

# Partial Differential Equations

## Problem sheet 1

### Answers

1. The Green's function for the IVP

$$y'' + p(x)y' + q(x)y = r(x), \quad \text{with } y(0) = 0 \text{ and } y'(0) = 0 \quad (1)$$

is given by

$$G(x, s) = \frac{y_1(s)y_2(x) - y_1(x)y_2(s)}{W[y_1, y_2](s)} \quad (2)$$

where  $y_1$  and  $y_2$  are independent solutions of the homogeneous equation

$$y'' + p(x)y' + q(x)y = 0 \quad (3)$$

(a) (2 marks) The functions

$$\tilde{y}_1(x) = ay_1(x) + by_2(x)$$

$$\tilde{y}_2(x) = cy_1(x) + dy_2(x)$$

are also solutions of (3). Give a condition on  $a, b, c$  and  $d$  which makes  $\tilde{y}_1(x)$  and  $\tilde{y}_2(x)$  independent solutions.

**Solution to 1(a):** Differentiating the equations for  $\tilde{y}_1(x)$  and  $\tilde{y}_2(x)$  one obtains

$$\tilde{y}_1'(x) = ay_1'(x) + by_2'(x)$$

$$\tilde{y}_2'(x) = cy_1'(x) + dy_2'(x)$$

Substituting this into  $W[\tilde{y}_1, \tilde{y}_2]$  gives after a small amount of algebra

$$W[\tilde{y}_1, \tilde{y}_2] = (ad - bc)(y_1y_2' - y_2y_1') = (ad - bc)W[y_1, y_2]$$

Since  $y_1$  and  $y_2$  are independent solutions we know that  $W[y_1, y_2](x) \neq 0$ . Hence

$$W[\tilde{y}_1, \tilde{y}_2] \neq 0 \Leftrightarrow ad - bc \neq 0$$

Note that this is simply the condition that the linear transformation relating  $\tilde{y}_1, \tilde{y}_2$  to  $y_1, y_2$  has non-zero determinant and is therefore invertible.

(b) (2 marks) Calculate the Green's function using independent functions  $\tilde{y}_1(x)$  and  $\tilde{y}_2(x)$  as given by the above equation and show that it is identical to the Green's function given by equation (2).

**Solution 1(b)** Using the above equations we calculate that

$$\tilde{y}_1(s)\tilde{y}_2(x) - \tilde{y}_1(x)\tilde{y}_2(s) = (ad - bc)\{y_1(s)y_2(x) - y_1(x)y_2(s)\}$$

We already know that

$$W[\tilde{y}_1, \tilde{y}_2](s) = (ad - bc)W[y_1, y_2](s)$$

So that using formula (2) but with  $\tilde{y}_1$  and  $\tilde{y}_2$  gives:

$$\begin{aligned}\frac{\tilde{y}_1(s)\tilde{y}_2(x) - \tilde{y}_1(x)\tilde{y}_2(s)}{W[\tilde{y}_1, \tilde{y}_2](s)} &= \frac{(ad - bc)\{y_1(s)y_2(x) - y_1(x)y_2(s)\}}{(ad - bc)W[y_1, y_2](s)} \\ &= \frac{y_1(s)y_2(x) - y_1(x)y_2(s)}{W[y_1, y_2](s)} \quad \text{Since } (ad - bc) \neq 0 \\ &= G(x, s)\end{aligned}$$

So that the Green's function constructed from  $\tilde{y}_1$  and  $\tilde{y}_2$  is **identical** to that constructed from  $y_1$  and  $y_2$ .

(c) (6 marks) Calculate the Green's function for the IVP

$$y'' - 4y' + 4y = r(x) \quad \text{with } y(0) = 0 \text{ and } y'(0) = 0 \quad (4)$$

and use this to obtain the solution to the above problem when  $r(x) = xe^{2x}$ .

**Solution 1(c)**

The homogeneous equation is

$$y'' - 4y' + 4y = 0$$

with *characteristic equation*

$$m^2 - 4m + 4 = 0$$

This is just  $(m - 2)^2 = 0$  so that the solution is  $m = 2$  (twice). We may therefore take  $y_1(x) = e^{2x}$  and  $y_2(x) = xe^{2x}$  as independent solutions of the homogeneous equation. The Wronskian is given by

$$W[y_1, y_2](x) = \begin{vmatrix} e^{2x} & xe^{2x} \\ 2e^{2x} & e^{2x} + 2xe^{2x} \end{vmatrix} = e^{4x}$$

Hence substituting into equation (2) gives:

$$G(x, s) = \frac{e^{2s}xe^{2x} - e^{2x}se^{2s}}{e^{4s}} = xe^{2x}e^{-2s} - e^{2x}se^{-2s}$$

Note (although we will not use this fact below) that we can write the above as  $G(x, s) = (x - s)e^{2(x-s)}$  so that  $G(x, s)$  can be written as a function of  $(x - s)$ .

To solve the equation  $y'' - 4y' + 4y = xe^{2x}$  with  $y(0) = 0$  and  $y'(0) = 0$  we use the Green's function calculated above and  $r(x) = xe^{2x}$

$$\begin{aligned}y(x) &= \int_{s=0}^x G(x, s)r(s)ds \\ &= \int_{s=0}^x (xe^{2x}e^{-2s} - e^{2x}se^{-2s})se^{2s}ds \\ &= xe^{2x} \int_{s=0}^x sds - e^{2x} \int_{s=0}^x s^2ds \\ &= xe^{2x}(x^2/2) - e^{2x}(x^3/3) \\ &= (x^3/6)e^{2x}\end{aligned}$$

So that the required solution is  $y(x) = (x^3/6)e^{2x}$ .

2. Let  $\lambda_n, n = 1, 2, 3, \dots$  be the eigenvalues, and  $y_n(x)$  be the corresponding eigenfunctions, for the boundary value problem

$$\frac{d^2y}{dx^2} = \lambda y, \quad y(0) = 0, \quad \frac{dy}{dx}(1) = 0, \quad x \in [0, 1] \quad (5)$$

- (a) (4 marks) **Without calculating** the  $y_n(x)$  use integration by parts to show that

$$\int_0^1 (y_m(x)y_n''(x) - y_n(x)y_m''(x)) dx = 0 \quad (6)$$

Hence show (**again without calculating the  $y_n(x)$** ) that

$$\int_0^1 y_m(x)y_n(x)dx = 0, \quad \text{for } m \neq n \quad (7)$$

**Solution 2(a)**

$$\begin{aligned} \int_0^1 (y_m(x)y_n''(x) - y_n(x)y_m''(x)) dx &= [y_m(x)y_n'(x) - y_n(x)y_m'(x)]_0^1 \\ &\quad + \int_0^1 (-y_m'(x)y_n'(x) + y_n'(x)y_m'(x)) dx \\ &= [y_m(x)y_n'(x) - y_n(x)y_m'(x)]_0^1 \\ &= 0 \quad \text{Since } y'(1) = 0 \text{ and } y(0) = 0 \end{aligned}$$

Which gives equation (6). We now use the fact that  $y_m''(x) = \lambda_m y_m(x)$  and  $y_n''(x) = \lambda_n y_n(x)$ . Substituting into (6) we get

$$\begin{aligned} 0 &= \int_0^1 (y_m(x)y_n''(x) - y_n(x)y_m''(x)) dx \\ &= \int_0^1 (\lambda_n y_m(x)y_n(x) - \lambda_m y_n(x)y_m(x)) dx \\ &= (\lambda_n - \lambda_m) \int_0^1 y_m(x)y_n(x)dx \end{aligned}$$

Now for  $m \neq n$  we have  $\lambda_n - \lambda_m \neq 0$  hence we must have

$$\int_0^1 y_m(x)y_n(x)dx = 0, \quad \text{for } m \neq n$$

which is equation (7).

- (b) (6 marks) Calculate the eigenvalues and eigenfunctions for the BVP given by equation (5). Substitute in the values you have obtained for  $y_n(x)$  and  $y_m(x)$  into (7) to obtain an “orthogonality condition” for certain trigonometric functions.

**Solution 2(b)** To calculate the eigenvalues and eigenfunctions we need to solve the BVP

$$y'' - \lambda y = 0, \quad y(0) = 0, \quad y'(1) = 0$$

For this ODE there are three cases:  $\lambda > 0$ ,  $\lambda = 0$  and  $\lambda < 0$ .

Case (i)  $\lambda > 0$  For this case we write  $\lambda = \mu^2$  (where  $\mu \neq 0$ ). Then the general solution is  $y(x) = Ae^{\mu x} + Be^{-\mu x}$ . The boundary condition  $y(0) = 0$  gives  $0 = A + B$ , so that  $B = -A$ . Hence we may write  $y(x) = A(e^{\mu x} - e^{-\mu x})$ . To apply the boundary condition  $y'(1) = 0$  we must differentiate  $y(x)$  to obtain  $y'(x) = A\mu(e^{\mu x} + e^{-\mu x}) = 2A\mu \cosh(\mu x)$ . Hence  $y'(1) = 0$  gives  $0 = 2A\mu \cosh(\mu)$ . Since  $\mu \neq 0$  this is only possible if  $A = 0$ . Hence the boundary conditions give  $y(x) = 0$  for all  $x$ . So case (i) only gives the *trivial solution*.

Case (ii)  $\lambda = 0$ . In this case the DE becomes  $y'' = 0$  with general solution  $y(x) = A + Bx$ . The boundary condition  $y(0) = 0$  gives  $0 = A$  so that  $y(x) = Bx$ . Differentiating gives  $y'(x) = B$  so the boundary condition  $y'(0) = 0$  gives  $B = 0$  and again we get the *trivial solution*.

Case (iii)  $\lambda < 0$ . In this case we take  $\lambda = -\mu^2$  (with  $\mu \neq 0$ ). Then the DE becomes  $y'' + \mu^2 y = 0$  with general solution  $y(x) = A \cos(\mu x) + B \sin(\mu x)$ . The boundary condition  $y(0) = 0$  gives  $0 = A$  and therefore  $y(x) = B \sin(\mu x)$ . Differentiating gives  $y'(x) = B\mu \cos(\mu x)$ . We now apply the boundary condition  $y'(1) = 0$ . This gives  $0 = B\mu \cos(\mu)$ . For non-trivial solutions this is only possible if

$$\cos(\mu) = 0 \quad \Rightarrow \quad \mu = \frac{\pi}{2} + k\pi, \quad k \in \mathbb{Z}$$

Hence the eigenvalues are given by

$$\lambda_k = - \left( \frac{(2k+1)\pi}{2} \right)^2 = - \frac{(2k+1)^2 \pi^2}{4}, \quad k \in \mathbb{Z}$$

The corresponding eigenfunctions are given by

$$y_k(x) = B \sin \left( \frac{(2k+1)\pi x}{2} \right), \quad k \in \mathbb{Z}$$

Note however that  $k = 0$  and  $k = -1$  give the same eigenvalues and eigenfunctions (at the expense of changing the sign of  $B$ ). Similarly  $k = 1$  and  $k = -2$  give the same eigenvalues and eigenfunctions, and in general  $n = n$  and  $k = -(n+1)$  give the same eigenvalues and eigenfunctions for  $n \in \mathbb{N}$ . Remember that Sturm-Liouville theory tells us that there is a minimum value of (with my convention)  $-\lambda$ . So that the **distinct** set of eigenvalues and eigenfunctions are given by

$$\lambda_n = - \frac{(2n+1)^2 \pi^2}{4} \quad n \in \mathbb{N}$$

$$y_n(x) = B \sin \left( \frac{(2n+1)\pi x}{2} \right), \quad n \in \mathbb{N}$$

Substituting the expressions for  $y_n(x)$  into equation (7) and ignoring the arbitrary constant  $B$  we get:

$$\int_0^1 \sin \left( \frac{(2m+1)\pi x}{2} \right) \sin \left( \frac{(2n+1)\pi x}{2} \right) dx = 0 \quad \text{for } m, n \in \mathbb{N} \text{ and } m \neq n$$

This is the required “orthogonality condition”.

Note without restricting  $m, n$  to be in  $\mathbb{N}$  rather than  $\mathbb{Z}$  the above result would not be true. For example the above would not be true if we chose  $m = 1$  and  $n = -2$  since it would be the integral of  $\sin^2(3\pi x/2)$  which integrates to  $1/2$ , rather than 0.

3. The function  $f(x)$  is  $2\pi$ -periodic and is defined for  $-\pi \leq x \leq \pi$  by

$$f(x) = \begin{cases} -1 & -\pi \leq x < 0 \\ 1 & 0 \leq x < \pi \end{cases} \quad (8)$$

- (a) (2 marks) Plot the graph of the function  $f(x)$  for  $-3\pi \leq x \leq 3\pi$  using Python.  
 (b) (2 marks) Explain why  $f(x)$  has a Fourier Sine series (i.e. the cosine terms in the Fourier series all vanish).

**Solution 3(b)** The function is  $2\pi$ -periodic so according to the theory of Fourier series one can write it as

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} a_n \cos(nx) + \sum_{n=1}^{\infty} b_n \sin(nx)$$

where

$$a_n(x) = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos(nx) dx, \quad n \in \mathbb{N}$$

$$b_n(x) = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx) dx, \quad n \in \mathbb{N}^*$$

To calculate  $a_n$  we note that in our case we have  $f(-x) = -f(x)$  and  $\cos(-nx) = \cos(nx)$  so that the integrand  $f(x) \cos(nx)$  is an **odd** function. Hence the integral on  $-\pi \leq x \leq 0$  exactly cancels the integral on  $0 \leq x \leq \pi$  and thus  $a_n = 0$ .

Alternatively just calculate the integral to show  $a_n = 0$  for  $n \in \mathbb{N}$ .

- (c) (6 marks) Calculate the Fourier series for  $f(x)$ .

**Solution 2(c)** We are left with the calculation of the  $b_n$ . Since  $f(x)$  has a different formula for  $x < 0$  and  $x > 0$  we need to split the integral for  $b_n$  into two parts as shown below:

$$\begin{aligned} b_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx) dx \\ &= \frac{1}{\pi} \int_{-\pi}^0 f(x) \sin(nx) dx + \frac{1}{\pi} \int_0^{\pi} f(x) \sin(nx) dx \\ &= \frac{1}{\pi} \int_{-\pi}^0 -\sin(nx) dx + \frac{1}{\pi} \int_0^{\pi} \sin(nx) dx \\ &= \frac{1}{\pi} \left[ \frac{1}{n} \cos(nx) \right]_{-\pi}^0 + \frac{1}{\pi} \left[ -\frac{1}{n} \cos(nx) \right]_0^{\pi} \\ &= \frac{1}{n\pi} (\cos(0) - \cos(-n\pi)) - \frac{1}{n\pi} (\cos(n\pi) - \cos(0)) \end{aligned}$$

Using the fact that  $\cos(0) = 1$ ,  $\cos(-n\pi) = \cos(n\pi) = (-1)^n$  we get

$$b_n = \frac{2}{n\pi} (1 - (-1)^n)$$

So that the Fourier series is given by

$$f(x) = \frac{2}{\pi} \sum_{n=1}^{\infty} \frac{(1 - (-1)^n)}{n} \sin(nx)$$

Noting that  $(1 - (-1)^n)$  is zero for  $n$  even and is equal to 2 for  $n$  odd. We may make sure that  $n$  is odd by writing  $n = 2m - 1$  for  $m \in \mathbb{N}^*$  which enables us to write the above in a slightly neater way as

$$f(x) = \frac{4}{\pi} \sum_{m=1}^{\infty} \frac{1}{2m-1} \sin((2m-1)x).$$