

PhysH308

Spinning *with* things!

Ted Brzinski, Nov. 7, 2024



Pre-registration ends this week!

There are a lot of upper-level options available this Spring!

Physics courses:

1. **PHYSH302: Advanced Quantum**, with Walter Smith
2. Bryn Mawr is offering **PHYSB309 Advanced E&M**
3. **PHYSH304: Computational Physics** with our new faculty member Vijay Singh (whose research is in computational biological physics).
4. **PHYSH353: Topics in Soft Matter Physics**, a special topics course with visiting faculty member Vianney Gimenez-Pinto.
5. Bryn Mawr is offering **PHYSB331: Advanced Experimental Physics**

Not listed as physics, but physics, Clyde Daly in Chemistry is teaching **CHEMH350: Topics in Computational Chemistry** (time TBD), which can use PHYSH214 as a prereq.

Astrophysics/astronomy courses:

6. **ASTR344: Topics in Astrophysics**: Gravitational Waves with Andrea Lommen.

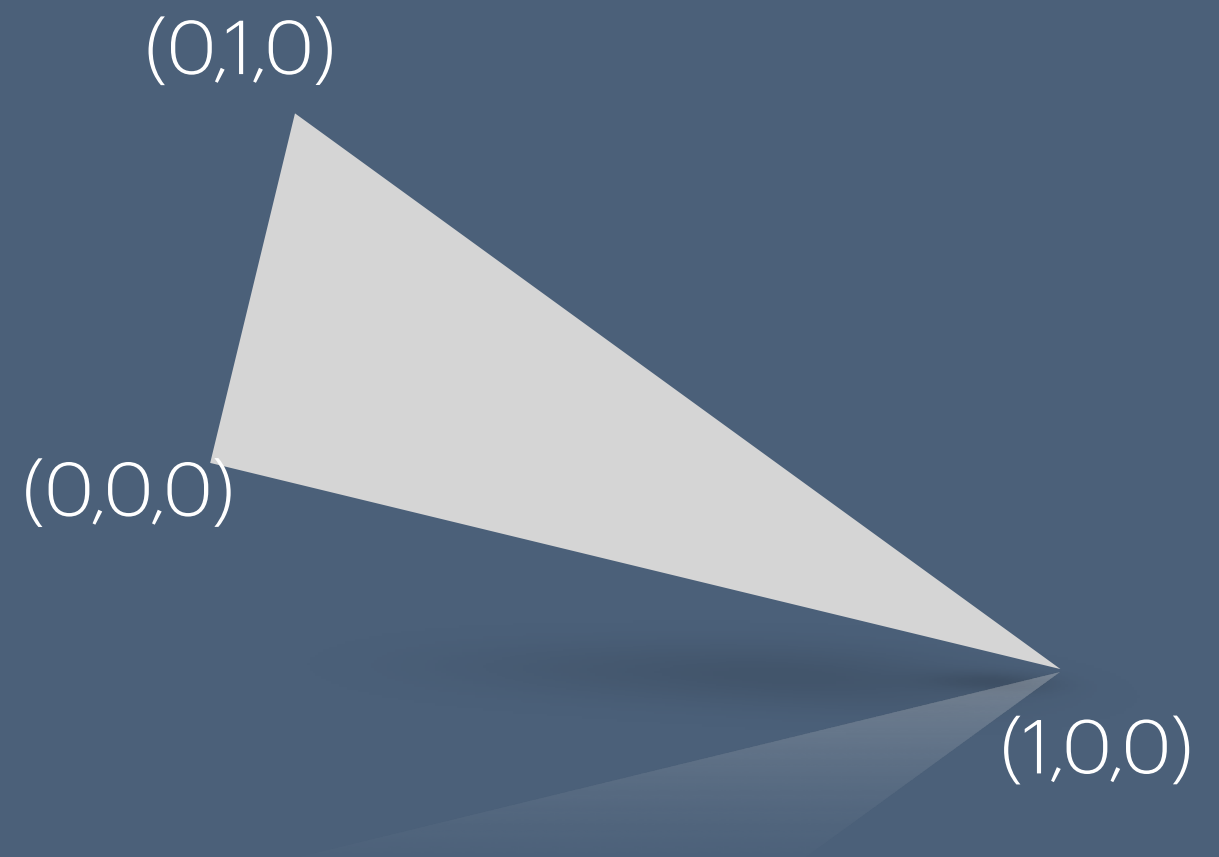
Problem 10.37

(From last week)

- Given a shape, find \mathbf{I} , diagonalize it, and find the principal axes
See that from \mathbf{I} in any coordinates you can find the principal axes

$$I_{xx} = m \sum (y^2 + z^2)$$

$$I_{xy} = I_{yx} = -m \sum xy$$



Why find principle axes?

- Diagonal matrices are easier to deal with!
- For rotation about these axes, $\vec{L} \parallel \vec{\omega}$ (no torque required to maintain $\hat{\omega}$) - this tells us about the (rotational translation) symmetries of the object!



Motion of a top

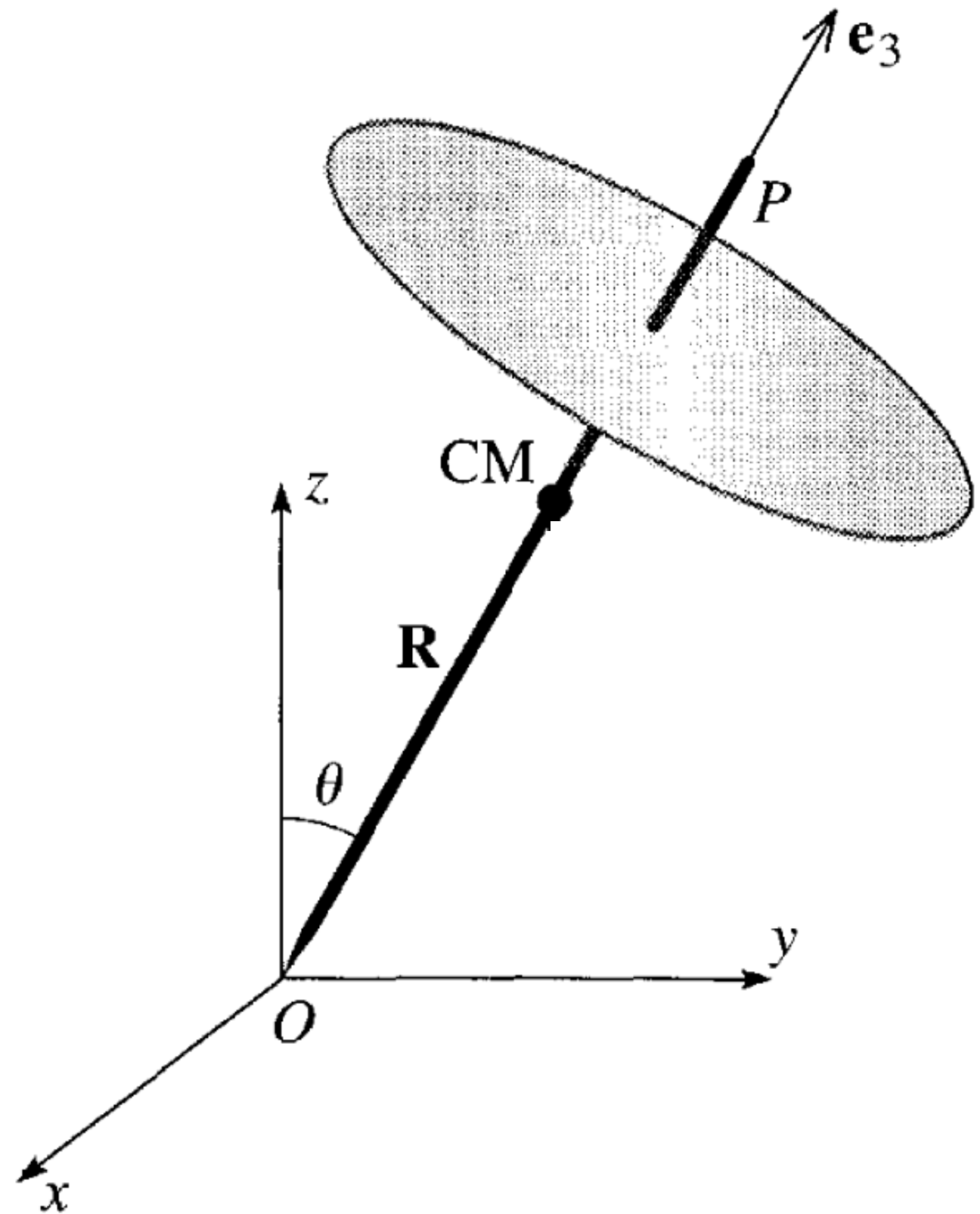
(First, without gravity)

Consider rotation of an axially symmetric about an axis of symmetry \hat{e}_3 through the center of mass:

$$\mathbf{I} = \begin{pmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_1 & 0 \\ 0 & 0 & \lambda_3 \end{pmatrix} \text{ and } \vec{\omega} = \begin{pmatrix} 0 \\ 0 \\ \omega \end{pmatrix}$$

$$\text{So } \vec{L} = \mathbf{I}\vec{\omega} = \lambda_3 \omega \hat{e}_3$$

Now let's add gravity!



Motion of a top

Now let's add gravity!

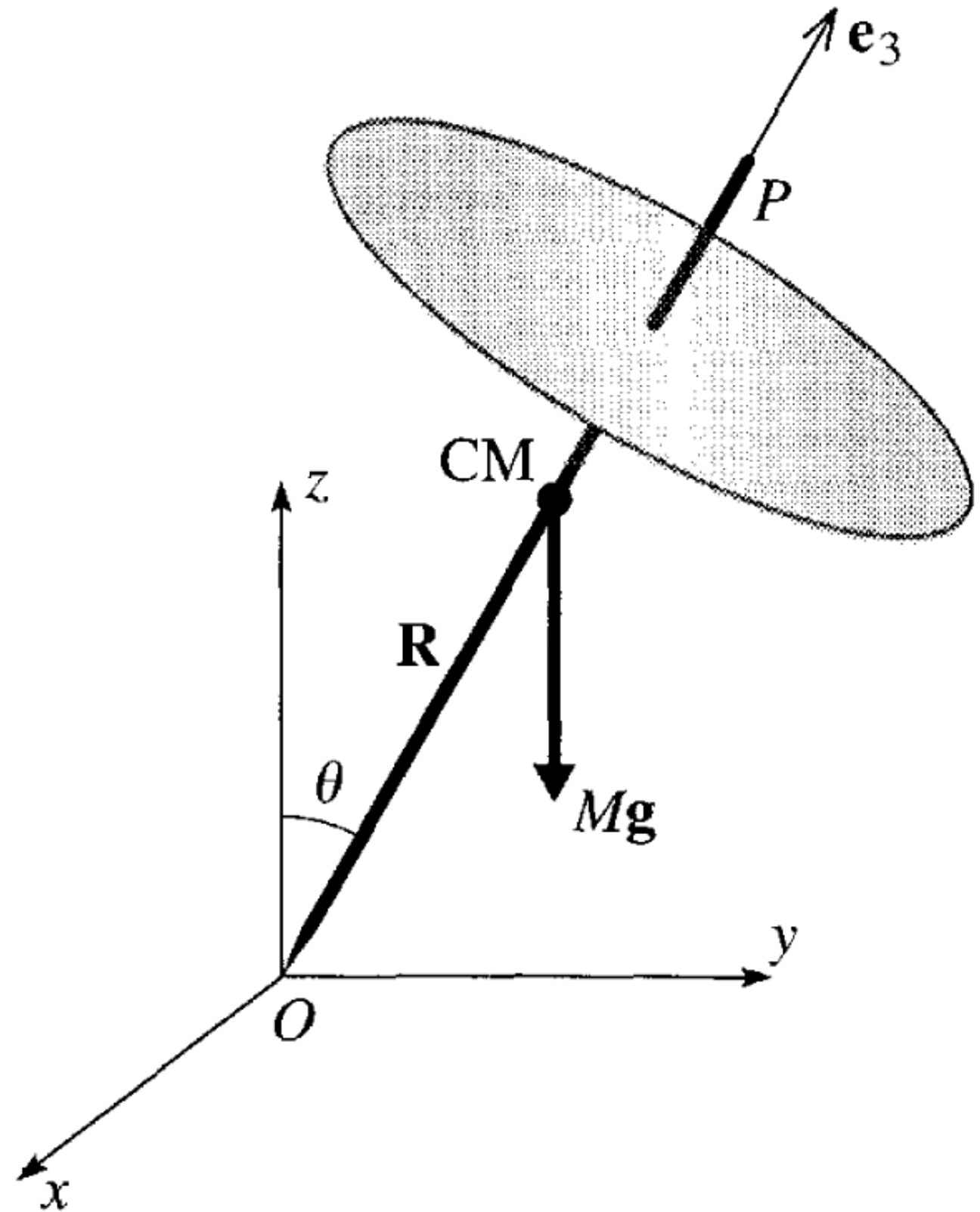
Gravity applies a force at the center of mass: $\vec{F}_{cm} = -Mg\hat{z}$

The result is a torque:

$$\vec{\Gamma} = \vec{F}_{cm} \times \vec{R}_{cm} = MgR_{cm}\hat{e}_3 \times \hat{z}$$

$$\text{Recall } \vec{\Gamma} = \dot{\vec{L}} = \lambda_3\omega\dot{\hat{e}}_3$$

$$\text{So } \dot{\hat{e}}_3 = \frac{MgR_{cm}}{\lambda_3\omega}\hat{e}_3 \times \hat{z}$$



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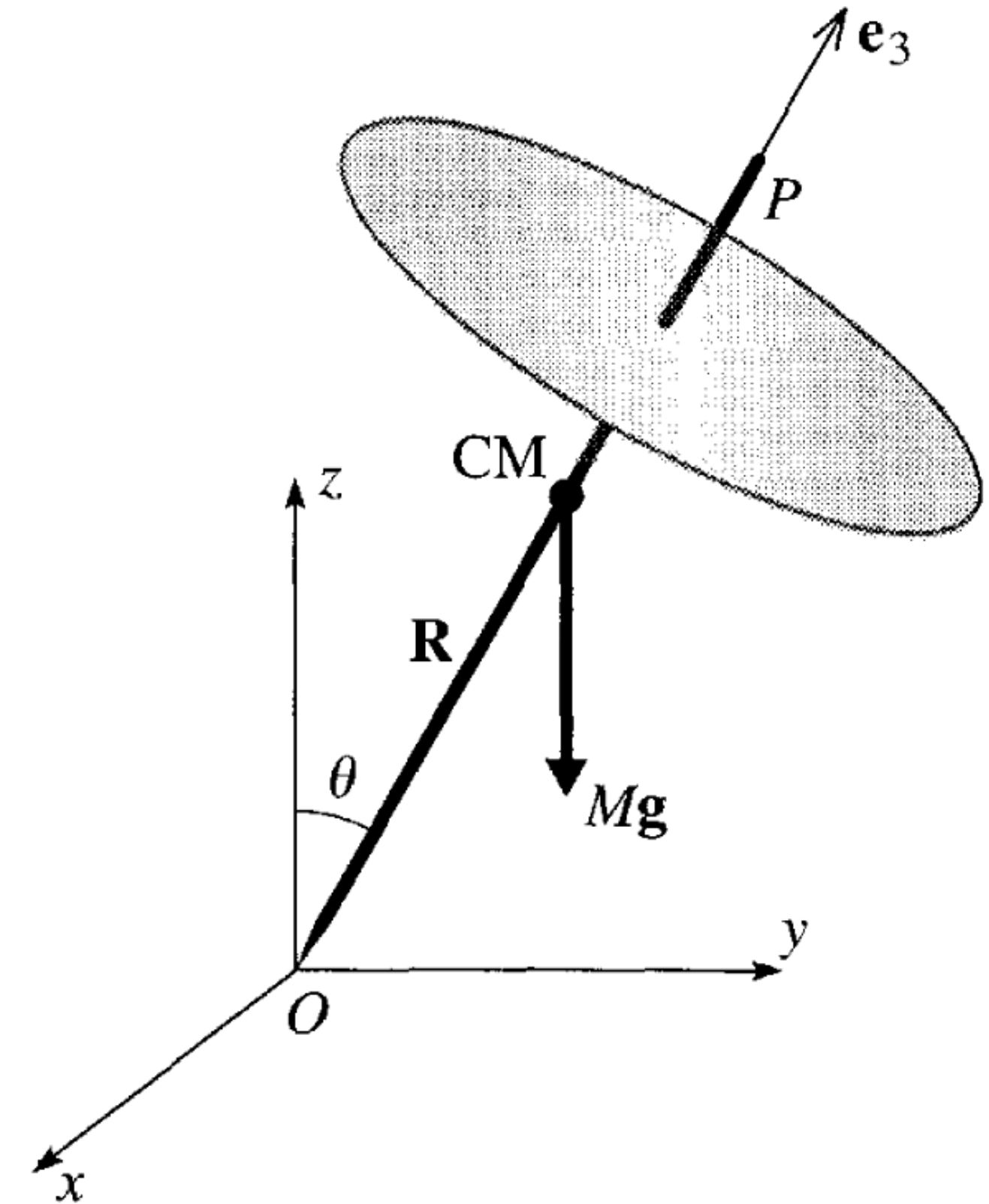
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$$\text{So } \dot{\hat{e}}_3 = \frac{MgR_{cm}}{\lambda_3 \omega} \hat{e}_3 \times \hat{z} = \hat{e}_3 \times \vec{\Omega}$$

Precession about \hat{z} with rate $\vec{\Omega}$!



Motion of a top

Now let's add gravity!

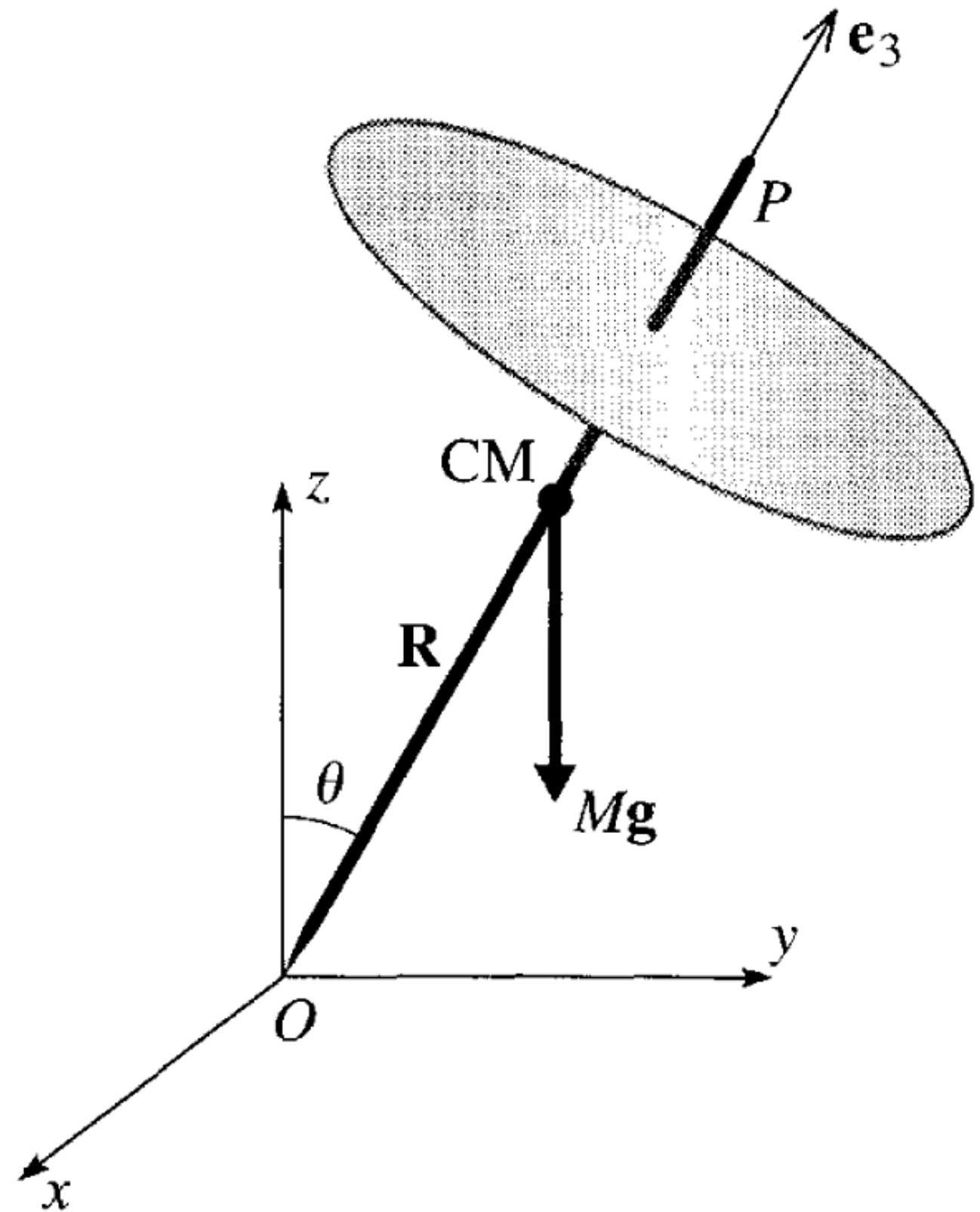
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Precession about \hat{z} with rate $\vec{\Omega}$!

You also can get small oscillations in $\hat{\theta}$ (*nutation*) which we'll explore Thursday

[Spintop demo!]

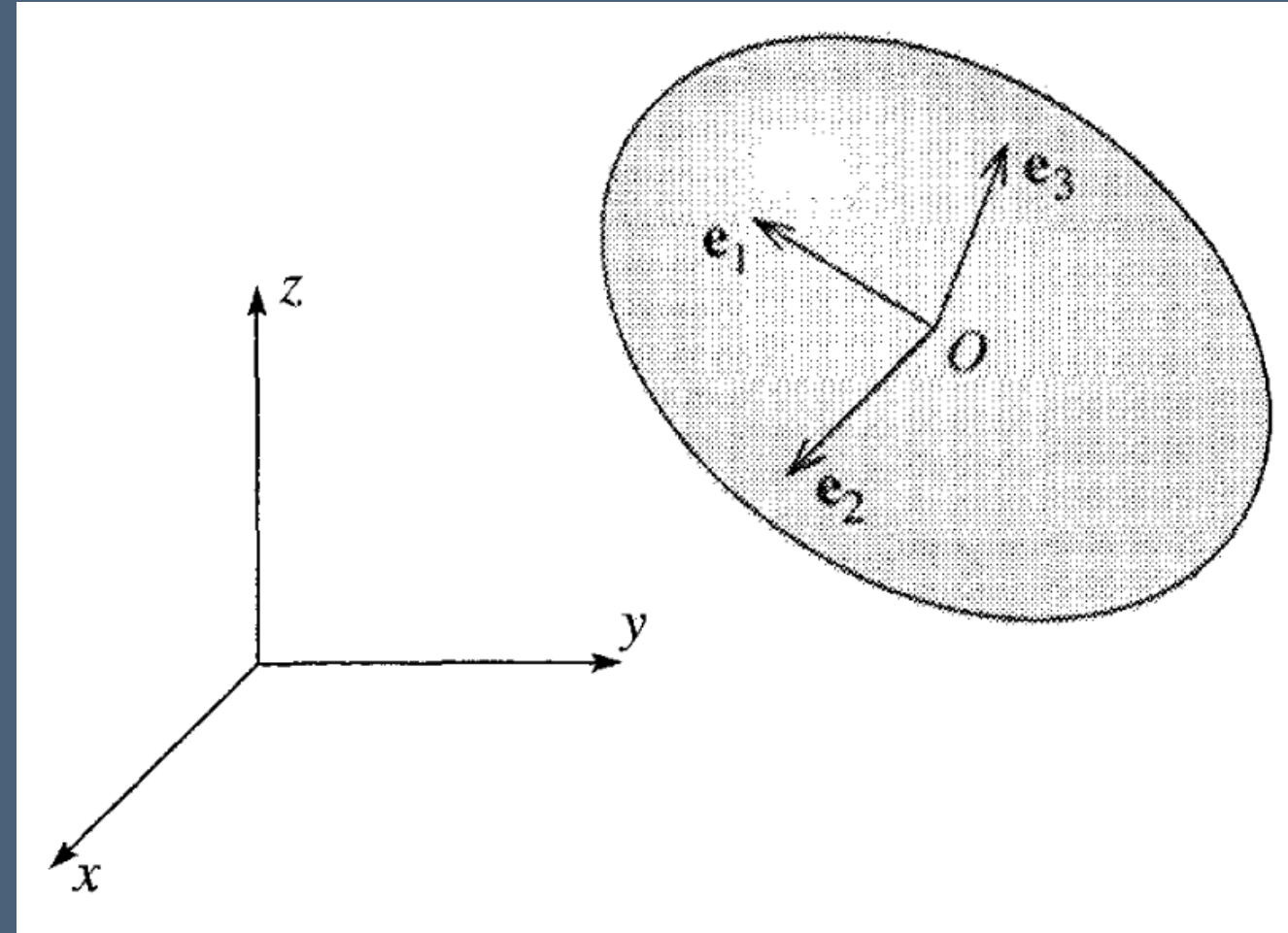
So how can we explain the Tippe Top? [DEMO 2]



Body vs space frame

- So far we've mostly discussed behavior in the "space frame" (inertial, "lab")
- Consider a "body frame"
 - Origin fixed to the body (usually CoM), axes = principal axes
 - Co-rotating with the object,
 $\vec{\Omega} = \vec{\omega}$

$$\vec{L} = \begin{pmatrix} \lambda_1 \omega_1 \\ \lambda_2 \omega_2 \\ \lambda_3 \omega_3 \end{pmatrix}$$



Body vs space frame

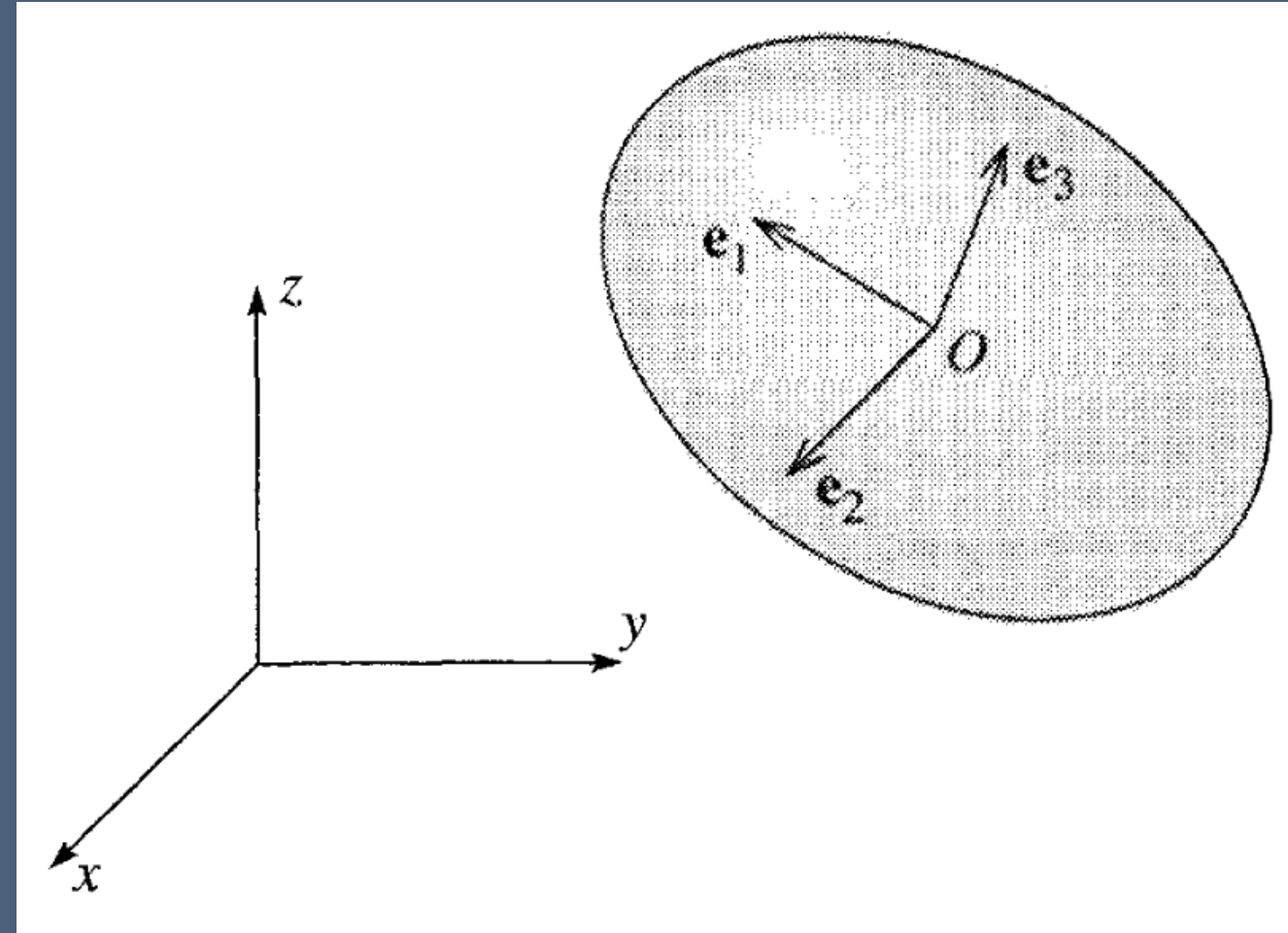
- $$\vec{L} = \begin{pmatrix} \lambda_1 \omega_1 \\ \lambda_2 \omega_2 \\ \lambda_3 \omega_3 \end{pmatrix}$$

- *But in a rotating frame, there is an inertial torque! (Ch. 9)*

- $$\vec{\Gamma} = \dot{\vec{L}}_{space} = \dot{\vec{L}}_{body} + \omega \times \vec{L}$$

- **Euler's Equations:**

$$\vec{\Gamma} = \dot{\vec{L}} + \omega \times \vec{L}$$



Body vs space frame

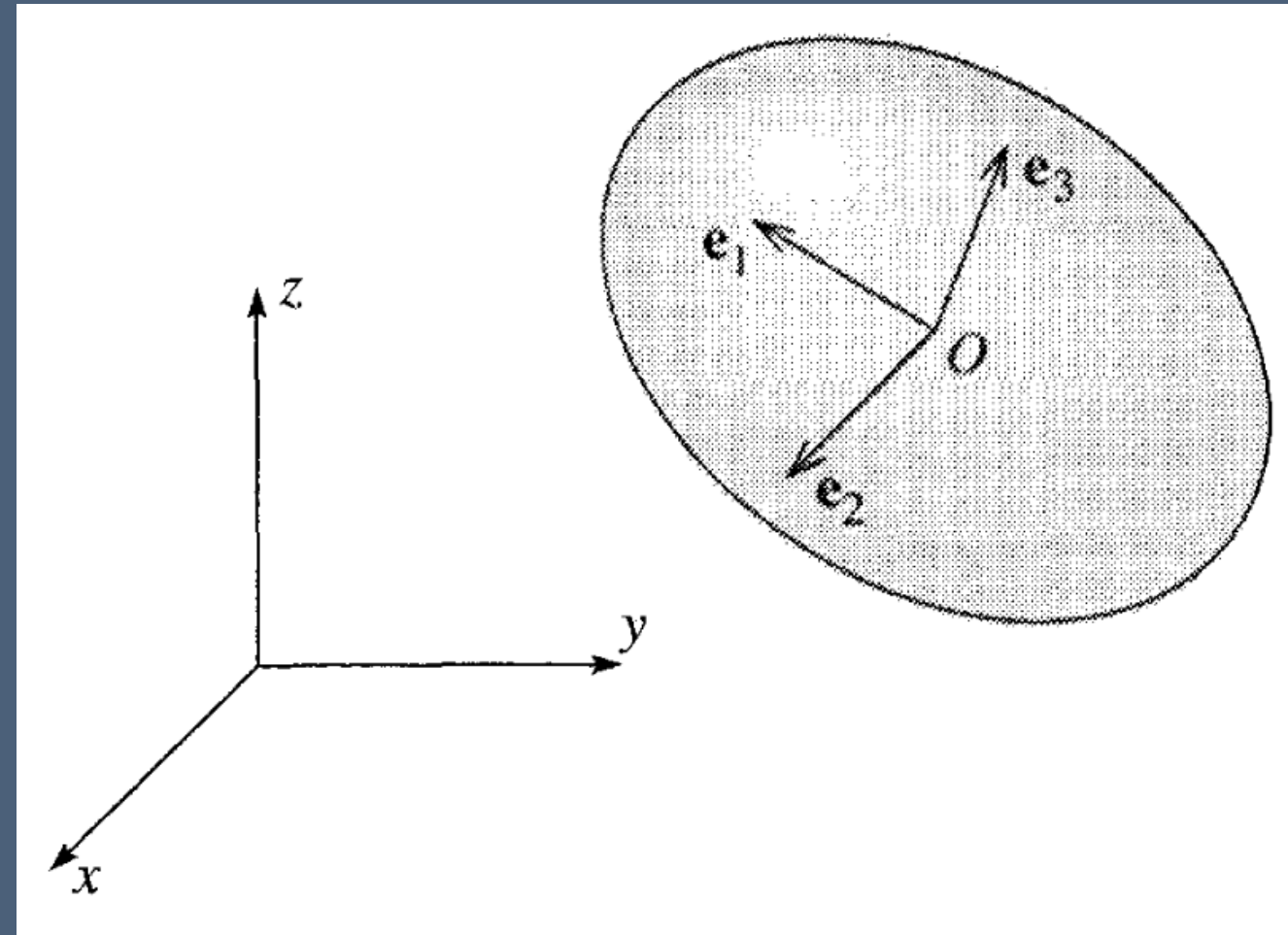
- **Euler's Equations:**

$$\vec{\Gamma} = \dot{\vec{L}} + \omega \times \vec{L}$$

- $\Gamma_i = \lambda_i \dot{\omega}_i - (\lambda_j - \lambda_k) \omega_j \omega_k$

- *With no external torque this becomes:*

$$\lambda_i \dot{\omega}_i = (\lambda_j - \lambda_k) \omega_j \omega_k$$



Rotation without torque

- **Euler's Equations:**

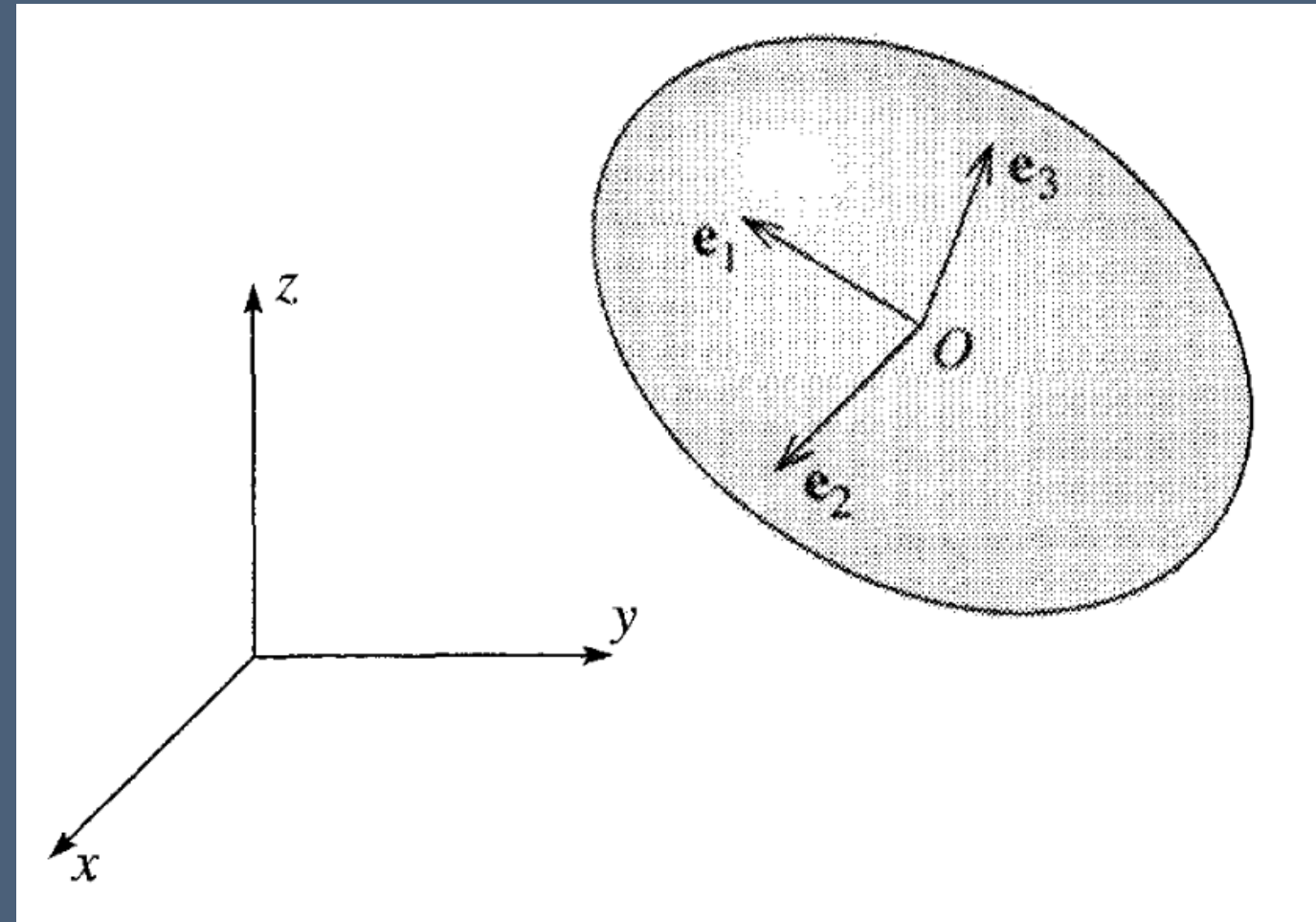
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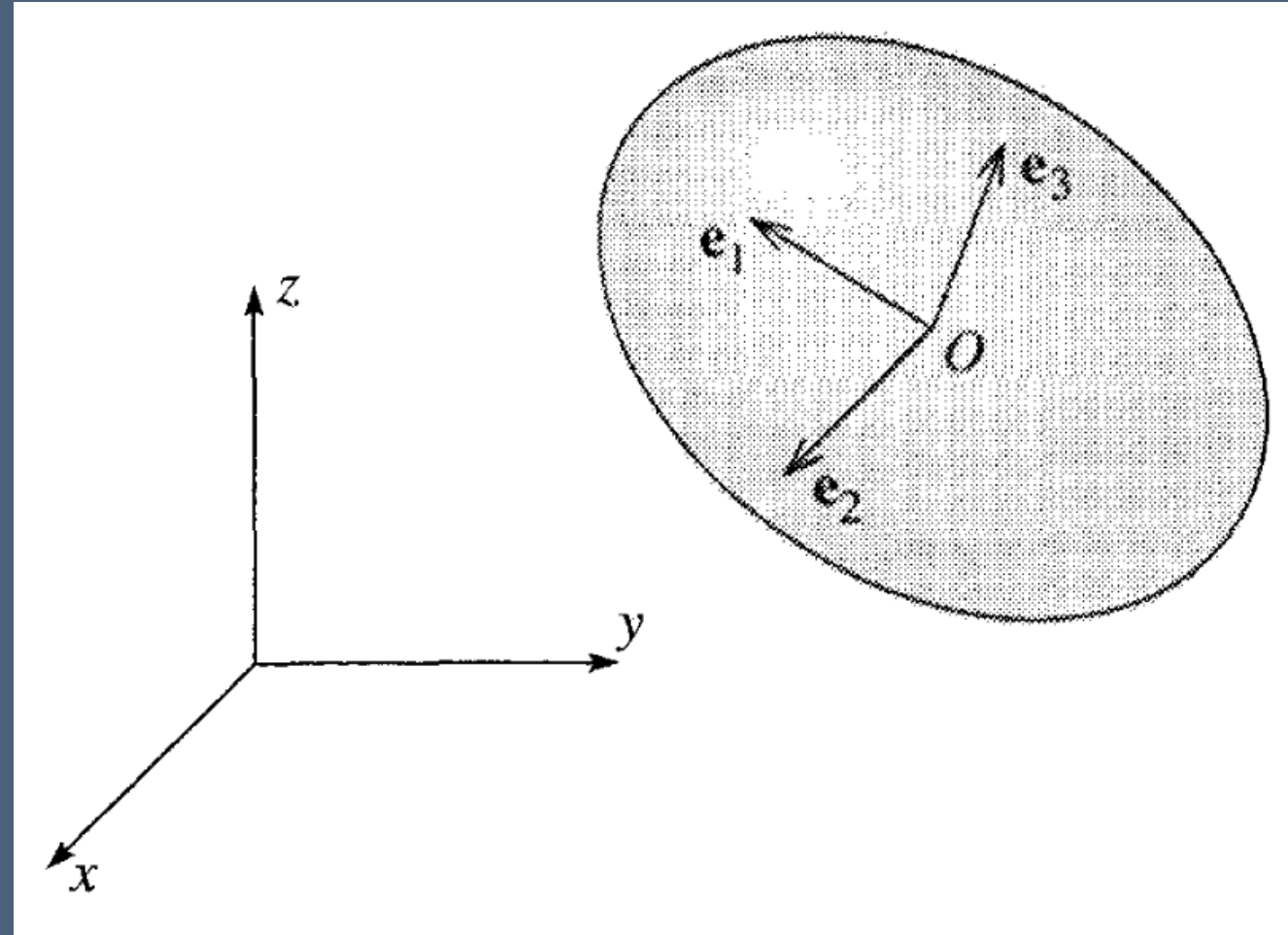
$$\lambda_i \dot{\omega}_i = (\lambda_j - \lambda_k) \omega_j \omega_k$$

- *Consider the case where*
 $\omega_3 \approx \text{constant} \gg \omega_1, \omega_2$



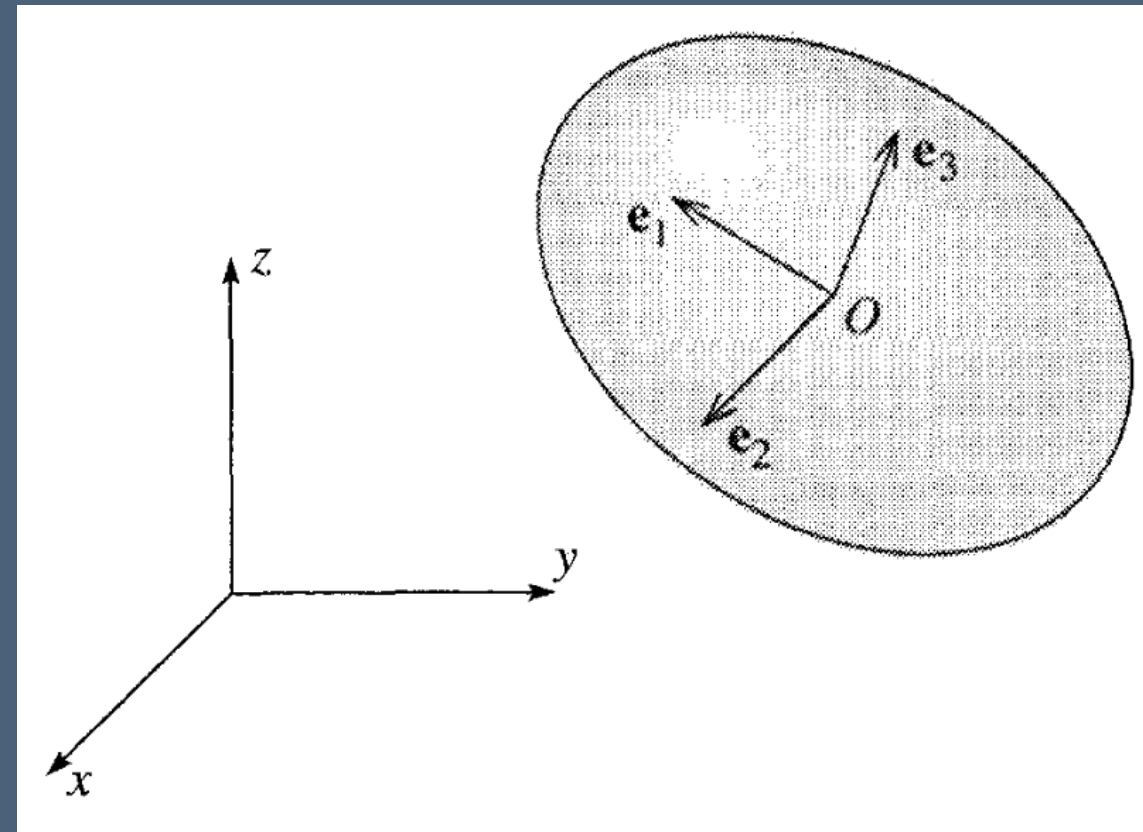
Rotation without torque

- $\lambda_i \dot{\omega}_i = (\lambda_j - \lambda_k) \omega_j \omega_k$
- Consider the case where $\omega_3 \approx \text{constant} \gg \omega_1, \omega_2$
- $\lambda_1 \dot{\omega}_1 = [(\lambda_2 - \lambda_3) \omega_3] \omega_2$
- $\lambda_1 \ddot{\omega}_1 = [(\lambda_2 - \lambda_3) \omega_3] \dot{\omega}_2$
- $\lambda_2 \dot{\omega}_2 = [(\lambda_3 - \lambda_1) \omega_3] \omega_1$



Rotation without torque

- $\lambda_i \dot{\omega}_i = (\lambda_j - \lambda_k) \omega_j \omega_k$
- Consider the case where $\omega_3 \approx \text{constant} \gg \omega_1, \omega_2$
- $\lambda_1 \dot{\omega}_1 = [(\lambda_2 - \lambda_3) \omega_3] \omega_2$
- $\ddot{\omega}_1 = - \left[(\lambda_3 - \lambda_2) (\lambda_3 - \lambda_1) \frac{\omega_3^2}{\lambda_1 \lambda_2} \right] \omega_1$



e.g. wobbly frisbee!

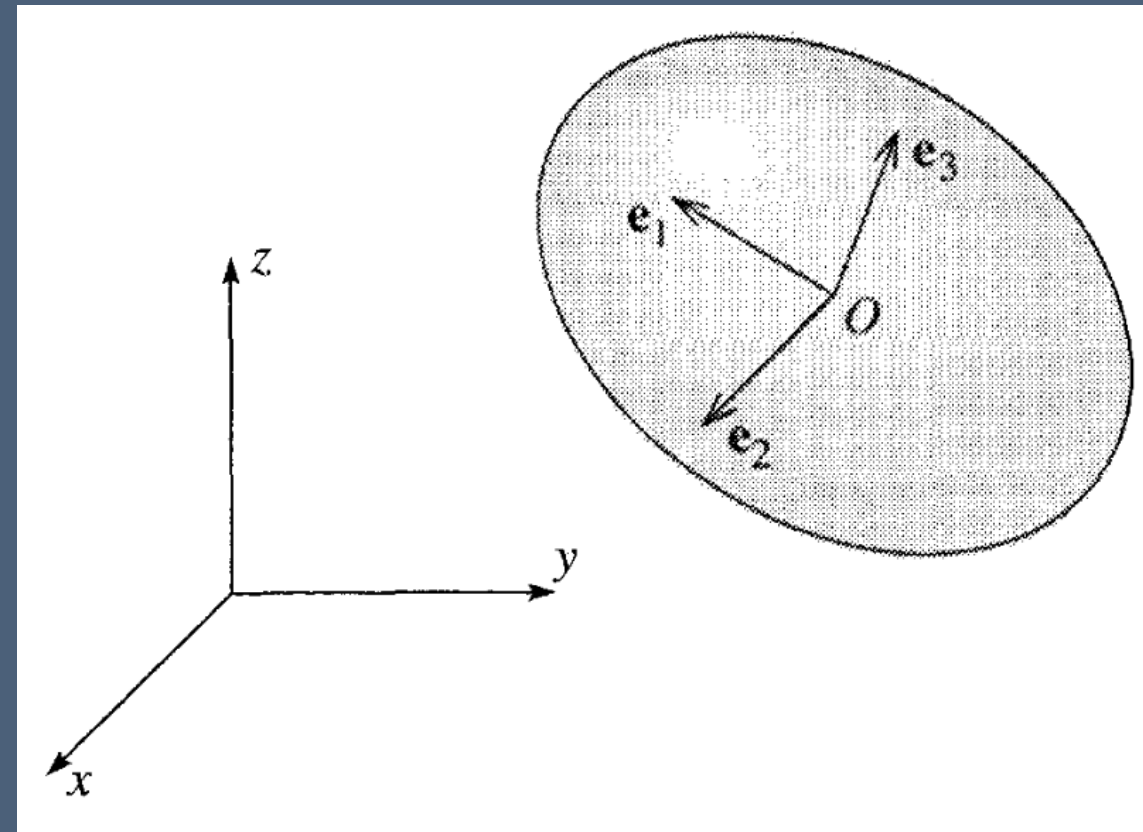


Intermediate Moment theorem!

AKA Tennis Racquet Thrm

$$\ddot{\omega}_1 = - \left[(\lambda_3 - \lambda_2) (\lambda_3 - \lambda_1) \frac{\omega_3^2}{\lambda_1 \lambda_2} \right] \omega_1$$

- *When is rotation stable?*
- *Unstable?*





Problem 10.42

Wobbly book frisbee!

- How wobbly is the spinning of a book about its shortest principal axis?

- Ex. 10.2 will be helpful, along with $\ddot{\omega}_1 = - \left[\frac{(\lambda_3 - \lambda_2)(\lambda_3 - \lambda_1)}{\lambda_1 \lambda_2} \omega_3^2 \right] \omega_1$

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