Neural Networks as Quantum States

NNs in Quantum Many-Body Problems

Deep Learning con Applicazioni

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- Quantum-many-body problems feature an exponential Hilbert-space growth
- Traditional variational ansätze reduce complexity sacrificing correlation and entanglement
- Neural-network quantum states can offer expressivity with a compact set of parameters

Neural Networks can represent ground states for a many-body quantum system¹

¹Giuseppe Carleo, Matthias Troyer, Solving the quantum many-body problem with artificial neural networks. *Science* 355, 602-606 (2017)

Project Outline Find the Ground State

- Restricted Boltzmann Machine as a neural network
- Sample configurations from RBM
- Train RBM to find the ground state

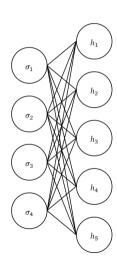




Restricted Boltzmann Machine

Generative Model

- Learns $p(\vec{\sigma})$ of the input data $\vec{\sigma}$
- Two layers: *N* visible and *M* hidden
- Parameters W: weights W_{ij} and biases b_i , c_j





Restricted Boltzmann Machine

Probability Distribution

Internal Energy:

$$E(\vec{\sigma}, \vec{h}) = -\sum_{i} b_{i}\sigma_{i} - \sum_{j} c_{j}h_{j} - \sum_{i,j} W_{ij}\sigma_{i}h_{j}$$

Probability Distribution:

$$p(\vec{\sigma}, \vec{h}) = \frac{1}{Z} e^{-E(\vec{\sigma}, \vec{h})}$$

We can obtain the wave function as unnormalized probability.



Restricted Boltzmann Machine

Application to Quantum States

Considering a spin configuration $\vec{\sigma}$ and a set of parameters \mathcal{W} :

$$\psi\left(\vec{\sigma},\mathcal{W}\right)=e^{\sum_{i}b_{i}\sigma_{i}+\sum_{j}c_{j}h_{j}+\sum_{i,j}W_{ij}\sigma_{i}h_{j}}$$

With no intra-layer connections:

$$\psi\left(ec{\sigma},\mathcal{W}
ight)=\mathrm{e}^{\sum_{i}b_{i}\sigma_{i}} imes\prod_{j=1}^{M}2cosh\left[c_{j}+\sum_{i}W_{ij}\sigma_{i}
ight]$$

From the RBM we can sample configurations $\vec{\sigma}$.



Update parameters W towards energy minimum.

$$E_0 \leq \frac{\langle \psi_{\mathcal{W}} | \hat{H} | \psi_{\mathcal{W}} \rangle}{\langle \psi_{\mathcal{W}} | \psi_{\mathcal{W}} \rangle}$$

The parameters W are optimized using a gradient descent method:

- Variational Monte Carlo (VMC)²
- Stochastic Reconfiguration (SR)³

²Moritz Reh, Markus Schmitt, Martin Gärttner, Optimizing design choices for neural quantum states, *Phys. Rev. B* 107, 195115 (2023)

³Becca F, Sorella S. *Quantum Monte Carlo Approaches for Correlated Systems*. Cambridge University Press; 2017



We can compute the gradients:

$$\begin{split} \nabla_{\mathcal{W}}\left(\mathcal{E}\right) &= \nabla_{\mathcal{W}} \langle \mathcal{E}_{\text{loc}} \rangle \\ &= \langle \mathcal{E}_{\text{loc}} \nabla_{\mathcal{W}} log(p_{\psi}) \rangle \\ &= 2 \Re \left[\langle \mathcal{E}_{\text{loc}} \nabla_{\mathcal{W}} log(\psi) \rangle - \langle \mathcal{E}_{\text{loc}} \rangle \langle \nabla_{\mathcal{W}} log(\psi) \rangle \right] \end{split}$$

And consequently update the parameters:

$$\Delta \mathcal{W} = -\eta \nabla_{\mathcal{W}} \left(E \right)$$



Parameters evolution in the variational space:

$$\Delta \mathcal{W} = -\eta \mathbf{S}^{-1} \vec{F}$$

where S^{-1} is the pseudo-inverse of the covariance matrix:

$$\mathbf{S}_{ij} = \langle O_i^* O_j \rangle - \langle O_i^* \rangle \langle O_j \rangle$$

and \vec{F} is the force vector:

$$F_i = \langle O_i^* E_{loc} \rangle - \langle O_i^* \rangle \langle E_{loc} \rangle$$

where O_i are the local operators defined as:

$$O_i = rac{\partial log\left(\psi\left(ec{\sigma},\mathcal{W}
ight)
ight)}{\partial W_{ij}}$$

- RBM
- Sampler
- Hamiltonian
- Optimizers





Code Structure

nngs/

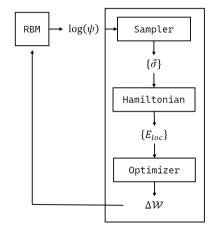
- nnqs.py

-> RBM

- hamiltonian.py -> Ising1D

- optimizer.py -> MRT2, GibbsSampler

sampler.py -> VMC, SR





Custom tf.module implementing the RBM.

Weights and biases are initialized randomly as complex values:

```
self.W = tf.Variable(
    tf.random.normal([self.N_visible, self.N_hidden], dtype=tf.complex64)
)
self.b = tf.Variable(tf.random.normal([self.N_visible], dtype=tf.complex64))
self.c = tf.Variable(tf.random.normal([self.N_hidden], dtype=tf.complex64))
```

$$\psi\left(ec{\sigma},\mathcal{W}
ight)= extbf{e}^{\sum_{i}b_{i}\sigma_{i}} imes\prod_{j=1}^{M}2cosh\left[c_{j}+\sum_{i}\textit{W}_{ij}\sigma_{i}
ight]$$



RBM - Logarithm of the Wave Function

$$\vec{\sigma} = (\sigma_1, \sigma_2, \sigma_3, ..., \sigma_N)$$
 NN log(ψ)

$$log(\psi) = \sum_{i} b_{i}\sigma_{i} + \sum_{j=1}^{M} log\left[2cosh\left(c_{j} + \sum_{i} W_{ij}\sigma_{i}\right)\right]$$



RBM - Logarithm of the Wave Function

```
sum_visible = tf.reduce_sum(
    a * spins, axis=1
)
w_h = b + tf.matmul(spins, W)
sum_hidden = tf.reduce_sum(
    tf.math.log(2.0 * (
        tf.math.cosh(b + tf.matmul(spins, W))
    )),axis=1
)
return sum_visible + sum_hidden
```

$$\sum_{i} b_{i}\sigma_{i}$$
 $\sum_{j=1}^{M} log \left[2cosh \left(c_{j} + \sum_{i} W_{ij}\sigma_{i}
ight)
ight]$



Sampler

Samples configurations from the RBM.

- Metropolis-Hastings MCMC using TensorFlow Probability
- Gibbs sampling Double-step method to sample visible and hidden variables⁴

Provides batches of configurations. Example:

⁴Francesco D'Angelo, Lucas Böttcher, Learning the Ising model with generative neural networks, *Phys. Rev. Research* 2, 023266 (2020)



Hamiltonian - 1D Ising

Simple 1D Ising Hamiltonian:

$$H = -J\sum_{i}\sigma_{i}\sigma_{i+1} - h\sum_{i}\sigma_{i}$$

Provides local_energy(samples) which computes E_{loc} for each configuration $\vec{\sigma}$.



Optimizers

Leverage <code>tf.GradientTape</code> to compute gradients of $\log\left(\psi\right)$

```
with tf.GradientTape() as tape:
    log_psi = self.wave_function.log_psi(samples)
grad_log_psi = tape.jacobian(log_psi, self.wave_function.trainable_variables)
```

VMC and SR methods obtain parameters update in different ways.

Update the parameters of the wave function in the same way:

```
for grad_val, var in grads_vars:
    var.assign_sub(self.learning_rate * grad_val)
```



Optimizers - VMC

gradients.append(vmc_grad)

 $\nabla_{\mathcal{W}}(E) = 2\Re\left[\langle E_{\text{loc}} \nabla_{\mathcal{W}} log(\psi) \rangle - \langle E_{\text{loc}} \rangle \langle \nabla_{\mathcal{W}} log(\psi) \rangle\right]$



The SR approach computes the variational increments as follows:



Optimizers - Stochastic Reconfiguration

$$F = \langle E_{loc} \nabla_{\mathcal{W}} log(\psi) \rangle - \langle E_{loc} \rangle \langle \nabla_{\mathcal{W}} log(\psi) \rangle$$
 # compute force vector mean_energy = tf.reduce_mean(local_energies)
$$F_{-} vec = (\\ tf.reduce_mean(local_energies * 0, axis=0) - mean_energy * 0_mean)$$



Optimizers - Stochastic Reconfiguration

$$\Delta W = S^{-1}F$$

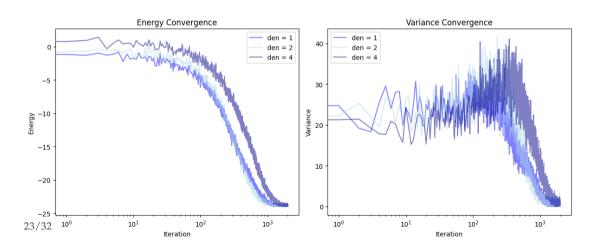
```
# solve for delta
P = tf.shape(S)[0]
S_reg = S + self.epsilon * tf.eye(P, dtype=S.dtype)
# dW = S^{-1} F
gradients = tf.linalg.solve(S_reg, tf.expand_dims(F_vec, 1))
```



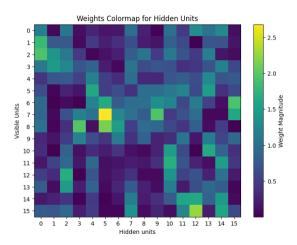


Results 1D Ising Model - 16 spins

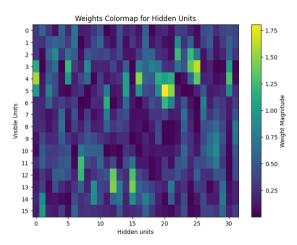
Using J = 1.0 and h = 0.5



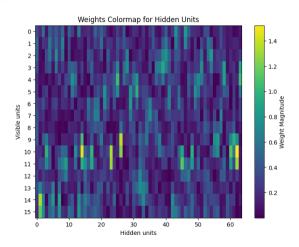








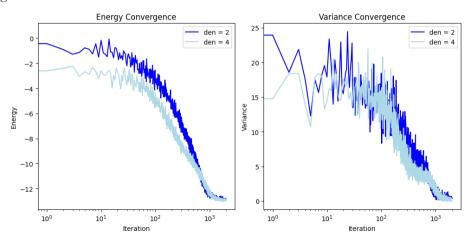




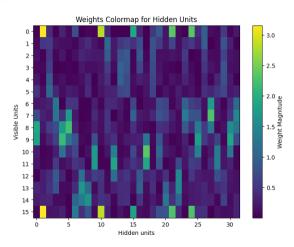


1D Ising Model - 16 spins

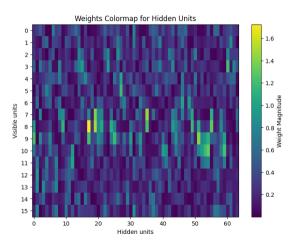
Using J = -1.0 and h = 0.5













Extensions

- Extend to 2D systems
- Other neural networks (FFNN, CNN)
- Unitary (and non) dynamics

Optimization

- Efficient SR update (single parameter updates)
- Change tensor framework (JAX, ...)

Other Projects

NetKet: The Machine-Learning toolbox for Quantum Physics





References

- Giuseppe Carleo, Matthias Troyer, Solving the quantum many-body problem with artificial neural networks. *Science* 355, 602-606 (2017).
- Moritz Reh, Markus Schmitt, Martin Gärttner, Optimizing design choices for neural quantum states, *Phys. Rev. B* 107, 195115 (2023).
- Becca F, Sorella S. *Quantum Monte Carlo Approaches for Correlated Systems*. Cambridge University Press; 2017.
- Francesco D'Angelo, Lucas Böttcher, Learning the Ising model with generative neural networks, *Phys. Rev. Research* 2, 023266 (2020).



Thanks for your attention!

Appendix





Appendix

Probabilities for the RBM

Thanks to no intra-layer connections, we can factorize the joint probability:

$$p(\vec{\sigma} \mid \vec{h}) = \prod_{i} p(\sigma_i \mid \vec{h})$$
 $p(\vec{h} \mid \vec{\sigma}) = \prod_{i} p(h_i \mid \vec{\sigma})$

where:

$$p(\sigma_i|h) = \sigma\left(b_i + \sum_j W_{ij}h_j\right)$$

$$ho(h_j|\sigma) = \sigma\left(c_j + \sum_i W_{ij}\sigma_i\right)$$

where $\sigma(x) = \frac{1}{1+e^{-x}}$ is the logistic or sigmoid function.



AppendixMetropolis-Hastings and MCMC sampling

Propose a single bit flip, and accept it with the Metropolis criterion:

$$\alpha = \min\left(1, \frac{p(\sigma')}{p(\sigma)}\right)$$

where σ' is the proposed state and σ is the current state.

Uses the tfp.mcmc module for efficient sampling.

- Requires target log(p)
- Requires new state proposal



We can sample the visible and hidden variables in a double-step Gibbs sampling:

Algorithm: For *k* iterations:

Step 1: Sample hidden units

$$p(h_j = 1|\vec{v}) = \sigma\left(c_j + \sum_i W_{ij}v_i\right)$$

Step 2: Sample visible units

$$p(v_i = 1|\vec{h}) = \sigma\left(b_i + \sum_j W_{ij}h_j\right)$$

where $\sigma(x) = \frac{1}{1+e^{-x}}$ is the sigmoid function.