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本 节 概 览

□ 学习内容

- □ 电场强度的分析和计算方法(续)
- ロ 电场线
- □ 电通量

仅供探讨,请勿上传网络

回顾:库仑定律



库仑定律的数学表达式

$$\vec{F} = k \frac{q_1 q_2}{r^3} \vec{r}$$

$$k = \frac{1}{4\pi\varepsilon_0}$$

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$$\varepsilon_0 = \frac{1}{4\pi k} = 8.85 \times 10^{-12} \frac{\text{C}^2}{\text{m}^2\text{N}}$$

$$\varepsilon_0$$
 为真空中的介电常数

$$\vec{E} = \frac{\vec{F}}{q_0}$$

电场强度
$$\vec{E} = \frac{\vec{F}}{q_0}$$
 点电荷电场 $\vec{E} = \frac{\vec{F}}{q_0} = \frac{1}{4\pi\varepsilon_0} \frac{q}{r^2} \vec{r}^0$

电场强度的计算步骤



1. 原则:利用点电荷的电场公式和叠加原理

$$d\vec{E} = \frac{1}{4\pi\varepsilon_0} \frac{dq}{r^2} \vec{r}^0 \longrightarrow \vec{E} = \int \frac{dq}{4\pi\varepsilon_0 r^2} \vec{r}^0$$

$$\mathrm{d}q = \begin{cases} \lambda \mathrm{d}l & (\text{带电体为线分布}) \\ \sigma \mathrm{d}S & (\text{带电体为面分布}) \end{cases} \qquad \bar{E} = \begin{cases} \int \frac{\lambda \mathrm{d}l}{4\pi\varepsilon_0 r^2} \vec{r}^0 \\ \int \frac{\sigma \mathrm{d}S}{4\pi\varepsilon_0 r^2} \vec{r}^0 \end{cases}$$
 (带电体为体分布)
$$Q \not = \begin{cases} \nabla \mathrm{d}S & (\text{带电体为体分布}) \\ \nabla \mathcal{L} \not= \nabla \mathcal{L} & \nabla \mathcal{L} \not= \nabla \mathcal{L} \end{cases}$$
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电场强度的计算步骤



2. 矢量积分化成标量积分,用分量式表示

$$\begin{cases} E_{x} = \int dE_{x} \\ E_{y} = \int dE_{y} \end{cases} \qquad \overrightarrow{E} = E_{x}\overrightarrow{i} + E_{y}\overrightarrow{j} + E_{z}\overrightarrow{k}$$

$$E_{z} = \int dE_{z}$$

- 3. 分析带电体的对称性,简化计算
- 4. 利用已知电场,化繁为简

均匀带电细圆环

均匀带电圆盘

无限大均匀带电平面

Page 4

$$E = \frac{1}{4\pi\varepsilon_0} \frac{qx}{(R^2 + x^2)^{3/2}} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1 - \frac{x}{(R^2 + x^2)^{1/2}}] \vec{i} \qquad \vec{E} = \frac{\sigma}{2\varepsilon_0} [1$$

dx





长为l均匀带电直线,电荷线密度为 λ

求:图示P点的电场强度



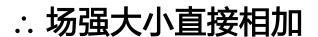
在坐标 x 处取一小段线元dx

$$dq = \lambda dx$$

该点电荷在 P 点的场强方向如图所示

大小为
$$dE = \frac{dq}{4\pi\varepsilon_0 r^2} = \frac{\lambda dx}{4\pi\varepsilon_0 (l+a-x)^2}$$





$$E = \int dE = \int_0^l \frac{\lambda dx}{4\pi\varepsilon_0 (l+a-x)^2} = \frac{\lambda}{4\pi\varepsilon_0} \left(\frac{1}{a} - \frac{1}{a+l} \right)$$

 χ



若P点与带电线不在同一直线上

$$dq = \lambda dx$$

$$\mathrm{d}E = \frac{1}{4\pi\varepsilon_0} \frac{\lambda \mathrm{d}x}{r^2}$$

$$dE_{r} = dE \cos \theta$$

$$dE_x = dE \cos \theta$$
 $dE_y = dE \sin \theta$

统一变量
$$x = a \tan(\theta - \frac{\pi}{2}) = -a \cot\theta$$

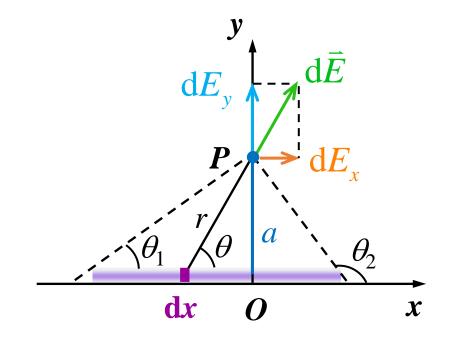
$$dx = a \csc^2 \theta d\theta$$

$$r^2 = a^2 + x^2 = a^2 \csc^2 \theta$$

$$dE_x = \frac{\lambda}{4\pi\varepsilon_a a} \cos\theta d\theta$$

$$dE_x = \frac{\lambda}{4\pi\varepsilon_0 a} \cos\theta d\theta \qquad dE_y = \frac{\lambda}{4\pi\varepsilon_0 a} \sin\theta d\theta$$

$$E_{x} = \int dE_{x} = \int_{\theta_{1}}^{\theta_{2}} \frac{\lambda}{4\pi\varepsilon_{0}a} \cos\theta d\theta = \frac{\lambda}{4\pi\varepsilon_{0}a} (\sin\theta_{2} - \sin\theta_{1})$$





计算得
$$E_x = \frac{\lambda}{4\pi\varepsilon_0 a} (\sin\theta_2 - \sin\theta_1)$$
 $E_y = \frac{\lambda}{4\pi\varepsilon_0 a} (\cos\theta_1 - \cos\theta_2)$

$$E_{y} = \frac{\lambda}{4\pi\varepsilon_{0}a}(\cos\theta_{1} - \cos\theta_{2})$$

(i) 无限长均匀带电直线 (l >> a)

$$\theta_1 = 0$$

$$\theta_2 = \pi$$

$$E_x = 0$$

$$E_{y} = \frac{\lambda}{2\pi\varepsilon_{0}a}$$

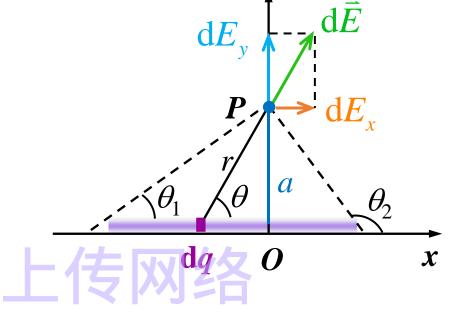
(ii) 半无限长均匀带电直线

$$\theta_{1} = \frac{\pi}{2}$$

$$\theta_{2} = \pi$$

$$E_{x} = \frac{\lambda}{4\pi\varepsilon_{0}a}$$

$$E_{y} = \frac{-\lambda}{4\pi\varepsilon_{0}a}$$







求面密度为 σ ,宽为d,长无限的平板中心轴线上方a 处P点的电场强度



可看作无数无限长带电线的集合

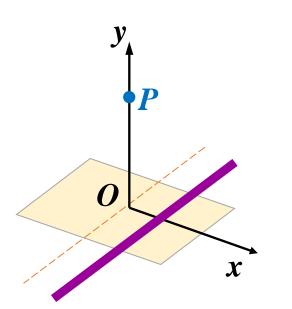
在x处取宽为dx的小窄条

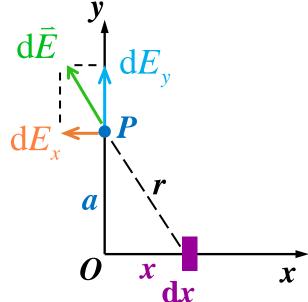
$$dE = \frac{\lambda}{2\pi\varepsilon_0 r} = \frac{\sigma \cdot dx}{2\pi\varepsilon_0 (a^2 + x^2)^{1/2}}$$

$$dE_x = dE \cdot (-\frac{x}{r})$$
 $dE_y = dE \cdot \frac{a}{r}$



$$E_{y} = \int dE_{y} = \frac{\sigma a}{2\pi\varepsilon_{0}} \int_{-d/2}^{d/2} \frac{dx}{a^{2} + x^{2}} = \frac{\sigma}{\pi\varepsilon_{0}} \tan^{-1}(\frac{d}{2a})$$



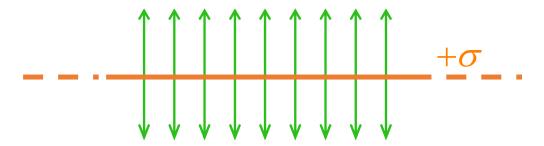




$$E = E_{y} = \frac{\sigma}{\pi \varepsilon_{0}} \tan^{-1}(\frac{d}{2a})$$

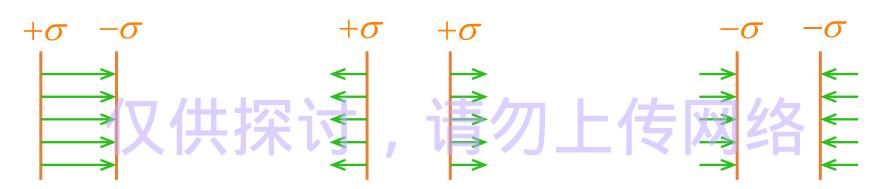
讨论

(i) 无限大均匀带电平面 $(d \to \infty)$



$$E = E_{y} = \frac{\sigma}{2\varepsilon_{0}}$$

(ii) 两个无限大均匀带电平面 (不考虑边缘效应)



10-3 电通量 高斯定理



电场线 (电力线)

1. 规定

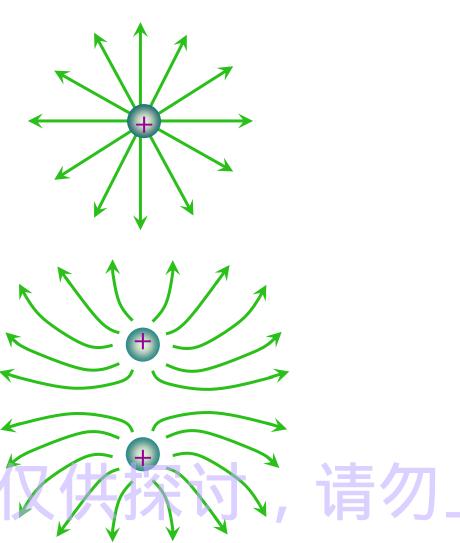
$$\left| \vec{E} \right| = E = dN / dS_{\perp}$$

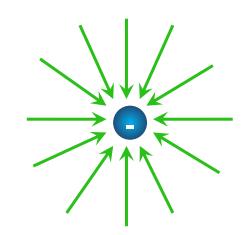
曲线上每一点切线方向为该点电场方向 通过垂直于电场方向单位面积电场线数为该点电场强度的大小

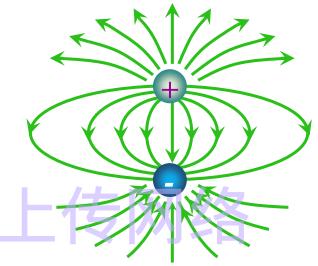
- 2. 电场线的性质
- ✓ 起始于正电荷(或无穷远处),终止于负电荷,不会在没有电荷处中断
- ✓ 不会形成闭合曲线,两条电场线不会相交
- ✓ 电场线一般不是电荷运动的轨迹
- 上传网络 ✓ 电场线密集处电场强, 电场线稀疏处电场弱

第10章:静电场

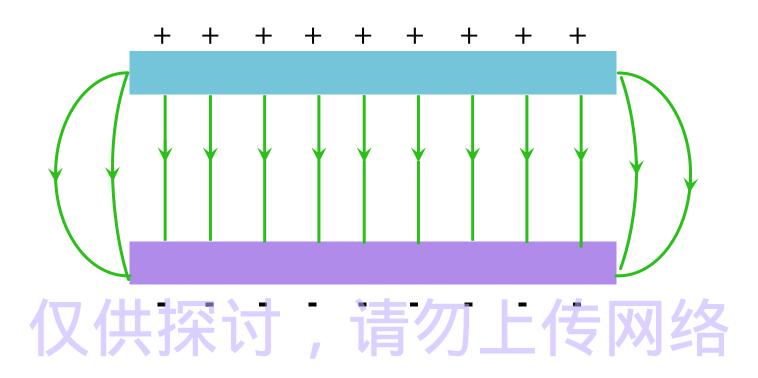
几种电荷分布的电场线图







带电平行板电容器的电场

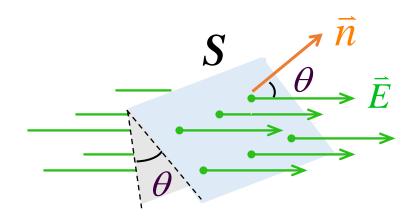




电通量

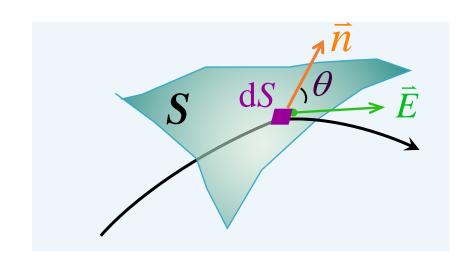
穿过某一有向曲面的电场线条数,用Φ。表示

电通量的计算公式



$$\Phi_e = ES\cos\theta = \vec{E}\cdot\vec{S}$$

均匀电场,S 法线方向与 电场强度方向成 θ 角



$$\Phi_e = \iint d\Phi_e = \iint \vec{E} \cdot d\vec{S}$$

请 电场不均匀,S为任意曲面



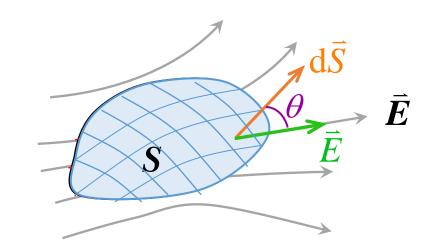
(讨论)

(i) 电通量是标量,有正负之分

- *θ* <90⁰,通量为正;
- $\theta = 90^{\circ}$,通量为零;
- $\theta > 90^{\circ}$,通量为负;
- (ii) 闭合曲面的电场强度通量

$$d\Phi_e = \vec{E} \cdot d\vec{S}$$

$$\Phi_e = \iint_S \vec{E} \cdot d\vec{S} = \iint_S E \cos \theta dS$$



规定:法线的正方向为指向闭合曲面的外侧





□ 学习内容

本

顾

- 电场强度的分析和计算方法(续)
- 电场线
- 电通量

- □ 课下任务
 - 口作业册"电场强度的计算电通量"与以《各