

Math 673

Multigrid Methods: A Mostly Matrix-Based Approach

Chapter 04: A Modern Energy-Norm Analysis of the Two-Grid Method

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Chapter 04

A Modern Energy-Norm Analysis of the Two-Grid Method

Introduction



Like in the last chapter, we will consider the finite element approximation of the model problem in 1D:

$$\begin{cases}
-u'' = f, & \text{in } \Omega = (0,1), \\
u = 0, & \text{on } \partial\Omega = \{0,1\}.
\end{cases}$$
(1)

We will again use a uniform mesh and prove that the two-grid algorithm converges, borrowing many of the results from the previous chapter. But, the analysis style in this chapter will foreshadow the more general, and more modern, style used in the later chapters. In particular, the convergence theory will depend upon two properties, an approximation property and a smoothing property.

The reader will quickly realize how much simpler and more powerful is this modern energy approach to convergence. It is nothing short of remarkable.



Recall we have defined the level-1 stiffness matrix, $A_1 = [a_{1,i,j}] \in \mathbb{R}^{n_1 \times n_1}$, by the following equations:

$$\mathbf{a}_{1,i,j} = \left(\psi_{1,j}',\psi_{1,i}'\right)_{L^2(0,1)} = \left(\psi_{1,i}',\psi_{1,j}'\right)_{L^2(0,1)}, \quad 1 \leq i,j \leq \mathit{n}_1.$$

This matrix is, recall, SPD, and is therefore orthogonally diagonalizable

$$\mathsf{A}_1 = \mathsf{V}_1 \mathsf{D}_1 \mathsf{V}_1^\top,$$

where V_1 is the orthogonal matrix containing the eigenvectors of A_1 ,

$$V_1 = \begin{bmatrix} \begin{matrix} & & & \\ \boldsymbol{v}_1^{(1)} & \boldsymbol{v}_1^{(2)} & \cdots & \boldsymbol{v}_1^{(n_1)} \\ & & & \end{matrix} \in \mathbb{R}^{n_1 \times n_1},$$

and D_1 is the diagonal matrix containing the the eigenvalues of A_1 ,

$$\mathsf{D}_1 = \mathrm{diag}\left(\lambda_1^{(1)}, \lambda_1^{(2)}, \cdots, \lambda_1^{(n_1)}\right) \in \mathbb{R}^{n_1 \times n_1}.$$



Also recall the eigenvalues, $\lambda_1^{(k)}$, and eigenvectors, $\mathbf{v}_1^{(k)}$ of A_1 are defined by

$$\lambda_1^{(k)} = \frac{4}{h_1} \sin^2 \left(\frac{k\pi h_1}{2} \right) = \frac{2}{h_1} \left(1 - \cos(k\pi h_1) \right). \tag{2}$$

and

$$\left[\mathbf{v}_{1}^{(k)}\right]_{i} = \mathbf{v}_{1,i}^{(k)} = \sin\left(k\pi x_{1,i}\right), \quad 1 \le i \le n_{1}, \tag{3}$$

respectively. Since A₁ is SPD, we can define its square root as

$$A_1^{1/2} = V_1 D_1^{1/2} V_1^\top,$$

where

$$\mathsf{D}_1^{1/2} = \mathrm{diag}\left(\sqrt{\lambda_1^{(1)}}, \sqrt{\lambda_1^{(2)}}, \cdots, \sqrt{\lambda_1^{(n_1)}}\right) \in \mathbb{R}^{n_1 \times n_1}.$$

Of course, $A_1^{1/2}$ is SPD, and $A_1^{1/2}A_1^{1/2} = A_1$.

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Now, recall, for any $\mathbf{v}_1 \in \mathbb{R}^{n_1}$,

$$\begin{aligned} \|\mathbf{v}_1\|_{A_1} &:= & \sqrt{(\mathbf{v}_1, \mathbf{v}_1)_{A_1}} \\ &= & \sqrt{(A_1 \mathbf{v}_1, \mathbf{v}_1)_1} \\ &= & \sqrt{\left(A_1^{1/2} \mathbf{v}_1, A_1^{1/2} \mathbf{v}_1\right)_1} \\ &= & \left\|A_1^{1/2} \mathbf{v}_1\right\|_1. \end{aligned}$$

In a similar way, we can define

$$\begin{aligned} \|\mathbf{v}_1\|_{A_1^2} &:= & \sqrt{(A_1^2\mathbf{v}_1, \mathbf{v}_1)_1} \\ &= & \sqrt{(A_1\mathbf{v}_1, A_1\mathbf{v}_1)_1} \\ &= & \|A_1\mathbf{v}_1\|_1. \end{aligned}$$

Lemma

The Approximation Property

Suppose that the two-level grids are uniform and nested and A_0 satisfies the Galerkin condition. Then,

$$\left\| (I_1 - \tilde{\Pi}_1) \mathbf{v}_1 \right\|_1 = \sqrt{2\Lambda_1^{-1}} \left\| (I_1 - \tilde{\Pi}_1) \mathbf{v}_1 \right\|_{A_1}, \tag{4}$$

for all $\mathbf{v}_1 \in \mathbb{R}^{n_1}$, where

$$\Lambda_1 := \frac{4}{h_1}$$
.



Proof.

The Approximation Property 000000000000000

Expand \mathbf{v}_1 in the basis of eigenvectors of A_1 :

$$\mathbf{v}_1 = \sum_{k=1}^{n_1} v_{1,k} \mathbf{v}_1^{(k)}.$$

Recall that

$$(I_1 - \tilde{\Pi}_1) \mathbf{v}_1^{(k)} = S_k \mathbf{v}_1^{(k)} + S_k \mathbf{v}_1^{(k')},$$

for $1 \le k \le n_0$. Similarly,

$$(I_1 - \tilde{\Pi}_1) \mathbf{v}_1^{(k')} = C_k \mathbf{v}_1^{(k)} + C_k \mathbf{v}_1^{(k')},$$

for $1 < k < n_0 + 1$.

Then

The Approximation Property 000000000000000

$$\begin{aligned} \left\| (\mathsf{I}_{1} - \tilde{\mathsf{\Pi}}_{1}) \mathbf{v}_{1} \right\|_{1}^{2} &= \left((\mathsf{I}_{1} - \tilde{\mathsf{\Pi}}_{1}) \mathbf{v}_{1}, (\mathsf{I}_{1} - \tilde{\mathsf{\Pi}}_{1}) \mathbf{v}_{1} \right)_{1} \\ &= \left(n_{1} + 1 \right) \sum_{k=1}^{n_{0}+1} \delta_{k} \gamma_{k}^{2}, \end{aligned}$$

where

$$\delta_k = \begin{cases} \frac{1}{2}, & k = n_0 + 1, \\ 1, & \text{otherwise}, \end{cases}$$

and

$$\gamma_k := \mathbf{v}_{1,k} \mathbf{S}_k + \mathbf{v}_{1,k'} \mathbf{C}_k.$$



Similarly,

$$\begin{split} \left\| (\mathsf{I}_{1} - \tilde{\mathsf{\Pi}}_{1}) \mathbf{v}_{1} \right\|_{\mathsf{A}_{1}}^{2} &= \left((\mathsf{I}_{1} - \tilde{\mathsf{\Pi}}_{1}) \mathbf{v}_{1}, (\mathsf{I}_{1} - \tilde{\mathsf{\Pi}}_{1}) \mathbf{v}_{1} \right)_{\mathsf{A}_{1}} \\ &= \left(\mathsf{A}_{1}^{1/2} (\mathsf{I}_{1} - \tilde{\mathsf{\Pi}}_{1}) \mathbf{v}_{1}, \mathsf{A}_{1}^{1/2} (\mathsf{I}_{1} - \tilde{\mathsf{\Pi}}_{1}) \mathbf{v}_{1} \right)_{1} \\ &= \left(n_{1} + 1 \right) \sum_{k=1}^{n_{0}+1} \delta_{k} \gamma_{k}^{2} \left(\frac{\lambda_{1}^{(k)} + \lambda_{1}^{(k')}}{2} \right) \\ &= \frac{2}{h_{1}} (n_{1} + 1) \sum_{k=1}^{n_{0}+1} \delta_{k} \gamma_{k}^{2} \\ &= \frac{1}{2} \mathsf{A}_{1} \left\| (\mathsf{I}_{1} - \tilde{\mathsf{\Pi}}_{1}) \mathbf{v}_{1} \right\|_{1}^{2}. \end{split}$$

Remark

The Approximation Property

We will prove a similar result later in the general FEM framework, where, in particular, the meshes need not be uniform. In that setting, we will seek to prove that there is a constant $C_1>0$ such that

$$\|(I_1 - \tilde{\Pi}_1)\mathbf{v}_1\|_1 \le C_1 \sqrt{\Lambda_1^{-1}} \|(I_1 - \tilde{\Pi}_1)\mathbf{v}_1\|_{A_1},$$
 (5)

for all $\mathbf{v}_1 \in \mathbb{R}^{n_1}$.

Theorem (The Approximation Property)

Suppose that the two-level grids are uniform and nested and A_0 satisfies the Galerkin condition. For any $\mathbf{v}_1 \in \mathbb{R}^{n_1}$,

$$\|(\mathbf{I}_1 - \tilde{\mathbf{\Pi}}_1)\mathbf{v}_1\|_{A_1} \le C_1 \sqrt{\Lambda_1^{-1}} \|\mathbf{v}_1\|_{A_1^2},$$
 (6)

where

The Approximation Property

$$\Lambda_1 := \frac{4}{h_1}$$
 and $C_1 = \sqrt{2}$.



Proof.

The Approximation Property

Let $\mathbf{v}_1 \in \mathbb{R}^{n_1}$ be arbitrary. Then

$$\begin{split} \left\| (I_{1} - \tilde{\Pi}_{1}) \mathbf{v}_{1} \right\|_{A_{1}}^{2} &= \left(A_{1} (I_{1} - \tilde{\Pi}_{1}) \mathbf{v}_{1}, (I_{1} - \tilde{\Pi}_{1}) \mathbf{v}_{1} \right)_{1} \\ &= \left((I_{1} - \tilde{\Pi}_{1}) \mathbf{v}_{1}, A_{1} (I_{1} - \tilde{\Pi}_{1}) \mathbf{v}_{1} \right)_{1} \\ &= \left((I_{1} - \tilde{\Pi}_{1}) \mathbf{v}_{1}, A_{1} \mathbf{v}_{1} \right)_{1} - \left((I_{1} - \tilde{\Pi}_{1}) \mathbf{v}_{1}, A_{1} \tilde{\Pi}_{1} \mathbf{v}_{1} \right)_{1}. \end{split}$$



The Approximation Property

The second term on the RHS is zero, as we now show:

$$\begin{split} \left((I_1 - \tilde{\Pi}_1) \boldsymbol{v}_1, \tilde{\Pi}_1 \boldsymbol{v}_1 \right)_{A_1} &= \left((I_1 - \tilde{\Pi}_1) \boldsymbol{v}_1, A_1 \tilde{\Pi}_1 \boldsymbol{v}_1 \right)_1 \\ &= \left((I_1 - \tilde{\Pi}_1) \boldsymbol{v}_1, A_1 (P_0 A_0^{-1} R_0 A_1) \boldsymbol{v}_1 \right)_1 \\ &= \left((I_1 - \tilde{\Pi}_1) \boldsymbol{v}_1, (A_1 P_0 A_0^{-1} R_0) A_1 \boldsymbol{v}_1 \right)_1 \\ &= \left((A_1 P_0 A_0^{-1} R_0)^\top (I_1 - \tilde{\Pi}_1) \boldsymbol{v}_1, A_1 \boldsymbol{v}_1 \right)_1 \\ &= \left(\tilde{\Pi}_1 (I_1 - \tilde{\Pi}_1) \boldsymbol{v}_1, \boldsymbol{v}_1 \right)_{A_1} \\ &= 0, \end{split}$$

since.

$$\tilde{\Pi}_1(I_1 - \tilde{\Pi}_1) = \tilde{\Pi}_1 - \tilde{\Pi}_1^2 = \tilde{\Pi}_1 - \tilde{\Pi}_1 = 0.$$

Therefore,

The Approximation Property 0000000000000000

$$\begin{split} \left\| (I_1 - \tilde{\Pi}_1) \boldsymbol{\nu}_1 \right\|_{A_1}^2 &= \left((I_1 - \tilde{\Pi}_1) \boldsymbol{\nu}_1, A_1 \boldsymbol{\nu}_1 \right)_1 \\ &\stackrel{\text{C.S.}}{\leq} \left\| (I_1 - \tilde{\Pi}_1) \boldsymbol{\nu}_1 \right\|_1 \left\| A_1 \boldsymbol{\nu}_1 \right\|_1 \\ &= \left\| (I_1 - \tilde{\Pi}_1) \boldsymbol{\nu}_1 \right\|_1 \left\| \boldsymbol{\nu}_1 \right\|_{A_1^2} \\ &\stackrel{(4)}{=} \sqrt{2 \Lambda_1^{-1}} \left\| (I_1 - \tilde{\Pi}_1) \boldsymbol{\nu}_1 \right\|_{A_1} \left\| \boldsymbol{\nu}_1 \right\|_{A_1^2}. \end{split}$$

The result follows.

Corollary

The Approximation Property

Suppose that the two-level grids are uniform and nested and A_0 satisfies the Galerkin condition. Then

$$\left((I_1-\tilde{\Pi}_1)\textbf{\textit{v}}_1,(I_1-\tilde{\Pi}_1)\textbf{\textit{v}}_1\right)_{A_1}=\left((I_1-\tilde{\Pi}_1)\textbf{\textit{v}}_1,\textbf{\textit{v}}_1\right)_{A_1},$$

or, equivalently,

$$\left(\big(\textbf{I}_1 - \tilde{\boldsymbol{\Pi}}_1 \big) \boldsymbol{v}_1, \tilde{\boldsymbol{\Pi}}_1 \boldsymbol{v}_1 \right)_{A_1} = 0.$$



Remark

The Approximation Property

Estimate (6) is called an approximation property with constant $C_1 > 0$.



In this section, we want to change our smoothing algorithm from damped Jacobi to Richardson's method. It turns out that this change is mostly cosmetic, as the underlying structure for the smoother will be essentially the same. Specifically, our Richardson smoother will have the following error transfer matrix:

$$K_1 = I_1 - \frac{1}{\Lambda_1} A_1, \quad \Lambda_1 = \frac{4}{h_1}.$$

Recall

$$\lambda_1^{(k)} = \frac{2}{h_1} \left(1 - \cos(k\pi h_1) \right).$$

So,

$$0 < \lambda_1^{(1)} < \lambda_1^{(2)} < \dots < \lambda_1^{(n_1)} = \rho(A_1) < \frac{4}{h_1} = \Lambda_1.$$

 Λ_1 is almost the spectral radius of A_1 ; the last estimate is asymptotically sharp. For damped Jacobi, recall that

$$\begin{aligned} \mathsf{K}_1 &= \mathsf{K}_1(\omega) &=& \mathsf{I}_1 - \omega \mathsf{D}^{-1} \mathsf{A}_1 \\ &=& \mathsf{I}_1 - \omega \frac{h_1}{2} \mathsf{A}_1. \end{aligned}$$

If we take $\omega = \frac{1}{2}$ in damped Jacobi, we get our new Richardson smoother.



The error transfer matrix, which, for Richardson's method is

$$\mathsf{K}_1 = \mathsf{I}_1 - \mathsf{\Lambda}_1^{-1} \mathsf{A}_1 = \mathsf{I}_1 - \frac{\mathit{h}_1}{4} \mathsf{A}_1,$$

has the following eigenvalues:

$$\mu_1^{(k)} \left(\frac{1}{2} \right) = \frac{1}{2} \cos(k\pi h_1) + 1 - \frac{1}{2}$$

$$= 1 - \frac{\lambda_1^{(k)}}{\Lambda_1}.$$

Consulting the figure on the next page, it follows easily that

$$\left|\mu_1^{(k)}\left(\frac{1}{2}\right)\right| = \mu_1^{(k)}\left(\frac{1}{2}\right) \le \mu_1^{(n_0+1)}\left(\frac{1}{2}\right) = \frac{1}{2}, \quad n_0 + 1 \le k \le n_1.$$

In other words the smoother reduces all high frequency modes of the error by at least half.



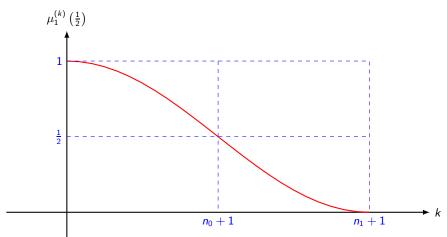


Figure: Eigenvalues of the error transfer matrix for the Richardson smoother.



The Smoothing Property



Before we establish the so-called $smoothing\ property$, we need a technical lemma.

Lemma

For any $m \in \mathbb{N}$,

$$\max_{0 \le x \le 1} x (1 - x)^{2m} \le \frac{1}{2m}.$$
 (7)

Proof.

Define, for any $m \in \mathbb{N}$,

$$f(x) = x(1-x)^{2m}, x \in [0,1].$$

Observe that

$$f(0) = f(1) = 0,$$

and, otherwise

$$f(x) > 0, \quad \forall x \in (0,1).$$

The derivative of f is

$$f'(x) = (x-1)^{2m} + x(2m)(x-1)^{2m-1},$$

and there is a single zero for f'(x) in (0,1), which we label x_0 . It is easy to show that

$$x_0=\frac{1}{2m+1}.$$

Then, the extreme value of f satisfies

$$f(x_0) = \frac{1}{2m+1} \left(\frac{2m}{2m+1}\right)^{2m} \le \frac{1}{2m+1} \le \frac{1}{2m}.$$

The result follows.

Theorem (The Smoothing Property)

Suppose that the two-level grids are uniform and nested and A₀ satisfies the Galerkin condition. Assume that smoothing is carried out by Richardson's method, or, equivalently, the damped Jacobi method with $\omega = \frac{1}{2}$. There is some constant $C_2 > 0$, such that

$$\|\mathsf{K}_{1}^{m_{1}}\mathbf{v}_{1}\|_{\mathsf{A}_{1}^{2}} \leq C_{2}\sqrt{\frac{\mathsf{A}_{1}}{m_{1}}}\,\|\mathbf{v}_{1}\|_{\mathsf{A}_{1}},$$
 (8)

The Smoothing Property 000000000

for all $\mathbf{v}_1 \in \mathbb{R}^{n_1}$, for any $m_1 \in \mathbb{N}$. In particular, we can take

$$C_2 = \sqrt{\frac{1}{2}}.$$



Proof.

First, observe that

$$\begin{aligned} \|\mathsf{K}_{1}^{m_{1}} \mathbf{v}_{1}\|_{\mathsf{A}_{1}^{2}}^{2} &= \|\mathsf{A}_{1} \mathsf{K}_{1}^{m_{1}} \mathbf{v}_{1}\|_{1}^{2} \\ &= (\mathsf{A}_{1} \mathsf{K}_{1}^{m_{1}} \mathbf{v}_{1}, \mathsf{A}_{1} \mathsf{K}_{1}^{m_{1}} \mathbf{v}_{1})_{1}. \end{aligned}$$

Using the eigenvector basis for A₁, let us write

$$\mathbf{v}_1 = \sum_{k=1}^{n_1} v_{1,k} \mathbf{v}_1^{(k)}.$$

Then

$$\begin{split} \|\mathsf{K}_{1}^{\mathit{m}_{1}}\,\boldsymbol{v}_{1}\|_{\mathsf{A}_{1}^{2}}^{2} & = & \frac{\mathit{n}_{1}+1}{2}\sum_{\mathit{k}=1}^{\mathit{n}_{1}}\left(\lambda_{1}^{(\mathit{k})}\mathit{v}_{1,\mathit{k}}\right)^{2}\left(\mu_{1}^{(\mathit{k})}\left(\frac{1}{2}\right)\right)^{2\mathit{m}_{1}} \\ & = & \Lambda_{1}\left(\frac{\mathit{n}_{1}+1}{2}\right)\sum_{\mathit{k}=1}^{\mathit{n}_{1}}\left(\frac{\lambda_{1}^{(\mathit{k})}}{\Lambda_{1}}\right)\left(1-\frac{\lambda_{1}^{(\mathit{k})}}{\Lambda_{1}}\right)^{2\mathit{m}_{1}}\lambda_{1}^{(\mathit{k})}\mathit{v}_{1,\mathit{k}}^{2} \\ & \leq & \Lambda_{1}\mathit{G}(\mathit{m}_{1})\left(\frac{\mathit{n}_{1}+1}{2}\right)\sum_{\mathit{k}=1}^{\mathit{n}_{1}}\lambda_{1}^{(\mathit{k})}\mathit{v}_{1,\mathit{k}}^{2} \\ & = & \Lambda_{1}\mathit{G}(\mathit{m}_{1})\left\|\boldsymbol{v}_{1}\right\|_{\mathsf{A}_{1}}^{2}, \end{split}$$

where

$$G(m) := \max_{1 \le k \le n_1} \left(\frac{\lambda_1^{(k)}}{\Lambda_1} \right) \left(1 - \frac{\lambda_1^{(k)}}{\Lambda_1} \right)^{2m}.$$

Upon rescaling and using the lemma we just proved, we have

$$G(m) \leq \max_{0 \leq x \leq 1} x (1-x)^{2m}$$

$$\leq \frac{1}{2m}.$$

Therefore

$$\|\mathsf{K}_{1}^{m_{1}} \mathbf{v}_{1}\|_{\mathsf{A}_{1}^{2}} \leq \sqrt{\frac{1}{2}} \sqrt{\frac{\mathsf{\Lambda}_{1}}{m_{1}}} \|\mathbf{v}_{1}\|_{\mathsf{A}_{1}},$$
 (9)

and the result follows with $C_2 = \sqrt{\frac{1}{2}}$.

The Smoothing Property 00000000

Remark

Estimate (8) is called a **smoothing property** with constant $C_2 > 0$.



Convergence in the Energy Norm

Now that we have the powerful smoothing and approximation properties, the energy-norm convergence of the two-grid method is simple matter.

Theorem (Convergence of the One-Sided Two-Grid Method)

Suppose that the two-level grids are uniform and nested and A₀ satisfies the Galerkin condition. Suppose $m_2 = 0$, and with $\omega = \frac{1}{2}$ (Richardson). Then the two grid method converges provided m₁ is sufficiently large. Moreover, we have the error estimate

 $\|\mathbf{e}_1^{\ell+1}\|_{A_1} \leq C_1 C_2 m_1^{-1/2} \|\mathbf{e}_1^{\ell}\|_{A_1},$

where $C_1, C_2 > 0$ are as given in Theorem 2 and Lemma 5, respectively. Since we can assume $C_1C_2=1$, it follows that the method converges if $m_1>2$.

Proof.

Recall, the error transfer matrix in this case is

$$\mathsf{E}_1 = (\mathsf{I}_1 - \tilde{\mathsf{\Pi}}_1) \mathsf{K}_1^{\mathit{m}_1}.$$

So

$$\begin{aligned} \left\| \boldsymbol{e}_{1}^{\ell+1} \right\|_{A_{1}} &= \left\| (I_{1} - \tilde{\Pi}_{1}) K_{1}^{m_{1}} \boldsymbol{e}_{1}^{\ell} \right\|_{A_{1}} \\ &\stackrel{(6)}{\leq} C_{1} \sqrt{\Lambda_{1}^{-1}} \left\| K_{1}^{m_{1}} \boldsymbol{e}_{1}^{\ell} \right\|_{A_{1}^{2}} \\ &\stackrel{(8)}{\leq} C_{1} \sqrt{\Lambda_{1}^{-1}} C_{2} \sqrt{\Lambda_{1}} m_{1}^{-\frac{1}{2}} \left\| \boldsymbol{e}_{1}^{\ell} \right\|_{A_{1}} \\ &= C_{1} C_{2} m_{1}^{-\frac{1}{2}} \left\| \boldsymbol{e}_{1}^{\ell} \right\|_{A_{1}} .\end{aligned}$$



Remark

In this simplified setting, we know that $C_1C_2=1$. For the more general FEM case that we will explore in the future – for which the h_1 -independent constants $C_1>0$ and $C_2>0$, may not be known explicitly – we will still be guaranteed that the method converges at a uniform, h_1 -independent rate provide $m_1\geq 1$ is large enough so that

$$0 < \mathit{C}_{1}\mathit{C}_{2}\mathit{m}_{1}^{-\frac{1}{2}} < 1.$$