Interacting Particle Systems Diploma Thesis Presentation

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Outline

- 1 Introduction and Some Background
- Independent Random Walks
- 3 Simple Exclusion Process
- 4 Zero Range Process



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- Introduction and Some Background
- 2 Independent Random Walks
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Interacting Particle Systems?

- Large-scale systems of components interacting with each other governed by stochastic dynamics
- Several applications on
 - natural sciences: reaction diffusion, gas particles systems...
 - social sciences: traffic flow, opinion dynamics, spread of epidemics...



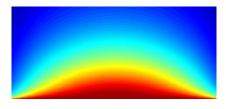
Introduction and Some Background Independent Random Walks Simple Exclusion Process Zero Range Process

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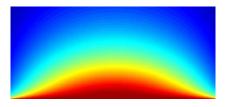
Deduce macroscopic behavior



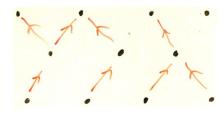


Goal?

Deduce macroscopic behavior



from microscopic interactions.





Introduction and Some Background Independent Random Walks Simple Exclusion Process Zero Range Process



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- Stochastic Process: family of random variables $(\eta_t)_{t\geq 0}$ in X



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- Markov Process: collection $(\mathbb{P}^{\eta}:\eta\in X)$ of probability measures **with**
- Markov Property: $\mathbb{P}^{\eta}(\eta_{t+s} \in A|\mathcal{F}_t) = \mathbb{P}^{\eta_t}(\eta_s \in A)$ "Memorylessness"



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- Markov Chain: Markov Process on countable now X characterized by transition rates $c(\eta, \eta')$



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An irreducible Markov chain with finite state space X has a unique invariant measure.

 Phase transitions, i.e. existence of more that one invariant measures, occur only in infinite systems.



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Space

- ullet \mathbb{Z}^d the d-dimensional integer lattice
- $\mathbb{T}_L = \mathbb{Z}_L = \{0,1,...,L-1\}$ the torus with L points and $\mathbb{T}_L^d = (\mathbb{T}_L)^d$
- \bullet L represents the inverse of the distance between the points of \mathbb{T}^d_L , namely the particle sites



Space and Stochastics

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- ullet L represents the inverse of the distance between the points of \mathbb{T}^d_L , namely the particle sites
- Transition probability p(x,y)=p(0,y-x)=:p(y-x) for some $p(\cdot)$ on \mathbb{Z}^d (elementary transition probability)



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Thus, there is only one invariant measure.



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Call a Poisson measure $\nu^L_{\rho(\cdot)}$ on \mathbb{T}^d_L associated to a fixed positive function $\rho:\mathbb{T}^d_L\to\mathbb{R}_+$, a probability on $\mathbb{N}^{\mathbb{T}^d_L}$, denoted by $\nu^L_{\rho(\cdot)}$, satisfying that:



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- under $\nu^L_{\rho(\cdot)}$ the random variables $(\eta(x):x\in\mathbb{T}^d_L)$ must be independent,
- $\bullet \ \ \text{for every fixed site} \ x \in \mathbb{T}^d_L,$

$$\nu_{\rho(\cdot)}^L(\eta(x) = k) = p_{\rho(x),k}.$$



Theorem

If particles are initially distributed according to a Poisson measure associated to a constant function equal to α then the distribution at time t is exactly the same Poisson measure ν_{α}^{L} .



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Note that

$$E_{\nu_{\alpha}^{L}}(\eta(x)) = \sum_{k>0} e^{-\alpha} \frac{\alpha^{k}}{k!} k = \alpha.$$

The Poisson measures are in this way naturally parametrized by the density of particles.

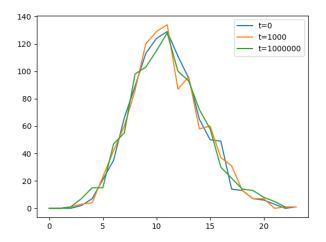


Figure: Distribution at time t of IRW on \mathbb{T}^1_{1000} with $\alpha=10$





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• Is that the same profile?



- ullet We have different space scales, \mathbb{T}^d and $L^{-1}\mathbb{T}^d_L$.
- Why not distinguish between two different time scales?
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- ullet A microscopic time t and a macroscopic time infinitely large with respect to t.
- Let $m := \sum xp(x) \in \mathbb{R}^d$, the expectation of $p(\cdot)$.
- Introducing time scale tL:

$$\lim_{L \to \infty} p_{tL}^L([uL]) = \rho_0(u - mt) =: \rho(t, u)$$

ullet Now the profile did change. We observe a new macroscopic profile: the original one translated by mt.



ullet Scaling limits o Hydrodynamic description

$$\partial_t \rho + m \cdot \nabla \rho = 0.$$



 $\bullet \ \, \mathsf{Scaling} \ \, \mathsf{limits} \to \mathsf{Hydrodynamic} \ \, \mathsf{description} \\$

$$\partial_t \rho + m \cdot \nabla \rho = 0.$$

- However, if the random walk is not asymmetric, then again the profile remains the same.
- Still, if we consider a larger time scale, times of order L^2 , even when m=0, we can observe an interesting time evolution.



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Model

- First of all, SEP allows at most one particle per site.
- State space: $\{0,1\}^{\mathbb{T}^d_L}$
- ullet $\eta^{x,y}$: configuration from η letting a particle jump from x to y

$$\eta^{x,y}(z) = \begin{cases} \eta(z) & \text{if } z \neq x, y, \\ \eta(x) - 1 & \text{if } z = x, \\ \eta(y) + 1 & \text{if } z = y \end{cases}$$



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$$\eta(x)(1-\eta(y))p(y-x)$$



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Again, the hydrodynamic description given by the heat equation

$$\partial_t p = (1/2)\Delta \rho$$



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Main quantity of interest: the position of the $m{\rm th}$ particle from the left at time t

$$x_m(t)$$
, with $x_m(0) = m$.



Marginal particle

Behavior of $|x_1|$, i.e. the distance that the marginal particle has covered on a given time t?



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Assume that

$$E(|x_1(t)|) \sim ct^a$$

which transforms to

$$\log E(|x_1(t)|) \sim a \log t + \log c.$$



Expected value

Performing experiments for $\gamma = 0.5$:

$$E(|x_1(t)|) \sim \frac{t}{2} = \gamma t$$



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It is true that, in general,

$$E(|x_1(t/\gamma)|) \sim t$$

for every $\gamma \in (0,1]$.



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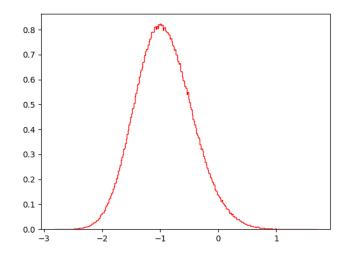
Therefore, what about

$$\frac{|x_1(t)| - \gamma t}{t^{0.6}}$$

?



Tracy-Widom distribution!







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So, distribution of the largest eigenvalue λ_{max} ?



Appearance of the distribution

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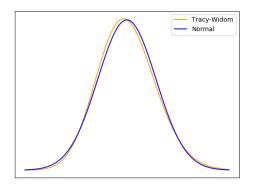
- Two decades later, in 1992, Tracy and Widom specified that distribution!
- But... in 1999, the same distribution was found in the length of the longest increasing subsequence of random permutations.
- It started to appear in models all over physics and mathematics.
- Especially in systems with interacting components.



 Universality: diverse microscopic effects → same collective behavior



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- Tracy-Widom complex cousin of the familiar bell curve...





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- Tracy-Widom from strongly correlated variables, such as interacting species, stock prices, matrix eigenvalues...
- Tracy-Widom universally proved to hold for certain classes of random matrices.
- Looser handle in counting problems, random walk problems, growth models...



Maybe an explanation?

ullet Asymmetric character of the distribution o some kind of universal phase transition



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 \bullet Asymmetric character of the distribution \to some kind of universal phase transition

In general,

- left tail: all components act in concert, (unstable)
- right tail: the components act alone. (stable)



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Model and Hydrodynamics

- ullet State space is $\mathbb{N}^{\mathbb{T}^d_L}$
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- \bullet Transition probability $p(\cdot,\cdot)$ on \mathbb{Z}^d



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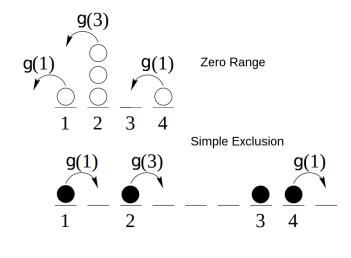
Non-linear (parabolic) hydrodynamic description

$$\partial_t \rho = (1/2)\Delta(\Phi(\rho))$$

where $\Phi(\cdot)$ is the inverse of $\rho(\cdot)$ until a critical density ρ_c .



Duality







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- Specifically, the jump rates will be given by

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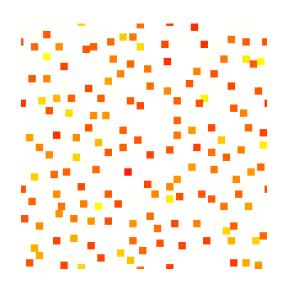
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• For b > 2, there is a critical density

$$\rho_c = \frac{1}{b-2}$$







Order of time needed to reach equilibrium?



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Performing experiments, for 3 initial states,

totally asymmetric zero range process:

$$T_{eq} = \mathcal{O}(L^2)$$

• symmetric zero range process:

$$T_{eq} = \mathcal{O}(L^3)$$

(diffusion without a drift)



Tagged Particles

Three classes \rightarrow a jump occurs from the tagged particle's site x,

- First class: the tagged will jump,
- Random: the tagged will jump with probability $1/\eta(x)$,
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Distance covered by the tagged particle in a given time: (space scale L, time scale L^2)

- First class: $X_{tag}(tL^{-2})/L = \mathcal{O}(1)$,
- Random: $X_{tag}(tL^{-2})/L = \mathcal{O}(\sqrt{t})$,
- Second class: $X_{tag}(tL^{-2})/L = \mathcal{O}(\sqrt{t})$.



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Thank you very much!

