. Kasuga,<sup>1</sup> K. Kobayashi,<sup>1</sup> Y. Kobayashi,<sup>1</sup> Y. Koshibo,<sup>1</sup> M. Miura,<sup>1</sup> M. Nakahata,<sup>1</sup> S. Nakayama,<sup>1</sup> A. Okada,<sup>1</sup> K. Okumura,<sup>1</sup> N. Sakurai,<sup>1</sup> M. Shiozawa,<sup>1</sup> Y. Suzuki,<sup>1</sup> Y. Takeuchi,<sup>1</sup> Y. Totsuka,<sup>1</sup> S. Yamada,<sup>1</sup> M. Earl,<sup>2</sup> A. Habig,<sup>2</sup> E. Kearns,<sup>2</sup> M. D. Messier,<sup>2</sup> K. Scholberg,<sup>2</sup> J. L. Stone,<sup>2</sup> L. R. Sulak,<sup>2</sup> C. W. Walter,<sup>2</sup> M. Goldhaber,<sup>3</sup> T. Barszczak,<sup>4</sup> D. Casper,<sup>4</sup> W. Gajewski,<sup>4</sup> P. G. Halverson,<sup>4,\*</sup> J. Hsu,<sup>4</sup> W. R. Kropp,<sup>4</sup> L. R. Price,<sup>4</sup> F. Reines,<sup>4</sup> M. Smy,<sup>4</sup> H. W. Sobel,<sup>4</sup> M. R. Vagins,<sup>4</sup> K. S. Ganezer,<sup>5</sup> W. E. Keig,<sup>5</sup> R. W. Ellsworth,<sup>6</sup> S. Tasaka,<sup>7</sup> J. W. Flanagan,<sup>8,†</sup> A. Kibayashi,<sup>8</sup> J. G. Learned,<sup>8</sup> S. Matsuno,<sup>8</sup> V. J. Stenger,<sup>8</sup> D. Takemori,<sup>8</sup> T. Ishii,<sup>9</sup> J. Kanazaki,<sup>9</sup> T. Kobayashi,<sup>9</sup> S. Mine,<sup>9</sup> K. Nakamura,<sup>9</sup> K. Nishikawa,<sup>9</sup> Y. Oyama,<sup>9</sup> A. Sakai,<sup>9</sup> M. Sakuda,<sup>9</sup> O. Sasaki,<sup>9</sup> S. Echigo,<sup>10</sup> M. Kohama,<sup>10</sup> A. T. Suzuki,<sup>10</sup> T. J. Haines,<sup>11,‡</sup> E. Blaufuss,<sup>12</sup> B. K. Kim,<sup>12</sup> R. Sanford,<sup>12</sup> R. Svoboda,<sup>12</sup> M. L. Chen,<sup>13</sup> Z. Conner,<sup>13,‡</sup> J. A. Goodman,<sup>13</sup> G. W. Sullivan,<sup>13</sup> J. Hill,<sup>14</sup> C. K. Jung,<sup>14</sup> K. Martens,<sup>14</sup> C. Mauger,<sup>14</sup> C. McGrew,<sup>14</sup> E. Sharkey,<sup>14</sup> B. Viren,<sup>14</sup> C. Yanagisawa,<sup>14</sup> W. Doki,<sup>15</sup> K. Miyano,<sup>15</sup> H. Okazawa,<sup>15</sup> C. Saji,<sup>15</sup> M. Takahata,<sup>15</sup> Y. Nagashima,<sup>16</sup> M. Takita,<sup>16</sup> T. Yamaguchi,<sup>16</sup> M. Yoshida,<sup>16</sup> S. B. Kim,<sup>17</sup> M. Etoh,<sup>18</sup> K. Fujita,<sup>18</sup> A. Hasegawa,<sup>18</sup> T. Hasegawa,<sup>18</sup> S. Hatakeyama,<sup>18</sup> T. Iwamoto,<sup>18</sup> M. Koga,<sup>18</sup> T. Maruyama,<sup>18</sup> H. Ogawa,<sup>18</sup> J. Shirai,<sup>18</sup> A. Suzuki,<sup>18</sup> F. Tsushima,<sup>18</sup> M. Koshiba,<sup>19</sup> M. Nemoto,<sup>20</sup> K. Nishijima,<sup>20</sup> T. Futagami,<sup>21</sup> Y. Hayato,<sup>21,§</sup> Y. Kanaya,<sup>21</sup> K. Kaneyuki,<sup>21</sup> Y. Watanabe,<sup>21</sup> D. Kielczewska,<sup>22,¶</sup> R. A. Doyle,<sup>23</sup> J. S. George,<sup>23</sup> A. L. Stachyra,<sup>23</sup> L. L. Wai,<sup>23,||</sup> R. J. Wilkes,<sup>23</sup> and K. K. Young<sup>23</sup> (Super-Kamiokande Collaboration)

# ATM observation Angular distribution

- Prof. Kajita's presentation at Neutrino 1998
- The disappearance of solar neutrinos is not “astrophysics”
- Neutrino particle physics need to be modified!



19



# ATM observation Angular distribution

We present an analysis of atmospheric neutrino data from a 33.0 kton yr (535-day) exposure of the Super-Kamiokande detector. The data exhibit a zenith angle dependent deficit of muon neutrinos which is inconsistent with expectations based on calculations of the atmospheric neutrino flux. Experimental biases and uncertainties in the prediction of neutrino fluxes and cross sections are unable to explain our observation. The data are consistent, however, with two-flavor  $\nu_\mu \leftrightarrow \nu_\tau$  oscillations with  $\sin^2 2\theta > 0.82$  and  $5 \times 10^{-4} < \Delta m^2 < 6 \times 10^{-3}$  eV $^2$  at 90% confidence level. [S0031-9007(98)06975-0]

$$P_{a \rightarrow b} = \sin^2 2\theta \sin^2 \left( \frac{1.27 \Delta m^2 (\text{eV}^2) L (\text{km})}{E_\nu (\text{GeV})} \right),$$

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## Evidence for Oscillation of Atmospheric Neutrinos

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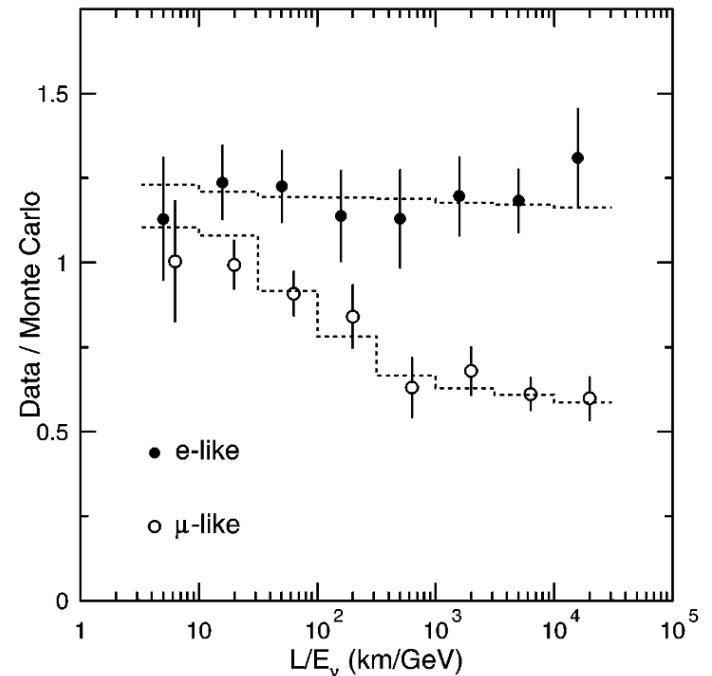


FIG. 4. The ratio of the number of FC data events to FC Monte Carlo events versus reconstructed  $L/E_\nu$ . The points show the ratio of observed data to MC expectation in the absence of oscillations. The dashed lines show the expected shape for  $\nu_\mu \leftrightarrow \nu_\tau$  at  $\Delta m^2 = 2.2 \times 10^{-3}$  eV $^2$  and  $\sin^2 2\theta = 1$ . The slight  $L/E_\nu$  dependence for  $e$ -like events is due to contamination (2–7%) of  $\nu_\mu$  CC interactions.