## **Neural Networks as Quantum States**

NNs in Quantum Many-Body Problems

Deep Learning con Applicazioni

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- Quantum-many-body problems feature an exponential Hilbert-space growth
- Traditional variational ansätze reduce complexity sactificing correlation and entanglement
- Neural-network quantum states can offer expressivity with a compact set of parameters

Neural Networks can represent ground states for a many-body quantum system<sup>1</sup>

<sup>&</sup>lt;sup>1</sup>Giuseppe Carleo, Matthias Troyer, Solving the quantum many-body problem with artificial neural networks. *Science* 355, 602-606 (2017)

## Project Outline Find the Ground State

- Restricted Boltzmann Machine as a neural network
- Sample configurations from RBM
- Train RBM to find the ground state

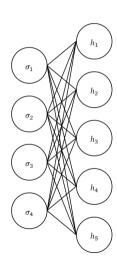




### **Restricted Boltzmann Machine**

Generative Model

- Learns  $p(\vec{\sigma})$  of the input data  $\vec{\sigma}$
- Two layers: *N* visible and *M* hidden
- Parameters W: weights  $W_{ij}$  and biases  $b_i$ ,  $c_j$





#### **Restricted Boltzmann Machine**

**Probability Distribution** 

Internal Energy:

$$E(\vec{\sigma}, \vec{h}) = -\sum_{i} b_i \sigma_i - \sum_{j} c_j h_j - \sum_{i,j} W_{ij} \sigma_i h_j$$

Probability Distribution:

$$p(\vec{\sigma}, \vec{h}) = \frac{1}{Z} e^{-E(\vec{\sigma}, \vec{h})}$$

where *Z* is the partition function:

$$Z = \sum_{ec{\sigma}, ec{h}} e^{-E(ec{\sigma}, ec{h})}$$



#### **Restricted Boltzmann Machine**

Application to Quantum States

Considering a spin configuration  $\vec{\sigma}$  and a set of parameters  $\mathcal{W}$ :

$$\psi\left(\vec{\sigma},\mathcal{W}\right)=e^{\sum_{i}b_{i}\sigma_{i}+\sum_{j}c_{j}h_{j}+\sum_{i,j}W_{ij}\sigma_{i}h_{j}}$$

With no intra-layer connections:

$$\psi\left(ec{\sigma},\mathcal{W}
ight)=\mathrm{e}^{\sum_{i}b_{i}\sigma_{i}} imes\prod_{j=1}^{M}2cosh\left[c_{j}+\sum_{i}W_{ij}\sigma_{i}
ight]$$

From the RBM we can sample configurations  $\vec{\sigma}$ .



Update parameters W towards energy minimum.

$$E_0 \leq \frac{\langle \psi_{\mathcal{W}} | \hat{H} | \psi_{\mathcal{W}} \rangle}{\langle \psi_{\mathcal{W}} | \psi_{\mathcal{W}} \rangle}$$

The parameters W are optimized using a gradient descent method:

- Variational Monte Carlo (VMC)<sup>2</sup>
- Stochastic Reconfiguration (SR)<sup>3</sup>

<sup>&</sup>lt;sup>2</sup>Moritz Reh, Markus Schmitt, Martin Gärttner, Optimizing design choices for neural quantum states, *Phys. Rev. B* 107, 195115 (2023)

<sup>&</sup>lt;sup>3</sup>Becca F, Sorella S. *Quantum Monte Carlo Approaches for Correlated Systems*. Cambridge University Press; 2017



We can compute the gradients:

$$\begin{split} \nabla_{\mathcal{W}}\left(\mathcal{E}\right) &= \nabla_{\mathcal{W}} \langle \mathcal{E}_{\text{loc}} \rangle \\ &= \langle \mathcal{E}_{\text{loc}} \nabla_{\mathcal{W}} log(p_{\psi}) \rangle \\ &= 2 \Re \left[ \langle \mathcal{E}_{\text{loc}} \nabla_{\mathcal{W}} log(\psi) \rangle - \langle \mathcal{E}_{\text{loc}} \rangle \langle \nabla_{\mathcal{W}} log(\psi) \rangle \right] \end{split}$$

And consequently update the parameters:

$$\Delta \mathcal{W} = -\eta \nabla_{\mathcal{W}} \left( E \right)$$



Parameters evolution in the variational space:

$$\Delta \mathcal{W} = -\eta \mathbf{S}^{-1} \vec{F}$$

where  $S^{-1}$  is the pseudo-inverse of the covariance matrix:

$$\mathbf{S}_{ij} = \langle O_i^* O_j \rangle - \langle O_i^* \rangle \langle O_j \rangle$$

and  $\vec{F}$  is the force vector:

$$F_i = \langle O_i^* E_{loc} \rangle - \langle O_i^* \rangle \langle E_{loc} \rangle$$

where  $O_i$  are the local operators defined as:

$$O_i = rac{\partial log\left(\psi\left(ec{\sigma},\mathcal{W}
ight)
ight)}{\partial W_{ij}}$$

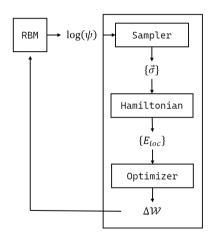
- RBM
- Sampler
- Hamiltonian
- Optimizers





Code Structure

nnqs/
- nnqs.py -> RBM
- hamiltonian.py -> Ising1D
- optimizer.py -> MRT2,
 GibbsSampler
- sampler.py -> VMC, SR





Custom tf.module implementing the RBM.

Weights and biases are initialized randomly as complex values:

```
self.W = tf.Variable(
    tf.random.normal([self.N_visible, self.N_hidden], dtype=tf.complex64)
)
self.b = tf.Variable(tf.random.normal([self.N_visible], dtype=tf.complex64))
self.c = tf.Variable(tf.random.normal([self.N_hidden], dtype=tf.complex64))
```

$$\psi\left(ec{\sigma},\mathcal{W}
ight)= extbf{e}^{\sum_{i}b_{i}\sigma_{i}} imes\prod_{j=1}^{M}2cosh\left[c_{j}+\sum_{i}\textit{W}_{ij}\sigma_{i}
ight]$$



RBM - Logarithm of the Wave Function

$$\vec{\sigma} = (\sigma_1, \sigma_2, \sigma_3, \dots, \sigma_N)$$
 NN  $NN$ 

$$log(\psi) = \sum_{i} b_{i}\sigma_{i} + \sum_{j=1}^{M} log\left[2cosh\left(c_{j} + \sum_{i} W_{ij}\sigma_{i}\right)\right]$$



RBM - Logarithm of the Wave Function

```
sum_visible = tf.reduce_sum(
    a * spins, axis=1
)
w_h = b + tf.matmul(spins, W)
sum_hidden = tf.reduce_sum(
    tf.math.log(2.0 * (
        tf.math.cosh(b + tf.matmul(spins, W))
    )),axis=1
)
return sum_visible + sum_hidden
```

$$\sum_{i} b_{i}\sigma_{i}$$
  $\sum_{j=1}^{M} log \left[ 2cosh \left( c_{j} + \sum_{i} W_{ij}\sigma_{i} 
ight) 
ight]$ 



Sampler

Samples configurations from the RBM.

- Metropolis-Hastings MCMC using TensorFlow Probability
- Gibbs sampling Double-step method to sample visible and hidden variables<sup>4</sup>

Provides batches of configurations. Example:

<sup>&</sup>lt;sup>4</sup>Francesco D'Angelo, Lucas Böttcher, Learning the Ising model with generative neural networks, *Phys. Rev. Research* 2, 023266 (2020)



Hamiltonian - 1D Ising

Simple 1D Ising Hamiltonian:

$$H = -J\sum_{i}\sigma_{i}\sigma_{i+1} - h\sum_{i}\sigma_{i}$$

Provides local\_energy(samples) which computes  $E_{loc}$  for each configuration  $\vec{\sigma}$ .



Optimizers

Leverage <code>tf.GradientTape</code> to compute gradients of  $\log\left(\psi\right)$ 

```
with tf.GradientTape() as tape:
    log_psi = self.wave_function.log_psi(samples)
grad_log_psi = tape.jacobian(log_psi, self.wave_function.trainable_variables)
```

VMC and SR methods obtain parameters update in different ways.

Update the parameters of the wave function in the same way:

```
for grad_val, var in grads_vars:
    var.assign_sub(self.learning_rate * grad_val)
```



Optimizers - VMC

gradients.append(vmc\_grad)

 $\nabla_{\mathcal{W}}(E) = 2\Re\left[\langle E_{\text{loc}} \nabla_{\mathcal{W}} log(\psi) \rangle - \langle E_{\text{loc}} \rangle \langle \nabla_{\mathcal{W}} log(\psi) \rangle\right]$ 



The SR approach computes the variational increments as follows:



Optimizers - Stochastic Reconfiguration

$$F = \langle E_{loc} \nabla_{\mathcal{W}} log(\psi) \rangle - \langle E_{loc} \rangle \langle \nabla_{\mathcal{W}} log(\psi) \rangle$$
 # compute force vector mean\_energy = tf.reduce\_mean(local\_energies) 
$$F_{-} vec = ( \\ tf.reduce_mean(local_energies * 0, axis=0) - mean_energy * 0_mean)$$



Optimizers - Stochastic Reconfiguration

$$\Delta W = S^{-1}F$$

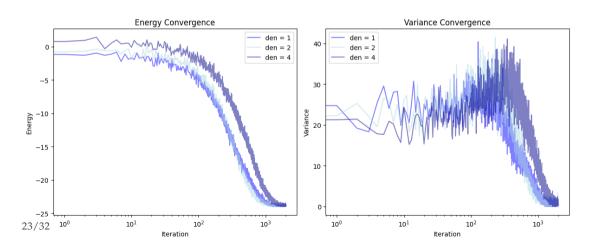
```
# solve for delta
P = tf.shape(S)[0]
S_reg = S + self.epsilon * tf.eye(P, dtype=S.dtype)
# dW = S^{-1} F
gradients = tf.linalg.solve(S_reg, tf.expand_dims(F_vec, 1))
```



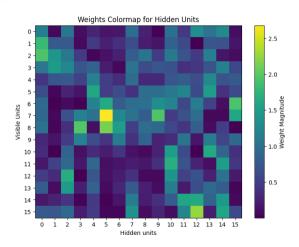


## Results 1D Ising Model - 16 spins

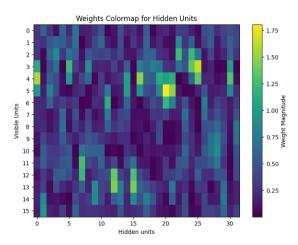
Using J = 1.0 and h = 0.5



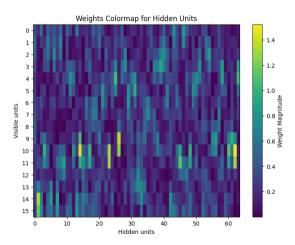








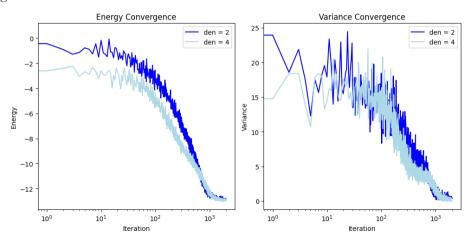




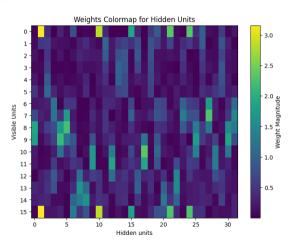


1D Ising Model - 16 spins

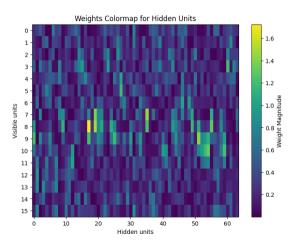
Using J = -1.0 and h = 0.5













#### **Extensions**

- Extend to 2D systems
- Other neural networks (FFNN, CNN)
- Unitary (and non) dynamics

#### Optimization

- Efficient SR update (single parameter updates)
- Change tensor framework (JAX, ...)

#### **Other Projects**

NetKet: The Machine-Learning toolbox for Quantum Physics





#### References

- Giuseppe Carleo, Matthias Troyer, Solving the quantum many-body problem with artificial neural networks. *Science* 355, 602-606 (2017).
- Moritz Reh, Markus Schmitt, Martin Gärttner, Optimizing design choices for neural quantum states, *Phys. Rev. B* 107, 195115 (2023).
- Becca F, Sorella S. *Quantum Monte Carlo Approaches for Correlated Systems*. Cambridge University Press; 2017.
- Francesco D'Angelo, Lucas Böttcher, Learning the Ising model with generative neural networks, *Phys. Rev. Research* 2, 023266 (2020).



# Thanks for your attention!

## Appendix





## **Appendix**

Probabilities for the RBM

Thanks to no intra-layer connections, we can factorize the joint probability:

$$p(\vec{\sigma} \mid \vec{h}) = \prod_{i} p(\sigma_i \mid \vec{h})$$
  $p(\vec{h} \mid \vec{\sigma}) = \prod_{i} p(h_i \mid \vec{\sigma})$ 

where:

$$p(\sigma_i|h) = \sigma\left(b_i + \sum_j W_{ij}h_j\right)$$

$$p(h_j|\sigma) = \sigma\left(c_j + \sum_i W_{ij}\sigma_i\right)$$

where  $\sigma(x) = \frac{1}{1+e^{-x}}$  is the logistic or sigmoid function.



Explicit form of the partial derivatives of  $\log (\psi)$ 



# **Appendix**Metropolis-Hastings and MCMC sampling

Propose a single bit flip, and accept it with the Metropolis criterion:

$$\alpha = \min\left(1, \frac{p(\sigma')}{p(\sigma)}\right)$$

where  $\sigma'$  is the proposed state and  $\sigma$  is the current state.

Uses the tfp.mcmc module for efficient sampling.

- Requires target *log(p)*
- Requires new state proposal



#### **Appendix** Gibbs sampling

We can sample the visible and hidden variables in a double-step Gibbs sampling:

```
def sample(self, wave_function):
    v = tf.cast(self.current_state, tf.float32)
    for _ in range(self.k):
        v_complex = tf.cast(v, wave_function.W.dtvpe)
        v_W = tf.matmul(v_complex, wave_function.W)
        p_h = tf.sigmoid(tf.math.real(wave_function.b + v_W))
        h = tf.cast(tf.random.uniform(tf.shape(p_h), dtype=tf.float32) < p_h, tf.
    float32)
        h_complex = tf.cast(h, wave_function.W.dtype)
        h_Wt = tf.matmul(h_complex, tf.transpose(wave_function.W))
        p_v = tf.sigmoid(tf.math.real(wave_function.a + h_Wt))
        v = tf.cast(tf.random.uniform(tf.shape(p_v), dtype=tf.float32) < p_v, tf.
    float32)
    self.current_state.assign(tf.cast(v. tf.int32))
return v
```