

## FORCED HARMONIC MOTION

### READINGS

Marion and Thornton, "Classical Dynamics", Chapter 3; or other Physics 381 dynamics textbook.

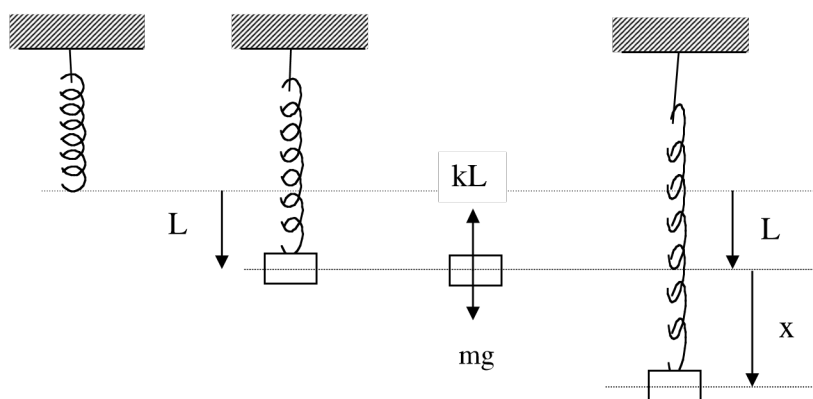
### PURPOSE

To study the resonant response of a system of a weight suspended from a spring where the system is driven up and down harmonically while subject to damping forces, and to determine both the relationships among frequency, amplitude and phase, and the effects of damping.

### THEORY

Consider a spring with force constant  $k$ . When a weight  $mg$  is hung from the spring, it stretches by an amount  $L$ . When the weight is motionless, we know by Newton's second law,  $\sum F = ma$  and Hooke's law,  $F = -kx$  that  $mg - kL = 0$ . When the spring is displaced a small distance  $x_0$  from  $L$  and released, the weight undergoes periodic vertical vibrations in  $x$  about the equilibrium length  $L$  described by Newton's second law,

$$-k(L + x) + mg = ma$$



Using  $kL = mg$  from above, we find

$$-kx = m \frac{d^2x}{dt^2} \tag{1}$$

If the weight is released at  $t = 0$ , the solution is

$$x = x_0 \cos(\omega_0 t) \tag{2}$$

where  $\omega_0 = \sqrt{k/m}$ .  $\omega_0$  is the natural or resonant frequency of the system and is expressed in radians/s. [Check that Eq. (2) is the solution by substituting it into Eq. (1).] Note that at  $t = 0, x = 0$  and  $v = dx/dt = 0$ .

Equation (2) tells us that the weight will continue to oscillate forever. This is because the system is assumed not to lose its starting energy. The term "damping" describes the condition where energy is lost. Damping forces act antiparallel to the velocity and are frequently found to obey Stokes' law,  $F = -Rv = -Rdx/dt$ , where  $R$  is a constant. The equation of motion is then

$$m \frac{d^2x}{dt^2} = -kx - R \frac{dx}{dt} \quad (3)$$

This equation was solved in the previous lab and we note that when released the weight will oscillate with decreasing amplitude until it comes to rest at  $x = 0$ . If the damping is too large, it won't even oscillate but instead will slowly approach its equilibrium position.

For this experiment we are interested in a more complicated case where the top of the spring is moved up and down at an angular frequency  $\omega$  by an external agent. This periodic displacement of the top of the spring translates into a periodic force on the weight

$$F = F_0 \sin \omega t \quad (4)$$

and the equation of motion becomes

$$m \frac{d^2x}{dt^2} = F_0 \sin \omega t - kx - R \frac{dx}{dt} \quad (5)$$

When the force is initially switched on, the system undergoes some transient motion which is eventually damped out and the system settles into a steady state motion given by

$$x = A \sin(\omega t + \phi) \quad (6)$$

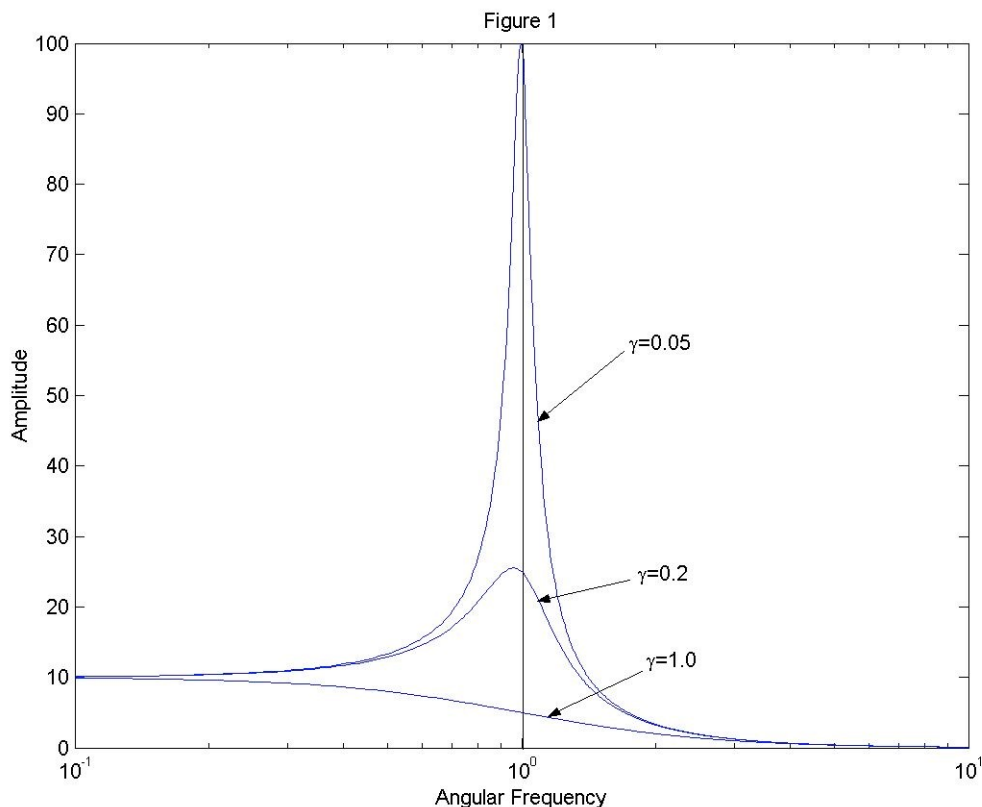
where the amplitude  $A$  is given by

$$A = \frac{F_0/m}{\sqrt{(\omega_0^2 - \omega^2)^2 + 4\gamma^2\omega^2}} \quad (7)$$

and the phase shift is

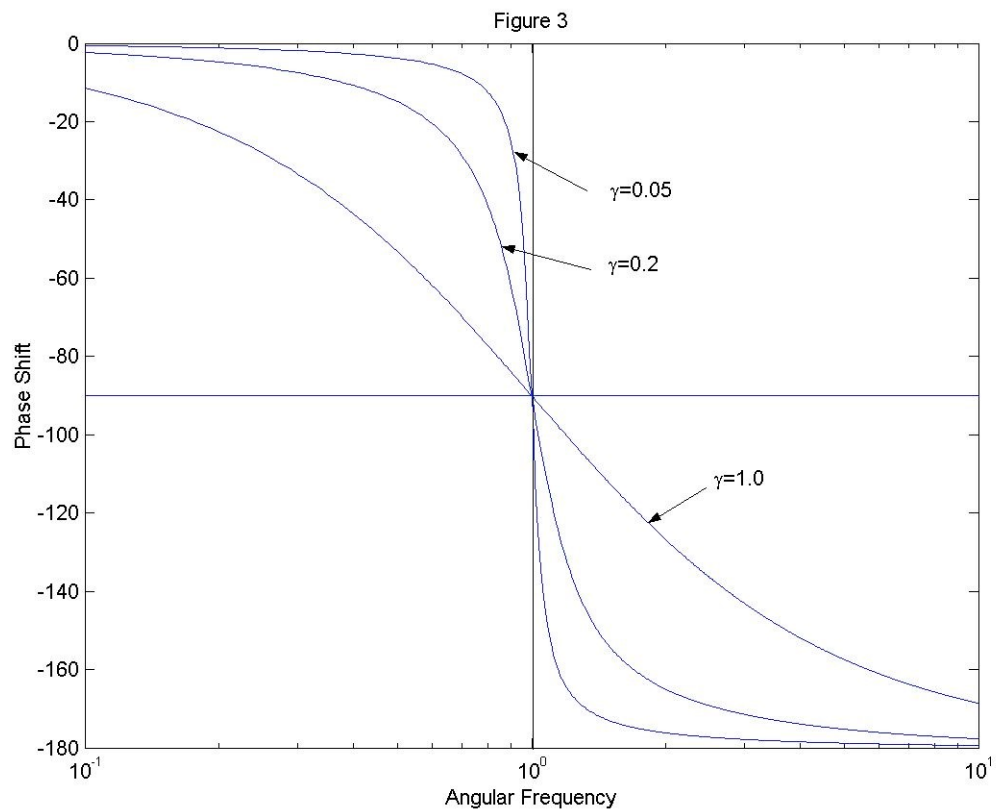
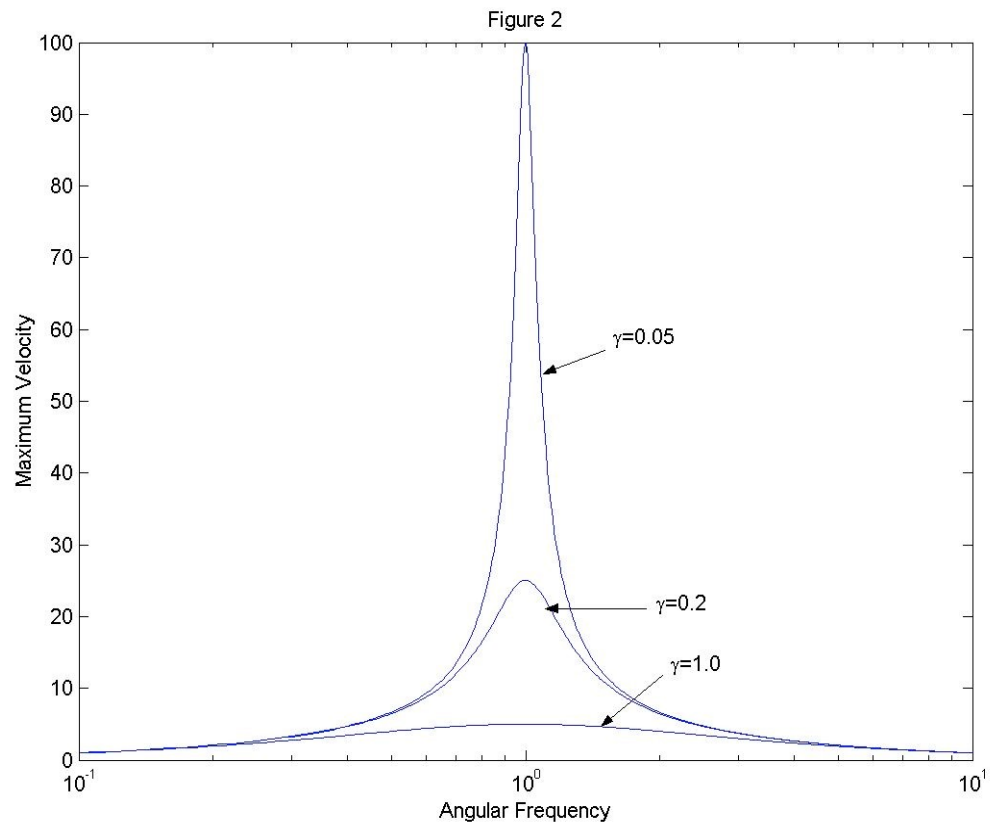
$$\tan \phi = \frac{2\gamma\omega}{\omega^2 - \omega_0^2} \quad (8)$$

with  $\omega_0 = \sqrt{k/m}$  and  $\gamma = R/2m$ .  $\gamma$  is called the damping constant. [Check that Eq. (6) is a solution to Eq. (5) by substituting it into Eq. (5).] Equation (7) is plotted in Fig. 1 (for  $\omega_0 = 1$ ).  $A$  is resonant, that is it has a maximum, at  $\omega = \omega_0$ , because the denominator in Eq. (7) reaches a minimum (actually the minimum occurs when  $\omega = \sqrt{\omega_0^2 - 2\gamma^2} \equiv \omega_R$ , but for this experiment  $\gamma \ll \omega_0$  and we will not be able to detect this shift). Note that if there were no damping ( $\gamma = 0$ ),  $A$  would be infinite at  $\omega = \omega_0$ . In practice there always is some damping and the amplitude remains finite.



Another quantity of interest is  $v_{max} = \omega A$ , which is obtained directly from Eq. (6) since  $v = dx/dt$ . As shown in Fig. 2,  $v_{max}$  is symmetric about  $\omega_0$  if it is plotted on a log scale (i.e.  $v_{max}$  vs.  $\log \omega$ ).

The phase shift  $\phi$  (see Fig. 3) gives the phase difference between the driving force and the displacement of the mass. This force reaches a maximum at a time  $t_0$  given by  $\omega t_0 = \pi/2$  (i.e.,  $\sin \omega t_0 = 1$ ), while  $x$  reaches a maximum at a time  $t_1$  given by  $\omega t_1 + \phi = \pi/2$ . So  $t_1 = t_0 - \phi/\omega$ . Because  $\phi$  is negative as shown in Fig. 3,  $t_1 > t_0$  and the displacement lags the force (i.e., reaches a maximum after the force does), we see that for very low frequencies  $x$  and  $F$  are almost in phase. But for higher frequencies  $x$  lags until it is exactly out of phase ( $180^\circ$ ) at very high frequencies.



We can likewise derive the phase shifts between the force and velocity and find that  $v$  leads  $F$  by  $\pi/2$  for low  $\omega$ , is exactly in phase with  $F$  at  $\omega = \omega_0$ , then lags by  $\pi/2$  for large  $\omega$ .

These phase relations are important for determining how much work  $F$  does on the mass. Since we are considering the steady state solution, the energy in the system does not change with time. So, the work done by the external force in one cycle equals the energy dissipated by the damping force. This is easily calculated since the power  $P$  supplied by  $F$  is

$$P = \frac{\Delta W}{\Delta t} = \frac{F(t)\Delta x}{\Delta t} = F(t)v(t)$$

The average over one cycle is

$$\langle P \rangle = \frac{1}{\tau} \int_0^\tau P(t) dt = \frac{1}{\tau} \int_0^\tau F(t)v(t) dt$$

where  $\tau = 2\pi/\omega$  is the period. Using Eqs. (4) and (6) we find

$$\langle P \rangle = \frac{A\omega}{2} F_0 \cos \theta \quad (9)$$

where  $\theta = \phi + \pi/2$  is the phase difference between  $F$  and  $v$ .  $\cos \theta$  is called the power factor. Using Eqs. (7) and (8) we can show

$$\langle P \rangle = \gamma m (A\omega)^2 = \gamma m v_{max}^2 \quad (10)$$

We see from Fig. (2) that at resonance the power supplied is a maximum and has the value (for  $\gamma \ll \omega_0$ )

$$\langle P \rangle_{max} = \frac{F_0^2}{4\gamma m} \quad (11)$$

The Quality Factor, or  $Q$ , is a very important parameter used to describe a resonance. It is defined to be

$$Q = \frac{\omega_R}{2\gamma} \quad (12)$$

and is also (for  $\gamma \ll \omega_0$ )

$$Q = \frac{2\pi(\text{maximum energy stored at resonance})}{(\text{energy dissipated at resonance in one cycle})}$$

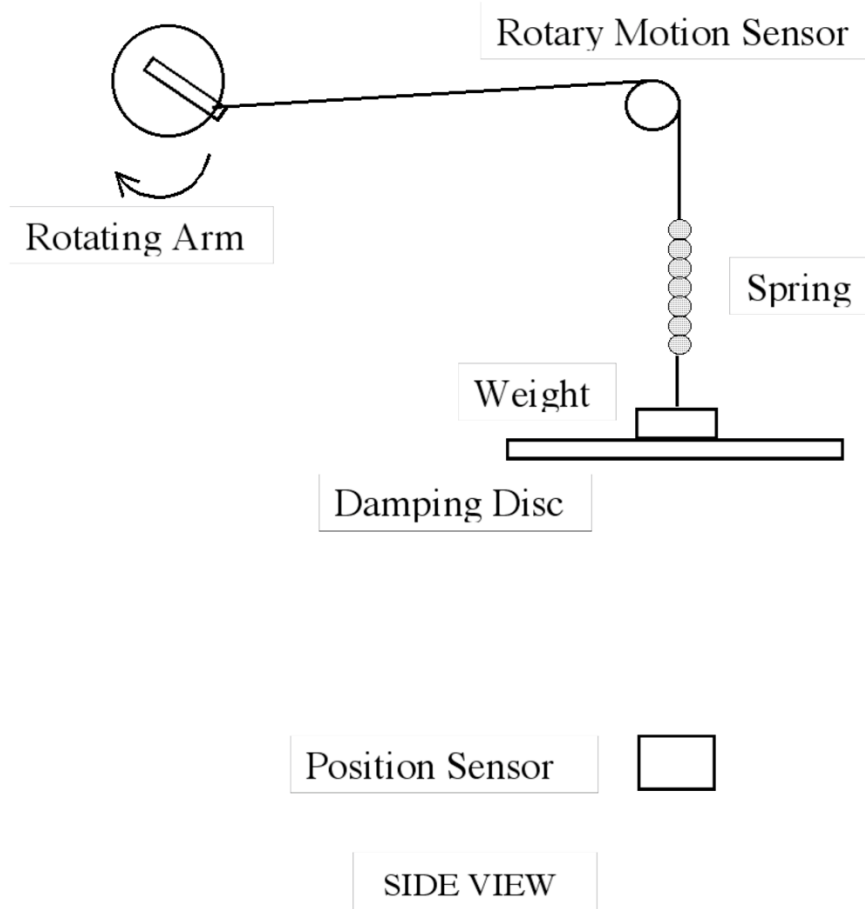
An alternative (equivalent) definition is that is more useful experimentally is (for  $\gamma \ll \omega_0$ )

$$Q = \frac{\omega_0}{2\Delta\omega} \quad (13)$$

where  $2\Delta\omega = \omega_+ - \omega_-$  and  $\omega_{\pm}$  are the frequencies on either side of  $\omega_0$  where  $\langle P \rangle = \frac{1}{2} \langle P \rangle_{max}$ .  $2\Delta\omega$  is called the full width at half power of the resonance. Experimentally one usually does not measure  $\langle P \rangle$  to get  $2\Delta\omega$ . Instead, one can use the data for  $v_{max}$  since by Eq. (10),  $v = \frac{1}{\sqrt{2}} v_{max} = 0.707 v_{max}$ . So  $2\Delta\omega$  is obtained from Fig. 2 by determining the two frequencies where the velocity drops to 71% of its maximum value.

## APPARATUS

The apparatus you will use is shown schematically below:



The rotating arm is driven by an electric motor whose speed is controlled by the voltage supplied by a variable power supply. Record the voltage applied to the motor at each speed setting so that a repeat of a setting is possible.

The rotary motion sensor shows the movement of the top of the spring, so its oscillations have the same phase as the external force applied to the system of the mass on spring.

The position sensor gives the position of the mass on the spring. Both angle sensor and position sensor should be read by Logger Pro as in the previous lab. From the Logger Pro displays, find the period of the oscillation and the amplitude  $A$  and phase  $\phi$  of the oscillation of the mass; measure the phase as the difference between the position and the rotary motion sensor oscillations.

## PROCEDURE

### Lab Activity for Week 1 (Forced Harmonic Oscillator with a Larger Damping)

#### 1. Measurement and comparison of oscillation frequencies to determine the natural frequency

##### **[Do not turn on the motor yet]**

A. [Case 1. Without damping disc]: Remove the damping disc, set the oscillator in motion, record the motion using Logger Pro, and export the data to a 'txt' file. Measure the period of the system using MATLAB. Time over 10 periods to reduce starting and stopping errors and repeat the measurement several times so that you have an estimate of your error. The available drive angular frequency ranges from about 10 rad/s to about 3 rad/s. If necessary, adjust the mass to set the angular frequency to about 5-6 rad/s.

[Case 2. With the equivalent mass of a large damping disc (but without the actual damping disc)]: Measure the mass of a large damping disc made out of cardboard, and add the equivalent mass (instead of adding the disc) to the oscillator. Repeat the procedure as in Case 1 to measure the frequency.

[Case 3. With a large damping disc]: Remove the equivalent mass and insert the large damping disc. Repeat the procedure as in Case 1 to measure the frequency.

**Question:** Compare the three measured frequencies. Are they the same? You will conduct experiments on forced harmonic oscillator with the large damping disc inserted. You will need to know the natural frequency  $\omega_0$  (as defined in the Theory section) for this experimental system. Does any of the three measured frequencies corresponds to the natural frequency? If you think it is, which one is the natural frequency? If you think

it is not, explain how to estimate the natural frequency more accurately and present the result.

## **2. Measurement and characterization of a forced harmonic oscillator with a large damping disc**

**[Now perform the following experiments with the motor turned on]**

B. Make sure that a large damping disc is installed. Run trial experiments with varying the motor frequency to roughly find the resonance frequency. Then adjust the amplitude of the external force (i.e., by changing the length of the rotating arm) appropriately so that spring is stretched somewhat extensively but not to the extent that the mass hits the floor or spring force is in nonlinear regime. Using Logger Pro, record the motions of oscillation (both angle and position sensors) as the motor frequency is varied (i.e., by changing the input voltage) over its entire range; and export the data to a 'txt' file at each frequency. Measure the amplitude and phase of the oscillation at each motor frequency using MATLAB. For each frequency setting, measure the period of the motor. Make your measurements very carefully, taking closely spaced frequency points near the peak of the resonance where the amplitude and phase are changing rapidly so that your data can be well compared with theoretical curves as in Figs. 1-3 in the Theory section of this manual. For each measurement allow time for the system to reach steady state after you changed the rotation speed.

C. Use MATLAB to make plots of  $A$  vs.  $\omega$ ;  $v_{max}$  vs.  $\omega$ ; and  $\phi$  vs.  $\omega$ . From your graph of  $v_{max}$  determine  $\omega_0$  and estimate your error. [ $v_{max}$  is calculated from the experimental data using the equation  $v_{max} = \omega A$ .] Compare this value of  $\omega_0$  with the natural frequency you determined in part A. How well do they agree? Also, from your plot of  $v_{max}$ , estimate  $2\Delta\omega$  from the points where  $v = 0.7 \times peak(v_{max})$ . Estimate the error. Use Eq. (13) to determine  $Q$  and Eq. (12) to determine  $\gamma$  and  $R$ . Estimate your errors.

D. Using the values  $\omega_0$  and  $\gamma$  you determined in part C, use MATLAB to calculate the theoretical frequency dependence of  $v_{max}$  and  $\phi$ . Add these curves to the plots you prepared in part C; and compare the experimental and the theoretical curves.

E. In part D you used a rather crude procedure to determine the parameters  $\gamma$  and  $\omega_0$ . What, in general terms, would be a better way to fit the theory to the experiment?



**Lab Activity for Week 2 (Forced Harmonic Oscillations for Smaller Damping)**

F. Repeat Parts A-E for a small paper plate instead of the large damping disc made out of cardboard. In this case, make sure you carefully adjust length of the rotating arm as instructed in Part B so that the extension of the spring at resonance is still in the measurable range in linear regime.

G. Now repeat Part F without even the small paper plate. In this case, the damping (intrinsic to spring, pulley, etc.) will be very small and the oscillation amplitude will be dramatically amplified at the resonance. Try to make the length of oscillation arm short enough for the motion to be reliably measurable at resonance.

H. Compare the results of the three damping cases by summarizing them in a table. The table should include the following information for each case: amplitude of angular motion detected by the rotation sensor (this represents the force amplitude  $F_0$ ); mass of the oscillator (i.e.,  $m$ ); natural frequency (i.e.,  $\omega_0$ ); and damping constant (i.e.,  $\gamma$ ).

I. Plot the normalized amplitude of displacement oscillation vs. normalized angular frequency (i.e.,  $f(\bar{\omega})$  vs.  $\bar{\omega}$  : see the lecture note), both experimental data and theoretical curves, for all the three damping cases in the same figure to compare them.

J. Plot the phase difference between the displacement and force vs. normalized angular frequency (i.e.,  $\phi(\bar{\omega})$  vs.  $\bar{\omega}$  : see the lecture note), both experimental data and theoretical curves, for all the three damping cases in the same figure to compare them.

K. Repeat Part I for the normalized amplitude of velocity oscillation (i.e.,  $g(\bar{\omega})$  vs.  $\bar{\omega}$  : see the lecture note), both experimental data and theoretical curves, for all the three damping cases in the same figure to compare them.

**Lab Report Submission Instruction**

Submit a single lab report that includes all the results for A-K by the due date set for the week-2 lab.

The total lab grade will be 40 points.

20 points: Introduction / Experimental Methods / Parts A-E

20 points: Parts F-K / Conclusions

**HINTS FOR USING LOGGER PRO**

To setup the program, go to:

Menu EXPERIMENT

- CONNECT INTERFACE
- CONNECT ON PORT: select COM1

Button LABPRO

- Drag MOTION DETECTOR to DIG/SONIC2 box
- Drag Rotary Motion to DIG/SONIC1 box
- CLOSE

Menu EXPERIMENT

- DATA COLLECTION ...
- Tab COLLECTION
- LENGTH 10 seconds
- Sample at Time Zero: off
- SAMPLING RATE about 20 samples per sec
- DONE

**MATLAB NOTE:**

The `atan()` function returns values in a range of angles that is wrong for our definition of the phase shift (Eq. (8)). To make the angles in array  $\phi = \text{atan}(\dots)$  fall in the correct range, use:

$k = \text{find}(\phi > 0)$ ;  $\phi(k) = \phi(k) - \pi$ ;

Explain in your report the difference between ranges of  $\phi$ .