

thesis

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The following equalities are of use, all of which can be verified through calculation and using the property $e^{\hat{A}}\hat{B}e^{-\hat{A}} = \hat{B} + [\hat{A}, \hat{B}] + \frac{1}{2!}[\hat{A}, [\hat{A}, \hat{B}]] + \dots$

$$\hat{D}^*(\alpha)\hat{D}(\alpha) = 1 \quad (1)$$

$$\begin{aligned} \hat{D}^*(\alpha)\hat{a}\hat{D}(\alpha) &= \hat{a} + \alpha \\ \hat{D}(\alpha)\hat{a}\hat{D}^*(\alpha) &= \hat{a} - \alpha \end{aligned} \quad (2)$$

$$\hat{S}^*(\xi)\hat{S}(\xi) = \hat{S}(-\xi)\hat{S}(\xi) = 1 \quad (3)$$

$$\begin{aligned} \hat{S}^*(\xi)\hat{a}\hat{S}(\xi) &= \hat{a} \cosh r - \hat{a}^* e^{i\theta} \sinh r \\ \hat{S}(\xi)\hat{a}\hat{S}^*(\xi) &= \hat{a}^* \cosh r - \hat{a} e^{-i\theta} \sinh r \end{aligned} \quad (4)$$

Let us write $\hat{a} = A + i\lambda B$, $\lambda \in \mathbb{R}$,¹. Let us also call $[\hat{A}, \hat{B}] = i\hat{C}$ and $\hat{A}, \hat{B} - 2\langle\hat{A}\rangle\langle\hat{B}\rangle$. Using \hat{F}, \hat{C} , we can rewrite the uncertainty relation as $(\Delta\hat{A})^2(\Delta\hat{B})^2 = \frac{1}{4}(\langle\hat{F}\rangle^2 + \langle\hat{C}\rangle^2)^2$. Applying this formula to \hat{a} , we see

$$\begin{aligned} \Delta\hat{A}^2 &= \frac{1}{2}(\lambda_i\langle\hat{F}\rangle + \lambda_r\langle\hat{C}\rangle) \\ \Delta\hat{B}^2 &= \frac{1}{|\lambda|^2}\Delta\hat{A}^2 \\ \lambda_i\langle\hat{C}\rangle - \lambda_r\langle\hat{F}\rangle &= 0 \end{aligned} \quad (5)$$

Combining we see our familiar uncertainty relation

$$\Delta\hat{A}^2\Delta\hat{B}^2 \geq \frac{1}{4}(\langle\hat{F}\rangle^2 + \langle\hat{C}\rangle^2) \quad (6)$$

¹This can be thought of breaking \hat{a} into hermitian and anti-hermitian components. For our usual operators \hat{q}, \hat{p} , the hermitian component is \hat{q} and the antihermitian is $\frac{\partial}{\partial x}$

²As a sanity check, our usual operators for position and momentum \hat{q}, \hat{p} follow the commutation relation $[\hat{q}, \hat{p}] = i\hbar \implies (\Delta\hat{q})^2(\Delta\hat{p})^2 \geq \frac{\hbar^2}{4}$ following our formula.