

Addition of Angular Momentum

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October 2021

1 Introduction

suppose we have two operators (J_i^1 and J_i^2) which satisfy the $\text{su}(2)$ algebra defined on 2 vector spaces V_1 and V_2 respectively.

We define J_i on $V_1 \otimes V_2$ by :

$J_i = J_i^1 \otimes 1 + 1 \otimes J_i^2$ now this new operator defined on $V_1 \otimes V_2$ also satisfies $\text{su}(2)$ algebra (PS. here 1 denotes the identity element not the number 1).

NOTE: If we had some arbitrary coefficients, then it wouldn't have satisfied the $\text{su}(2)$ algebra. check it for yourself. HINT!: write $J_i = aJ_i^1 \otimes 1 + b1 \otimes J_i^2$, now use the commutator relation $[J_i, J_j] = i\hbar\epsilon_{ijk}J_k$ and you'll find that for it to satisfy $\text{su}(2)$ algebra, the coefficients must be 1.

2 Spin 1/2

We know that: $S^2|s, m\rangle = \hbar^2 s(s+1)|s, m\rangle$ and $S_z|s, m\rangle = \hbar m|s, m\rangle$ and we know that these two eigenstates form an orthonormal eigenbase.

here $s = 1/2, m = \pm 1/2$

we define $S_i = S_i^1 \otimes 1 + 1 \otimes S_i^2$ on $V_1 \otimes V_2$ which is spanned by

$$\{|1/2, 1/2\rangle_1 \otimes |1/2, 1/2\rangle_2, |1/2, 1/2\rangle_1 \otimes |1/2, -1/2\rangle_2, |1/2, -1/2\rangle_1 \otimes |1/2, 1/2\rangle_2, |1/2, -1/2\rangle_1 \otimes |1/2, -1/2\rangle_2\}$$

which leads to 5 possible multiplets:

- 1) for $s = 0$ singlets.
- 2) Two doublets of $s = 1/2$.
- 3) one doublet of $s = 1/2$ and 2 singlets of $s = 0$.
- 4) one triplet of $s = 1$ and one singlet of $s = 0$.

5) one $s = 3/2$ multiplet.

only the 4) option is a viable option. Let that sink in and think as to why this is the only correct option, remember s takes half integer and integer $2s+1$ values and m ranges between $+s$ and $-s$.

Thus, we can write: $(s = 1/2) \otimes (s = 1/2) = (s = 1) \oplus (s = 0)$.

now, we can write the basis as:

$$\begin{aligned} |1, 1\rangle &= |1/2, 1/2\rangle_1 \otimes |1/2, 1/2\rangle_2 \\ |1, 0\rangle &= a|1/2, 1/2\rangle_1 \otimes |1/2, -1/2\rangle_2 + b|1/2, -1/2\rangle_1 \otimes |1/2, 1/2\rangle_2 \\ |1, -1\rangle &= |1/2, -1/2\rangle_1 \otimes |1/2, 1/2\rangle_2 \\ |0, 0\rangle &= c|1/2, 1/2\rangle_1 \otimes |1/2, -1/2\rangle_2 + d|1/2, -1/2\rangle_1 \otimes |1/2, 1/2\rangle_2 \end{aligned}$$

After a careful observation you'll see that $m=0$ has two states $m=0$ and $s=m=0$ and these states are entangled.

We define $S_+ = S_-^1 + S_-^2$ analogous to J_\pm operators.
similarly, we can define $S^2 = (S^1 + S^2)^2 = (S^{(1)})^2 + (S^{(2)})^2 + 2S^1 \cdot S^2$

Here, I made use of the following result:

$$\begin{aligned} J^{(1)} \cdot J^{(2)} &= \sum_{i=1}^3 J_i^{(1)} \otimes J_i^{(2)} \\ &= 1/2(J_+^1 \otimes J_-^2 + J_-^1 \otimes J_+^2) + J_z^1 \otimes J_z^2 \end{aligned}$$