

# Summaries of Papers for FELIX Internship 2026

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## I. KIRILYUK 2010 ULTRAFAST OPTICAL MANIPULATION OF MAGNETIC ORDER

[1]

## II. GIDDING 2023 DYNAMIC SELF-ORGANISATION AND PATTERN FORMATION BY MAGNON-POLARONS

In this article [2], it is shown that some peculiar patterns arise when a sample is hit by a pump pulse to change its magnetic order.

Realy precise switching of magnetic order(spins) is important for low energy cost data storage. The abstract states that it is a known fact that when a sample is hit by a ultra short pump pulse the resulting magenitization is chaotic. This would be due to internal instabilities in the sample. However it turns out that the bahavior of the magnetic reagion is not necceraly chaotic at all, some patterns will arise in these cases. It is also well understood that a spatially-localised perturbation creates propagating waves with wave vectors determined by the profile of the excitation.

I now wonder what could be the cause of these peculier paterns and how they might be used in the future. I also still don't know what magnon-polarons are.

## III. KWAAITAAL 2024 EPSILON-NEAR-ZERO REGIME ENABLES PERMANENT ULTRAFAST ALL-OPTICAL REVERSAL OF FERROELECTRIC POLARIZATION

In this article [3], it is shown that ultrafast excitation under epsilon-near-zero (ENZ) conditions can permanently reverse ferroelectric polarization between stable states.

ENZ materials have a dielectric constant  $\epsilon \approx 0$  whch enhances light-matter interactions. This article shows that in ENZ conditions it is possible to achieve permanent all optical switching of an order parameter. This means that only light is used to achieve a switching of an order parameter, the order parameter in this article is the ferroelectric polarization. After switching from one order to the other the polarization remains stable and thus permanent.

## IV. DAVIES 2024 PHONONIC SWITCHING OF MAGNETIZATION BY THE ULTRAFAST BARNETT EFFECT

In this article [4], it is shown that spontanious magnetization can be achieved using the ultrafast Barnett effect.

This is done through the resonant excitation of circularly polarized optical phonons in a paramagnetic substrate. The Barnett effect describes how an inertial body with zero net magnetic momentum can aquire magnetization when mechanically spinning. When the substrate is circularly polarized it generates a magnetic field that can permanently and selectively change the magnetization of the upper layer. This effect only happens when the laser frequnicy is in resonance with the phononmodiof the substrate.

## V. STUPAKIEWICZ 2021 ULTRAFAST PHONONIC SWITCHING OF MAGNETIZATION

In this article [5], it is shown that certain patterns arise when we magnetize a sample using a laser.

The patterns rotates as the polarization direction of the laser rotates. The collective excitation modes (like magnons and phonons) define the energy range that determines all the important and intriguing thermodynamic and macroscopic properties of solids, such as electric, magnetic or crystallographic order, and the superconducting transition temperature. Control of the crystal structure of materials is the core aim of the field of straintonics. For the experiment they used an accumulation of pump pulses to not damage the sample but still get results. The multi pulse approach seems to only have the effect of a slight growth of the domains. The best magnetic switching occurs at wavelengths of  $\lambda = 14\mu m$  which also shows a big phonon response. It seems tha the LO phonon are more responsible for the magnetic switching than the TO phonons, but i don't know what they mean, further investigation required.

In the future, ultrafast modification of the crystal field environment, and thus of magnetocrystalline anisotropy, may become the most universal way to manipulate magnetization. Magneto-elastic interactions are present in all materials and thus can be used everywhere, for example in antiferromagnets.

## VI. BEAUREPAIRE 1996 ULTRAFAST SPIN DYNAMICS IN FERROMAGNETIC NICKEL

In this article [6], the relaxation processes of electrons and spins systems following the absorption of femtosecond optical pulses in ferromagnetic nickel have been studied. They have used pump probe techniques and have shown that the magnetization drops rapidly after just a few pico-seconds.

They talk about the kerr effect, I should investigate what it is exactly. The magneto optical kerr effect is known as MOKE. The aim of this paper is to study both electronic and spin dynamics after excitation of a Ni film with 60 fs pulses. The delays between pump and probe are achieved using a modified Michelson interferometer (should learn what this is). The signals are recorded using a boxcar and a lock-in synchronous detection. The information about the spin dynamics is contained in the time evolution of the hysteresis loops recorded for each time delay  $\delta t$ . hysteresis loop basically means that the output depends on the history of the input.

This work basically showed that using optical and magneto-optical techniques can be used to measure extremely fast events. The experiment showed that for the first few picoseconds the dynamics of spin and electron temperatures are different.

## VII. KIMEL 2005 NON-THERMAL OPTICAL CONTROL OF MAGNETIZATION IN FERROMAGNETIC SEMICONDUCTORS

[7]

## VIII. BIGOT 2009 COHERENT ULTRAFAST MAGNETISM INDUCED BY FEMTOSECOND LASER PULSES

[8]

## IX. FORST 2011 NONLINEAR PHONONICS AS AN ULTRAFAST ROUTE TO LATTICE CONTROL

[9]

## X. NOVA 2017 AN EFFECTIVE MAGNETIC FIELD FROM OPTICALLY DRIVEN LATTICE VIBRATIONS

[10]

## XI. MAGNETISM: FROM FUNDAMENTALS TO NANOSCALE DYNAMICS

[11]

## XII. INTRODUCTION TO SOLID STATE PHYSICS

[12]

## XIII. NONLINEAR OPTICS

[13]

## XIV. MAGNETIC DOMAINS: THE ANALYSIS OF MAGNETIC MICROSTRUCTURES

[14]

## XV. QUANTUM OPTICS

[15]

## XVI. KIMEL 2019 NONTHERMAL OPTICAL CONTROL OF MAGNETISM AND ULTRAFAST SPINTRONICS

[16]

## XVII. MISHRA 2021 PLASMON-ENHANCED ULTRAFAST DEMAGNETIZATION IN MAGNETOPHOTONIC NANOSTRUCTURES

[17]

## XVIII. SAVOINI 2018 TRACING THE ULTRAFAST MAGNETIC RESPONSE WITH RESONANT X-RAY DIFFRACTION

[18]

## XIX. SCHUBERT 2017 SUB-CYCLE CONTROL OF TERAHERTZ HIGH-HARMONIC GENERATION BY DYNAMICAL BLOCH OSCILLATIONS

[19]

**XX. KALASHNIKOVA 2018 ULTRAFAST LATTICE CONTROL OF MAGNETIC ANISOTROPY IN ORTHOFERRITES**

[20]

**XXI. HENIGHAN 2016 GENERATION MECHANISM OF THZ-FREQUENCY COHERENT ACOUSTIC PHONONS IN FE BY ULTRAFAST OPTICAL EXCITATION**

[21]

**XXII. MISHRA 2020 ULTRAFAST DEMAGNETIZATION AND SPIN TRANSPORT IN CO/PT MULTILAYERS**

[22]

**XXIII. CIUCIULKAITE 2021 COHERENT CONTROL OF OPTICAL PHONONS IN IRON GARNET FILMS**

[23]

**XXIV. MISHRA 2019 DYNAMIC REGIMES OF MULTI-SHOT ALL-OPTICAL SWITCHING IN FERRIMAGNETIC ALLOYS**

[24]

**XXV. SAVOINI 2020 SPIN-LATTICE RELAXATION AT ULTRAFAST TIMESCALES INVESTIGATED VIA RESONANT X-RAY SCATTERING**

[25]

**XXVI. MICROMAGNETIC STUDIES OF LASER-INDUCED MAGNETIZATION DYNAMICS IN FEPT-C FILMS**

In this article [26], they have simulated laser-induced magnetization dynamics using a hybrid Monte Carlo micromagnetic method. The results show that the magnetization dynamics includes an ultrafast demagnetization, a slower magnetization recovery, and a long-timescale magnetization reversal or the continuing recovery, depending on the magnitude of laser fluence and the external magnetic field.

They have only studied linearly polarized light, but the experimental and simulated results closely match.

This shows that hybrid Monte carlo micromagnetic simulations are a good way to go about simulating these kinds of effects. They have used time-resolved magneto-optical Kerr effect (TR-MOKE) to derive the spin temperature profile directly. In the TR-MOKE measurement, a pump pulse excites the sample, and then the reflected probe beam is split into two orthogonal polarized components, denoted by signal A and signal B. The difference between their intensities, IA - IB, is due to the rotation of plane of polarization, which is induced by the magneto-optical Kerr effect.

**XXVII. MICROMAGNETIC MODELING OF ALL OPTICAL SWITCHING OF FERROMAGNETIC THIN FILMS: THE ROLE OF INVERSE FARADAY EFFECT AND MAGNETIC CIRCULAR DICHROISM**

In this article [27], a all optical switching (AOS) experiment has been simulated. They investigated the Inverse Faraday Effect (IFE) and the Magnetic Circular Dichroism (MCD) as effects of the magnetic switching. They found that the IFE produced domain wall movements the most accurately, but for local inversion of the initial magnetic state no big difference has been found.

The simulation used a micromagnetic model based on the Landau–Lifshitz–Bloch equation coupled to the heat transport. The temperature evolution in the system under the action of the laser pulses is described by the two Temperature Model (2TM), with  $T_e$  the electrons temperature and  $T_l$  the lattice temperature. So they simulated ultra short laser pulses hitting ferromagnetic thin films. The power of the laser s describe by:

$$P(r, t) = P_0 \exp \left[ -\frac{r^2}{d_0^2/(4 \ln(2))} \right] \exp \left[ -\frac{(t - t_0)^2}{\tau_L^2/(4 \ln(2))} \right]$$

With  $r = \sqrt{x^2 + y^2}$ : the distance from the laser spot center,  $d_0$ : the Full Width at Half Maximum (FWHM) of laser spot,  $\tau_L$ : the FWHM of the pulse duration. The maximum power of the laser is:

$$P_0 = \frac{F}{t_{FM} \tau_L}$$

where  $F$  is the laser fluence and  $t_{FM}$  is the thickness of the ferromagnetic sample. Then as the laser heats up the sample, we use the 2TM with temperature dynamics given by:

$$\begin{aligned} C_e \frac{\partial T_e}{\partial t} &= -k_e \nabla^2 T_e - g_{el}(T_e - T_l) + \eta(\mathbf{m}, \sigma)(1 - R)P(\mathbf{r}, t) \\ C_l \frac{\partial T_l}{\partial t} &= -k_l \nabla^2 T_l - g_{el}(T_l - T_e) - \frac{C_l}{\tau_{sub}}(T_l - T_{sub}) \end{aligned}$$

With  $C_e = \gamma_e T_e$  the electron heat capacity(which is linear in  $T_e$ ),  $C_l$  the lattice heat capacity(which is constant above the debye temperature),  $k_i$  the thermal conductivity,  $g_{el}$

the electron-lattice coupling,  $\mu(m, \sigma)$  the polarization- and magnetization-dependent absorption factor (MCD),  $R$  the reflectivity and  $\tau_{sub}$  the heat transfer time to the substrate. As the laser pulse usually heats the sample close or even over its Curie temperature  $T_C$ , the Landau–Lifshitz–Gilbert equation (LLG) cannot be used to describe the magnetization dynamics. Instead, the Landau–Lifshitz–Bloch (LLB) equation must be employed. In this case, the dynamics of the normalized magnetization is given by:

$$\frac{d\vec{m}}{dt} = -\gamma'_0 \vec{m} \times \vec{H}_{\text{eff}} - \gamma'_0 \frac{\alpha_{\perp}}{m^2} \vec{m} \times (\vec{m} \times (\vec{H}_{\text{eff}} + \vec{H}_{\text{th}}^{\perp})) \\ + \gamma'_0 \frac{\alpha_{\parallel}}{m^2} (\vec{m} \cdot \vec{H}_{\text{eff}}) \vec{m} + \vec{H}_{\text{th}}^{\parallel}$$

where the normalized magnetization is:

$$\vec{m}(\vec{r}, t) = \frac{\vec{M}(r, t)}{M_s^0}$$

With  $M_s^0$  is the saturation magnetization at  $T = 0K$ . The longitudinal and transverse damping parameters in the LLB model are given by:

$$\alpha_{\parallel} = \frac{2\lambda T}{3T_C m_e^2}, \quad \alpha_{\perp} = \lambda \left(1 - \frac{T}{3T_C}\right),$$

with  $\lambda$  the microscopic coupling to the thermal bath,  $T$  the local temperature (from the 2TM),  $T_C$  the Curie temperature and  $m_e(T)$  the equilibrium magnetization. The effective field entering the LLB equation consists of several contributions:

$$\vec{H}_{\text{eff}} = \vec{H}_{\text{ext}} + \vec{H}_{\text{ani}} + \vec{H}_{\text{ex}} + \vec{H}_{\text{dem}} + \vec{H}_{\text{long}},$$

where  $\vec{H}_{\text{ext}}$  is the external field,  $\vec{H}_{\text{ani}}$  the uniaxial anisotropy field,  $\vec{H}_{\text{ex}}$  the exchange field,  $\vec{H}_{\text{dem}}$  the demagnetizing field, and  $\vec{H}_{\text{long}}$  the longitudinal relaxation field. The longitudinal field is expressed as:

$$\vec{H}_{\text{long}} = \frac{1}{2\chi_{\parallel}} \left(1 - \frac{m^2}{m_e^2(T)}\right) \vec{m},$$

with  $\chi_{\parallel}$  the longitudinal susceptibility. The transverse dynamics is driven by the transverse susceptibility:

$$\vec{H}_{\perp} = \frac{1}{\chi_{\perp}} (\vec{m} \cdot \hat{e}_{\perp}) \hat{e}_{\perp},$$

where  $\chi_{\perp}$  is the transverse susceptibility and  $\hat{e}_{\perp}$  is the local transverse direction of deviation. Finally, the model includes the thermal stochastic fields in longitudinal and transverse directions:

$$\langle H_{\text{th},i}^{\perp}(t) \rangle = 0, \quad \langle H_{\text{th},i}^{\perp}(t) H_{\text{th},j}^{\perp}(t') \rangle = \frac{2\alpha_{\perp} k_B T}{\gamma'_0 M_s^0 V} \delta_{ij} \delta(t - t'), \\ \langle H_{\text{th}}^{\parallel}(t) \rangle = 0, \quad \langle H_{\text{th}}^{\parallel}(t) H_{\text{th}}^{\parallel}(t') \rangle = \frac{2\alpha_{\parallel} k_B T}{\gamma'_0 M_s^0 V} \delta(t - t'),$$

with  $k_B$  the Boltzmann constant and  $V$  the computational cell volume.

The optical excitation produced by the circularly polarized laser pulse introduces two additional mechanisms into the micromagnetic dynamics: the Inverse Faraday Effect (IFE) and the Magnetic Circular Dichroism (MCD). The IFE produces an effective magnetic field acting along the laser propagation direction, while the MCD modifies the local absorption depending on the magnetization orientation. The IFE field is written as:

$$\vec{B}_{\text{MO}}(\vec{r}, t) = (\sigma_{\pm}) B_0^{\text{MO}} f_{\text{MO}}(\vec{r}, t) \hat{u}_z, \\ \vec{H}_{\text{MO}}(\vec{r}, t) = \frac{1}{\mu_0} \vec{B}_{\text{MO}}(\vec{r}, t),$$

where  $\sigma_{\pm} = \pm 1$  denotes the laser helicity (right- or left-handed circular polarization),  $B_0^{\text{MO}}$  the maximum optically induced magnetic flux density, and  $f_{\text{MO}}(\vec{r}, t)$  a spatiotemporal envelope similar to the laser profile:

$$f_{\text{MO}}(\vec{r}, t) = \begin{cases} \exp\left[-\frac{r^2}{d_0^2/(4\ln 2)}\right] \\ \times \exp\left[-\frac{(t-t_0)^2}{\tau_L^2/(4\ln 2)}\right], & t < t_0, \\ \exp\left[-\frac{r^2}{d_0^2/(4\ln 2)}\right] \\ \times \exp\left[-\frac{(t-t_0)^2}{(\tau_L + \tau_d)^2/(4\ln 2)}\right], & t \geq t_0. \end{cases}$$

Here  $d_0$  is the laser spot FWHM,  $\tau_L$  is the pulse duration FWHM, and  $\tau_d$  accounts for a possible delay of the optically induced field decay. The IFE field therefore introduces a helicity-dependent torque which can move domain walls or contribute to magnetization reversal. The Magnetic Circular Dichroism (MCD) modifies the absorption of the laser power depending on the relative orientation between the magnetization and the light helicity. This is described by:

$$\eta(\vec{m}, \sigma) = A_{\text{LP}} \left[1 - \sigma_{\pm} m_z \text{MCD}(\%) \right],$$

where  $A_{\text{LP}}$  is the absorption coefficient for linearly polarized light and  $\text{MCD}(\%)$  is the dichroism amplitude. The absorbed power density entering the 2TM heat equation then becomes  $\eta(\vec{m}, \sigma)(1 - R)P(\vec{r}, t)$ , introducing a helicity-dependent local heating. Finally, both optical effects enter the Landau–Lifshitz–Bloch dynamics through the effective field:

$$\vec{H}_{\text{eff}} = \vec{H}_{\text{ext}} + \vec{H}_{\text{ani}} + \vec{H}_{\text{ex}} + \vec{H}_{\text{dem}} + \vec{H}_{\text{long}} + \vec{H}_{\text{MO}},$$

thus ensuring that the laser–material interaction is fully included in both the thermal and magnetic evolution of the system.

## XXVIII. ALL-OPTICAL SPIN SWITCHING PROBABILITY IN [TB/CO] MULTILAYERS

## XXIX. MICROMAGNETICS AT FINITE TEMPERATURE

This article [29], shows the development of hybrid Monte carlo (HMC) micromagnetic simulations. The HMC micromagnetics is a self-consistent method for the magnetic studies at finite temperature.

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