# Appendix: CMB Dipole Constraints in RSVP (Entropic Redshift Form)

## A.1 Fields, Normalization, and ACDM Dictionary

The RSVP framework employs the following plenum fields:

- Scalar capacity ≡ (drives outward expansion),
- Matter density  $\delta \rho_m$  (drives inward gravitational attraction),
- Bulk flow **u** (peculiar velocity).

The entropic redshift potential is defined as:

$$\Upsilon \equiv \delta \Phi - \beta(\eta) \, \varphi_m$$

$$\varphi_m(k,\eta) = \frac{4\pi G a^2(\eta) \,\bar{\rho}_m(\eta)}{k^2} \,\mathcal{T}_m(k,\eta) \,\delta_m(k,\eta)$$

where  $\delta\Phi$  is the scalar potential perturbation,  $\varphi_m$  is the dimensionless matter potential,  $\beta(\eta)$  is a time-dependent coupling,  $a(\eta)$  is the scale factor,  $\bar{\rho}_m(\eta)$  is the mean matter density,  $\mathcal{T}_m(k,\eta)$  is the matter transfer function, and  $\delta_m(k,\eta)$  is the matter density contrast.

**Convention.** The potential  $\Upsilon$  is normalized such that the instantaneous Sachs–Wolfe contribution at decoupling (denoted by subscript  $_*$ ) is:

$$\left(\frac{\Delta T}{T}\right)_{\rm SW} = \frac{1}{3}\Upsilon_*$$

For consistency with  $\Lambda$ CDM, we set  $\beta(\eta_*) = \frac{4\pi G \, a^2(\eta_*) \, \bar{\rho}_m(\eta_*)}{k^2} \mathcal{T}_m(k, \eta_*)$  at last scattering, ensuring  $\Upsilon_* = \delta \Phi_* - \alpha_m \delta \rho_m$  with  $\alpha_m = \beta(\eta_*)$ .

# A.2 Large-Angle Anisotropy in RSVP Form

For a sightline  $\hat{\mathbf{n}}$ , the CMB temperature anisotropy is:

$$\frac{\Delta T}{T}(\hat{\mathbf{n}}) = \underbrace{\hat{\mathbf{n}} \cdot \frac{\mathbf{u}_0}{c}}_{\text{kinematic dipole } \varepsilon_{\text{kin}} \sim 10^{-3}} + \underbrace{\frac{1}{3}\Upsilon_*(\hat{\mathbf{n}})}_{\text{entropic SW}} + \underbrace{2\int_{\eta_*}^{\eta_0} \dot{\Upsilon} \, d\eta}_{\text{entropic ISW}}$$

The intrinsic dipole amplitude, after subtracting the kinematic contribution, is:

$$\left| \left( \frac{\Delta T}{T} \right)_{\ell=1}^{\text{int}} \right| \equiv \varepsilon_{\text{int}} \text{few} \times 10^{-5}$$

### A.3 Super-Horizon Gradient Bound

A beyond-horizon inhomogeneity is modeled as a nearly uniform gradient:

$$\Upsilon(\mathbf{x}, \eta) \simeq \Upsilon_0(\eta) + \mathbf{G}(\eta) \cdot \mathbf{x}, \quad \|\mathbf{G}\|_{R_*} \ll 1$$

where  $R_*$  is the radius of the observable patch at decoupling. The intrinsic dipole contribution is:

$$D_{ ext{int}} pprox rac{1}{3} \| \mathbf{G}_* \| R_* + \mathcal{O} \left( \int \dot{\Upsilon} \, d\eta 
ight)$$

The gradient bound is:

$$\boxed{ \|\nabla \Upsilon_* \| R_* 3\varepsilon_{\mathrm{int}} \|} \iff \boxed{ \|\mathbf{G}_* \| \frac{3\varepsilon_{\mathrm{int}}}{R_*} }$$

# A.4 Effective Potential and "Falling Outward"

The effective potential governing outward and inward accelerations is:

$$ig| \mathbf{a}_{ ext{eff}} = -
abla \Phi_{ ext{eff}}, \quad \Phi_{ ext{eff}} \equiv \Phi - \gamma(\eta) arphi_m$$

where  $\gamma(\eta)$  is a time-dependent coupling. The dipole bound implies:

$$\boxed{ |\delta\Phi|_* \frac{\varepsilon_{\rm int}}{\alpha_\Phi} = \varepsilon_{\rm int}, \quad |\delta\rho_m|_* \frac{\varepsilon_{\rm int}}{\alpha_m} }$$

Under the chosen calibration,  $\alpha_{\Phi} = 1$  at decoupling, ensuring consistency with the normalization.

# A.5 Bulk-Flow Convergence (Vector Test)

The RSVP bulk-flow estimator over a sphere of radius R is:

$$\mathbf{u}_0^{\mathrm{RSVP}}(R) := \arg\min_{\mathbf{u}} \sum_{i: r_i < R} w_i \left( z_i^{\mathrm{obs}} - z_i^{\mathrm{RSVP}}(\mathbf{u}) \right)^2, \quad w_i \propto \frac{1}{\sigma_{S,i}}$$

where  $z_i^{\text{obs}}$  is the observed redshift,  $z_i^{\text{RSVP}}$  is the model-predicted redshift, and  $\sigma_{S,i}$  is the uncertainty. The estimator converges as:

$$\angle \left( \mathbf{u}_0^{\mathrm{RSVP}}(R), \mathbf{d}_{\mathrm{CMB}} \right) \rightarrow 0, \quad \| \mathbf{u}_0^{\mathrm{RSVP}}(R) \| \rightarrow c \varepsilon_{\mathrm{kin}}$$

# A.6 Long-Mode Consistency (Semantic-Slicing Gauge)

Super-horizon adiabatic modes in RSVP correspond to a semantic-slicing gauge redefinition. The leading dipole cancels, and any residual dipole is linked to the quadrupole via  $\Upsilon$  evolution at horizon entry, analogous to the Grishchuk–Zel'dovich relation. The observed quadrupole ( $\sim 10^{-5}$ ) further constrains the bound in A.3.

#### A.7 Bottom Line

The potential  $\Upsilon$  encodes both outward expansion ( $\Xi$ ) and inward gravitational pull ( $\delta \rho_m$ ) in the CMB. The residual dipole limit is:

$$\Delta \Upsilon_* \equiv \|\nabla \Upsilon_*\| R_* \text{few} \times 10^{-5}$$

Bulk-flow alignment and the small quadrupole amplitude indicate that homogeneity and isotropy extend at least one horizon length beyond the observable universe, disfavoring adjacent "bubble" domains with radically different large-scale parameters.