

Problem 9.1

a) We know the comoving wave number at horizon from eq 9.63 will be:

$$k_{eq} = a_{eq} H = H_0 \sqrt{2 \frac{\rho_m}{\rho_r}} = 0.02 h \text{ Mpc}^{-1}$$

and $\frac{1}{k_{eq}}$ is the actual comoving horizon.

With equation 9.32 we know $C_S = \frac{c}{\sqrt{3}}$

so really we only have to multiply $a_{eq} H$ by a factor of $\frac{1}{\sqrt{3}}$

$$\frac{1}{k_{eq,cs}} = \frac{1}{\sqrt{3} a_{eq} H} = \frac{50}{\sqrt{3}} \text{ Mpc } h^{-1} = 289 \text{ Mpc } h^{-1}$$

b) So we know the linear growth function:

$$D(a) = \frac{\delta(a)}{\delta(1)} = \frac{a g(a)}{g(1)}$$

We know that $g(a) = 1$ for an Einstein-de Sitter Universe, so we

can say $g(a) = 1$ and $g(1) = 1$:

$$D(a) = a = \frac{\delta(a)}{\delta(1)} \quad a = \frac{1}{z} \quad z = 2500 \text{ at matter radiation equality}$$

$$\frac{1}{2500} = \frac{\delta(a)}{2 \times 10^{-5}} \Rightarrow \delta(a) = \frac{2 \times 10^{-5}}{2500} = \boxed{5.71 \times 10^{-9}}$$

Problem 9.4

$$\frac{dn}{d \ln M} d \ln M = \frac{\rho_{M,0}}{M} \left| \frac{dF(M)}{d \ln M} \right| d \ln M. \quad (9.90)$$

After taking the derivative of Eq. (9.87) analytically, and including the miraculous factor of 2, we get (see Problem 9.4) the Press-Schechter mass function:

$$\frac{dn}{d \ln M} = \sqrt{\frac{2}{\pi}} \frac{\rho_{M,0}}{M} \frac{\delta_c}{\sigma} \left| \frac{d \ln \sigma}{d \ln M} \right| e^{-\delta_c^2 / (2\sigma^2)}. \quad (9.91)$$

$$F(M) \equiv \int_{\delta_c}^{\infty} P(\delta) d\delta = \frac{1}{\sqrt{2\pi}\sigma(M)} \int_{\delta_c}^{\infty} e^{-\delta^2 / (2\sigma(M)^2)} d\delta \quad (9.87)$$

$$\equiv \frac{1}{2} \text{erfc} \left(\frac{\nu_c}{\sqrt{2}} \right),$$

where erfc is the complementary error function (see e.g. Wikipedia) and

$$\nu_c \equiv \frac{\delta_c}{\sigma(M)} \quad (9.88)$$

$$\frac{d}{dz} \text{erf } z = \frac{2}{\sqrt{\pi}} e^{-z^2}.$$

$$F(M) = \frac{1}{2} \operatorname{erfc}\left(\frac{V_c}{\sqrt{2}}\right) = \frac{1}{2} \operatorname{erfc}\left(\frac{s_c}{\sqrt{2}\sigma(M)}\right)$$

$$\frac{dF(M)}{d\sigma(M)} = \frac{1}{2} \frac{2}{\sqrt{\pi}} e^{-\frac{s_c^2}{2\sigma^2}} \left(-\frac{s_c}{\sqrt{2}\sigma^2(M)}\right)$$

$$d \ln a = \frac{1}{a} da \Rightarrow d \ln \sigma = \sigma(M) d \ln \sigma$$

$$\frac{dF(M)}{\sigma d \ln \sigma} = \frac{1}{\sqrt{\pi}} e^{-\frac{s_c^2}{2\sigma^2}} \left(-\frac{s_c}{\sqrt{2}\sigma^2}\right)$$

$$dF(M) = -\frac{1}{\sqrt{2\pi}} \frac{s_c}{\sigma} e^{-\frac{s_c^2}{2\sigma^2}} d \ln \sigma$$

Now taking 9.90

$$\frac{dn}{d \ln M} d \ln M = \frac{P_{M,0}}{M} \left| \frac{dF(M)}{d \ln M} \right| d \ln M$$

and substituting $dF(M)$

$$\frac{dn}{d \ln M} = \frac{P_{M,0}}{M} \left| \frac{-\frac{1}{\sqrt{2\pi}} \frac{s_c}{\sigma} e^{-\frac{s_c^2}{2\sigma^2}} d \ln \sigma}{d \ln M} \right|$$

$$\frac{dn}{d \ln M} = \frac{P_{M,0}}{M} \frac{1}{\sqrt{2\pi}} \frac{s_c}{\sigma} \left| \frac{d \ln \sigma}{d \ln M} \right| e^{-\frac{s_c^2}{2\sigma^2}}$$

and introducing a magic factor of two like the book suggests:

$$\frac{dn}{d \ln M} = \frac{P_{M,0}}{M} \frac{2}{\sqrt{\pi}} \frac{s_c}{\sigma} \left| \frac{d \ln \sigma}{d \ln M} \right| e^{-\frac{s_c^2}{2\sigma^2}}$$

Additional Problem #2

looking at 4.3.54 and 4.3.55 in the Beuermann book

$$\partial_n(s\rho) = -3\mathcal{H}(s\rho + \rho\rho) + 3\rho'(p+p) - \partial_i q^i \quad \left. \begin{array}{l} \text{continuity equation} \\ \text{(4.3.49)} \end{array} \right\}$$

Remember from lecture:

$$\partial_n(s\rho) = \partial_n(\bar{\rho} \cdot \overset{\text{overdots}}{\rho}) = \partial_n(\bar{\rho}) \cdot \rho + \bar{\rho} \partial_n(\rho) = \bar{\rho}' \cdot \rho + \bar{\rho} \rho'$$

$$\text{Remember } \bar{\rho}' = -3\mathcal{H}(\bar{\rho} + \bar{p}) \quad \text{for radiation } p_r = \frac{1}{3}\rho_r \Rightarrow \bar{\rho}' = -3\mathcal{H}\left(\frac{4}{3}\rho_r\right) = -4\mathcal{H}\rho_r$$

$$\partial_n(\delta\bar{\rho}) = -4H\bar{J}_r \cdot \delta + \bar{J}\delta'$$

$$\partial_n(\delta\bar{\rho}) = -3\mathcal{H}(\delta\rho + \delta p) + 3\delta'(\bar{\rho} + \bar{p}) - \partial_i q^i$$

$$-4H\bar{J}_r \cdot \delta + \bar{J}\delta' = -3\mathcal{H}(\delta\rho + \delta p) + 3\delta'(\bar{\rho} + \bar{p}) - \partial_i q^i \quad \xrightarrow{q^i = (\bar{\rho} + \bar{p})V^i = \frac{4}{3}\bar{J}V^i}$$

$$-4H\bar{J}_r \cdot \delta + \bar{J}\delta' = -3\mathcal{H}(\bar{\rho} \cdot \delta) - \mathcal{H}\delta\bar{p} + 4\delta'\bar{J} - \frac{4}{3}\bar{J}\partial_i V^i$$

$$\bar{J}\delta' = 4\delta'\bar{J} - \frac{4}{3}\bar{J}\nabla V$$

$$\delta' = 4\delta' - \frac{4}{3}\nabla V$$