

## Part-I 第一卷

Q1 題 1 (8 points 8 分) Solution 解:

For resistance  $R$ , only half of the area is conducting, as the other half is blocked by medium-2. Let the voltage between the plates be  $V$ , then the electric field

先求電阻  $R$ 。介質-2 不導電，所以只有一半的導電板導電。令兩板之間的電壓為  $V$ ，則電場為

$$E = V/d, (1 \text{ point } 1 \text{ 分})$$

The current density is 電流密度為  $J = \sigma_1 E$ , (1 point 1 分)

The current is 電流為  $I = \frac{1}{2} J a^2 = \frac{a^2}{2} \sigma_1 \frac{V}{d}$ . (1 point 1 分)

So the resistance is 因此電阻為  $R = \frac{V}{I} = \frac{2d}{\sigma_1 a^2}$ . (1 point 1 分)

For capacitance, it can be treated as two capacitors in parallel. The capacitance on the right is 再求電容。總電容可當作是左右兩個電容并聯。右邊的電容為

$$C_1 = \frac{\epsilon_0 \epsilon_1 a^2}{2d}. (1 \text{ point } 1 \text{ 分})$$

For the left half, let the electric displacement be  $D_2$  which is the same throughout the region. 在左半邊，令電位移為  $D_2$ ，電位移處處相同。 (1 point 1 分)

The total free charge on the left half is

左半邊的總自由電荷為

$$Q = \frac{a^2}{2} D_2. (1 \text{ point } 1 \text{ 分})$$

The total voltage between the two plates is

兩導電板之間的總電壓為

$$V = \frac{1}{2} E_1 d + \frac{1}{2} E_2 d = \frac{d}{2\epsilon_0} D_2 \left( \frac{1}{\epsilon_1} + \frac{1}{\epsilon_2} \right). (1 \text{ point } 1 \text{ 分})$$

So 因此  $C_2 = \frac{\varepsilon_0 a^2}{d} \frac{\varepsilon_1 \varepsilon_2}{\varepsilon_1 + \varepsilon_2},$

$$C = C_1 + C_2 = \frac{\varepsilon_0 a^2}{d} \left( \frac{\varepsilon_1}{2} + \frac{\varepsilon_1 \varepsilon_2}{\varepsilon_1 + \varepsilon_2} \right). \text{ (1 point 1 分)}$$

Q2 題 2 (10 points 10 分) solution 解:

(a) Let the mass of the blackhole be  $M$ , then

令黑洞的質量為  $M$ ，則

$$\frac{GM}{r^2} = \frac{v^2}{r}, (2 \text{ points } 2 \text{ 分})$$

So 因此  $\sqrt{\frac{GM}{r}} = v$ , (1 point 1 分)

$$n = -\frac{1}{2}. (1 \text{ point } 1 \text{ 分})$$

(b) Let the areal mass density be  $\sigma$ , then

令質量面密度為  $\sigma$ ，則

$$\frac{G\pi\sigma r^2}{r^2} = \frac{v^2}{r}, (2 \text{ points } 2 \text{ 分})$$

So 因此  $\sqrt{G\pi\sigma r} = v$ , (1 point 1 分)

$$n = \frac{1}{2}. (1 \text{ point } 1 \text{ 分})$$

(c) Anything reasonable is fine. It DOES NOT have to be dark matter.

任何有一定理由的解釋都行，不一定非暗物質不可。(2 points 2 分)

Q3 題 3 (10 points 10 分) Solution 解:

Method 1: Using conservation of energy

方法-1：利用能量守恒

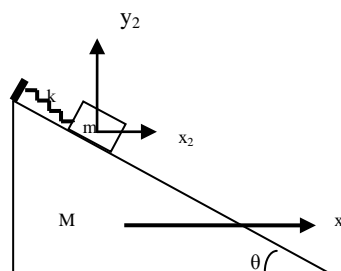
Equations 方程:

$$Mx_1 = -mx_2, \quad y_2 = (x_2 - x_1) \tan \theta = (1 + \frac{m}{M})x_2 \tan \theta$$

(1 point 1 分)

All coordinates are in the rest frame. 所有坐標取

自靜止參照系。



The total kinetic energy of the system is

總動能為

$$T = \frac{1}{2}M\dot{x}_1^2 + \frac{1}{2}m\dot{x}_2^2 + \frac{1}{2}m\dot{y}_2^2 \quad (1 \text{ point } 1 \text{ 分})$$

$$= \frac{1}{2}M \frac{m^2}{M^2} \dot{x}_2^2 + \frac{1}{2}m\dot{x}_2^2 + \frac{1}{2}m \tan^2 \theta \cdot \dot{x}_2^2 \left(1 + \frac{m}{M}\right)^2 \quad (1 \text{ point } 1 \text{ 分})$$

The total potential energy 總勢能為:

$$V = \frac{1}{2}k[(x_2 - x_1)^2 + y_2^2] - mgy_2 \quad (1 \text{ point } 1 \text{ 分})$$

$$= \frac{1}{2}k\left[\left(1 + \frac{m}{M}\right)^2 x_2^2 + \left(1 + \frac{m}{M}\right)^2 x_2^2 \tan^2 \theta\right] - mg\left(1 + \frac{m}{M}\right)x_2 \tan \theta$$

$$= \frac{1}{2}kx_2^2 \left(1 + \frac{m}{M}\right)^2 \cdot \frac{1}{\cos^2 \theta} - mg\left(1 + \frac{m}{M}\right)x_2 \tan \theta \quad (1 \text{ point } 1 \text{ 分})$$

Using 利用  $ax^2 - bx = a\left(x - \frac{b}{2a}\right)^2 - \frac{b^2}{4a}$ , we get 得

$$V = \frac{1}{2}k\left(1 + \frac{m}{M}\right)^2 \cdot \frac{1}{\cos^2 \theta} \left(x_2 - \frac{mMg \sin 2\theta}{k(M+m)}\right)^2 - \frac{(mg \sin \theta)^2}{2k} \quad (1 \text{ point } 1 \text{ 分})$$

Make a transfer of coordinate 取坐標變換

$$x \equiv x_2 - \frac{mMg \sin 2\theta}{k(M+m)}, \quad (1 \text{ point } 1 \text{ 分})$$

We reach an expression for the total energy that is of the form of  $T + V = ax^2 + b\dot{x}^2 + c$  which must be constant by energy conservation. Let  $x = A \cos(\omega t + \phi)$ , we get

$$T + V = aA^2 \cos^2(\omega t + \phi) + bA^2 \omega^2 \sin^2(\omega t + \phi) + c = A^2(a - b\omega^2) \cos^2(\omega t + \phi) + bA^2 \omega^2 + c.$$

For the total energy to be constant the  $\cos^2(\omega t + \phi)$  term must be zero all the time, which leads to  $\omega^2 = a/b$ .

我們得到總能量的表達式為  $T + V = ax^2 + b\dot{x}^2 + c$ 。因能量守恒總能量應為常數。令

$x = A \cos(\omega t + \phi)$ ，得

$$T + V = aA^2 \cos^2(\omega t + \phi) + bA^2 \omega^2 \sin^2(\omega t + \phi) + c = A^2(a - b\omega^2) \cos^2(\omega t + \phi) + bA^2 \omega^2 + c$$

作為常數，上式中  $\cos^2(\omega t + \phi)$  項必須為零。因此  $\omega^2 = a/b$ 。

Therefore, the oscillation frequency  $\omega$  in this case is

由上式得系統的頻率 $\omega$ 為

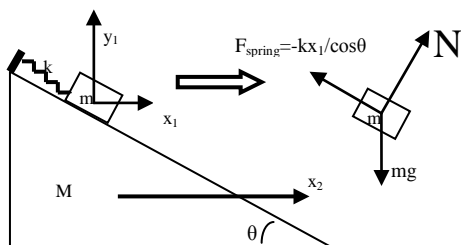
$$\omega^2 = \frac{k(1 + \frac{m}{M})^2 / \cos^2 \theta}{\frac{m}{M^2} + m + m(1 + \frac{m}{M}) \tan^2 \theta} = \frac{k}{m} \frac{1 + \frac{m}{M}}{\cos^2 \theta + (1 + \frac{m}{M}) \sin^2 \theta} = \frac{k}{m} \left( \frac{M + m}{M + m \sin^2 \theta} \right). \quad (1 \text{ point } 1 \text{ 分})$$

分)

For  $\theta = 0$ , we have 當  $\theta = 0$ , 得  $\omega^2 = k \left( \frac{1}{m} + \frac{1}{M} \right)$  (1 point 1 分),

For  $\theta = 90^\circ$  we have 當  $\theta = 90^\circ$ , 得  $\omega^2 = \frac{k}{m}$ . (1 point 1 分)

Method 2: Analytical Mechanics 方法-2：分析力學



Force figure 力圖 (2 points 2 分)

Equations 方程

$$m \ddot{x}_1 = -\frac{kx_1}{\cos \theta} \cos \theta - N \sin \theta + m \ddot{x}_2,$$

$x_1$  is in the frame on the slope.  $x_1$  是相對與斜面的橫坐標。 (1 point 1 分)

$$y_1 = x_1 \tan \theta$$

$$m \ddot{x}_1 \tan \theta = N \cos \theta - mg - kx_1 \tan \theta \quad (1 \text{ point } 1 \text{ 分})$$

$$M \ddot{x}_2 = -(N \sin \theta + kx_1) \quad (1 \text{ point } 1 \text{ 分})$$

Process 解方程過程

Step 1: Eliminate  $\ddot{x}_2$  步驟-1：消去  $\ddot{x}_2$

$$m \ddot{x}_1 = -kx_1 - N \sin \theta - \frac{m}{M} (N \sin \theta + kx_1)$$

$$m \ddot{x}_1 \tan \theta = N \cos \theta - mg - kx_1 \tan \theta$$

Step 2: Eliminate  $N$  步驟-2: 消去  $N$

$$m \ddot{x}_1 + (1 + \frac{m}{M})kx_1 = -N \sin \theta (1 + \frac{m}{M})$$

$$m \ddot{x}_1 \tan \theta + mg + kx_1 \tan \theta = N \cos \theta$$

Which means 整理後得

$$m \cos \theta \ddot{x}_1 + \cos \theta (1 + \frac{m}{M})kx_1 = -\sin \theta (1 + \frac{m}{M})[m \tan \theta \ddot{x}_1 + mg + kx_1 \tan \theta] \quad (2 \text{ points } 2 \text{ 分})$$

By assuming  $x_1 = Ae^{i\omega t}$ , the oscillation frequency  $\omega$  is obtained

設解  $x_1 = Ae^{i\omega t}$ , 得頻率  $\omega$

$$\omega^2 = \frac{k(1 + \frac{m}{M})}{m} \frac{\cos \theta + \frac{\sin^2 \theta}{\cos \theta}}{\cos \theta + (1 + \frac{m}{M}) \frac{\sin^2 \theta}{\cos \theta}} = \frac{k}{m} \left( \frac{M + m}{M + m \sin^2 \theta} \right). \quad (1 \text{ point } 1 \text{ 分})$$

For  $\theta = 0$ , we have 當  $\theta = 0$ , 得  $\omega^2 = k \left( \frac{1}{m} + \frac{1}{M} \right)$  (1 point 1 分),

For  $\theta = 90^\circ$  we have 當  $\theta = 90^\circ$ , 得  $\omega^2 = \frac{k}{m}$ . (1 point 1 分)

Q4 題 4 (10 points 10 分) Solution 解:

(a) Let the magnetic field be  $\vec{B}(z, t) = \vec{B}_0 e^{i(\tilde{k}z - \omega t)}$ . The k-vector of the wave is  $\vec{k} = \tilde{k}\vec{z}_0$ .

Using the equation  $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$ ,

令磁場的表達式為  $\vec{B}(z, t) = \vec{B}_0 e^{i(\tilde{k}z - \omega t)}$ 。k-矢量為  $\vec{k} = \tilde{k}\vec{z}_0$ 。利用方程  $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

we get 得  $\vec{B}(z, t) = \frac{\vec{k} \times \vec{E}(z, t)}{\omega}$ , (1 point 1 分)

$$= \frac{\tilde{k}}{\omega} E_0 (\vec{z}_0 \times \vec{x}_0) e^{i(\tilde{k}z - \omega t)} = \frac{\tilde{k}}{\omega} E_0 \vec{y}_0 e^{i(\tilde{k}z - \omega t)}$$

So 因此  $\vec{B}_0 = \frac{\tilde{k}}{\omega} E_0 \vec{y}_0$ . (1 point 1 分)

(b) Note that 利用  $\tilde{k} = \frac{\omega}{c} \sqrt{\varepsilon} + i \frac{\sigma}{2} \sqrt{\frac{\mu_0}{\varepsilon \varepsilon_0}} \equiv k_R + i k_I$ .

So 得  $\langle \vec{S} \rangle = \frac{1}{\mu_0} \langle (\vec{E} \times \vec{B}) \rangle = \frac{1}{2\mu_0} \text{Re}(\vec{E} \times \vec{B}^*)$  (1 point 1 分)

$$= \frac{E_0^2}{2\mu_0 \omega} \vec{z}_0 \text{Re}(k_R - i k_I) e^{-2k_I z} = \frac{1}{2\mu_0} \frac{k_R}{\omega} E_0^2 \vec{z}_0 e^{-2k_I z} = \frac{1}{2\mu_0 c} E_0^2 \vec{z}_0 e^{-2k_I z} \quad (1 \text{ point } 1 \text{ 分})$$

$$(c) \quad q = -\frac{d \langle S \rangle}{dz} = \frac{k_I k_R}{\mu_0 \omega} E_0^2 e^{-2k_I z} = \frac{\sigma}{2} \frac{\frac{\omega}{c} \sqrt{\varepsilon} \sqrt{\frac{\mu_0}{\varepsilon \varepsilon_0}}}{\mu_0 \omega} E_0^2 e^{-2k_I z} = \frac{\sigma}{2} E_0^2 e^{-2k_I z}. \quad (2 \text{ points } 2 \text{ 分})$$

$$(d) \quad \text{Joule Heat 焦爾熱} = \langle \vec{J} \cdot \vec{E} \rangle = \frac{1}{2} \text{Re}(\vec{J} \cdot \vec{E}^*) = \frac{\sigma}{2} \text{Re}(\vec{E} \cdot \vec{E}^*) = \frac{\sigma}{2} E_0^2 e^{-2k_I z} \quad (2 \text{ points } 2 \text{ 分})$$

(e) The energy loss of EM wave is equal to the Joule Heat.

電磁波能量的損失等于焦爾熱。 (2 points 2 分)

Q5 (12 points) 題 5 (12 分) Solution 解：

(a) First law 熱力學第一定律:  $dU = -PdV + \delta Q$ .

The equation of adiabatic processes is 絕熱過程的方程式為

$$\delta Q = dU + PdV = d(3PV) + PdV = 4PdV + 3VdP = 0 \quad (1 \text{ point } 1 \text{ 分})$$

$$\Rightarrow PV^{4/3} = \text{Constant}$$

$$\Rightarrow PV^{4/3} = \text{常數} \quad (1 \text{ point } 1 \text{ 分})$$

(b) In the Carnot cycle 在卡諾循環過程中:

$$(P_1, V_1) \xrightarrow{\text{Isothermal}} (P_1, V_2) \xrightarrow{\text{Adiabatic}} (P_2, V_3) \xrightarrow{\text{Isothermal}} (P_2, V_4) \xrightarrow{\text{Adiabatic}} (P_1, V_1).$$

Isothermal = 等溫 ; Adiabatic = 絕熱

The heat supplied to the gas during the first isothermal process is

第一個等溫過程中吸收的熱量為

$$Q_1 = 3P_1V_2 - 3P_1V_1 + \int_{V_1}^{V_2} P_1 dV = 4P_1(V_2 - V_1). \quad (1 \text{ point } 1 \text{ 分})$$

(c) Similarly, the heat supplied to the gas during the second isothermal process is

第二個等溫過程中吸收的熱量為

$$Q_2 = 4P_2(V_4 - V_3). \quad (1 \text{ point } 1 \text{ 分})$$

(d) From (a) we have 由(a)得:

$$\begin{aligned} P_1V_2^{4/3} &= P_2V_3^{4/3} \\ P_2V_4^{4/3} &= P_1V_1^{4/3} \end{aligned} \quad (1 \text{ point } 1 \text{ 分})$$

By definition 由定義

$$\frac{T_1}{T_2} = -\frac{Q_1}{Q_2} = -\frac{P_1(V_2 - V_1)}{P_2(V_4 - V_3)} = -\frac{P_1^{1/4}P_2^{3/4}(V_3 - V_4)}{P_2(V_4 - V_3)} = \frac{P_1^{1/4}}{P_2^{1/4}}. \quad (2 \text{ points } 2 \text{ 分})$$

Therefore, one may define the absolute temperature by  $T = AP^{1/4}$ , where  $A$  is an arbitrary constant. Since  $T = 1$  when  $P = 1$ ,  $T = P^{1/4}$ .

因此我們可以定義溫度為  $T = AP^{1/4}$ , 其中  $A$  為任意常數。由  $P = 1$  時  $T = 1$ , 得  $T = P^{1/4}$ 。

(1 point 1 分)

(e) The internal energy is then 內能為  $U = 3T^4V$ . (1 point 1 分)

Hence the heat capacity is 因此熱容量為  $C_V = \left( \frac{\partial U}{\partial T} \right)_V = 12T^3V$ . (1 point 1 分)

(f) The entropy is 熵為  $S = \int_0^T C_V \frac{dT}{T} = 12V \int_0^T T^2 dT = 4T^3V = 4P^{3/4}V$ . (2 points 2 分)



## Part-II 第二卷

Q1 (16 points) 題 1 (16 分) Solution 解：

- (a) Consider a thin layer of gas of unit area and thickness  $dr$ . The pressure at  $r$  should be larger than the pressure at  $r + dr$  in order to balance the gravity  $\frac{GMmn}{r^2}$ . So we have  $\frac{dP}{dr} = -\frac{GM_s mn}{r^2}$ , where  $n$  is the molecular number density of the gas.

考慮離太陽  $r$  處一厚度為  $dr$  的單位面積氣體。在  $r$  處的氣壓應比在  $r + dr$  處的大一點，從而平衡太陽的引力  $\frac{GMmn}{r^2} dr$ 。因此有  $\frac{dP}{dr} = -\frac{GM_s mn}{r^2}$ ，其中  $n$  為氣體的分子數密度。(1 point 1 分)

We also have the ideal gas law  $P = nkT_0$ . Replace  $P$  with  $n$  we get  $\frac{dn}{n} = -\frac{GM_s m}{kT_0} \frac{dr}{r^2}$ .

另有理想氣體方程  $P = nkT_0$ 。將  $P$  用  $n$  代入，得  $\frac{dn}{n} = -\frac{GM_s m}{kT_0} \frac{dr}{r^2}$ 。(1 point 1 分)

Finally,  $\rho = \rho_0 e^{\alpha/r}$ , where  $\alpha = \frac{GM_s m}{kT_0}$ .

最後得  $\rho = \rho_0 e^{\alpha/r}$ ，其中  $\alpha = \frac{GM_s m}{kT_0}$ 。(1 point 1 分)

- (b) When  $r \rightarrow \infty$ ,  $\rho \rightarrow \rho_0$  instead of zero. That means the gas ball is infinitely large, which is unphysical.

當  $r \rightarrow \infty$ ， $\rho \rightarrow \rho_0$  而不是零。這意味著氣體球是無限大的，不符合實際情況。(2 points 2 分)

- (c) The amount of energy per second through any concentric sphere shells should be constant. That is,  $J_0 = 4\pi r^2 I$ .

每秒鐘穿過任意一個同心圓殼的能量應為常數。所以  $J_0 = 4\pi r^2 I$ 。(1 point 1 分)

Then 從而得  $\frac{J_0}{4\pi r^2} = I$ 。(2 points 2 分)

- (d)  $I(r) = -\sigma \frac{dT}{dr}$ , so 因此  $\frac{dT}{dr} = -\frac{J_0}{4\pi\sigma r^2}$ 。(1 point 1 分)

Then 由此得  $T = \frac{J_0}{4\pi\sigma r}$ .

The integral constant should be zero, as  $T$  should be zero at large distance.

因為在無窮遠處溫度須為零，所以由積分產生的常數必須為零。(2 points 2 分)

(e) Again 重力和壓力平衡,  $\frac{dP}{dr} = -\frac{GM_s mn}{r^2}$ .

But now  $P = nkT(r) = \frac{J_0}{4\pi\sigma r} kn$ . Replace  $n$  by  $P$  in the first equation, we

$$\text{get } \frac{dP}{dr} = -\frac{4\pi GM_s m\sigma P}{kJ_0 r},$$

現在  $P = nkT(r) = \frac{J_0}{4\pi\sigma r} kn$ 。將  $P$  代入  $n$ ，得  $\frac{dP}{dr} = -\frac{4\pi GM_s m\sigma P}{kJ_0 r}$  (1 point 1 分)

which leads to  $P = P_0 \left(\frac{r}{r_0}\right)^{-\beta}$ , where  $\beta = \frac{4\pi GM_s m\sigma}{kJ_0}$ .

從而得  $P = P_0 \left(\frac{r}{r_0}\right)^{-\beta}$ ，其中  $\beta = \frac{4\pi GM_s m\sigma}{kJ_0}$ 。(1 point 1 分)

$$\rho = \frac{4\pi\sigma mr P_0}{kJ_0} \left(\frac{r}{r_0}\right)^{-\beta}. \quad (1 \text{ point } 1 \text{ 分})$$

This time  $P$  and  $\rho$  go to zero at large  $r$ . 現在的  $P$  和  $\rho$  在無窮遠處為零。

(f) From the surface temperatures of the planets we know today we estimate that  $r_0$  is about the radius of the orbit of Mars.

由現在各行星的表面溫度我們可以推測大概和火星的軌道相近。(2 points 2 分)

Q2 (16 points) 題 2 (16 分) Solution 解：

(a) The number of electrons crossing the junction per second is  $I/e$ .

每秒鐘通過界面的電子數為  $I/e$ 。(1 point 1 分)

On average, there are  $(\alpha - 0.5)I/e$  electrons flip their spins.

平均有  $(\alpha - 0.5)I/e$  的電子的自旋反轉。(1 point 1 分)

The net angular momentum change per second is then  $(\alpha - 0.5)\frac{I\hbar}{4\pi e} \times 2$ .

每秒鐘角動量的淨變化為  $(\alpha - 0.5) \frac{Ih}{4\pi e} \times 2$  。 (1 point 1 分)

This is equal to the torque so  $\tau = (\alpha - 0.5) \frac{h}{2\pi e} I$  .

角動量的淨變化等于力矩  $\tau = (\alpha - 0.5) \frac{h}{2\pi e} I$  。 (1 point 1 分)

(b) The equation of motion is  $J \frac{d^2\theta}{dt^2} = -\eta \frac{d\theta}{dt} - \kappa\theta$  .

導線扭擺的運動方程為  $J \frac{d^2\theta}{dt^2} = -\eta \frac{d\theta}{dt} - \kappa\theta$  。 (1 point 1 分)

Let  $\theta(t) = \theta_0 e^{i\tilde{\omega}t}$  , where  $\tilde{\omega}$  is complex,

令  $\theta(t) = \theta_0 e^{i\tilde{\omega}t}$  , 其中  $\tilde{\omega}$  是複數 , we get 得 (1 point 1 分)

$$\tilde{\omega}^2 - i\gamma\tilde{\omega} - \omega_0^2 = 0, \quad \text{where 其中 } \omega_0^2 \equiv \kappa/J, \quad \gamma \equiv \eta/J. \text{ (1 point 1 分)}$$

Let  $\tilde{\omega} = \omega_R + i\omega_I$  and solving the equation, we get  $\theta = \theta_0 e^{-\omega_I t} e^{i\omega_R t}$  ,

令  $\tilde{\omega} = \omega_R + i\omega_I$  , 并解上述方程 , 得  $\theta = \theta_0 e^{-\omega_I t} e^{i\omega_R t}$  , (1 point 1 分)

Where 其中  $\omega_I = \gamma/2$  , (1 point 1 分) ,  $\omega_R = \sqrt{\omega_0^2 + \gamma^2/4}$  . (1 point 1 分)

(c) This is a forced oscillation with the force given by  $\tau(t) = (\alpha - 0.5) \frac{h}{2e} I_0 e^{i\omega t} = \tau_0 e^{i\omega t}$  .

這是個受迫振動問題 , 驅動力矩為  $\tau(t) = (\alpha - 0.5) \frac{h}{2e} I_0 e^{i\omega t} = \tau_0 e^{i\omega t}$  。 (1 point 1 分)

The equation of motion is 運動方程為  $J \frac{d^2\theta}{dt^2} = -\eta \frac{d\theta}{dt} - \kappa\theta + \tau(t)$  . (1 point 1 分)

Let 令  $\theta(t) = \theta_0 e^{i\omega t}$  , (1 point 1 分)

we get the oscillation amplitude  $\theta_0 = \frac{\tau_0 / J}{\omega_0^2 - \omega^2 + i\gamma\omega}$  .

得振動幅度為  $\theta_0 = \frac{\tau_0 / J}{\omega_0^2 - \omega^2 + i\gamma\omega}$  。 (1 point 1 分)

The speed of the side wing is 邊翼的速度為  $v(t) = i\omega d\theta_0 e^{i\omega t}$  , (1 point 1 分)

and the electromotive potential is 電動勢為  $\xi(t) = BLv(t) = i\omega dBL\theta_0 e^{i\omega t}$  。 (1 point 1 分)

Q3 (18 points) 題 3 (18 分) Solution 解：

- (a) The magnetic dipole experiences a torque  $\vec{\tau} = \vec{m} \times \vec{B}_0$  which is always perpendicular to the  $\vec{S} \sim \vec{z}_0$  plane. The torque will turn the direction of  $\vec{S}$  so  $\vec{S}$  rotates around  $\vec{B}_0$  at constant angular speed.

磁偶極子受到的力矩為  $\vec{\tau} = \vec{m} \times \vec{B}_0$ ，其方向始終與  $\vec{S} \sim \vec{z}_0$  平面垂直。力矩改變  $\vec{S}$  的方向，因此  $\vec{S}$  繞著  $\vec{B}_0$  以勻角速度旋轉。(1 point 1 分)

Let the angle between  $\vec{S}$  and  $\vec{z}_0$  be  $\theta$ . The torque is  $mB_0 \sin \theta = \mu SB_0 \sin \theta$ , while the change of angular momentum over time  $\delta t$  is  $\delta S = S \sin \theta \delta \phi$ .

令  $\vec{S}$  與  $\vec{z}_0$  之間的夾角為  $\theta$ 。則力矩的大小為  $mB_0 \sin \theta = \mu SB_0 \sin \theta$ ，而角動量的變化為  $\delta S = S \sin \theta \delta \phi$ 。(1 point 1 分)

Since 既然  $S \sin \theta \delta \phi = \delta S = \mu SB_0 \sin \theta \delta t$  (1 point 1 分)

We have 我們得  $\omega_0 = \frac{\delta \phi}{\delta t} = \mu B_0$ . (1 point 1 分)

- (b) In the reference frame rotating at angular velocity  $\omega_0 \vec{z}_0 = -\mu B_0 \vec{z}_0$ , the spin appears stationary.

在以角速度  $\omega_0 \vec{z}_0 = -\mu B_0 \vec{z}_0$  旋轉的參照系裏，自旋是不動的。(2 points 2 分)

- (c) The effective B-field is 有效磁場為  $\vec{B}_\omega = -\frac{\vec{\omega}}{\mu}$ . (2 points 2 分)

- (d) In the rotating frame of  $-\omega_1 \vec{z}_0$ ,  $\vec{B}_1$  also appears static.

在以角速度  $-\omega_1 \vec{z}_0$  旋轉的參照系裏， $\vec{B}_1$  是不動的。(1 point 1 分)

Let it be along the X' axis in the rotating frame, the total B-field is

$$\vec{B} = (B_0 - \frac{\omega_1}{\mu}) \vec{z}_0' + B_1 \vec{x}_0'$$

令  $\vec{B}_1$  在旋轉參照系裏沿 X' 方向，則總磁場為  $\vec{B} = (B_0 - \frac{\omega_1}{\mu}) \vec{z}_0' + B_1 \vec{x}_0'$ 。(2 points 2

分)

- (e) In this case, only  $B_1 \vec{x}_0'$  remains. The spin will rotate around the  $\vec{x}_0'$  axis at angular speed  $\omega_1 = \mu B_1$ ,

這時的磁場只剩下  $B_1 \vec{x}_0'$ 。自旋繞其以角速度  $\omega_1 = \mu B_1$  旋轉， (2 points 2 分)

and the time to flip the spin is 倒轉自旋所需的時間為  $t = \frac{1}{2} \frac{2\pi}{\omega_1} = \frac{\pi}{\mu B_1}$ . (2 points 2 分)

- (f) In this case the spin will rotate around the total B-field given by (c) at angular

speed  $\omega = \mu \sqrt{\left(B_0 - \frac{\omega}{\mu}\right)^2 + B_1^2}$ .

這時的磁場由(c)給出。自旋的角速度為  $\omega = \mu \sqrt{\left(B_0 - \frac{\omega}{\mu}\right)^2 + B_1^2}$  (3 points 3 分)

~~~~~ End 完 ~~~~~