

Effective approximations of multiscale PDE based on limited information

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Joint work with Claude Le Bris & Frédéric Legoll

Multiscale Seminar, Laboratoire Navier

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General introduction

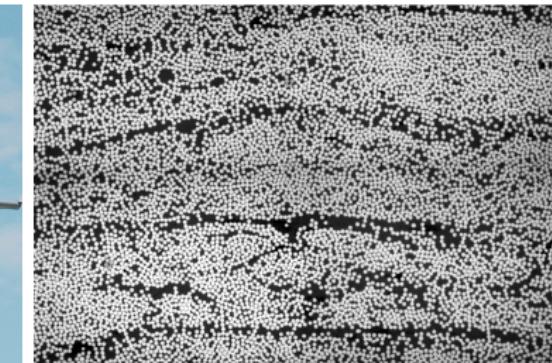
Multiscale systems

Multiscale system are characterized by the presence of multiple scales of interests that interact or influence one another.

They may be found in various scientific areas: engineering, medicine, physics, ...



airplane wing $\approx 10\text{m}$



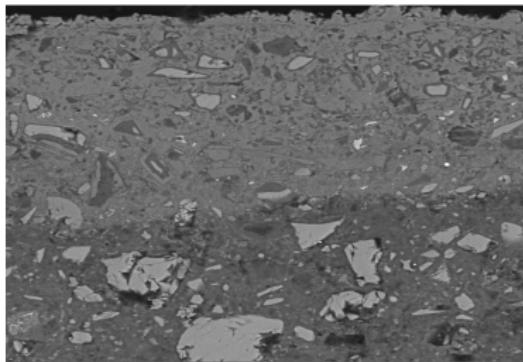
v.s. carbon fibers $\approx 10^{-6}\text{m}$

Figure: Composite material used in the aeronautics industry.

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bridge $\approx 10^3$ m v.s. mineral aggregate $\approx 10^{-5}$ m

Figure: Concrete: a multiscale material.

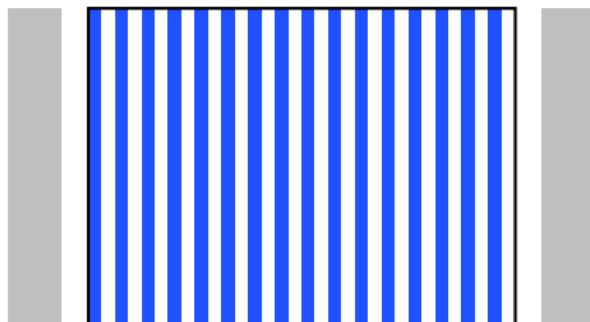
Inverse multiscale problems

- Consider the problem

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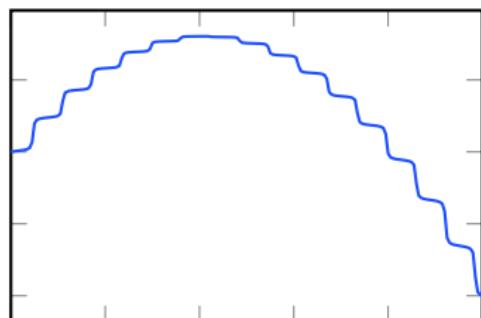
where A_ε is oscillating at a small length scale $\varepsilon \ll |\Omega|$.

- Applications: heat transfer in thermal engineering, (simplification of) elastic problem in mechanics, ...



$$\xrightarrow{\varepsilon} \quad \xleftarrow{L}$$

■ $A_\varepsilon = 0.023 \text{ W.m}^{-1}.\text{K}^{-1}$
□ $A_\varepsilon = 380 \text{ W.m}^{-1}.\text{K}^{-1}$



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Question

Based on measurements about the system, can A_ε be reliably reconstructed ? If not, what coefficients may be ?

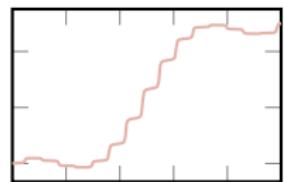
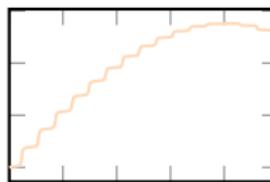
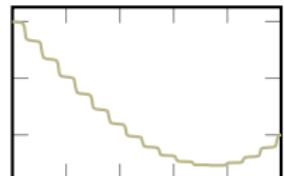
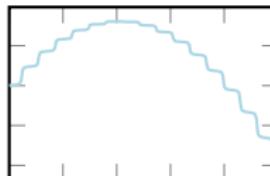
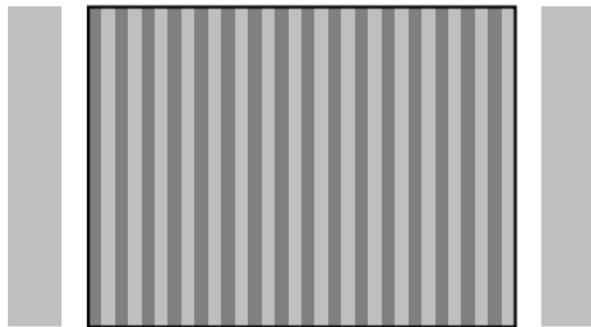
Limited information

Experimental settings:

- few knowledge on microstructure,
- only couples (*configuration, system responses*) are available.

Settings with limited information:

- No assumptions on microstructure (non periodic case, ε small but not infinitely small, ...),
- Coarse measurements,
- Noisy measurements,
- Quantitative restrictions (limited budget of measurements).



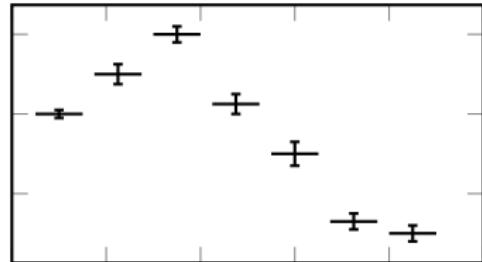
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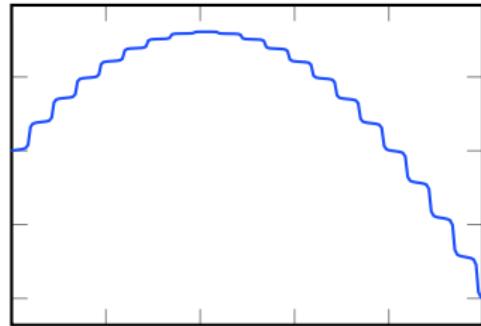
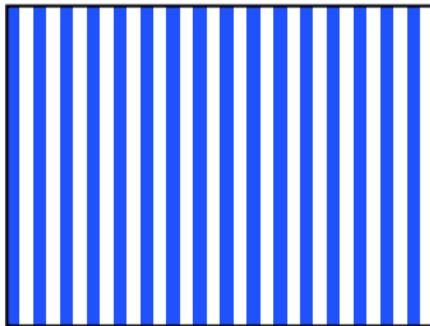
Homogenization (see e.g. [BLP78]¹) builds PDE with slowly varying coefficient that accurately approximate the oscillating PDE.

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Oscillating System

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Homogenized System



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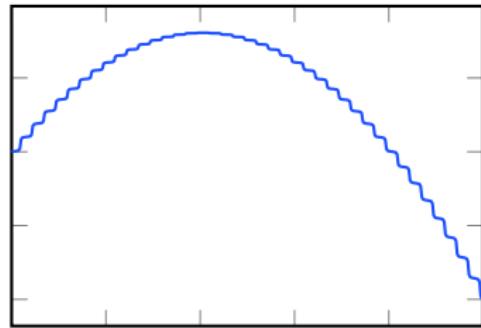
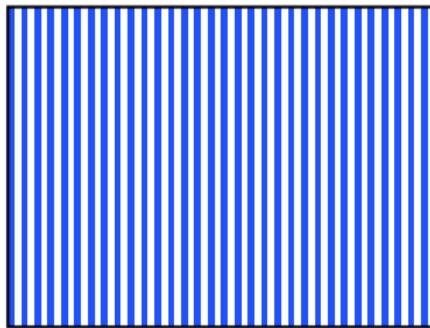
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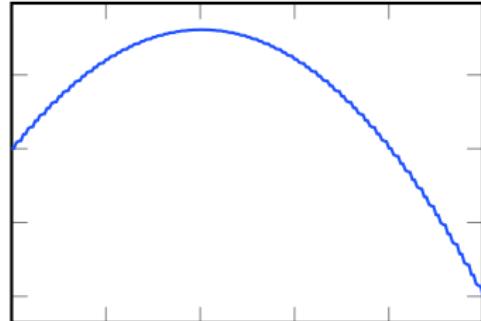
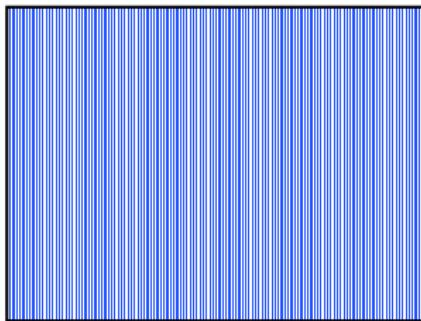
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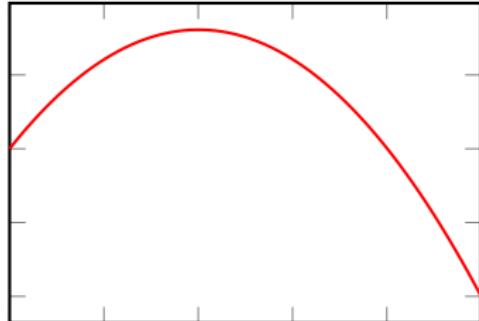
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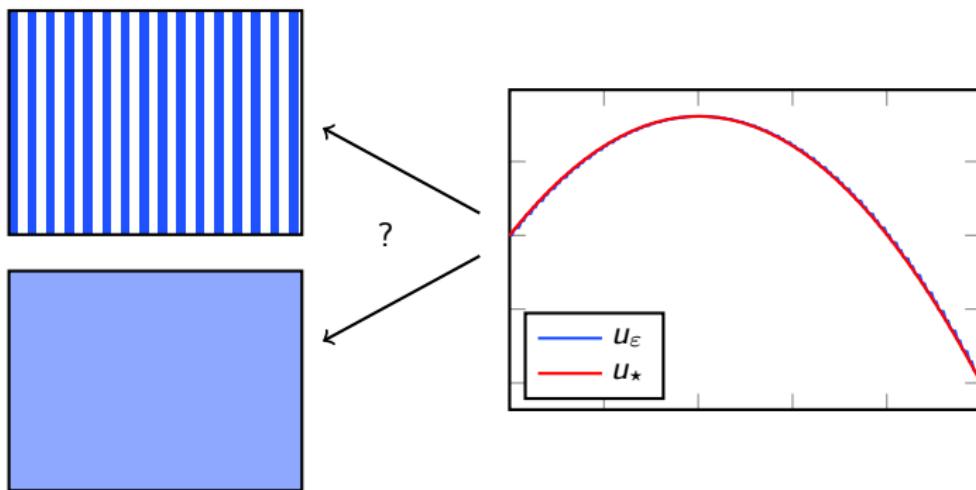
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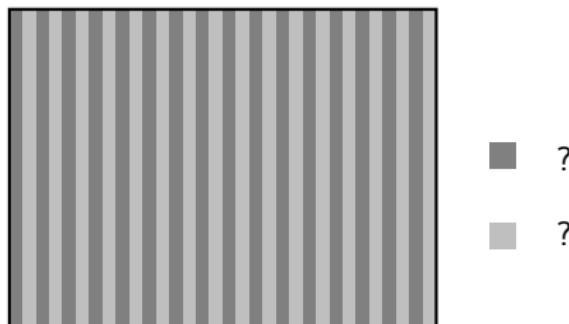
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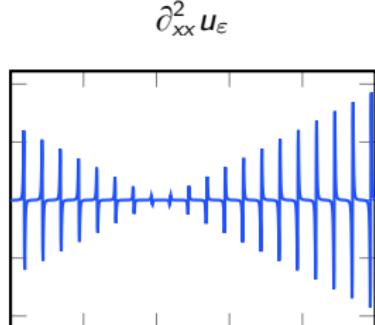
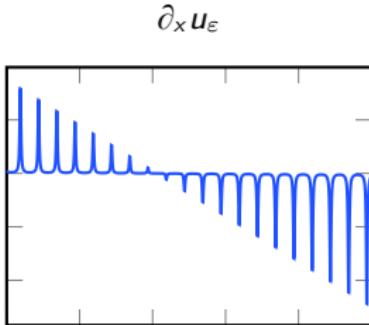
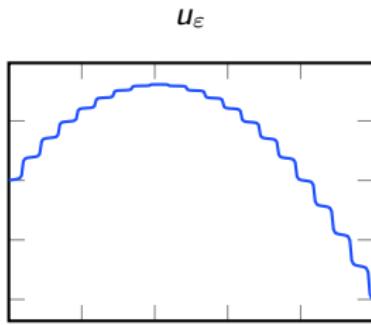
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Issue: how to proceed in contexts of limited information?

Effective coefficients

Effective coefficients are coefficients varying at the macroscopic scale that encapsulate the fine-scale features of highly oscillatory coefficients.

Example. Homogenization assesses the existence of an *effective coefficient A_** such that

$$\mathcal{L}_\varepsilon : f \longrightarrow u_\varepsilon(f) \quad \text{sol. to } -\operatorname{div}(A_\varepsilon \nabla u_\varepsilon) = f$$

$$\downarrow_{\circlearrowleft}^{\circlearrowright}$$

$$\mathcal{L}_{A_*} : f \longrightarrow u_*(f) \quad \text{sol. to } -\operatorname{div}(A_* \nabla u_*) = f$$

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with respect to the $L^2(\Omega)$ norm.

In the periodic case $A_\varepsilon(x) = A_{\text{per}}\left(\frac{x}{\varepsilon}\right)$, with A_{per} Q -periodic:

- in dimension $d = 1 \rightsquigarrow A_\star = \left(\int_Q \frac{1}{A_{\text{per}}} \right)^{-1}$,
- in dimension $d \geq 2 \rightsquigarrow A_\star = \int_Q A_{\text{per}}(\nabla \omega + \text{Id})$ where ω is a corrector defined through a PDE involving A_{per} .

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with respect to the $L^2(\Omega)$ norm.

Two major limitations of homogenization:

- No formula for A_* unless **strong assumptions** on A_ε (e.g. periodicity $A_\varepsilon(x) = A_{\operatorname{per}}(\frac{x}{\varepsilon})$).
- Valid only in the regime of separated scale (i.e. $\varepsilon \rightarrow 0$).

Objective

Based on available observables, define an effective coefficient \bar{A} such that, for any f , the solution $u_\varepsilon(f)$ to

$$-\operatorname{div}(A_\varepsilon \nabla u_\varepsilon) = f$$

are satisfactorily approximated by the solution $\bar{u} = u(\bar{A}, f)$ to the coarse problem

$$-\operatorname{div}(\bar{A} \nabla \bar{u}) = f.$$

Part I • Construct \bar{A} in the whole set $\mathbb{R}_{\text{sym}}^{d \times d}$.

Part II • Identify \bar{A} in the vicinity of a known coefficient \bar{A}_0 .

Part I

Effective modeling from boundary aggregated measurements

A first formulation [CRAS2013]², [COCV2018]³

For any $g \in L^2(\partial\Omega)$, consider the solution $u_\varepsilon = u_\varepsilon(g)$ to

$$-\operatorname{div}(A_\varepsilon \nabla u_\varepsilon) = 0 \text{ in } \Omega, \quad (A_\varepsilon \nabla u_\varepsilon) \cdot n = g \text{ on } \partial\Omega. \quad (1)$$

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For $\bar{A} \in \mathbb{R}_{\text{sym}}^{d \times d}$ a *constant* symmetric coefficient, consider $\bar{u} = u(\bar{A}, g)$ the solution to

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The quality of the effective coefficient \bar{A} can be quantified through the functional

$$\sup_{\|g\|_{L^2(\partial\Omega)}=1} \|u_\varepsilon(g) - u(\bar{A}, g)\|_{L^2(\Omega)}.$$

The strategy consists in **minimizing** the **worst case scenario** by looking at the optimization problem

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Issue : Using the **full solutions** u_ε **in the whole domain** Ω as observables is **disproportionate** to estimate a $d \times d$ constant symmetric matrix, and **irrealistic** from an experimental point of view.

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Practical observables

Only *coarser* observables are usually acquirable, such as the energy

$$\mathcal{E}(A_\varepsilon, g) = \frac{1}{2} \int_{\Omega} A_\varepsilon \nabla u_\varepsilon \cdot \nabla u_\varepsilon - \int_{\partial\Omega} g u_\varepsilon(g) = -\frac{1}{2} \int_{\partial\Omega} g u_\varepsilon(g). \quad (3)$$

Motivation:

- $\mathcal{E}(A_\varepsilon, g)$ passes to the **homogenized limit**:

$$\mathcal{E}(A_\varepsilon, g) \xrightarrow{\varepsilon \rightarrow 0} \mathcal{E}(A_*, g) \text{ in } \mathbb{R},$$

where $\mathcal{E}(A_*, g) = \frac{1}{2} \int_{\Omega} A_* \nabla u_* \cdot \nabla u_* - \int_{\partial\Omega} g u_*$ and where u_* denotes the homogenized solution.

- $\mathcal{E}(A_\varepsilon, g)$ is an **integrated quantity** (scalar !), thus it provides no direct insights about the microscale.

A new formulation

For $\bar{A} \in \mathbb{R}_{\text{sym}}^{d \times d}$ a *constant* symmetric coefficient, denote $\bar{u} = u(\bar{A}, g)$ the solution to

$$-\operatorname{div}(\bar{A} \nabla \bar{u}) = 0 \text{ in } \Omega, \quad (\bar{A} \nabla \bar{u}) \cdot n = g \text{ on } \partial\Omega.$$

To assess the quality of the effective coefficient \bar{A} , we use the functional

$$\sup_{\|g\|_{L^2(\partial\Omega)}=1} \|u_\varepsilon(g) - u(\bar{A}, g)\|_{L^2(\Omega)}^2 \rightarrow \sup_{\|g\|_{L^2(\partial\Omega)}=1} |\mathcal{E}(A_\varepsilon, g) - \mathcal{E}(\bar{A}, g)|^2.$$

Our strategy consists in **minimizing** the **worst case scenario** by looking at the optimization problem

$$\inf_{\substack{\bar{A} \in \mathbb{R}_{\text{sym}}^{d \times d} \\ \alpha \leq \bar{A} \leq \beta}} \sup_{\|g\|_{L^2(\partial\Omega)}=1} |\mathcal{E}(A_\varepsilon, g) - \mathcal{E}(\bar{A}, g)|^2.$$

Theoretical analysis

In the limit of vanishing ε , the problem leads to the homogenized diffusion coefficient as shown by the following proposition.

$$I_\varepsilon = \inf_{\substack{\bar{A} \in \mathbb{R}_{\text{sym}}^{d \times d} \\ \alpha \leq \bar{A} \leq \beta}} \sup_{g \in L^2(\partial\Omega)} |\mathcal{E}(A_\varepsilon, g) - \mathcal{E}(\bar{A}, g)|^2$$
$$\|g\|_{L^2(\partial\Omega)} = 1$$

Proposition (Asymptotic consistency, periodic case)

For any sequence of quasi-minimizer $(\bar{A}_\varepsilon^\#)_{\varepsilon > 0}$, i.e. sequence such that

$$I_\varepsilon \leq J_\varepsilon(\bar{A}_\varepsilon^\#) \leq I_\varepsilon + \text{err}(\varepsilon), \quad (4)$$

the following convergence holds :

$$\lim_{\varepsilon \rightarrow 0} \bar{A}_\varepsilon^\# = A_\star. \quad (5)$$

Computational procedure

We apply a **an iterative algorithm** to solve

$$\inf_{\substack{\bar{A} \in \mathbb{R}_{\text{sym}}^{d \times d} \\ \alpha \leq \bar{A} \leq \beta}} \sup_{g \in L^2(\partial\Omega)} |\mathcal{E}(A_\varepsilon, g) - \mathcal{E}(\bar{A}, g)|^2.$$
$$\|g\|_{L^2(\partial\Omega)} = 1$$

Given an iterate \bar{A}^n ,

- ① Define g^n , the argsup to

$$\sup_{\substack{g \text{ s.t. } \|g\|_{L^2(\partial\Omega)} = 1}} \left(\mathcal{E}(A_\varepsilon, g) - \mathcal{E}(\bar{A}^n, g) \right)^2.$$

In practice, $\sup_{g \in L^2(\Omega)} \rightarrow \sup_{g \in V_P}$ on $V_P = \text{Span}\{P \text{ loadings}\}$, with $P \approx 3$.

This step requires computing P solutions to a coarse PDE in order to get the energy $\mathcal{E}(\bar{A}^n, \cdot)$.

We next solve a $P \times P$ eigenvalue problem.

- ② Define \bar{A}^{n+1} , the optimizer to

$$\inf_{\substack{\bar{A} \in \mathbb{R}_{\text{sym}}^{d \times d}}} \left(\mathcal{E}(A_\varepsilon, g^n) - \mathcal{E}(\bar{A}, g^n) \right)^2.$$

In practice, we perform a gradient descent.

The gradient can be expressed with solutions computed in previous step, hence no additionnal costs.

Choice of the loadings

We identify P appropriate loadings $(g_i)_{1 \leq i \leq P}$ such that

$$\sup_{g \in L^2(\partial\Omega)} |\mathcal{E}(A_\varepsilon, g) - \mathcal{E}(\bar{A}, g)| \approx \sup_{\substack{g \in \text{Span } (g_i) \\ 1 \leq i \leq P}} |\mathcal{E}(A_\varepsilon, g) - \mathcal{E}(\bar{A}, g)|.$$

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Rayleigh quotient: we optimize

$$\sup_{\|g\|_{L^2(\partial\Omega)}=1} |\mathcal{E}(A_\varepsilon, g) - \mathcal{E}(\bar{A}, g)| = \sup_{g \in L^2(\partial\Omega)} \left| \frac{\int_{\partial\Omega} g (\mathcal{T}_\varepsilon - \mathcal{T}_{\bar{A}}) g}{\int_{\partial\Omega} g^2} \right|,$$

where

$$\mathcal{T}_\varepsilon : g \in L^2(\partial\Omega) \longrightarrow \gamma(u_\varepsilon(g)) \quad \text{with } u_\varepsilon(g) \text{ sol. to (1),}$$

$$\mathcal{T}_{\bar{A}} : g \in L^2(\partial\Omega) \longrightarrow \gamma(u(\bar{A}, g)) \quad \text{with } u(\bar{A}, g) \text{ sol. to (2).}$$

Thus, we seek the eigenmode of $\mathcal{T}_\varepsilon - \mathcal{T}_{\bar{A}}$ with largest (unsigned) eigenvalue.

Choice of the loadings

We identify P appropriate loadings $(g_i)_{1 \leq i \leq P}$ such that

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with $\mathcal{T} : g \in L^2(\partial\Omega) \longrightarrow \gamma(w(g))$ where $w(g)$ is solution to

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Practical choice: We select the $P \gtrapprox P_d = \frac{d(d+1)}{2}$ first eigenmodes of \mathcal{T} .

Numerical results (periodic)

In 2D ($\Omega =]0, 1[^2$), we consider the coefficient

$$A_\varepsilon(x, y) = A^{\text{per}}\left(\frac{x}{\varepsilon}, \frac{y}{\varepsilon}\right) = \begin{pmatrix} 22 + 10 \times (\sin(2\pi \frac{x}{\varepsilon}) + \sin(2\pi \frac{y}{\varepsilon})) & 0 \\ 0 & 12 + 2 \times (\sin(2\pi \frac{x}{\varepsilon}) + \sin(2\pi \frac{y}{\varepsilon})) \end{pmatrix},$$

for which

$$A_* \approx \begin{pmatrix} 19.3378 & 0 \\ 0 & 11.8312 \end{pmatrix}.$$

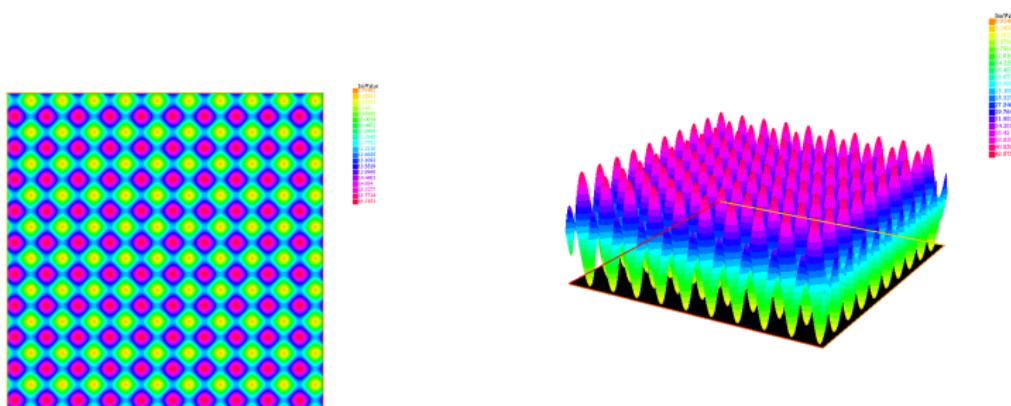


Figure: Components 11 and 22 of coefficient A_ε .

Numerical results (periodic)

$$\frac{|\bar{A} - A_*|_2}{|A_*|_2}$$

$$\text{Err}_{\varepsilon, Q}(\bar{A}) = \sup_{g \in \text{Span}(g_1, \dots, g_Q)} \left(\frac{\|u_\varepsilon(g) - u(\bar{A}, g)\|_{L^2(\Omega)}}{\|u_\varepsilon(g)\|_{L^2(\Omega)}} \right)$$

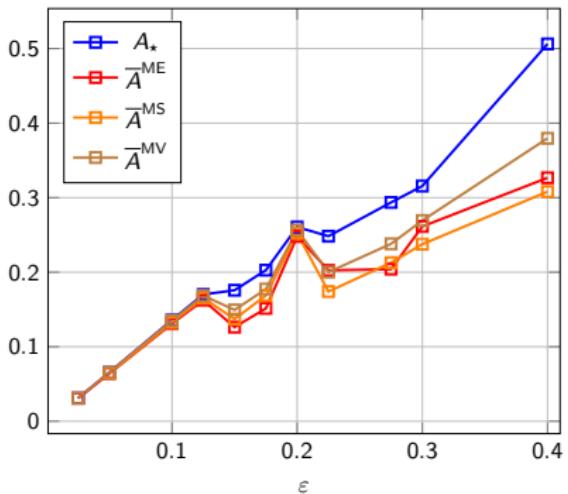
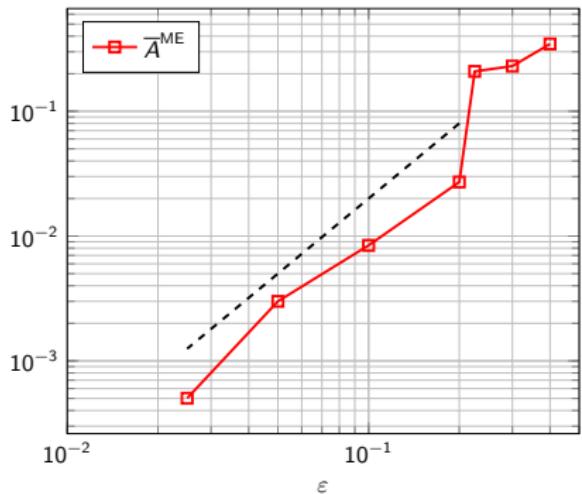


Figure: (left) Error between coefficients A_* and $\bar{A}_{\varepsilon, P}^{\text{ME}}$.

(right) Criterion $\text{Err}_{\varepsilon, Q}(\bar{A})$ for $\bar{A} \in \{A_*, \bar{A}_{\varepsilon, P}^{\text{MV}}, \bar{A}_{\varepsilon, P}^{\text{ME}}, \bar{A}_{\varepsilon, P}^{\text{MS}}\}$ (with $Q = 11$).

Numerical results (stochastic)

We now use a non periodic coefficient (random checkerboard),

$$A_\varepsilon(x, y, \omega) = a^{\text{sto}}\left(\frac{x}{\varepsilon}, \frac{y}{\varepsilon}, \omega\right) = \left(\sum_{k \in \mathbb{Z}^2} X_k(\omega) \mathbb{1}_{k+Q}(x, y)\right) \text{Id},$$

with X_k i.i.d random variables such that $\mathbb{P}(X_k = \gamma_1) = \mathbb{P}(X_k = \gamma_2) = \frac{1}{2}$ and $(\gamma_1, \gamma_2) = (4, 16)$.

We have

$$A_* = \sqrt{\gamma_1 \gamma_2} \text{ Id.}$$

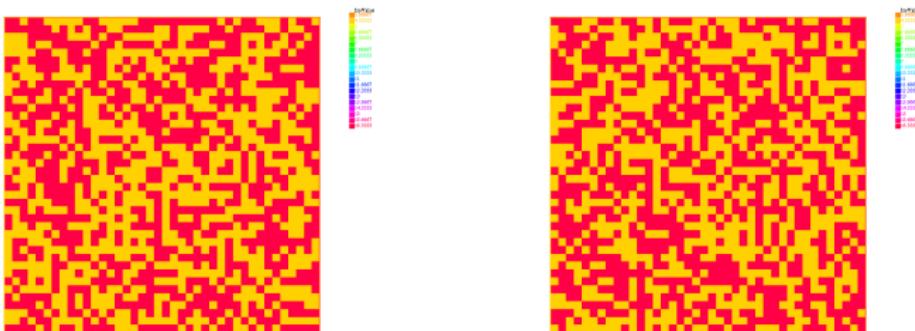
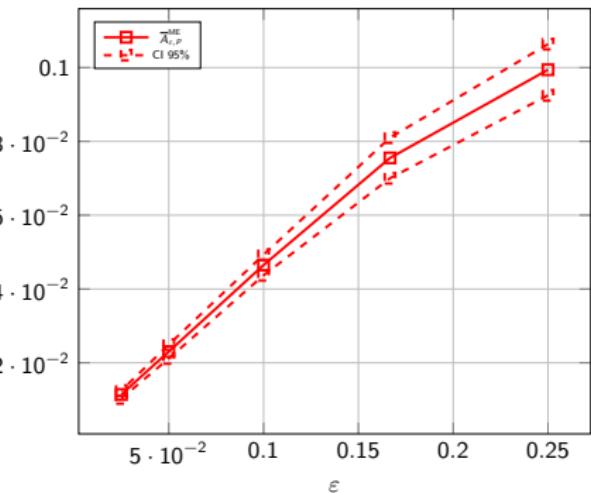


Figure: Two realizations of coefficient A_ε .

Our strategy rewrites $I_\varepsilon = \inf \sup |\mathbb{E}(\mathcal{E}(A_\varepsilon(\cdot, \omega), f)) - \mathcal{E}(\bar{A}, f)|$. Confidence intervals are computed from 40 realizations of the expectation (itself approximated by its empirical mean using 40 realizations of the coefficient a^{sto}).

Numerical results (stochastic)

$$\frac{|\bar{A} - A_\star|_2}{|A_\star|_2}$$



$$\text{Err}_{\varepsilon, Q}^{\mathbb{E}}(\bar{A}) = \sup_{g \in \text{Span}(g_1, \dots, g_Q)} \left(\frac{\|\mathbb{E}(u_\varepsilon(g)) - u(\bar{A}, g)\|_{L^2(\Omega)}}{\|\mathbb{E}(u_\varepsilon(g))\|_{L^2(\Omega)}} \right)$$

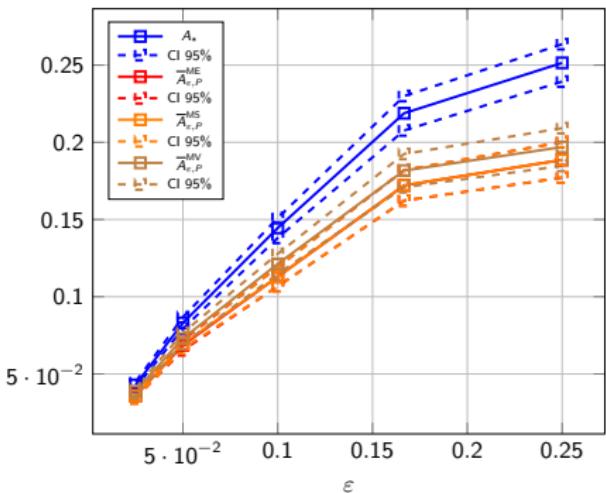


Figure: (left) Error between coefficients A_\star and $\bar{A}_{\varepsilon, P}^{\text{ME}}$.

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Noise

Consider a multiplicative Gaussian noise in the energy:

$$\mathcal{E}(A_\varepsilon, g; \sigma) = (1 + \sigma\eta) \mathcal{E}(A_\varepsilon, g).$$

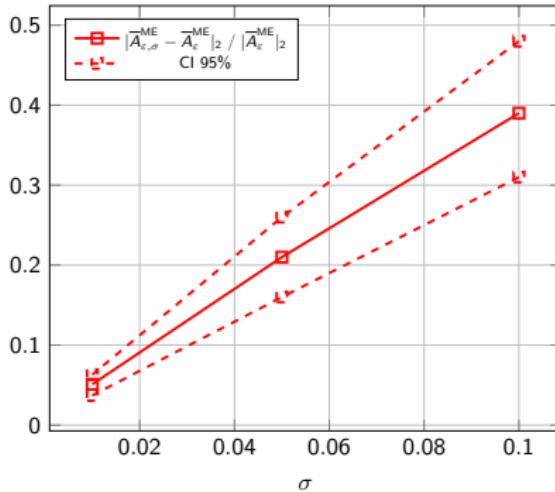


Figure: Error $\frac{|\bar{A}_{\varepsilon,\sigma}^{\text{ME}} - \bar{A}_\varepsilon^{\text{ME}}|_2}{|\bar{A}_\varepsilon^{\text{ME}}|_2}$ as a function of the noise magnitude σ (for $\varepsilon = 0.025$).

Part II

Perturbative reconstruction of effective coefficients

Perturbative reconstruction of effective coefficients

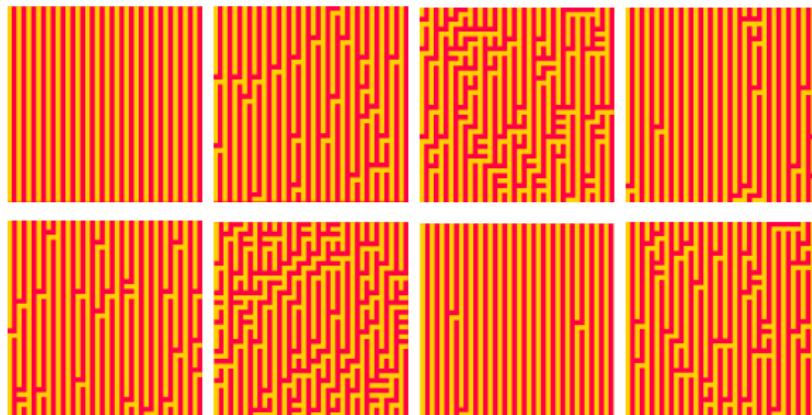
- **Assumption:** The effective coefficient lies in neighbourhood of a known coefficient \bar{A}_0 .
- **Example:** Randomly defectuous (periodic) material

$$A_{\varepsilon,\eta}(x, \omega) = A_\varepsilon^{\text{per}}(x) + b_\eta(\omega) C_\varepsilon^{\text{per}}(x),$$

with $C_\varepsilon^{\text{per}}$ possibly not negligible, but

$$A_{*,\eta} = \bar{A}_0 + \eta \bar{A}_1 + o(\eta),$$

where \bar{A}_0 is known (e.g. given as an industrial reference).



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where \bar{A}_0 is known (e.g. given as an industrial reference).

- **Issue:** Computing the effective coefficient using previous methods for many realizations ω and different defect rates η may end up having a **prohibitive computational costs...**

Question

How can we use the *a priori* knowledge of \bar{A}_0 to guide and speed up the optimization ?

Perturbative development

Consider the problem

$$-\operatorname{div}(A_\varepsilon \nabla u_\varepsilon) = f \text{ in } \Omega, \quad \text{and} \quad u_\varepsilon = 0 \text{ on } \partial\Omega,$$

and its approximation by

$$-\operatorname{div}(\bar{A} \nabla \bar{u}) = f \text{ in } \Omega, \quad \text{and} \quad \bar{u} = 0 \text{ on } \partial\Omega.$$

Exact	Perturbative expansion
\bar{A}	$\bar{A}_0 + \eta \bar{B}$

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\bar{A}	$\bar{A}_0 + \eta \bar{B}$
$u(\bar{A}, f)$	$u_0 + \eta v$

where $u_0 = u(\bar{A}_0, f)$ and $v = v(\bar{A}_0, \bar{B}, f)$ is solution to

$$\begin{cases} -\operatorname{div}(\bar{A}_0 \nabla v) = \operatorname{div}(\bar{B} \nabla u_0) & \text{in } \Omega, \\ v = 0 & \text{on } \partial\Omega. \end{cases}$$

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Linearity implies that $v = \sum_{ij} \bar{B}_{ij} v_{ij}(\bar{A}_0, f)$ with $v_{ij} = v_{ij}(\bar{A}_0, f)$ solution to

$$\begin{cases} -\operatorname{div}(\bar{A}_0 \nabla v_{ij}) = \operatorname{div}(E_{ij} \nabla u_0) & \text{in } \Omega, \\ v_{ij} = 0 & \text{on } \partial\Omega. \end{cases} \quad (6)$$

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$u(\bar{A}, f)$	$u_0 + \eta v$
$\mathcal{E}(\bar{A}, f)$	$\mathcal{E}(\bar{A}_0, f) + \eta \sum_{ij} \bar{B}_{ij} \mathcal{F}_{ij}(\bar{A}_0, f)$

where $\mathcal{E}(\bar{A}_0, f) = -\frac{1}{2} \int_\Omega f u_0$ and

$$\begin{aligned}\mathcal{F}_{ij}(\bar{A}_0, f) &= -\frac{1}{2} \int_\Omega f v_{ij} \\ &= \frac{1}{2} \int_\Omega \partial_i u_0 \partial_j u_0.\end{aligned}$$

Perturbative development

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$\mathcal{E}(\bar{A}, f)$	$\mathcal{E}(\bar{A}_0, f) + \eta \sum_{ij} [\bar{B}]_{ij} \mathcal{F}_{ij}(\bar{A}_0, f)$

We formulate the optimization problem

$$\inf_{\substack{\bar{B} \in \mathbb{R}_{\text{sym}}^{d \times d}, \\ f \in L^2(\Omega)}} \left(\mathcal{E}(A_{\varepsilon, \eta}, f) - \mathcal{E}(\bar{A}_0, f) - \sum_{1 \leq i, j \leq d} [\bar{B}]_{ij} \mathcal{F}_{ij}(\bar{A}_0, f) \right)^2.$$
$$\alpha \leq \bar{A}_0 + \bar{B} \leq \beta \|f\|_{L^2(\Omega)} = 1.$$

Implementation aspects

- **Offline stage:**

- Compute $u(\bar{A}_0, f)$.
- Compute $\mathcal{E}(\bar{A}_0, f)$ and $\mathcal{F}_{ij}(\bar{A}_0, f)$ for any $1 \leq i \leq j \leq d$.

↳ computing $P \approx 3$ solutions to a coarse PDE and $P(1 + \frac{d(d+1)}{2})$ domain integrals.

- **Online stage:** We apply a **gradient descent**. Given an iterate \bar{B}^n ,

- Define f^n , the argsup to

$$\sup_{f \text{ s.t. } \|f\|_{L^2(\Omega)} = 1} \left(\mathcal{E}(A_\varepsilon, f) - \mathcal{E}(\bar{A}_0, f) - \sum_{ij} [\bar{B}^n]_{ij} \mathcal{F}_{ij}(\bar{A}_0, f) \right)^2.$$

In practice, $\sup_{f \in L^2(\Omega)} \rightarrow \sup_{f \in V_P}$ on $V_P = \text{Span}\{P \text{ loadings}\}$, with $P \approx 3$.

This step amounts to solving a $P \times P$ eigenvalue problem.

- Define

$$\bar{B}^{n+1} = \bar{B}^n - \mu \nabla_{\bar{B}} J_\varepsilon^n(\bar{B}^n),$$

with

$$J_\varepsilon^n(\bar{B}) = \left(\mathcal{E}(A_\varepsilon, f^n) - \mathcal{E}(\bar{A}_0, f^n) - \sum_{ij} [\bar{B}]_{ij} \mathcal{F}_{ij}(\bar{A}_0, f^n) \right)^2.$$

↳ no additional computations of coarse PDE !

Numerical results

- Do not damage drastically the quality.
- Reduction of computational costs (by a factor of ≈ 80 to 400).

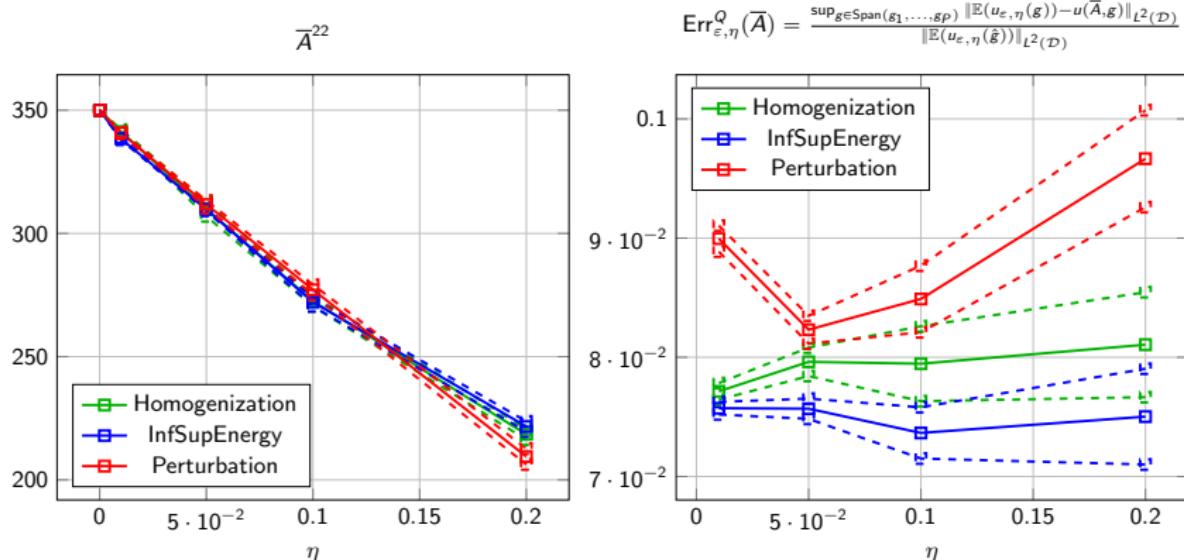


Figure: (left) Component 22 for various approximations of the effective coefficient.
(right) Criterion $\text{Err}_{\varepsilon, Q}^E(\bar{A})$ for different constant coefficients (with $Q = 9$ and $\varepsilon = 0.025$).

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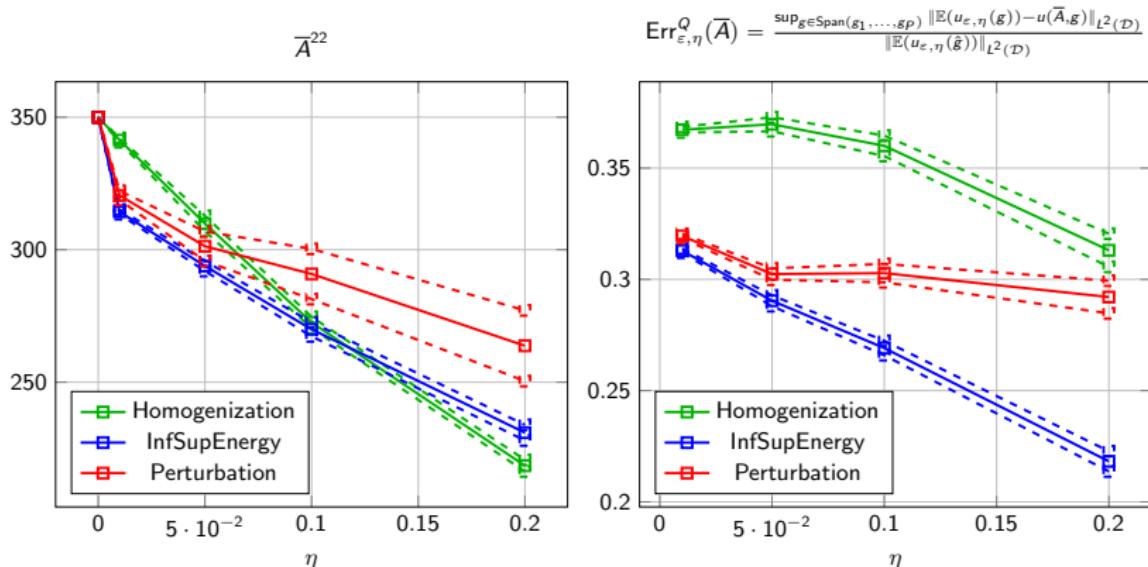


Figure: (left) Component 22 for various approximations of the effective coefficient.
(right) Criterion $\text{Err}_{\varepsilon,Q}^E(\bar{A})$ for different constant coefficients (with $Q = 9$ and $\varepsilon = 0.1$).

Conclusion and perspectives

Our strategies

- aim at determining effective approximation for multiscale PDEs through coarse coefficients,
- are designed for context with limited information is available,
- are inspired by homogenization theory and consistent with it (numerically and theoretically),
- can be extended outside the regime of separated scale.

Perspectives

- Application to real experimental data.
- Extension to non-constant effective coefficients.
- Convergence analysis of \bar{A} to A_* .

Thank you !

Sketch of proof

Three ingredients:

- Optimization over **compact set** $\mathcal{S}_{\alpha,\beta}$ $\implies \overline{A}_\varepsilon^\#$ converges to $\overline{A}_\#$ up to an extraction.

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- Optimization over **compact set** $\mathcal{S}_{\alpha,\beta} \implies \overline{A}_\varepsilon^\#$ converges to $\overline{A}_\#$ up to an extraction.
- **Homogenization** $\implies \mathcal{E}_\varepsilon(g) \xrightarrow{\varepsilon \rightarrow 0} \mathcal{E}_\star(g) \implies I_\varepsilon \xrightarrow{\varepsilon \rightarrow 0} 0 \implies \mathcal{E}_\varepsilon(g) \xrightarrow{\varepsilon \rightarrow 0} \mathcal{E}_\#(g).$

Polarization relation implies that for any $f, g \in L^2(\partial\Omega)$:

$$\int_{\partial\Omega} f \ u(A_\star, g) = \int_{\partial\Omega} f \ u(\overline{A}_\#, g).$$

Thus,

$$u(A_\star, g) = u(\overline{A}_\#, g) \text{ in } L^2(\partial\Omega).$$

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Thus,

$$u(A_\star, g) = u(\overline{A}_\#, g) \text{ in } L^2(\partial\Omega).$$

- Identification of particular couples $(f_i, g_i)_{1 \leq i \leq \frac{d(d+1)}{2}}$ to get

$$A_\star = A_\#.$$

A different perspective on noise

Motivation: anticipate on reproducibility errors during model deployment.

Idea: treat \bar{A} as a random field and optimize upon its mean.

Formulation: consider the problem

$$\inf_{\bar{A} \in \mathcal{S}_{\alpha, \beta}} \sup_{g \in L^2(\partial\Omega), \|g\|_{L^2(\partial\Omega)} = 1} \left| \mathcal{E}(A_\varepsilon, g) - \mathbb{E} (\mathcal{E}(\bar{A} + \sigma\eta, g)) \right|^2,$$

where η is a Gaussian variable.

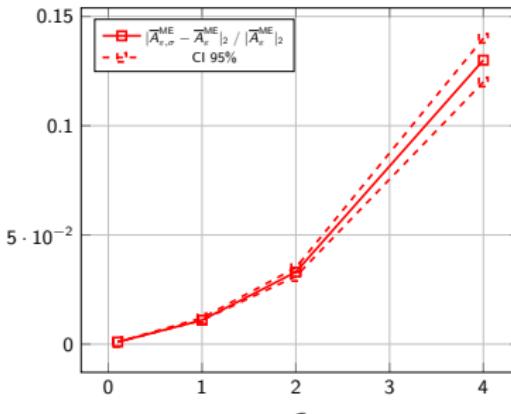


Figure: Error between $\bar{A}_{\varepsilon, \sigma}^{\text{ME}}$ and $\bar{A}_\varepsilon^{\text{ME}}$ as a function of the noise magnitude σ (for $\varepsilon = 0.05$).

Part III

Efficient selection of effective coefficients

Framework

- **Setting:** we are given

- a list of candidate coefficients $\mathcal{A} = \{\bar{A}_1, \dots, \bar{A}_N\}$.
- a list of admissible loadings $\mathcal{F} = \{f_1, \dots, f_P\}$.
- a measurement operator $\mathcal{O} : \mathcal{A} \times \mathcal{F} \rightarrow \mathbb{R}$ or $L^2(\Omega)$ (e.g. $\mathcal{O}(A_\varepsilon, f) = u_\varepsilon(f)$ or $\mathcal{E}(A_\varepsilon, f)$).

- **Challenge:**

- Expensive measurement costs \implies budget of $Q \ll P$ measurements.
- (Unknown) decomposition of \mathcal{F} into $\mathcal{F}_{\text{disc}}$ and $\mathcal{F}_{\text{non-disc}}$ such that

$$\text{card}(\mathcal{F}_{\text{disc}}) \ll \text{card}(\mathcal{F}),$$

and for any $f \in \mathcal{F}_{\text{non-disc}}$ and any $\bar{A}, \bar{B} \in \mathcal{A}^2$,

$$\|\mathcal{O}(A_\varepsilon, f) - \mathcal{O}(\bar{A}, f)\|_{\mathcal{O}} \approx \|\mathcal{O}(A_\varepsilon, f) - \mathcal{O}(\bar{B}, f)\|_{\mathcal{O}}.$$

Objective

Select the coefficient in \mathcal{A} that provides the best effective approximation of the underlying system while simultaneously minimizing the number of measurement operations $\mathcal{O}(A_\varepsilon, f)$ with $f \in \mathcal{F}$.

Iterative algorithm: each step k selects a loading f^k in \mathcal{F} and update the choice of the best coefficient \bar{A}^k in \mathcal{A} .

Core components:

- **Discrimination rate:** it estimates how a loadings is discriminative w.r.t observable \mathcal{O} .
- **Effectiveness score:** it assesses the quality of a coefficient as an effective coefficient for the system.

Strategy

Iterative algorithm: each step k selects a loading f^k in \mathcal{F} and update the choice of the best coefficient \bar{A}^k in \mathcal{A} .

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$\Delta_1(f)$	$\max_{(\bar{A}, \bar{B}) \in \mathcal{A}^2} \text{Err}(\bar{A}, \bar{B}, f)$	$\text{Err}(\bar{A}, \bar{B}, f) = \frac{\ \mathcal{O}(\bar{A}, f) - \mathcal{O}(\bar{B}, f)\ _{\mathcal{O}}}{\ \mathcal{O}(\bar{A}, f)\ _{\mathcal{O}}}$
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$\Delta_2^k(f)$	$\max_{(\bar{A}, \bar{B}) \in \mathcal{A}^2} \text{Err}^k(\bar{A}, f) - \text{Err}^k(\bar{B}, f) $	$\text{Err}^k(\bar{A}, f) = \frac{\ \mathcal{O}(\bar{A}^k, f) - \mathcal{O}(\bar{A}, f)\ _{\mathcal{O}}}{\ \mathcal{O}(\bar{A}^k, f)\ _{\mathcal{O}}}$
$\Delta_{\varepsilon}(f)$	$\max_{(\bar{A}, \bar{B}) \in \mathcal{A}^2} \text{Err}_{\varepsilon}(\bar{A}, f) - \text{Err}_{\varepsilon}(\bar{B}, f) $	$\text{Err}_{\varepsilon}(\bar{A}, f) = \frac{\ \mathcal{O}(\bar{A}_{\varepsilon}, f) - \mathcal{O}(\bar{A}, f)\ _{\mathcal{O}}}{\ \mathcal{O}(\bar{A}_{\varepsilon}, f)\ _{\mathcal{O}}}$

Strategy

Iterative algorithm: each step k selects a loading f^k in \mathcal{F} and update the choice of the best coefficient \bar{A}^k in \mathcal{A} .

Core components:

- **Discrimination rate:** it estimates how a loadings is discriminative w.r.t observable \mathcal{O} .

$\Delta_1(f)$	$\max_{(\bar{A}, \bar{B}) \in \mathcal{A}^2} \text{Err}(\bar{A}, \bar{B}, f)$	$\text{Err}(\bar{A}, \bar{B}, f) = \frac{\ \mathcal{O}(\bar{A}, f) - \mathcal{O}(\bar{B}, f)\ _{\mathcal{O}}}{\ \mathcal{O}(\bar{A}, f)\ _{\mathcal{O}}}$
$\Delta_2^k(f)$	$\max_{(\bar{A}, \bar{B}) \in \mathcal{A}^2} \text{Err}^k(\bar{A}, f) - \text{Err}^k(\bar{B}, f) $	$\text{Err}^k(\bar{A}, f) = \frac{\ \mathcal{O}(\bar{A}^k, f) - \mathcal{O}(\bar{A}, f)\ _{\mathcal{O}}}{\ \mathcal{O}(\bar{A}^k, f)\ _{\mathcal{O}}}$
$\Delta_3(f)$	$\max_{(\bar{A}, \bar{B}) \in \mathcal{A}^2} \text{Err}_\varepsilon(\bar{A}, f) - \text{Err}_\varepsilon(\bar{B}, f) $	$\text{Err}_\varepsilon(\bar{A}, f) = \frac{\ \mathcal{O}(A_\varepsilon, f) - \mathcal{O}(\bar{A}, f)\ _{\mathcal{O}}}{\ \mathcal{O}(A_\varepsilon, f)\ _{\mathcal{O}}}$

- **Effectiveness score:** it assesses the quality of a coefficient as an effective coefficient for the system.

↳ e.g.

$$\gamma^k(\bar{A}) = \max_{f \in \{f_1, \dots, f_k\}} \text{Err}_\varepsilon(\bar{A}, f).$$

Strategy

Selection algorithm

Initialization:

Select $f^1 \in \mathcal{F}$ by solving

$$f^1 = \arg \max_{f \in \mathcal{F}} \max_{(\bar{A}, \bar{B}) \in \mathcal{A}^2} \frac{\|\mathcal{O}(\bar{A}, f) - \mathcal{O}(\bar{B}, f)\|_{\mathcal{O}}}{\|\mathcal{O}(\bar{A}, f)\|_{\mathcal{O}}}.$$

Iterate k :

- ① Compute discrimination rate $\Delta^k(f)$ for any f in $\mathcal{F}^k = \mathcal{F} \setminus \{f^p\}_{p=1,\dots,k-1}$.
- ② Select $f^k \in \arg \max_{f \in \mathcal{F}^k} \Delta^k(f).$
- ③ Measure $\mathcal{O}(A_\varepsilon, f^k).$
- ④ Define the effective coefficient

$$\bar{A}^k \in \arg \min_{\bar{A} \in \mathcal{A}} \gamma^k(\bar{A}).$$

Numerical results

- **Microstructure.** Consider

$$A_\varepsilon(x) = \begin{cases} \gamma_1 + \gamma_2 \sin\left(\frac{2\pi x}{\varepsilon}\right) & \text{if } x \in D_1, \\ \gamma_3 & \text{if } x \in D_2. \end{cases}$$

with

$$\gamma_3 = a_\star$$

the limit in the sense of homogenization of
 $x \mapsto \gamma_1 + \gamma_2 \sin\left(\frac{2\pi x}{\varepsilon}\right)$.

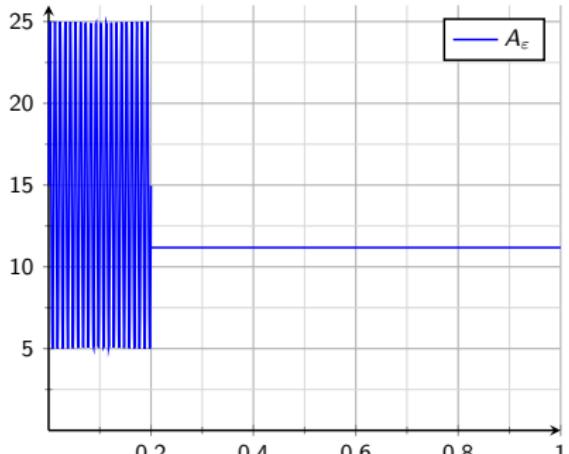


Figure: Coefficient A_ε with $\varepsilon = 0.01$.

Numerical results

- Microstructure.
- Set \mathcal{F} . Consider

$$f_n = \mathbb{1}_{\left(\frac{n-1}{N}, \frac{n}{N}\right)}$$

and

$$\mathcal{F} = \underbrace{\{f_n \text{ s.t. } \text{Supp}(f_n) \subset D_1\}}_{\mathcal{F}_{\text{disc}}} \cup \underbrace{\{f_n \text{ s.t. } \text{Supp}(f_n) \subset D_2\}}_{\mathcal{F}_{\text{non-disc}}}.$$

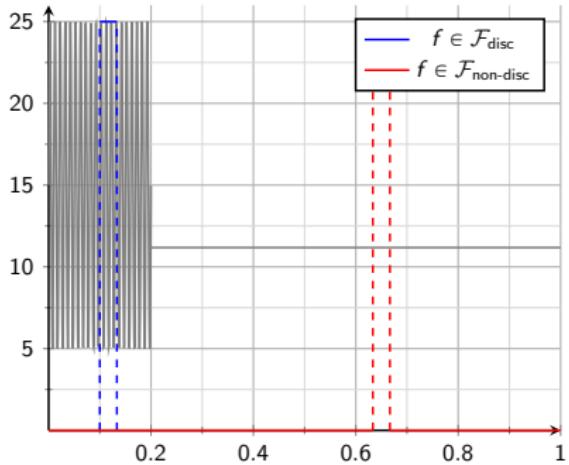


Figure: Admissible loadings.

Numerical results

Microstructure.

Set \mathcal{F} .

Set \mathcal{A} . Consider

$$\bar{A}(x) = \begin{cases} \bar{A}^1 & \text{if } x < 0.2, \\ \bar{A}^2 & \text{if } x > 0.2. \end{cases}$$

with \bar{A}^1 and \bar{A}^2 are constants.

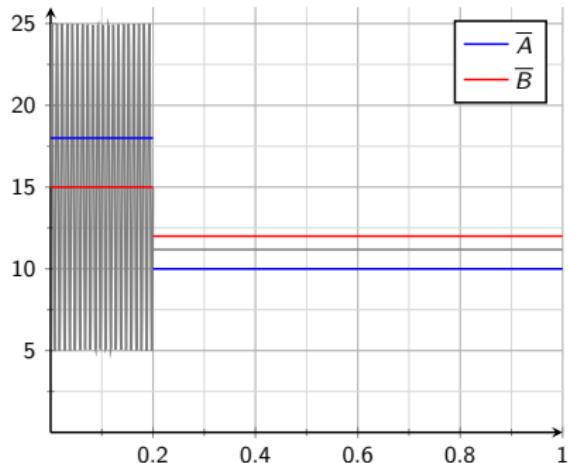


Figure: Effective coefficients.

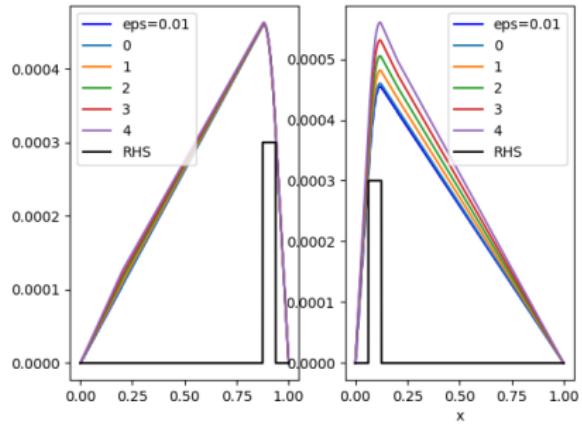
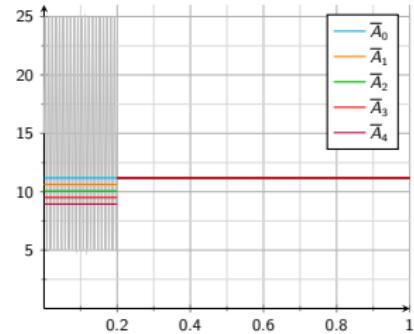
Numerical results: case 1

Setting:

$$\begin{aligned}\mathcal{A}_1 = \{\bar{A}_0 &= (a_*, a_*), \\ \bar{A}_1 &= (0.95a_*, a_*), \\ \bar{A}_2 &= (0.9a_*, a_*), \\ \bar{A}_3 &= (0.85a_*, a_*), \\ \bar{A}_4 &= (0.8a_*, a_*)\},\end{aligned}$$

and

$$\frac{\text{card}(\mathcal{F}_{\text{disc}})}{\text{card}(\mathcal{F})} \approx 0.2.$$



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and

$$\frac{\text{card}(\mathcal{F}_{\text{disc}})}{\text{card}(\mathcal{F})} \approx 0.2.$$

Conclusions:

- Δ_1 , Δ_2 and Δ_ε perform the same selection.
- Loadings in $\mathcal{F}_{\text{disc}}$ are first identified.
- $\bar{A}_0 = (a_*, a_*)$ is the best coefficient at each step.

Step	Loadings	Best Coefficient
1	$f_1 \in \mathcal{F}_{\text{disc}}$	\bar{A}_0
2	$f_2 \in \mathcal{F}_{\text{disc}}$	\bar{A}_0
3	$f_3 \in \mathcal{F}_{\text{disc}}$	\bar{A}_0
4	$f_4 \in \mathcal{F}_{\text{disc}}$	\bar{A}_0
5	$f_{12} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0
6	$f_{13} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0
7	$f_{14} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0
8	$f_{15} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0
9	$f_{16} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0
10	$f_{17} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0
11	$f_{18} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0
12	$f_{19} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0
13	$f_{20} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0
14	$f_{21} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0
15	$f_{22} \in \mathcal{F}_{\text{non-disc}}$	\bar{A}_0

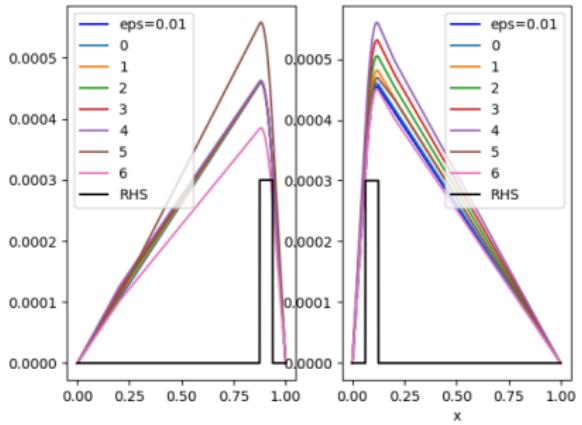
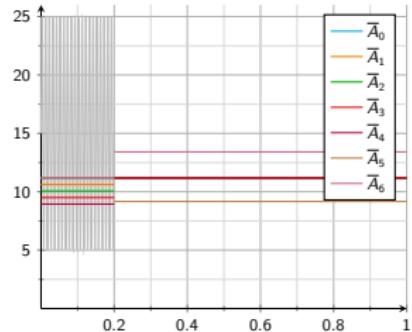
Numerical results: case 2

Setting:

$$\begin{aligned}\mathcal{A}_2 = \mathcal{A}_1 \cup \{\bar{A}_5 = (a_*, 0.82a_*), \\ \bar{A}_6 = (a_*, 1.2a_*)\},\end{aligned}$$

and

$$\frac{\text{card}(\mathcal{F}_{\text{disc}})}{\text{card}(\mathcal{F})} \approx 0.2.$$



Numerical results: case 2

Setting:

$$\mathcal{A}_2 = \mathcal{A}_1 \cup \{\bar{A}_5 = (a_*, 0.82a_*), \\ \bar{A}_6 = (a_*, 1.2a_*)\},$$

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Conclusions:

- Δ_2 better replicates Δ_ε than Δ_1 .
- $\bar{A}_0 = (a_*, a_*)$ is the best coefficient at each step.
- Δ_1 selects loadings that do not precisely discriminate \bar{A}_0 from other coefficients in \mathcal{A}_1 .

Step	Δ_ε	Δ_2
1	$f_1 \in \mathcal{F}_{\text{disc}}$	$f_{29} \in \mathcal{F}_{\text{non-disc}}$
2	$f_2 \in \mathcal{F}_{\text{disc}}$	$f_1 \in \mathcal{F}_{\text{disc}}$
3	$f_3 \in \mathcal{F}_{\text{disc}}$	$f_2 \in \mathcal{F}_{\text{disc}}$
4	$f_4 \in \mathcal{F}_{\text{disc}}$	$f_3 \in \mathcal{F}_{\text{disc}}$
5	$f_{29} \in \mathcal{F}_{\text{non-disc}}$	$f_{28} \in \mathcal{F}_{\text{non-disc}}$
6	$f_{28} \in \mathcal{F}_{\text{non-disc}}$	$f_4 \in \mathcal{F}_{\text{disc}}$
7	$f_{12} \in \mathcal{F}_{\text{non-disc}}$	$f_{27} \in \mathcal{F}_{\text{non-disc}}$
8	$f_{27} \in \mathcal{F}_{\text{non-disc}}$	$f_{12} \in \mathcal{F}_{\text{non-disc}}$
9	$f_{26} \in \mathcal{F}_{\text{non-disc}}$	-
10	$f_{25} \in \mathcal{F}_{\text{non-disc}}$	-
11	$f_{24} \in \mathcal{F}_{\text{non-disc}}$	-
12	$f_{23} \in \mathcal{F}_{\text{non-disc}}$	-
13	$f_{22} \in \mathcal{F}_{\text{non-disc}}$	-
14	$f_{21} \in \mathcal{F}_{\text{non-disc}}$	-
15	$f_{20} \in \mathcal{F}_{\text{non-disc}}$	-
16	$f_{19} \in \mathcal{F}_{\text{non-disc}}$	-
17	$f_{18} \in \mathcal{F}_{\text{non-disc}}$	-
17	$f_{17} \in \mathcal{F}_{\text{non-disc}}$	-

Numerical results: case 2

Setting:

$$\mathcal{A}_2 = \mathcal{A}_1 \cup \{\bar{A}_5 = (a_*, 0.82a_*), \\ \bar{A}_6 = (a_*, 1.2a_*)\},$$

and

$$\frac{\text{card}(\mathcal{F}_{\text{disc}})}{\text{card}(\mathcal{F})} \approx 0.2.$$

Conclusions:

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Step	Δ_ε	Δ_1
1	$f_1 \in \mathcal{F}_{\text{disc}}$	$f_{29} \in \mathcal{F}_{\text{non-disc}}$
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3	$f_3 \in \mathcal{F}_{\text{disc}}$	$f_{27} \in \mathcal{F}_{\text{non-disc}}$
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5	$f_{29} \in \mathcal{F}_{\text{non-disc}}$	$f_{25} \in \mathcal{F}_{\text{non-disc}}$
6	$f_{28} \in \mathcal{F}_{\text{non-disc}}$	$f_{24} \in \mathcal{F}_{\text{non-disc}}$
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17	$f_{18} \in \mathcal{F}_{\text{non-disc}}$	$f_{13} \in \mathcal{F}_{\text{non-disc}}$
17	$f_{17} \in \mathcal{F}_{\text{non-disc}}$	$f_1 \in \mathcal{F}_{\text{non-disc}}$

Convergence analysis

Work with A. Cohen.

Schrödinger equation