First proposal for the REWOLF project

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Abstract

In this proposal, we show that by using standard instrumentation we can measure a possible fluctuation of the group velocity of light as predicted by some theories. We also introduce the idea to integrate this apparatus in a existing long baseline experiment to reach the finest results.

1 General idea

The speed of light can be decomposed in three components.

- The first is defined as the maximum speed, and leads to the Special Relativity. This speed (c_{rel}) appears in the Lorentz transformation and is the one that defines the relationship between the mass and the energy.
- The second is the magnetic wave speed, which is commonly called phase velocity (c_{Φ})
- The last is the speed of propagation of a photon (in the quantic mean of impulsion). It is called the group velocity (c_{group}) , and it is defined in average as $c_{group} = L/T$ where L is the distance and T the transit time through L.

While c_{rel} is not expected to fluctuate and is strongly tested, some theories predict that c_{Φ} and c_{group} can fluctuate during the propagation in vacuum.

The objective of the REWOLF project (name?) is to detect a possible group velocity fluctuation of light in vacuum, larger than a possible phase velocity fluctuation. It is important to note that, even if a limit based on cosmological observations was put, among the current experiments, no dedicated experiment has yet reached the predictions from recent theoretical models.

In the considered theories, this fluctuation is expected to be stochastic and can be studied with the arrival time of a short light pulse after propagation in the vacuum. The longer the propagation path, the wider the timing enlargement of a pulse is.

In the first family of theories (Th 1), the group velocity fluctuations are predicted if the photon propagation is considered as a succession of interactions with pairs of ephemeral (virtual) fermions. This kind of model is developed in [1] and independently in [2]. More generally, from such definition of the vacuum filled with virtual fermions, one can demonstrate that the values of the parameters μ_0 and ϵ_0 and the group speed of light appear naturally at the correct scale. Moreover, the equation $\langle c_{group} \rangle = c_{Phi} = \sqrt{1/(\mu_0 \epsilon_0)}$ where $\langle c_{group} \rangle$ is the mean value of the group velocity, can be deduced independently of the Maxwell equations.

This description of the vacuum predicts fluctuations of the photons' transit time of the order of $\sigma_t(\text{FWHM}) \approx 120 \, [\text{as}] \cdot (L[\text{m}])^{-1/2}$ (attosecond, 10^{-15} s) at the distance L. The experimental signature of such an effect is an enlargement of the light pulse temporal envelop during its propagation. This enlargement increases proportionally to the square root of the path length in the vacuum.

There are two things to note in this model. First, the lifetime of the virtual pair is inversely proportional to the square root of the total pair energy. This order of magnitude of 120 as is given for a peaked energy distribution, considering averaged parameters instead of their full distribution (as life time of pairs, Heisenberg principle...). The second one is that since no final interpretation is given to this energy distribution, its average can be bigger (with weak effect on the fluctuations order of magnitude), or smaller (a non relativistic level of energy bring the fluctuation FWHM at $\sigma_t \approx 240 \, [{\rm as}].(L[{\rm m}])^{-1/2}$.

In this model, this phenomenon is totally independent of the light wavelength and the phase as it relies only on a photon interaction probability with the ephemeral pairs.

In the second family of theories (Th 2), the effects are at smaller scale, but for some of them and under certain conditions, they are reachable in our proposed experiment. These theories are based on physics beyond the standard model, like Quantum Gravity [3] and have a sensibly different behaviour as a function of distance and energy. Assuming additional dimensions, the Quantum Gravity theory implies a dispersion in function of distance of $[s]/[m]^n$ where $n \geq 1$. If a speed of light fluctuation is detected, it is essential to measure the spreading distribution at different distances to differentiate the theories. Note that in these theories we can have a velocity fluctuation dependency on energy, that make interesting to do a spectral measurement of the arriving pulse.

There is a third family of theories (Th 3) concerning space/time fluctuations at Planck scale. That implies a too small scale to be detected by the proposed method, or that the effect impacts also the phase speed which can be measured much more precisely thanks to interferometry [4]. It is important to note that any model that implies a space/time fluctuation as an effect on c_{Phi} and c_{group} at the same scale, an interferometry system is more adapted, due to its very accurate precision.

The current limit put on this fluctuation effect is $\sigma_t(\text{FWHM}) \approx 700 as. m^{-1/2}$ [5]. The purpose of this project is to push this limit down to the theoretical predictions, and, possibily, to distinguish different theoretical fluctuations predictions.

In the case of 100 m light path and 2 fs detector resolution the current cosmological limit can be reached in laboratories.

The final goal of this experiment is to use ultrashort pulses of 5 fs FWHM and a path of 1 to 10 km. According to the predictions of the foregoing first family of theory, the time width of the pulse is expected to be spreaded to 9 fs (FWHM). It is possible to see this enlargement using detectors with a resolution of about 2 fs.

The best method to characterize both the initial and final pulse must be an intensity auto-correlation system. This type of detector is described in the section 2.3 We will see in the next parts a description of the different phases of the experiment and the suggested technical solutions.

2 Experimental setup description.

The setup has three main components: the light sources, the propagation path and the detector. We plan several phases of the experiment:

Phase-0) 6 months: Laboratory tests of the source and the detector. At this phase we verify the source properties and accuracy of the detector.

Phase-1) 8 months: Laboratory measurements with a multi-reflections cell [6]. In this phase we plan to reach the current cosmological limit of the light transit time fluctuations. The light path should reach the 100 m scale. It was calculated that an accurate measurement can overtake this limit by about 30% (fig. 3).

Phase-2) 10 months: Using the knowlege out of the phase-0 and phase-1 (source, mirrors and detector), the final experimental setup will use a km-scale path. Here we plan to reach the model limit for the light transit dispersion. In this proposal we want to investigate two possibilities - an improved Herriot cell, or to integrate the setup with a long baseline detector. In Italy, a very good candidate could be the VIRGO facilities ¹.

¹https://wwwcascina.virgo.infn.it/

The research for a future integration in a long baseline detector will be under preparation and discussion during the full developpement process. This element is for us a very important point, because, further a better sensitivity, such an integration can bring to a better perenity of the experiment, thanks to the dependency to a much bigger collaboration.

In the further sections the technical solutions of the experiment components will be described.

2.1 Source

In the modern ultrashort pulse applications, the Titanium-Sapphire lasers are used thanks to the handling simplicity compared to the previously used dye lasers. Ti-Sapphire lasers able to produce pulses with a width below 10 fs. The typical pulse energy is 2-6 nJ and the rate is about 80 MHz. The wavelength has a range of about 700-1000 nm². The divergence is < 2 mrad and the beam diameter is < 2.5 mm. For the pulse measurement an energy of several nJ may not be enough. In this case an amplification should be used. Ti:Sapphire multipass amplifier with the dispersion compensation is able to provide < 25 fs pulses with energy of 1 mJ, but reduce the pulse rate to the kHz scale, keeping the divergence < 2 mrad and increasing the beam diameter to 15 mm. Stretching of this pulses is possible to < 7 fs with hollow fiber compressors. 3 . The state of art for the compression is 4 fs from the amplificator. At this order of magnitude the measurement method gives the real limitation on the time resolution. 4

2.2 Light path

Considering the phase-1 of the proposal we need a Herriott cell at meter scale with about 100 reflections on mirrors. To ensure < 2/3 energy loss at the end of the path > 0.99 reflectivity coefficient is required. This requirement is possible with gold mirrors with a roughness of less than 100 Å. The state of art is the 1 Å scale 5 . This kind of mirrors are commercially available. The time dispersion due to thermal fluctuations of the mirrors are expected to be negligible comparing to the dispersion due to the mirror roughness, but it should be tested in the phase-0.

The air group velocity dispersion is of the order of $1 \text{ fs}^2/\text{m}$ (at 1 atm) [7] Which represents, for a pulse of 5 fs width, a dispersion of 0.2 fs/m. This dispersion is directly proportional to the pressure. That means that to reach a negligible total dispersion (less than 10 as), the pressure should be less that 0.5 mbar for the 100 m scale experiment, and less than 0.05 mbar for the km scale. These orders of magnitude are quite easy to reach in both cases.

According to the experience from the phase-1, the extension to the phase-2 is straightforward. Due to the length of several kilometers, 0.1 μ rad divergence of the light source is needed to keep cm scale beam radius. This constraint will require an additional optical collimator.

Note that in the case of the integration in a long baseline experiment like VIRGO, the pressure inside the detector is very low $(10^{-11} \text{ mbar for VIRGO})$, and is far better that the requirements for the proposed experiment. Moreover, the use of such a facility would strongly reduce the number of reflections, which will decrease the systematic errors sources and the complexity.

For a long baseline detector, a dedicated study will be required to determine the compatibility of our detector set up with the facility. The avaibility of experienced peoples in our Institute in the VIRGO collaboration allowed us to get a first contact, and will help to clarify this part.

2.3 Detector

For this experiment a direct spectrometry or phase measurement method cannot be used, because they rely on some assumptions about the pulse properties and use the Fourier limit to deduce its shape.

²Rainbow oscillator www.femtolasers.com,Octavius-85www.thorlabs.de

 $^{^3}$ The data are taken for CEP Pro amplifier + Kaleidoscope compressor www.femtolasers.com

 $^{^{4}} MOSAIC \qquad OS \qquad compressor \qquad \text{http://pdf.directindustry.com/pdf/femtolasers-produktions-gmbh/mosaic-os-octave-spanning-gdd-module/36035-231646.html}$

⁵http://www.thorlabs.de/NewGroupPage9.cfm?objectgroup_id=7028

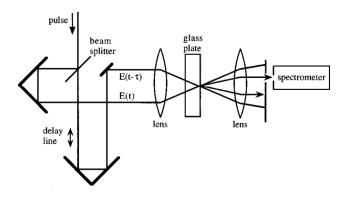


Figure 1: FROG system scheme. The pulse is splitted, one component is delayed. A beam is produced by harmonic generation autocorrelation in a crystal. The generated beam is detected by a spectrometer.

Sensitivity index for an optimal detector resolution of 0.2 fs

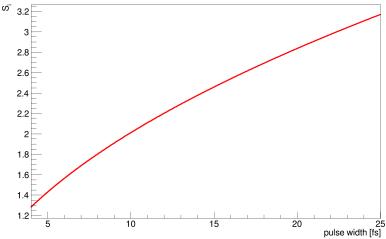


Figure 2: Sensitivity index for a resolution of 0.2 fs in function of the laser pulse width.

On the other hand an autocorrelation system is based on the effective interaction of the photon with a cristal, which allows to determine precisely the electric field of the pulse without any assumptions.

For this alternative there are several kinds of autocorrelator detectors, and FROG (Frequency Resolved Optical Gating) which satisfy the requirements.

In general the measure is done by splitting the pulse beam in two parts and sending them on a non-linear crystal. The crystal, thanks to the second or third harmonic generation effect, will produce an output beam proportional to the cumulative intensity of the input beams. By adding a delay to one of the incoming beam (changing one path length), it becomes possible to measure the full autocorrelation of the incoming pulse.

By using a surface third harmonic generation FROG it is possible to characterize pulses with 5 fs width and an intensity of few nJ [8]. The time resolution of such a system can reach the sub fs order of magnitude (0.2 fs).

2.4 Perspectives and conclusions

If we consider a light pulse as a gaussian photon distribution, the sensitivity index of this experiment can be calculated by

Sensitivity index for a laser between 4 fs and 20 fs

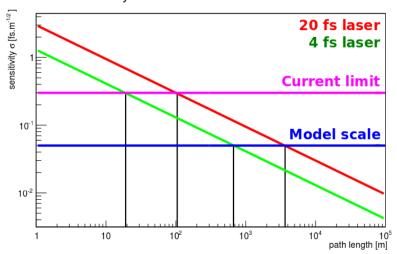


Figure 3: Sensitivity for different cases of laser pulse width (detector resolution of 0.2 fs). The current limit is given by cosmologic analysis, the Model scale is for the model order of magnitude defined in the Th. 1.

$$S_i = \sqrt{(R + P_W)^2 - P_W^2}$$

where R is the total resolution of the experiment, and P_W the pulse width. That allows to get the total sensitivity S:

$$S = S_i / \sqrt{L}$$

Given that the FROG sensitivity is not the most heavy part in the total cost, we consider now that the priviliged solution is a 0.2 fs resolution. It permits to evaluate the sensitivity index, represented in the figure 2 and the total sensitivity, represented in the figure 3.

Therefore we can set the goals of the experiment:

In the most pessimistic solution for the laser, we can push back the current limit with a 100 m path and the detection of the fluctuations are allowed for about 3 km to 6 km. We could exclude the fluctuation hypothesis at 3 sigmas at a path of 4 km and a path of 10 km can exclude it at 5 sigma (considering a fluctuation of $50 \text{ as.m}^{-1/2}$).

In the case of a fluctuation detection, to exclude the linear dependance of the fluctuations with distance (instead a square root dependence, $S = S_i/L$), which expected in some alternative models (although at smaller scale) 18 km measurement (equivalent to 3 round trip in the Virgo tunnel) is enough to reach 5 sigma difference.

In the most optimistic case of a laser of 4 fs, we could push back the current limit at the ten meter scale (easy to reach in phase-1), and gain one order of magnitude for the fluctuation hypothesis detection or exclusion.

Therefore it is clear that this research program is very ambitious. Even if it uses common and well known technologies (almost all of the material is commercially available for industrial uses), and it will need a good experimental mastery and a good experimental artifact control and analysis.

Nevertheless, for us, these skills seem to be reachable. The team that we are forming combine a theoritical knowledge, an reliable experimental experience and a steady electronics/DAQ. We are convinced that, thanks to such an experiment, we will be able either to detect a group velocity fluctuation, or to set a new strong limit that confirms that the group velocity and show that the phase velocity are indifferentiable at a scale of the attosecond per square meter (and much further for a direct proportionality with the distance). In

any case this result would be very interesting for the physics community, so we wish to get the opportunity offered by your funding proposal to realize this experiment.

References

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