Topological Quantum Computing

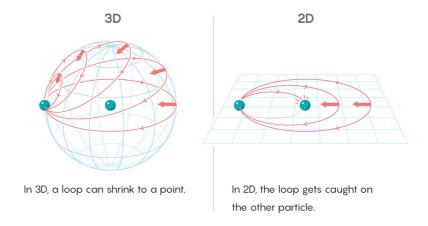
Xiangdong Zeng June 25, 2025

Anyons

- 3D case
 - Paths are topologically equivalent:

$$|\Psi(\lambda_2)
angle=|\Psi(\lambda_1)
angle=|\Psi(0)
angle$$

- ullet Encircle = exchange $^{ exttt{2}}$: $|\Psi(\lambda_2)
 angle=R^2\,|\Psi(0)
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- $R=\pm 1$: bosons and fermions



Anyons

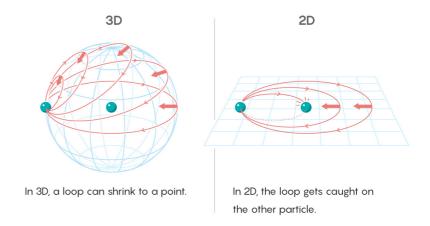
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- $R=\pm 1$: bosons and fermions
- 2D case
 - No longer topologically trivial:

$$|\Psi(\lambda_2)
angle
eq |\Psi(\lambda_1)
angle = |\Psi(0)
angle$$

- lacksquare $R=\mathrm{e}^{i heta}$: abelian anyons
- lacktriangleq R is a unitary matrix: non-abelian anyons



Topological states

- Symmetry-protected topological (SPT) states
 - Topological only given that some protecting symmetry
 - Do not support anyons as intrinsic quasiparticle excitations, unless ...
 - Having defects
 - Majorana zero modes
 - Examples: integer quantum Hall states, topological insulators

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 - Examples: integer quantum Hall states, topological insulators
- (Intrinsic) topological order
 - Exhibit long-range entanglement
 - Give rise to ground state degeneracy and topological entanglement entropy
 - Examples: fractional quantum Hall states, spin liquids

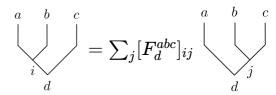
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lacktriangle F-move: basis transformation in \mathbb{V}_d^{abc}



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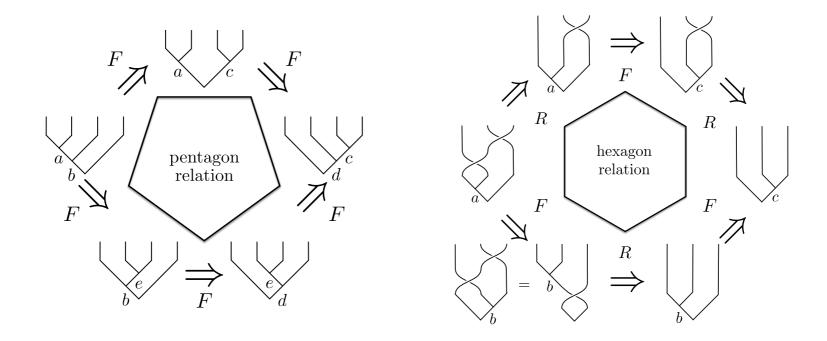
lacktriangleright F-move: basis transformation in \mathbb{V}^{abc}_d

$$\left[egin{array}{c} a & b & c \ igcup_{i} & igcup_{d} & igcup_{d}$$

• *R*-move: exchanging two anyons

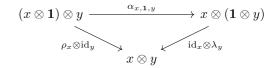
$$\stackrel{a}{\smile} \stackrel{b}{\smile} = R_c^{ab} \stackrel{a}{\smile} \stackrel{b}{\smile}$$

Anyon models: consistency relation

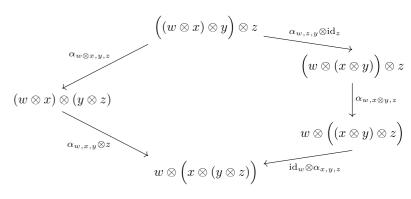


Mathematical structure (1)

- A monoidal (tensor) category is a category C with:
 - Tensor product: a bi-functor \otimes : $\mathcal{C} \times \mathcal{C} \to \mathcal{C}$
 - Unit object: $\mathbf{1} \in \mathcal{C}$
 - Natural isomorphisms:
 - lacksquare Associator: $lpha_{x,y,z}\colon (x\otimes y)\otimes z\stackrel{\sim}{ o} x\otimes (y\otimes z)$
 - lacksquare Left unitor: $\lambda_x \colon \mathbf{1} \otimes x \stackrel{\sim}{ o} x$
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 - The following two diagrams commute



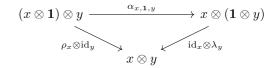
Triangle equation



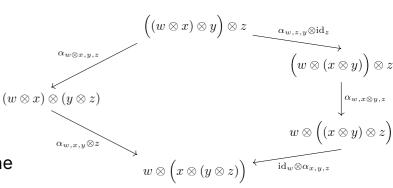
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- MacLane's coherence theorem:
 - Every monoidal category is equivalent to a strict one
 - " $\stackrel{\sim}{\rightarrow}$ " becomes "="



Triangle equation



Pentagon equation

Mathematical structure (2)

- A braided monoidal category is a monoidal category C with:
 - lacksquare Braiding: $\sigma_{x,y}\colon x\otimes y\stackrel{\sim}{ o} y\otimes x$
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$$x \otimes (y \otimes z) \xrightarrow{\alpha_{x,y,z}^{-1}} (x \otimes y) \otimes z \xrightarrow{\sigma_{x,y} \otimes \mathrm{id}_z} (y \otimes x) \otimes z$$

$$\downarrow^{\alpha_{y,x,z}} \downarrow \qquad \qquad \downarrow^{\alpha_{y,x,z}}$$

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 - With finitely many simple objects and finite dimensional spaces of morphisms
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 - Unit object 1 is simple
- A modular tensor category is
 - A braided monoidal fusion category

$$lacksquare$$
 With S matrix $s_{ij}=$ non-degenerate

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- Why Fibonacci?
 - $\begin{array}{ll} \bullet & \tau \otimes \tau \otimes \tau = \mathbf{1} + 2 \cdot \tau, \\ & \tau \otimes \tau \otimes \tau \otimes \tau = 2 \cdot \mathbf{1} + 3 \cdot \tau, \\ & \tau \otimes \tau \otimes \tau \otimes \tau \otimes \tau = 3 \cdot \mathbf{1} + 5 \cdot \tau \cdots \end{array}$

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• *F*-matrix (basis transformation):

•
$$F=F^ au_{ au au au}=egin{pmatrix}\phi^{-1}&\phi^{-1/2}\\phi^{-1/2}&-\phi^{-1}\end{pmatrix}$$
 , where $\phi=rac{1+\sqrt{5}}{2}$

R-matrix (braiding):

$$lacksquare R = egin{pmatrix} R_{ au au}^1 & 0 \ 0 & R_{ au au}^ au \end{pmatrix} = egin{pmatrix} e^{4\pi i/5} & 0 \ 0 & e^{-3\pi i/5} \end{pmatrix}$$

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- Ising anyon model
- Anyons: $\mathbf{1}, \psi, \sigma$
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- Same as Ising CFT
- F-matrix:

$$lacksquare F = F^{\sigma}_{\sigma\sigma\sigma} = rac{1}{\sqrt{2}} egin{pmatrix} 1 & 1 \ 1 & -1 \end{pmatrix}$$

■ *R*-matrix:

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Quantum computation with anyons (1)

- Advantages:
 - Protect against local perturbations
 - Gates are free from control errors since they depend only on topological characteristics

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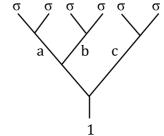
- Protect against local perturbations
- Gates are free from control errors since they depend only on topological characteristics
- Qubits: basis in fusion space

$$ullet \ |00
angle = |\sigma\sigma; {f 1}
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$$ullet |01
angle = |\sigma\sigma; \mathbf{1}
angle |\sigma\sigma; \psi
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•
$$|10\rangle = |\sigma\sigma;\psi\rangle |\sigma\sigma;\psi\rangle |\sigma\sigma;\mathbf{1}\rangle$$

$$lacksquare |11
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	a	b	c
00 }	1	1	1
10 }	ψ	ψ	1
01 }	1	ψ	ψ
11 }	ψ	1	ψ

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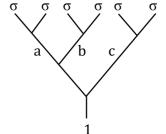
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- Measurements: fusion outcomes
 - Detecting energy difference between 1 (vacuum) and ψ (a massive particle)



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11 }	ψ	1	ψ

Quantum computation with anyons (2)

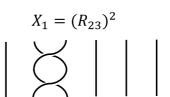
- Gates: braiding
- Single-qubit gates:

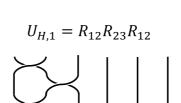
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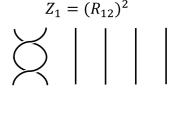
$$lacksquare Z_1=(R_{12})^2=R^2\otimes \mathbb{I}$$

•
$$U_{H,1} = R_{12}R_{23}R_{12} = RF^{-1}RFR \otimes \mathbb{I}$$

- R_{ij} : exchange of anyons i and j
- Two-qubit gate:
 - $U_{CZ} = (R_{12})^{-1} R_{34} (R_{56})^{-1}$







$$U_{CZ} = R_{12}^{-1} R_{56}^{-1} R_{34}$$

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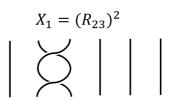
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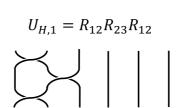
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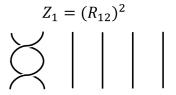
Non-Clifford gate (necessary for universality):

$$\bullet \ \ U = \begin{pmatrix} 1 & 0 \\ 0 & e^{i\Delta Et} \end{pmatrix}$$

lacktriangleright Lift the degeneracy of the fusion channels by ΔE so it's not topologically protected







$$U_{CZ} = R_{12}^{-1} R_{56}^{-1} R_{34}$$

Majorana zero modes: Majorana fermions

- Creation and annihilation operators:
 - lacksquare Bosons: $[b,b^\dagger]=bb^\dagger-b^\dagger b=1$
 - lacksquare Fermions: $\{f,f^{\dagger}\}=ff^{\dagger}+f^{\dagger}f=1$
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- Rewrite as:
 - $lacksquare f = rac{1}{2}(\gamma_1 + i\gamma_2) \implies \gamma_1 = f + f^\dagger, \, \gamma_2 = i(f f^\dagger)$
 - ullet $\gamma_i=\gamma_i^\dagger,\,\{\gamma_i,\gamma_j\}=2\delta_{ij},\,\gamma_i^2=1$



Ettore Majorana

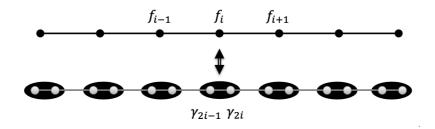
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- Majorana Fermion is its own anti-particle
 - Neutrino?



Ettore Majorana

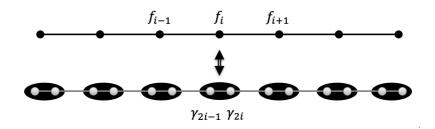
Majorana zero modes: Kitaev chain (1)



$$\blacksquare \quad \text{Hamiltonian: } H = \sum_{j=1}^L \left[-t(f_j^\dagger f_{j+1} + f_{j+1}^\dagger f_j) - \mu \left(f_j^\dagger f_j - \frac{1}{2} \right) + \left(\Delta_p f_j f_{j+1} + \Delta_p^* f_{j+1}^\dagger f_j^\dagger \right) \right]$$

- ullet t: tunnelling amplitude, μ : chemical potential, $\Delta_p=|\Delta_p|e^{i heta}$: superconducting pairing potential
- lacksquare Global symmetry: fermion parity $\mathcal{P} = \exp(i\pi\sum_j f_j^\dagger f_j)$

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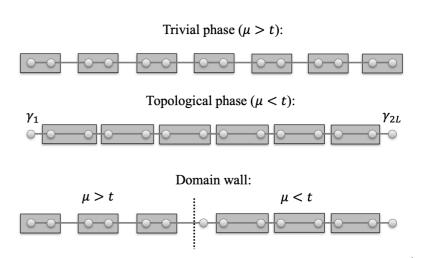
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$$\blacksquare \quad \text{Use Majorana operators: } H = \sum_{j=1}^L \frac{i}{2} \left[-\mu \gamma_{2j-1} \gamma_{2j} + \left(t + |\Delta_p|\right) \gamma_{2j} \gamma_{2j+1} + \left(-t + |\Delta_p|\right) \gamma_{2j-1} \gamma_{2j+2} \right]$$

Nearest and third-nearest neighbor interaction

Majorana zero modes: Kitaev chain (2)

- lacksquare Chemical potential term dominates: $\mu\gg t,\, |\Delta_p|$
 - $lacksquare H = -rac{i\mu}{2}\sum_{j=1}^L \gamma_{2j-1}\gamma_{2j}$
 - Product state, (topologically) trivial phase

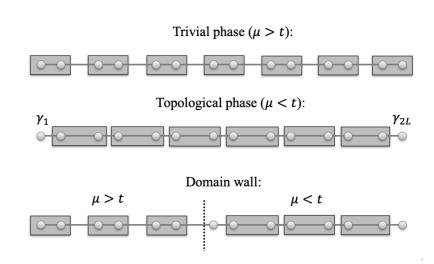


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- lacktriangle Kinetic/pairing term dominates: $t=|\Delta_p|\gg \mu$

$$lacksquare H=it\sum_{j=1}^L\gamma_{2j}\gamma_{2j+1}=2t\sum_{j=1}^{L-1}\left(ilde{f}_j^\dagger ilde{f}_j-rac{1}{2}
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- lacksquare Missing fermion: $d=e^{-i heta/2}(\gamma_1+i\gamma_{2L})/2$
 - γ_1 and γ_{2L} : Majorana zero modes
 - $[d^{\dagger}d, H] = 0$: two-fold degeneracy
- Topological phase

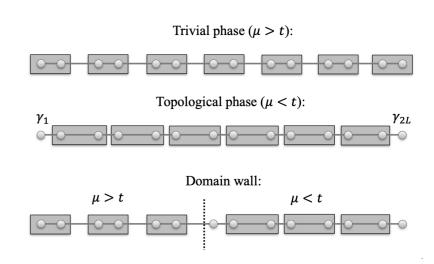


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$$lacksquare H=it\sum_{j=1}^L\gamma_{2j}\gamma_{2j+1}=2t\sum_{j=1}^{L-1}\left(ilde{f}_j^\dagger ilde{f}_j-rac{1}{2}
ight)$$

- lacksquare Missing fermion: $d=e^{-i heta/2}(\gamma_1+i\gamma_{2L})/2$
 - γ_1 and γ_{2L} : Majorana zero modes
 - $[d^{\dagger}d, H] = 0$: two-fold degeneracy
- Topological phase
- Domain wall separating the phases also binds a localized Majorana mode

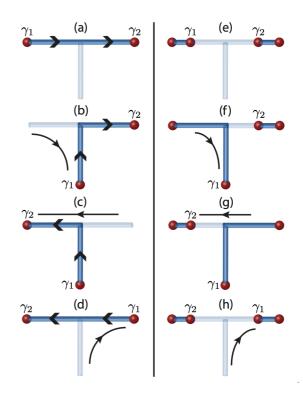


- Anyons:
 - 1: vacuum
 - σ : Majorana zero mode γ_1, γ_{2L}
 - ψ : fermionic excitation $ilde{f}_j^\dagger$

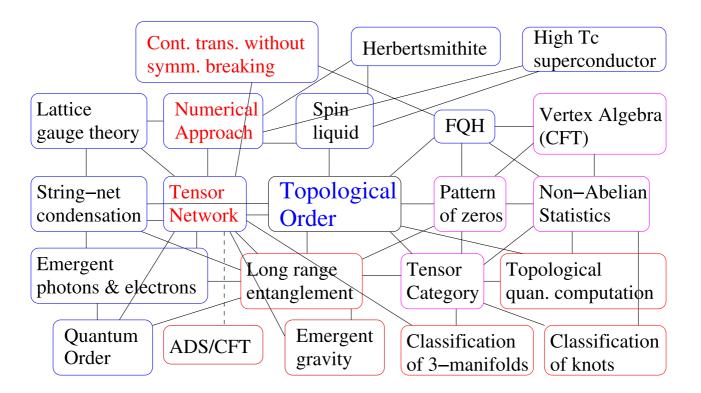
- Anyons:
 - 1: vacuum
 - σ : Majorana zero mode γ_1, γ_{2L}
 - ψ : fermionic excitation \tilde{f}_i^{\dagger}
- Fusion rules:
 - $d^\dagger d = (1+i\gamma_1\gamma_{2L})/2$: occupied or not
 - $\sigma \otimes \sigma \rightarrow \mathbf{1} + \psi$: eigenvalue 0 or 1
- Fusion channel states:
 - $ullet i\gamma_1\gamma_{2L}\ket{\sigma\sigma;\mathbf{1}}=-\ket{\sigma\sigma;\mathbf{1}},\ i\gamma_1\gamma_{2L}\ket{\sigma\sigma;\psi}=+\ket{\sigma\sigma;\psi}$

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- Qubits:
 - ullet Different ψ -parity sectors: two Kitaev chains are required
 - $\bullet \hspace{0.3cm} |0\rangle = |\sigma\sigma; \mathbf{1}\rangle \hspace{0.1cm} |\sigma\sigma; \mathbf{1}\rangle \hspace{0.1cm}, \hspace{0.1cm} |1\rangle = |\sigma\sigma; \psi\rangle \hspace{0.1cm} |\sigma\sigma; \psi\rangle$

- Anyons:
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- Qubits:
 - ullet Different ψ -parity sectors: two Kitaev chains are required
 - $ullet \ |0
 angle = |\sigma\sigma; \mathbf{1}
 angle \ |\sigma\sigma; \mathbf{1}
 angle \ , \ |1
 angle = |\sigma\sigma; \psi
 angle \ |\sigma\sigma; \psi
 angle$
- Braiding: T-junction



The paradigm of topological order



From Xiao-Gang Wen (arXiv:0903.1069)

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