# Temperature Effect on Rubber Balloon

#### Abstract

Rubber balloon is inflated by the pressure jump across the membrane due to inner gas. If the inner gas is lighter than air, the rubber balloon will produce buoyance to lift a payload. So the large rubber balloon can serve as a vehicle to transport scientific or military apparatus into high altitude. For rubber is material of entropy induced elasticity, the deformation of rubber balloon depends on temperature. The environmental temperature varies along altitude and changes significantly during the diurnal cycles, it is significant to study the thermoelastic effect of rubber on the ascending process of rubber balloon.

In this paper, we proposed a modified Mooney-Rivlin model of rubber which considered the dependence of elastic coefficients on temperature. We studied the ascending process of rubber balloon with payload using this modified model and compared its results with the normal one. We also found that the rubber balloon would keep its shape at high altitude during the diurnal cycles, which means the balloon could maintain nearly a constant altitude for a long period. The results given here show the temperature dependence of elasticity is not negligible and can be wisely utilized when we use rubber balloon as High Altitude Long Duration(HALD) aerostat.

 $Key\ words$ : Rubber balloons, Mooney–Rivlin model, thermoelasticity, HALD aerostat

### 1 Introduction

Rubber balloon plays an important role in scientific activities, such as high altitude atmosphere research, meteorology, stratospheric composition and astronomy. In fact, it also can serve as a vehicle to transport a stratospheric airship into the nearspace (20-100km). All these scientific investigations can be realized by rubber balloons filled with hydrogen or helium, see Figure 1.

The inflation of rubber balloon refers to the problem of large deformation[1, 2, 3, 4]. To analyse or simulate this deformation, it is important to choose the material law[4, 5]. Mooney-Rivlin model is always used. Results from this model are compared with experimental data[4]. Müller and Struchtrup simulate the process of inflating a single balloon by mouth[5].

In fact the coefficients of elasticity in Mooney-Rivlin model depend on temperature due to its entropy induced elasticity[6], this may influence the performance of rubber balloon.

In this paper, we proposed a modified Mooney-Rivlin model which considered the temperature dependence of rubber's elasticity. We studied the temperature effect of rubber's elasticity on the ascending process of rubber balloon. The results shows temperature has a significant effect on the performance of

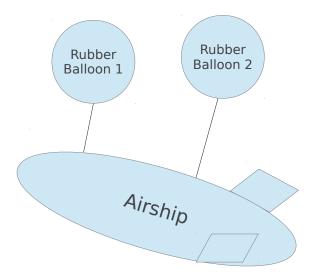


Figure 1: Airship with rubber balloon

rubber balloon especially at high altitude. Besides, we also found rubber balloon will keep its shape at high altitude when temperature changes, this means that the rubber balloon could stay aloft at a constant altitude over many diurnal cycles.

# 2 Mooney-Rivlin model

Since rubber is typical hyperelastic material, We choose the Mooney-Rivlin model as its constitute law. For a rubber balloon, the pressure equation is

$$\Delta p = p_{in} - p_{out} \tag{1}$$

The inner pressure  $p_{in}$  can be calculated by ideal gas relation:

$$p_{in} = \frac{nRT}{\frac{4}{3}\pi r^3} \tag{2}$$

n is the number of moles of gas in the balloon,  $R = 8.3144J \cdot K^{-1} \cdot mol^{-1}$  is the gas constant. The  $\Delta p - r$  characteristic, which describes the dependence of pressure difference across the spherical rubber balloon's membrane on its radius, is non-monotonic, see Figure 2. Its analytical form is

$$\Delta p(r) = 2s_{+} \frac{t_0}{r_0} \left( \frac{r_0}{r} - \left( \frac{r_0}{r} \right)^7 \right) \left( 1 - \frac{s_{-}}{s_{+}} \left( \frac{r}{r_0} \right)^2 \right) \tag{3}$$

From(1),(2), and (3), we obtain

$$p_{out}(h) - \frac{nRT(h)}{\frac{4}{3}\pi r^3} + 2s_{+}\frac{t_0}{r_0} \left(\frac{r_0}{r} - \left(\frac{r_0}{r}\right)^7\right) \left(1 - \frac{s_{-}}{s_{+}} \left(\frac{r}{r_0}\right)^2\right) = 0$$
 (4)

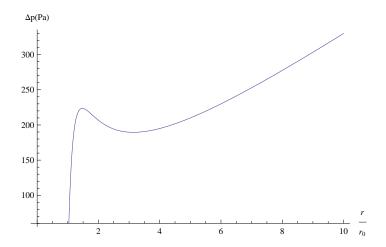


Figure 2: Pressure radius characteristic

 $t_0$  and  $r_0$  are the unstretched thickness and radius of the balloon.  $s_+$  and  $s_-$  are the elastic coefficients. We notice that the atmospherical pressure and temperature is functions of altitude. We now study the ascending process of a rubber balloon with payload. The mass of payload  $m_p = 1.5kg$ , the unstretched radius  $r_0$  and thickness  $t_0$  is given by the manufacturer, here, we set them 0.5m and 0.25mm. With  $r_0$ ,  $t_0$ , and the density of the rubber, we can work out the mass of the rubber balloon. We assume the radius at bursting to be 5m. By setting r to the initial inflation radius at h = 0, we can obtain n. Then we can solve equation (4) for r at each altitude h.

# 3 Modified Mooney-Rivlin model

### 3.1 Thermoelasticity

For rubber is the material with entropy-induced elasticity, the elastic coefficients  $s_+$  and  $s_-$  actually depend on temperature; their dependence on temperature is linear[6]. This means

$$\frac{s_{+}(T)}{s_{+}(T_{0})} = \frac{s_{-}(T)}{s_{-}(T_{0})} = \frac{T}{T_{0}}$$
(5)

So equation (4) needs some modification. It becomes to

$$p_{out} - \frac{nRT}{\frac{4}{3}\pi r^3} + 2s_+(T)\frac{t_0}{r_0} \left(\frac{r_0}{r} - \left(\frac{r_0}{r}\right)^7\right) \left(1 - \frac{s_-(T_0)}{s_+(T_0)} \left(\frac{r}{r_0}\right)^2\right) = 0 \quad (6)$$

Then the thermodynamic state of the inner gas is coupled with the strain or the stress of balloon. The behavior of rubber varies at different temperature, also at different altitude. The study here is to reveal how much of an effect this has on the ascending process.

#### 3.2 Comparison between Mooney-Rivlin models

The typical value given in [6] for  $s_+$  and  $s_-$ :

$$s_{+}(300K) = 300kPa, s_{-}(300K) = -30kPa \tag{7}$$

Figure 3 shows the resulting graph of balloon radius versus height. As well as the Mooney-Rivlin model ( $s_-$  and  $s_+$  are constant during ascending process) and the modified Mooney-Rivlin model( $s_-$  and  $s_+$  change with altitude during ascending process), the graph also includes a purely theoretical non-restoring model, in which the balloon exerts no retractive elastic force at all ( $\Delta p = 0$ , so the inner pressure is always equal to the atmospherical pressure).

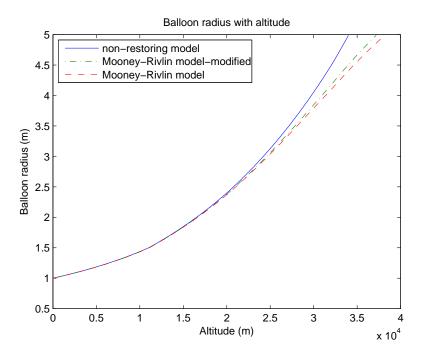


Figure 3: Balloon radius with altitude

We have assumed the burst radius of the balloon as 5m, This occurs at 34km for the non-restoring model, 38.2km for the Mooney-Rivlin model and 37.2km for the modified Mooney-Rivlin model with the provided parameters. The result here suggests that the existence of elasticity increases the ceiling of the balloon though very limited. The difference between Mooney-Rivlin models is slight.

Figure 4 shows the pressure jump across the balloon membrane for each model.

For the balloon with the provided parameters here, it can be seen that the membrane pressure is of the order of 140-300Pa throughout the flight envelope. This means that the difference in pressure between inside and outside is negligible at low altitudes(sea level pressure is about 101325 Pa). However, at high altitudes it may start to become significant. At 30 km the atmospheric pressure has fallen to 1197Pa; at 40 km it is 287Pa. This may lead some interesting phenomenon which we will see below.

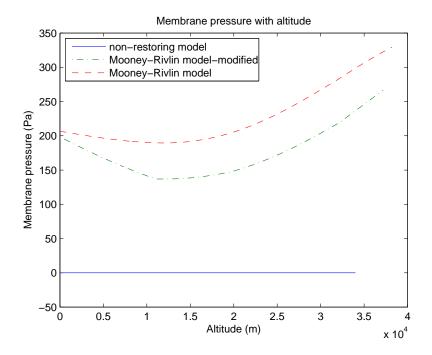


Figure 4: Membrane pressure with altitude

We also get that the membrane pressure in modified Mooney-Rivlin model is always smaller than in Mooney-Rivlin model. This is because the temperature is lower at high altitude than sea level, then the restoring force decreases due to the entropy induced elasticity.

Lift is calculated by buoyancy minus gravity. Buoyancy is proportional to the product the volume of the balloon and the density of the atmosphere, while gravity contains three parts(payload, balloon and inner gas). The difference is very obvious when we plot the lift curves of different models, see Figure 5.

Because there is no pressure jump across the membrane in the non-restoring model, the lift is constant. But this should be changed when we consider elasticity. The free lift decreases from 2.7kg to 1.0kg below as the balloon approaches its burst point in Mooney-Rivlin models. Though the balloon may still ascend, it is important when we need estimate whether the target altitude is achieved and assess the lift performance. We also get that the modified Mooney-Rivlin model has better lift performance than the normal one.

The ascent velocity was calculated based on the formula for terminal velocity:

$$L = F_d = C_d \frac{\rho}{2} A v^2, v = \sqrt{\frac{2L}{C_d \rho A}}$$
 (8)

here  $C_d=0.47$  for a sphere,  $A=\pi r^2$  is the cross-section area, and  $\rho$  is the atmospherical density. It is easy to know lift L and  $\rho$  are functions of altitude. Then we plot the ascent velocity with altitude for each models (Figure 6).

The ascent velocity generally increases with height (below 30km), but it decreases at high altitude (above 30km) in Mooney-Rivlin models. When we con-

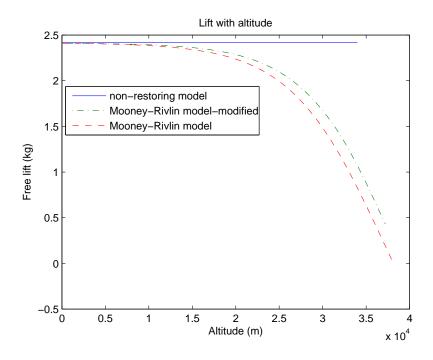


Figure 5: Lift with altitude  $\,$ 

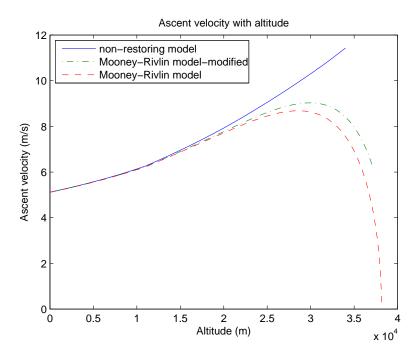


Figure 6: Ascent velocity with altitude

sider the temperature effect in Mooney-Rivlin model, we found that the velocity is larger due to the better lift performance.

## 3.3 Staying aloft

Now we turn to consider another interesting phenomenon when the temperature effect of elasticity is taken account.

If the rubber balloon stay aloft at very high altitude (nearspace), say 80km, the atmospherical pressure is approximately 0Pa. Then equation (6) changes to

$$-\frac{nRT}{\frac{4}{3}\pi r^3} + 2s_+(T)\frac{t_0}{r_0} \left(\frac{r_0}{r} - \left(\frac{r_0}{r}\right)^7\right) \left(1 - \frac{s_-(T_0)}{s_+(T_0)} \left(\frac{r}{r_0}\right)^2\right) = 0 \tag{9}$$

we can cancel T in each term, then we get:

$$-\frac{nR}{\frac{4}{3}\pi r^3} + 2\frac{s_+(T_0)}{T_0} \frac{t_0}{r_0} \left(\frac{r_0}{r} - \left(\frac{r_0}{r}\right)^7\right) \left(1 - \frac{s_-(T_0)}{s_+(T_0)} \left(\frac{r}{r_0}\right)^2\right) = 0 \tag{10}$$

It suggests that the balloon's radius r is independent with temperature, the balloon will keep its shape when the balloon's temperature changed, this is important when we consider a balloon that is capable of staying aloft for half a year at high altitude. At night the temperature of the lifting gas cools, while the temperature rises caused by solar heating during the day. These are called "supercool" and "superheat". From the equation above, though the temperature varies during the diurnal cycles, the balloon will keep its volume, so the residence altitude. For some aerostats need carry a significant amount of ballast to maintain constant altitude over several diurnal cycles, it will be a great advantage when we use rubber balloon as the High Altitude Long Duration aerostat.

## 4 Conclusion

The ascending process of rubber balloon has been studied. Due to the entropy induced elasticity, the elastic behavior of rubber varies at different temperature. We set rubber as material of a modified Mooney-Rivlin relation, and we found that the temperature effect of elasticity plays an important role in ascending process. We also show the interesting phenomenon that the balloon will keep its height when temperature changes at high altitude. This fact may be wisely utilized when we develop a High Altitude Long Duration aerostat.

But the dependence of elastic coefficients of rubber on temperature is rather complex, and the nature of them is still not clear[7]. In fact, from the experiment the coefficient  $s_{-}$  is always increasing more slowly than the absolute temperature[8]. Even more, the material of rubber balloons is not strictly a Mooney-Rivlin model. All these make it more difficult to predict the performance of rubber balloon, which needs more examination.

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