# MATRIX CHERNOFF INEQUALITY AND CLASSICAL & QUANTUM CORRELATIONS

Stephen Diadamo

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Seminar: Applications of Matrix Techniques in Information and Signal Theory

## **OVERVIEW**

## Overview of the talk

- 1. Matrix Chernoff Inequality
- 2. Quantum and Classical Correlations
- 3. Motivating Example
- 4. Total Bipartite Correlations

Classical Chernoff Bounds: Let  $X_1, X_2, ..., X_n$  be random variables such that  $a \le X_k \le b$  for each k. Let  $Y := \sum_k X_k$  and define  $\mu := \mathbb{E} Y$ . Then for  $\epsilon > 0$ ,

$$\mathbb{P}(Y \ge (1 + \epsilon)\mu) \le \exp\left(\frac{2\epsilon^2 \mu^2}{n(b - a)^2}\right)$$

$$\mathbb{P}(Y \le (1 - \epsilon)\mu) \le \exp\left(\frac{\epsilon^2 \mu^2}{n(b - a)^2}\right)$$

# Matrix Chernoff Inequality

- · Let  $X_1, X_2, ..., X_n$  be a finite set of independent, random, positive-semidefinite matrices of dimension d
- · We have that for each k,  $\lambda_{min}(X_k) \geq 0$ , and  $X_k$  is Hermitian
- · Assume for each k,  $\lambda_{max}(X_k) \leq L$
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**Goal**: Study the expectation and tail behaviour of  $\lambda_{\min}(Y)$  and  $\lambda_{\max}(Y)$  Why?

- $\|Y\| \geq \lambda_{\max}(Y)$
- $\cdot \lambda_{\min}(Y)$  tell us if Y is singular

**Theorem**: Consider  $X_1, X_2, ..., X_n$  and Y defined previously, then for

$$\cdot \mu_{\min} := \lambda_{\min}(\mathbb{E}\mathsf{Y})$$

$$\cdot \ \mu_{\max} := \lambda_{\max}(\mathbb{E} Y)$$

It holds:

$$\quad \cdot \ \mathbb{P}\{\lambda_{\min}(\mathsf{Y}) \leq (1-\epsilon)\mu_{\min}\} \leq \mathsf{d} \left[ \tfrac{e^{-\epsilon}}{(1-\epsilon)^{1-\epsilon}} \right]^{\mu_{\min}/L} \ \mathsf{for} \ \epsilon \in [0,1)$$

$$\quad \cdot \ \mathbb{P}\{\lambda_{\max}(Y) \geq (1+\epsilon)\mu_{\max}\} \leq d \left[\frac{e^{\epsilon}}{(1+\epsilon)^{1+\epsilon}}\right]^{\mu_{\max}/L} \ \text{for } \epsilon \geq 0.$$

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**Lemma:** For  $X \in \{X_k\}$  as above and  $\theta \in \mathbb{R}$ ,  $\theta > 0$ , it holds:

· Mgf: 
$$\Phi_X(\theta) := \mathbb{E} e^{\theta X} \le \exp\left(\frac{e^{\theta L} - 1}{L} \cdot \mathbb{E} X\right)$$

· Cgf: 
$$\Xi_{X}(\theta) := \log \mathbb{E} e^{\theta X} \leq \frac{e^{\theta L} - 1}{L} \cdot \mathbb{E} X$$

**Remark**: For  $X_1,...,X_N$  iid random Hermitian matrices of dimension d, such that  $0 \le X_i \le \mathbb{I}$ ,  $\lambda_{min}(\mathbb{E}X_i) \ge \mu$ , we have for  $Y := \frac{1}{N} \sum X_i$ , and  $\epsilon \in [0,1]$ 

$$\mathbb{P}(\lambda_{\min}(\mathbf{Y}) \le (1 - \epsilon)\mu) \le \operatorname{d} \exp(-\operatorname{N}\epsilon^2 \mu/2)$$

$$\mathbb{P}(\lambda_{\max}(Y) \ge (1+\epsilon)\mu) \le d \exp(-N\epsilon^2\mu/2)$$

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- Classical Correlations: Correlations governed by laws of classical mechanics
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**Goal**: Find an operational meaning for total correlation of a bipartite quantum system

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- · Question: What is the total correlation of this state?
- Claim: This state contains two bits of correlation, 1 bit of quantum correlation (entanglement) and 1 bit of classical correlation

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- · From site A, 1 bit of randomness is applied, that is,  $\mathbb{I}$  or  $\sigma_z$  is applied with equal probability
- ·  $|\Phi^{+}\rangle \rightarrow \rho = \frac{1}{2} |\Phi^{+}\rangle\langle\Phi^{+}| + \frac{1}{2} |\Phi^{-}\rangle\langle\Phi^{-}|$

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# Step Two: Erase the classical correlation

- From site A, another 1 bit of randomness is applied, that is,  $\mathbb{I}$  or  $\sigma_x$  is applied with equal probability
- $\cdot \ \rho \to {\textstyle \frac{1}{2}} \mathbb{I}_{\mathsf{A}} \otimes {\textstyle \frac{1}{2}} \mathbb{I}_{\mathsf{B}}$

**Theorem**: The total correlation, measured by the amount of local erasure, for a bipartite state  $\rho_{AB}$  is the quantum mutual information

$$I(A:B) = S(\rho_A) + S(\rho_B) - S(\rho_{AB})$$

Let A and B be finite dimensional Hilbert spaces.

**Definition**: COLUR Maps

For  $N \in \mathbb{N}$ , let the ensemble  $\{p_i, U_i \otimes V_i\}_{i=1}^N$  be such that for each i,  $U_i \otimes V_i$  is a unitary on  $A \otimes B$ . A randomizing map  $R : \mathcal{S}(A \otimes B) \to \mathcal{S}(A \otimes B)$ ,

$$\mathsf{R}:\rho\mapsto\sum_{\mathsf{i}=1}^{\mathsf{N}}\mathsf{p}_{\mathsf{i}}(\mathsf{U}_{\mathsf{i}}\otimes\mathsf{V}_{\mathsf{i}})\rho(\mathsf{U}_{\mathsf{i}}\otimes\mathsf{V}_{\mathsf{i}})^{\dagger}$$

is called a coordinated local unitary randomizing (COLUR) map (COLUR) if it is CPTP. If for each i,  $\mathbf{U}_i = \mathbb{I}$ , then  $\mathbf{R}$  is called a B-LUR map. Similarly, if  $\mathbf{V}_i = \mathbb{I}$  for each i, then  $\mathbf{R}$  is called A-LUR.

**Definition**:  $\epsilon$ -decorrelates

A COLUR map **R** is said to  $\epsilon$ -decorrelate a state  $\rho \in \mathcal{S}(A \otimes B)$  if there is a product state  $\omega_A \otimes \omega_B \in \mathcal{S}(A \otimes B)$  such that,

$$\|\mathbf{R}(\rho) - \omega_{\mathsf{A}} \otimes \omega_{\mathsf{B}}\|_1 \leq \epsilon$$

where  $\|\cdot\|_1$  is the trace norm of an operator.

**Definition**: Entropy Exchange

For a COLUR map **R** and a purification of state  $\rho \in \mathcal{S}(\mathsf{A} \otimes \mathsf{B})$ ,  $\rho_\mathsf{P} = |\psi\rangle\langle\psi|_\mathsf{ZAB}$ , with reference Hilbert space Z, we define the entropy exchange as,

$$\mathsf{S}_{\mathsf{e}}(\mathsf{R}, \rho_{\mathsf{A}}) := \mathsf{S}((\mathbb{I}_{\mathsf{Z}} \otimes \mathsf{R}) \left| \psi \middle
angle \psi \middle
angle_{\mathsf{ZAB}})$$

**Proposition 1**: For any COLUR map  $\mathbf{R}$  that  $\epsilon$ -decorrelates the state  $\rho^{\otimes n} \in \mathcal{S}(\mathbf{A}^{\otimes n} \otimes \mathbf{B}^{\otimes n})$ , the entropy exchange of  $\mathbf{R}$  relative to  $\rho^{\otimes n}$  has the lower bound

$$S_e(R, \rho^{\otimes n}) \ge n[I(A : B) - O(\epsilon)],$$

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$$S_e(R, \rho^{\otimes n}) \ge n[I(A : B) - O(\epsilon)],$$

## Proof Idea:

- · Using the locality of R, separate  $\rho^{\otimes n}$  into two parts  $R_A := \operatorname{Tr}_B \mathbf{R}(\rho^{\otimes n})$  and  $R_B = \operatorname{Tr}_A \mathbf{R}(\rho^{\otimes n})$
- · Using concavity of von Neumann entropy and Fannes inequality,  $S(R_A) + S(R_B) S(R(\rho^{\otimes n})) \le O(\epsilon)$
- · Introduce a purification of  $\rho = \operatorname{Tr}_{\mathbb{Z}}(|\psi\rangle\langle\psi|)$ ,  $\psi := |\psi\rangle\langle\psi|$ , and conclude using Araki-Lieb (triangle) inequality  $S_{e}(\mathbf{R}, \rho^{\otimes n}) = S(\mathbb{I}_{\mathbb{Z}}^{\otimes n} \otimes \mathbf{R})(\psi^{\otimes n}) \geq n[I(A:B) O(\epsilon)]$

**Proposition 2**: For any state  $\rho \in \mathcal{S}(A \otimes B)$  and  $\epsilon > 0$ , there exists for sufficiently large  $n \in \mathbb{N}$ , an A-LUR map,

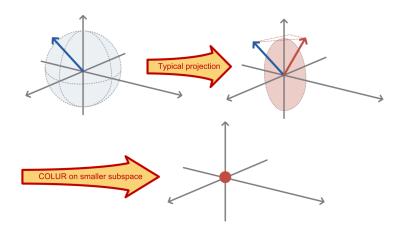
$$R: \mathcal{S}(A^{\otimes n} \otimes B^{\otimes n}) \to \mathcal{S}(A^{\otimes n} \otimes B^{\otimes n})\text{,}$$

$$\mathsf{R}: 
ho^{\otimes \mathsf{n}} \mapsto rac{1}{\mathsf{N}} \sum_{\mathsf{i}=\mathsf{1}}^{\mathsf{N}} (\mathsf{U}_{\mathsf{i}} \otimes \mathbb{I}) 
ho^{\otimes \mathsf{n}} (\mathsf{U}_{\mathsf{i}} \otimes \mathbb{I})^{\dagger}$$

which  $\epsilon$ -decorrelates  $\rho^{\otimes n}$  such that

$$\log N \le n[I(A:B) + O(\epsilon)]$$

## Proof Idea

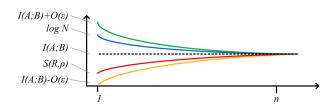


**Theorem** (restated more mathematically): The total correlations in a bipartite state  $\rho_{AB}$  measured in terms of how much noise needs to be added to turn  $\rho_{AB}$  into a product state is given by

$$\begin{split} \sup_{\epsilon>0} & \liminf_{n\to\infty} \frac{1}{n} \min\{S_e(R,\rho^{\otimes n}) \mid R, \ \epsilon\text{-decorr} \ \rho^{\otimes n}, \ \text{COLUR}\} \\ &= \sup_{\epsilon>0} \limsup_{n\to\infty} \frac{1}{n} \min\{\log N \mid R, \ \epsilon\text{-decorr} \ \rho^{\otimes n}, \ \text{A-LUR}\} \\ &= I(A:B) \end{split}$$

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## References

- 1. J. Tropp: An Introduction to Matrix Concentration Inequalities, Chapter 5
- http://www.markwilde.com/teaching/qinfomcgill/presentations/2-JanFlorjanczyk.pdf, Jan Florjanczyk
- 3. Quantum, classical, and total amount of correlations in a quantum state, Groisman, Popescu, Winter