

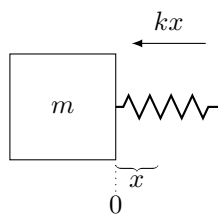
Harmonic Oscillator

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1 Simple Harmonic Oscillator

The simple harmonic oscillator comes about, most commonly, from a mass on spring. No friction.



Gives us the differential equation

$$m\ddot{x} - kx = 0$$

We solve this equation by solving the characteristic equation

$$r^2 - \frac{k}{m} = 0$$
$$r = \pm \sqrt{-\frac{k}{m}}$$

Since both k and m are positive

$$r = \pm i\sqrt{\frac{k}{m}}$$

and the general solution is

$$x = A \cos\left(\sqrt{\frac{k}{m}}t\right) + B \sin\left(\sqrt{\frac{k}{m}}t\right)$$

We let

$$\omega_0 = \sqrt{\frac{k}{m}}$$

This is both a useful simplification and a meaningful label, the *angular frequency* of the oscillation. Therefore we have

$$x = A \cos(\omega_0 t) + B \sin(\omega_0 t)$$

or

$$x = A \cos(\omega_0 t + \phi)$$

Both of these forms are useful, I will work with both.

1.1 Initial Conditions

We will solve for A and B when given initial conditions

$$x(0) = x_0, \quad \dot{x}(0) = \dot{x}_0$$

1.1.1 Form 1

$$\begin{aligned} x &= A \cos(\omega_0 t) + B \sin(\omega_0 t) \\ \dot{x} &= -A\omega_0 \sin(\omega_0 t) + B\omega_0 \cos(\omega_0 t) \end{aligned}$$

the initial conditions give

$$\begin{aligned} x(0) &= A = x_0 \\ \dot{x}(0) &= B\omega_0 = \dot{x}_0 \end{aligned}$$

So the particular solution is

$$x = x_0 \cos(\omega_0 t) + \frac{\dot{x}_0}{\omega_0} \sin(\omega_0 t)$$

1.1.2 Form 2

$$\begin{aligned} x &= A \cos(\omega_0 t + \phi) \\ \dot{x} &= -A\omega_0 \sin(\omega_0 t + \phi) \end{aligned}$$

the initial conditions give

$$\begin{aligned} x(0) &= A \cos \phi = x_0 \\ \dot{x}(0) &= -A\omega_0 \sin \phi = \dot{x}_0 \end{aligned}$$

We can get that

$$\cos \phi = \frac{x_0}{A}, \quad \sin \phi = -\frac{\dot{x}_0}{A\omega_0}$$

So

$$\begin{aligned}\cos^2 \phi + \sin^2 \phi = 1 &= \frac{x_0^2}{A^2} + \frac{\dot{x}_0^2}{A^2 \omega_0^2} \\ 1 &= \frac{1}{A^2} \left(x_0^2 + \frac{\dot{x}_0^2}{\omega_0^2} \right) \\ A^2 &= x_0^2 + \frac{\dot{x}_0^2}{\omega_0^2} \\ A &= \sqrt{x_0^2 + \dot{x}_0^2 / \omega_0^2}\end{aligned}$$

and

$$\tan \phi = -\frac{\dot{x}_0}{A \omega_0} \frac{A}{x_0} = -\frac{\dot{x}_0}{\omega_0 x_0}$$

So the particular solution is

$$x = \cos\left(\omega_0 t - \frac{\dot{x}_0}{\omega_0 x_0}\right) \sqrt{x_0^2 + \frac{\dot{x}_0^2}{\omega_0^2}}$$

This looks way worse, but it gives the actual amplitude of the wave, and the *actual phase*, whatever that's useful for. So it's physically more useful to look at, as an engineer or something, but it's harder to do math with.

If you let A' and B' be the constants for **Form 1** of the solution,

$$A' = x_0, \quad B' = \frac{\dot{x}_0}{\omega_0}$$

And then you can relate the two forms as

$$A^2 = (A')^2 + (B')^2, \quad \tan \phi = \frac{B'}{A'}$$

Which is an interesting relation.

1.1.3 Initial Conditions Sidenote

Say our initial conditions don't start at $t = 0$, say we have initial condition $x(t_0) = x_0$, we let $t = t' - t_0$ so that

$$\begin{aligned}x &= A \cos(\omega_0 t + \phi) && \text{becomes} \\ x &= A \cos(\omega_0 (t' - t_0) + \phi)\end{aligned}$$

Plugging in the initial condition will still give us the same values for constants A and ϕ . The fact that the initial conditions don't start at zero just doesn't mean the solution changes, the same things happen, just at a different time. So if need to know what happens at a later time, we just shift that time as well. If we wanted to know that happens at $t = t_1$, then we just plug in $t_1 + t_0$.

1.2 Energy