

$$\phi(r, \theta, z, t) = h(z, \theta, t) - r \quad (1)$$

$$\frac{D\phi}{Dt} = 0 \quad (2)$$

$$\left(\frac{\partial\phi}{\partial r}\right) \frac{dr}{dt} + \left(\frac{\partial\phi}{\partial\theta}\right) \frac{d\theta}{dt} + \left(\frac{\partial\phi}{\partial z}\right) \frac{dz}{dt} + \left(\frac{\partial\phi}{\partial t}\right) = 0 \quad (3)$$

$$\left(\frac{\partial(h-r)}{\partial r}\right) \frac{dr}{dt} + \left(\frac{\partial(h-r)}{\partial\theta}\right) \frac{d\theta}{dt} + \left(\frac{\partial(h-r)}{\partial z}\right) \frac{dz}{dt} + \left(\frac{\partial(h-r)}{\partial t}\right) = 0 \quad (4)$$

$$\left(\frac{\partial h}{\partial r}\right) u_r - u_r + \frac{u_\theta}{r} \left(\frac{\partial h}{\partial\theta}\right) + u_z \left(\frac{\partial h}{\partial z}\right) + \left(\frac{\partial h}{\partial t}\right) - u_r = 0 \quad (5)$$

$$\left(\frac{\partial h}{\partial r}\right) u_r + \frac{u_\theta}{r} \left(\frac{\partial h}{\partial\theta}\right) + u_z \left(\frac{\partial h}{\partial z}\right) + \left(\frac{\partial h}{\partial t}\right) = 2u_r \quad (6)$$

rearranging,

$$\left(\frac{\partial h}{\partial t}\right) + u_r \left(\frac{\partial h}{\partial r}\right) + \frac{u_\theta}{r} \left(\frac{\partial h}{\partial\theta}\right) + u_z \left(\frac{\partial h}{\partial z}\right) = 2u_r \quad (7)$$

we note the following ,

$$h(z, \theta, t) = \bar{h}(z) + h'(z, \theta, t) \quad (8)$$

$$u_r(h, z, \theta, t) = \bar{U}_r(r) + u'_r(h, z, \theta, t) \quad (9)$$

$$u_\theta(h, z, \theta, t) = u'_\theta(h, z, \theta, t) \quad (10)$$

(since the base flow is axisymmetric $\bar{U}_\theta = 0$)

$$u_z(h, z, \theta, t) = \bar{U}_z(r) + u'_z(h, z, \theta, t) \quad (11)$$

$$\left(\frac{\partial h(z, \theta, t)}{\partial t}\right) + u_r(h, z, \theta, t) \left(\frac{\partial h(z, \theta, t)}{\partial r}\right) +$$

$$\frac{u_\theta(h, z, \theta, t)}{r} \left(\frac{\partial h(z, \theta, t)}{\partial\theta}\right) + u_z(h, z, \theta, t) \left(\frac{\partial h(z, \theta, t)}{\partial z}\right) = 2(\bar{U}_r(r) + u'_r(h, z, \theta, t)) \quad (12)$$

$$\begin{aligned}
& \left(\frac{\partial \bar{h}}{\partial t} \right) + \left(\frac{\partial h'}{\partial t} \right) + \bar{U}_r \left(\frac{\partial \bar{h}(z)}{\partial r} \right) + \bar{U}_r \left(\frac{\partial h'}{\partial r} \right) + u'_r \left(\frac{\partial \bar{h}(z)}{\partial r} \right) + u'_r \left(\frac{\partial h'}{\partial r} \right) + \\
& \frac{u'_\theta}{r} \left(\frac{\partial \bar{h}}{\partial \theta} \right) + \frac{u'_\theta}{r} \left(\frac{\partial h'}{\partial \theta} \right) + \bar{U}_z \left(\frac{\partial \bar{h}}{\partial z} \right) + \bar{U}_z \left(\frac{\partial h'}{\partial z} \right) + u'_z \left(\frac{\partial \bar{h}}{\partial z} \right) + u'_z \left(\frac{\partial h'}{\partial z} \right) \\
& = 2(\bar{U}_r(r) + u'_r(h, z, \theta, t))
\end{aligned}$$

since \bar{h} is a function of z only, $\frac{\partial \bar{h}(z)}{\partial r} = \frac{\partial \bar{h}(z)}{\partial \theta} = 0$ and $\frac{\partial h'(z, \theta, t)}{\partial r} = 0$. We have,

$$\begin{aligned}
& \left(\frac{\partial \bar{h}}{\partial t} \right) + \left(\frac{\partial h'}{\partial t} \right) + \frac{u'_\theta}{r} \left(\frac{\partial h'}{\partial \theta} \right) \\
& + \bar{U}_z \left(\frac{\partial \bar{h}}{\partial z} \right) + \bar{U}_z \left(\frac{\partial h'}{\partial z} \right) + u'_z \left(\frac{\partial \bar{h}}{\partial z} \right) + u'_z \left(\frac{\partial h'}{\partial z} \right) = 2(\bar{U}_r(r) + u'_r(h, z, \theta, t)) \quad (13)
\end{aligned}$$

Equation for the mean is given by

$$\left(\frac{\partial \bar{h}}{\partial t} \right) + \bar{U}_r \left(\frac{\partial \bar{h}(z)}{\partial r} \right) + \bar{U}_z \left(\frac{d\bar{h}}{dz} \right) = 2\bar{U}_r(r) \quad (14)$$

or,

$$\left(\frac{\partial \bar{h}}{\partial t} \right) + \bar{U}_z \left(\frac{d\bar{h}}{dz} \right) = 2\bar{U}_r(r) \quad (15)$$

Subtracting the mean from the total flow, we get the perturbed equation of h :

$$\left(\frac{\partial h'}{\partial t} \right) + \frac{u'_\theta}{r} \left(\frac{\partial h'}{\partial \theta} \right) + \bar{U}_z \left(\frac{\partial h'}{\partial z} \right) + u'_z \left(\frac{dh'}{dz} \right) + u'_z \left(\frac{\partial h'}{\partial z} \right) = 2u'_r(h, z, \theta, t) \quad (16)$$

assuming small perturbations, we neglect the product of small quantities to yield,

$$\left(\frac{\partial h'(z, \theta, t)}{\partial t} \right) + \bar{U}_z(r) \left(\frac{\partial h'(z, \theta, t)}{\partial z} \right) + \underbrace{u'_z(h, z, \theta, t)}_{(17)} \left(\frac{d\bar{h}(z)}{dz} \right) = 2 \underbrace{u'_r(h, z, \theta, t)}_{(17)}$$

The above equation is nonlinear in h . We know that $h = \bar{h}(z) + h'(z, \theta, t)$

To linearize eq(17) equation, we need to expand the terms in the curly braces using Taylor series about \bar{h} .

$$u'_r(h, z, \theta, t) = u'_r(\bar{h}(z), z, \theta, t) + (h - \bar{h}) \left(\frac{\partial u'_r}{\partial h} \right)_{(\bar{h}, z, \theta, t)}$$

$$+(z - \bar{z}) \left(\frac{\partial u'_z}{\partial h} \right)_{(\bar{h}, z, \theta, t)} + (\theta - \bar{\theta}) \left(\frac{\partial u'_\theta}{\partial h} \right)_{(\bar{h}, z, \theta, t)} + \{O(h - \bar{h})^2\} + \dots \quad (18)$$

Neglecting higher order terms,

$$u'_r(h, z, \theta, t) = u'_r(\bar{h}(z), z, \theta, t) + (h - \bar{h}) \left(\frac{\partial u'_r}{\partial h} \right)_{(\bar{h}, z, \theta, t)} \quad (19)$$

but $h - \bar{h} = h'$ which is a function of z, θ, t .
i.e.,

$$u'_r(h, z, \theta, t) = u'_r(\bar{h}(z), z, \theta, t) + h'(z, \theta, t) \left(\frac{\partial u'_r}{\partial h} \right)_{(\bar{h}, z, \theta, t)} \quad (20)$$

Similarly,

$$u'_z(h, z, \theta, t) = u'_z(\bar{h}(z), z, \theta, t) + h'(z, \theta, t) \left(\frac{\partial u'_z}{\partial h} \right)_{(\bar{h}, z, \theta, t)} \quad (21)$$

substituting eq (20) and eq(21) into eq(17)

$$\begin{aligned} & \left(\frac{\partial h'(z, \theta, t)}{\partial t} \right) + \bar{U}_z(r) \left(\frac{\partial h'(z, \theta, t)}{\partial z} \right) + \\ & \left\{ u'_z(\bar{h}(z), z, \theta, t) + h'(z, \theta, t) \left(\frac{\partial u'_z}{\partial h} \right)_{(\bar{h}, z, \theta, t)} \right\} \left(\frac{d\bar{h}(z)}{dz} \right) = 2 \left\{ u'_r(\bar{h}(z), z, \theta, t) + h'(z, \theta, t) \left(\frac{\partial u'_r}{\partial h} \right)_{(\bar{h}, z, \theta, t)} \right\} \\ & \frac{\partial h'}{\partial t} + \bar{U}_z(r) \left(\frac{\partial h'}{\partial z} \right) + u'_z \frac{d\bar{h}(z)}{dz} + h' \left(\frac{\partial u'_z}{\partial h} \right)_{(\bar{h}, z, \theta, t)} \frac{d\bar{h}(z)}{dz} = 2u'_r + 2h' \left(\frac{\partial u'_r}{\partial h} \right)_{(\bar{h}, z, \theta, t)} \quad (22) \\ & \frac{\partial h'}{\partial t} + \bar{U}_z(r) \left(\frac{\partial h'}{\partial z} \right) + u'_z \frac{d\bar{h}(z)}{dz} + \underbrace{h' \left(\frac{\partial u'_z}{\partial h} \right)_{(\bar{h}, z, \theta, t)}}_{(23)} \left(\frac{d\bar{h}(z)}{dz} \right) = \underbrace{2u'_r + 2h' \left(\frac{\partial u'_r}{\partial h} \right)_{(\bar{h}, z, \theta, t)}}_{(23)} \end{aligned}$$

We can further linearize the above equation assuming small perturbation, so that the terms in the curly brackets are ignored.

Note: consider the term $\underbrace{h' \left(\frac{\partial u'_z}{\partial h} \right)_{(\bar{h}, z, \theta, t)}}_{(23)} \left(\frac{d\bar{h}(z)}{dz} \right)$. Even though the term in the curly bracket is small, $\frac{d\bar{h}(z)}{dz}$ need not be,

so that their product is finite. But we assume that $\frac{d\bar{h}(z)}{dz}$ is bounded, hence the entire term can be neglected.

$$\frac{\partial h'(z, \theta, t)}{\partial t} + \bar{U}_z(r) \left(\frac{\partial h'(z, \theta, t)}{\partial z} \right) + u'_z(\bar{h}(z), z, \theta, t) \left(\frac{d\bar{h}(z)}{dz} \right) = 2u'_r(\bar{h}(z), z, \theta, t) \quad (24)$$

The above equation is linear .