

Fundamentalne konstante

Planckova konstanta: $\hbar = 1.055 \cdot 10^{-34} \text{ Js}$

Brzina svjetlosti: $c = 2.998 \cdot 10^8 \text{ m/s}$

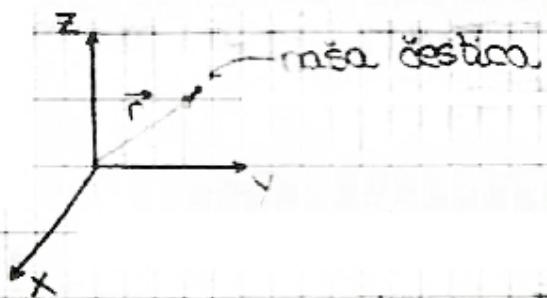
Konstanta gravitacije: $G = 6.672 \cdot 10^{-11} \text{ Nm}^2 \text{kg}^{-2}$

$$N = \text{kg m s}^{-2}$$

Izvedena Planckova duljina (najfundamentalija duljina za rad, manje od toga fizika ne može)

$$l_p = \hbar^2 c^3 G^8 = \dots = 1.6 \cdot 10^{-35} \text{ m}$$

Kinematika točke



Položaj: $x(t) \vec{i} + y(t) \vec{j} + z(t) \vec{k}$

$$\text{Brzina: } \vec{v}(t) = \frac{d\vec{r}(t)}{dt} = \dot{\vec{r}}(t) = \dot{x}(t)\vec{i} + \dot{y}(t)\vec{j} + \dot{z}(t)\vec{k} = v_x(t)\vec{i} + v_y(t)\vec{j} + v_z(t)\vec{k}$$

$$\text{Integralni zapisi: } \vec{s}(t) = \frac{d\vec{r}(t)}{dt} / dt$$

$$d\vec{r}(t) = \vec{v}(t) dt / \int_p^k$$

$$\int_p^k d\vec{r}(t) = \int_p^k \vec{v}(t) dt$$

$$p = t_0 \quad k = t$$

$$\vec{r} \Big|_{t_0}^t = \int_{t_0}^t \vec{v}(t) dt$$

$$\vec{r}(t) - \vec{r}(t_0) = \int_{t_0}^t \vec{v}(t) dt$$

$$\boxed{\vec{r}(t) = \vec{r}(t_0) + \int_{t_0}^t \vec{v}(t) dt}$$

Akceleracija: $\vec{a}(t) = \frac{d\vec{v}(t)}{dt} = \dot{\vec{v}}(t) = \dot{v}_x(t)\vec{i} + \dot{v}_y(t)\vec{j} + \dot{v}_z(t)\vec{k}$
 $= a_x(t)\vec{i} + a_y(t)\vec{j} + a_z(t)\vec{k}$

$$\vec{v}(t) = \vec{v}(t_0) + \int_{t_0}^t \vec{a}(t) dt$$

Newtonovi aksioni

Prvi: Kada na matičnu točku ne djeluje sila, ona ostaje na mjestu ili se giba po pravcu (princip trinostnosti)

Dруги: Vremenska promjena količine gibanja matične točke razmjera je sili koja na nju djeluje
 $\vec{p} = m\vec{v}$, $\vec{F} = \frac{d\vec{p}}{dt} = \frac{d(m\vec{v})}{dt} = m \frac{d\vec{v}}{dt} = m\vec{a}$

Treći: Ako jedno tijelo djeluje na drugo nekom silom, onda drugo djeluje na prvo silom istog intesa, ali drugog smjera.
 $\vec{F}_{12} = -\vec{F}_{21}$, $|\vec{F}_{12}| = |\vec{F}_{21}|$

Rješavanje jednogodišnjeg gibanja

Uz stacionarni, tj. u istom položaju i u istočnom vremenu sila je konstantna $\frac{d\vec{F}}{dt} = \vec{F}_0$, $m\vec{a} = \frac{d\vec{v}}{dt}$, $m\vec{v} = \frac{d\vec{r}}{dt}$, $m\vec{r} = \frac{d\vec{r}^2}{dt^2}$

$$\text{Brzina: } \frac{d\vec{v}}{dt} = \frac{\vec{F}_0}{m} \Rightarrow m\frac{d\vec{v}}{dt} = \vec{F}_0$$

$$\frac{d\vec{v}}{dt} = \frac{\vec{F}_0}{m} \quad | \cdot dt$$

$$d\vec{v} = \frac{\vec{F}_0}{m} dt \quad | \int$$

$$\vec{v}|_p = \frac{1}{m} \int_{t_0}^k \vec{F}_0 dt$$

$$\vec{v}_k - \vec{v}_p = \frac{1}{m} \vec{F}_0 |$$

$$\vec{v}_k - \vec{v}_p = \frac{1}{m} \vec{F}_0 (k - p)$$

$$\begin{aligned} p &= t_0 & k &= t \\ \vec{v}(+) &= \vec{v}(t_0) + \frac{\vec{F}_0}{m} (t - t_0) \end{aligned}$$

$$\text{Položaj: } \frac{d\vec{r}}{dt} = \vec{v}(+) \quad | \cdot dt$$

$$d\vec{r} = \vec{v}(+) dt \quad | \int$$

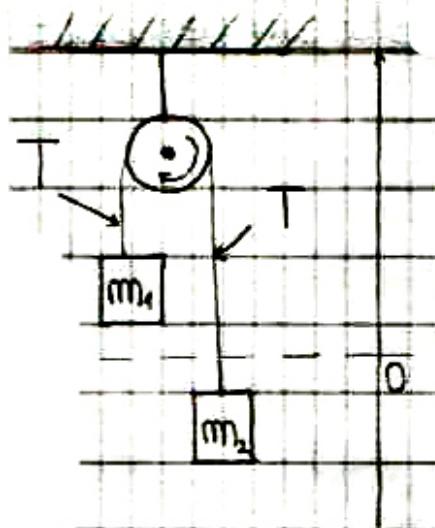
$$\vec{r} = \vec{r}(t_0) + \int_{t_0}^t \vec{v}(+) dt$$

$$\vec{r}(t) - \vec{r}(t_0) = \int_{t_0}^t \vec{v}(+) dt$$

$$\vec{r}(+) = \vec{r}(t_0) + \int_{t_0}^t \left[\vec{v}(t_0) + \frac{\vec{F}_0}{m} (t - t_0) \right] dt$$

$$\vec{r}(+) = \vec{r}(t_0) + \vec{v}(t_0)(t - t_0) + \frac{1}{2} \frac{\vec{F}_0}{m} (t - t_0)^2$$

Atwoodov podstavlj



$$\vec{r}_1(t) = \mathbf{x}_1(t) \hat{i}$$

$$\vec{r}_2(t) = \mathbf{x}_2(t) \hat{i}$$

$$\mathbf{x}_1(t) = -\mathbf{x}_2(t)$$

$$\dot{\mathbf{x}}_1(t) = -\dot{\mathbf{x}}_2(t)$$

$$\ddot{\mathbf{x}}_1(t) = -\ddot{\mathbf{x}}_2(t)$$

$$\left. \begin{aligned} m_1 \ddot{x}_1 &= T - m_1 g \\ m_2 \ddot{x}_2 &= T - m_2 g \end{aligned} \right\} -$$

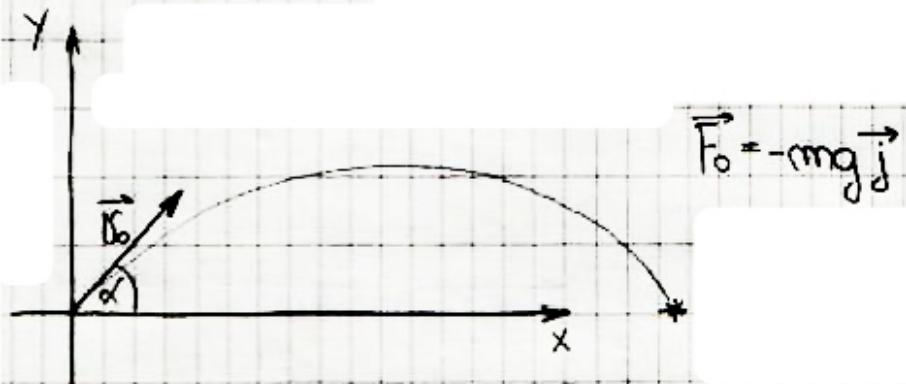
$$m_1 \ddot{x}_1 - m_2 \ddot{x}_2 = m_2 g - m_1 g$$

$$m_1 \ddot{x}_1 + m_2 \ddot{x}_1 = g(m_2 - m_1)$$

$$\ddot{x}_1 = g \left(\frac{m_2 - m_1}{m_1 + m_2} \right) \quad \text{akceleracija sistema}$$

3.

Kosi hitac



Početni uvjeti: $t_0=0$, $\vec{v}(t_0)=v_0(\cos\alpha\vec{i} + \sin\alpha\vec{j})$
 $\vec{r}(t_0)=0$

$$\vec{v}(t) = \vec{v}(t_0) + \frac{\vec{F}_0}{m}(t-t_0)$$

wrastimo u \vec{v}

$$\vec{v}(t) = v_0(\cos\alpha\vec{i} + \sin\alpha\vec{j}) - gt\vec{j} = \underbrace{v_0 \cos\alpha\vec{i}}_{\vec{v}_x} + \underbrace{(v_0 \sin\alpha - gt)\vec{j}}_{\vec{v}_y}$$

$$\vec{r}(t) = \vec{r}(t_0) + \vec{v}(t_0)(t-t_0) + \frac{\vec{F}_0}{2m}(t-t_0)^2$$

wrastimo u \vec{r}

$$\vec{r}(t) = v_0 t (\cos\alpha\vec{i} + \sin\alpha\vec{j}) - \frac{gt^2}{2}\vec{j} = \underbrace{v_0 t \cos\alpha\vec{i}}_{x(t)} + \underbrace{(v_0 t \sin\alpha - \frac{gt^2}{2})\vec{j}}_{y(t)}$$

$$t = \frac{x(t)}{v_0 \cos\alpha}$$

$$y(x) = v_0 \frac{x(t)}{v_0 \cos\alpha} \sin\alpha - \frac{g}{2} \left(\frac{x(t)}{v_0 \cos\alpha} \right)^2 = x \tan\alpha - \frac{g}{2} \frac{x^2}{v_0^2 \cos^2\alpha} \quad (u = \tan\alpha)$$

$$= xu - \frac{gx^2}{2v_0^2} (1+u^2)$$

Rješavanje jednadžbi gibanja

Kada je sila otpora razmjetna brzina, $\vec{F} = -\gamma \vec{v}$ (\vec{F} je obrnutog smjera od \vec{v})

$$\frac{d\vec{p}}{dt} = \vec{F} = m\dot{\vec{v}} = -\gamma \vec{v}$$

$$m(\dot{v}_x \vec{i} + \dot{v}_y \vec{j} + \dot{v}_z \vec{k}) = -\gamma(v_x \vec{i} + v_y \vec{j} + v_z \vec{k})$$

koefficijent

$$m\dot{v}_x = -\gamma v_x$$

$$m\dot{v}_y = -\gamma v_y$$

$$m\dot{v}_z = -\gamma v_z$$

Rješavamo sasno za x komponentu (analogno y i z)

Brzina:

$$m \frac{d\vec{v}_x}{dt} = -\gamma \vec{v}_x \quad / \cdot dt$$

$$m d\vec{v}_x = -\gamma \vec{v}_x dt \quad / : m \vec{v}_x$$

$$\frac{d\vec{v}_x}{\vec{v}_x} = -\frac{\gamma}{m} dt \quad / \int_p^k$$

$$\int_p^k \frac{1}{\vec{v}_x} d\vec{v}_x = -\frac{\gamma}{m} \int_p^k dt$$

$$\ln \vec{v}_x \Big|_p^k = -\frac{\gamma}{m} (k-p)$$

Položaj:

$$\vec{r}(+) = \vec{r}(t_0) + \int_t^{+} \vec{v}(+) dt$$

$$= \vec{r}(t_0) + \int_{t_0}^{+} [\vec{v}(t_0) e^{-\frac{\gamma}{m}(t-t_0)}] dt$$

$$= \vec{r}(t_0) + \vec{v}(t_0) \int_{t_0}^{+} \frac{m}{\gamma} e^{-\frac{\gamma}{m}(t-t_0)} \frac{\gamma}{m} dt$$

$$= \vec{r}(t_0) - \frac{m}{\gamma} \vec{v}(t_0) e^{-\frac{\gamma}{m}(t-t_0)}$$

$$= \vec{r}(t_0) + \frac{m}{\gamma} \vec{v}(t_0) (1 - e^{-\frac{\gamma}{m}(t-t_0)})$$

$$\ln \vec{v}_x(+) = \ln \vec{v}_x(t_0) - \frac{\gamma}{m} (t-t_0) / e^{\gamma / m} / i$$

$$\vec{v}(+) = \vec{v}(t_0) \cdot e^{-\frac{\gamma}{m}(t-t_0)}$$

4.

Rješavanje jednadžbi gibanja

Uz konstantnu silu i otpor (koji je proporcionalan brzini)

$$\frac{d\vec{v}}{dt} = m\vec{a} = \vec{F}_0 - \gamma \vec{v}$$

Opet rješavamo za x komponentu.

Brzina:

$$m\dot{v}_x = F_{0x} - \gamma v_x$$

$$m \frac{dv_x}{dt} = F_{0x} - \gamma v_x / \frac{dt}{m}$$

$$dv_x = \left(\frac{F_{0x}}{m} - \frac{\gamma v_x}{m} \right) dt$$

$$\frac{dv_x}{\frac{F_{0x}}{m} - \frac{\gamma v_x}{m}} = dt \quad | \int$$

$$\frac{m}{\gamma} \ln \left(\frac{F_{0x} - \gamma v_x}{F_{0x}} \right) = t + t_0$$

$$\ln \left(F_{0x} - \gamma v_x(t) \right) = \ln \left(F_{0x} - \gamma v_x(t_0) \right) - \frac{\gamma}{m} (t - t_0) / e^x$$

$$F_{0x} - \gamma v_x(t) = (F_{0x} - \gamma v_x(t_0)) e^{-\frac{\gamma}{m}(t-t_0)}$$

$$\vec{v}(t) = \frac{\vec{F}_{0x}}{\gamma} - \left(\frac{F_{0x}}{\gamma} - \frac{\gamma v_x(t_0)}{\gamma} \right) e^{-\frac{\gamma}{m}(t-t_0)}$$

$$\vec{r}(t) = \vec{r}(t_0) e^{-\frac{\gamma}{m}(t-t_0)} + \frac{\vec{F}_0}{\gamma} \left(1 - e^{-\frac{\gamma}{m}(t-t_0)} \right)$$

$$\vec{r}(t) = \vec{r}(t_0) + \frac{\vec{F}_0}{\gamma} (t - t_0) + \left(\vec{r}(t_0) - \frac{\vec{F}_0}{\gamma} \right) \frac{m}{\gamma} \left(1 - e^{-\frac{\gamma}{m}(t-t_0)} \right)$$

Rješavanje jednadžbi gibanja

Kada je sila otpora razmjerna kvadratu brzine

$$\frac{d\vec{p}}{dt} = m \dot{\vec{v}}_x = -K v_x^2$$

dizisa:

$$m d\vec{v}_x = -K v_x^2 \quad | \quad \frac{dt}{m}$$

$$d\vec{v}_x = -\frac{K}{m} v_x^2 dt \quad | : v_x^2$$

$$\frac{d\vec{v}_x}{v_x^2} = -\frac{K}{m} dt \quad | \int$$

$$-\frac{1}{v_x} \Big|_{t_0}^t = -\frac{K}{m} t \Big|_{t_0}$$

$$-\frac{1}{v_x(t)} + \frac{1}{v_x(t_0)} = -\frac{K}{m} (t-t_0)$$

$$v_x(t) = \frac{v_x(t_0)}{1 + \frac{K}{m} v_x(t_0)(t-t_0)}$$

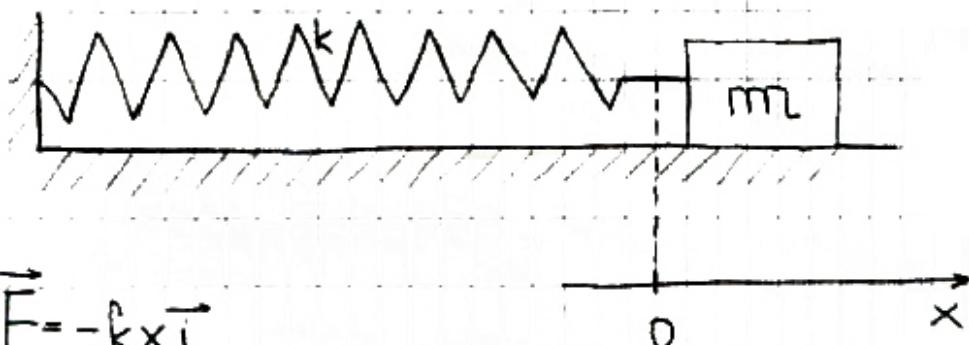
Položaj:

$$\vec{r}(t) = \vec{r}(t_0) + \frac{m}{K} \ln \left(1 + \frac{K}{m} v_x(t_0)(t-t_0) \right)$$

5.

Rješavanje jednadžbi gibanja

Uz harmoničku silu (sile opuge).



$$\vec{F} = -kx\vec{i}$$

← odstojan od ravn položaja (od 0)

$$\frac{d\vec{p}}{dt} = \vec{F} = -kx\vec{i}$$

dt

$$m\ddot{x}(t) = -kx(t)$$

$$\ddot{x}(t) = -\frac{k}{m}x(t)$$

$$\omega^2 = \frac{k}{m}$$

$$\boxed{\ddot{x}(t) + \omega^2 x(t) = 0} \quad \text{jednadžba harmoničkog oscilatora}$$

$$\text{opće rješenje: } x(t) = A \cos \omega t + B \sin \omega t \quad (1)$$

$$\dot{x}(t) = -A\omega \sin \omega t + B\omega \cos \omega t = \omega(-A \sin \omega t + B \cos \omega t) \quad (2)$$

$$\ddot{x}(t) = -\omega^2(A \cos \omega t + B \sin \omega t) = -\omega^2 x(t)$$

1. slučaj: Mirovanje u ravnotežnom položaju ($t_0 = 0, x_0 = x(t_0) = 0, \dot{x}(t_0) = 0$)

$$\text{uvrštavajući u (1): } x(t_0) = A \Rightarrow A = 0$$

$$\text{uvrštavajući u (2): } \dot{x}(t_0) = \omega B \Rightarrow B = 0$$

$$\text{uvrštavajući u (1): } x(t) = 0$$

2. slučaj: Mirovanje izuzv ravnotežnog položaja ($t_0 = 0, x_0 = x(t_0) \neq 0, \dot{x}(t_0) = 0$)

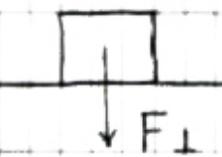
$$\text{uvrštavajući u (1): } x(t_0) = A \Rightarrow A \neq 0$$

$$\text{uvrštavajući u (2): } \dot{x}(t_0) = \omega B \Rightarrow B = 0$$

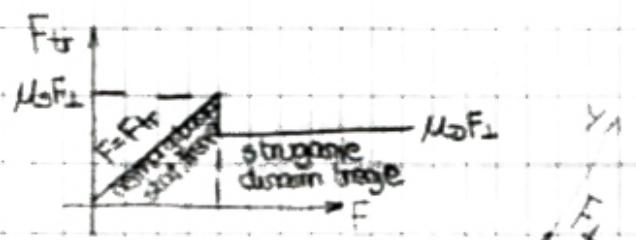
$$\text{uvrštavajući u (1): } x(t) = A \cos \omega t - x(t_0) \cos \omega t$$

3. slučaj: Početna brzina u ravnotežnom položaju ($t_0=0, x_0=0, \dot{x}_0=0$)
 uvrštanjem u (1): $x(t) = A \Rightarrow A = 0$
 uvrštanjem u (2): $\ddot{x}(t_0) = \omega^2 B \Rightarrow B = \frac{\ddot{x}(t_0)}{\omega}$
 uvrštanjem u (1): $x(t) = \frac{\ddot{x}(t_0)}{\omega} \sin \omega t$

Trenje

1) Staticko (u mirovanju)  $F_{tr} \leq \mu_s F_{\perp}$

2) Dinamičko (pri straganju) $F_{tr} = \mu_d F_{\perp}$



KOSINA:



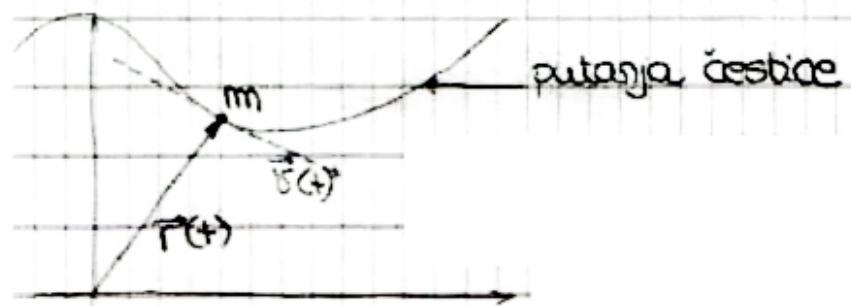
$$\frac{d\vec{p}}{dt} = m\vec{a} = \vec{F} = m\vec{g} + \vec{F}_{\perp} + \vec{F}_{tr} \quad (m\vec{g} = m\vec{g} (\vec{i} \sin \alpha - \vec{j} \cos \alpha), |\vec{F}_{\perp}| = |\vec{F}_{\perp}|, |\vec{F}_{tr}| = \pm \mu |\vec{F}_{\perp}|)$$

$$m(\ddot{\vec{x}}_i + \ddot{\vec{x}}_j) = mgs \sin \alpha \vec{i} - mg \cos \alpha \vec{j} + |\vec{F}_{\perp}| \vec{j} \pm \mu |\vec{F}_{\perp}| \vec{i} \\ = \vec{i} (mgs \sin \alpha \pm \mu |\vec{F}_{\perp}|) + \vec{j} (|\vec{F}_{\perp}| - mg \cos \alpha) \quad (\vec{j} = 0, |\vec{F}_{\perp}| = mg \cos \alpha)$$

$$m\ddot{x}_i = mgs \sin \alpha \pm \mu mg \cos \alpha$$

$a = g (\sin \alpha \pm \mu \cos \alpha)$	$+ \text{ (uzgorado)}$
	$- \text{ (nizbrdo)}$

Sila okomita na brzinu (Gibanje po krivulji)



$$\frac{d\vec{p}}{dt} = m\vec{a} = \vec{F}$$

$$\vec{a} = \frac{d\vec{v}}{dt} = \frac{d}{dt}(\vec{r} \cdot \hat{\vec{r}}) = \frac{d(\vec{r} \cdot \hat{\vec{r}})}{dt} = \frac{d\vec{r}}{dt} \hat{\vec{r}} + \vec{r} \frac{d\hat{\vec{r}}}{dt}$$

okomit na
brzine

Rastavili smo akceleraciju na 2 vektora: $\vec{a} = \vec{a}_{tang} + \vec{a}_{\perp}$

Na tangencijsku i transverzalnu.

Tangencijska je u smjeru tangente putanje i govori nam o promjeni brzine, tj. koliko čestica obzira gubojuci se duž vlastite putanje.

Transverzalna (radijalna, centralna...) je okomita na smjer gibajuća i govori o promjeni smjera brzine.

TVRDNJA: Ako je sila okomita na brzinu ($\vec{F} \cdot \vec{v} = 0$) onda se iznos brzine ne mijenja ($\frac{d\vec{v}}{dt} = 0$).

DOKAZ:

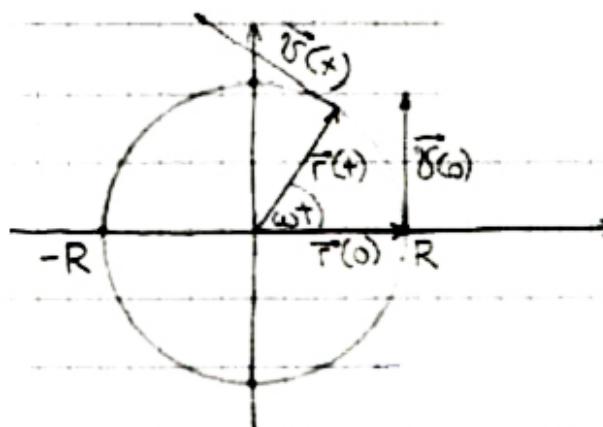
$$\vec{a} = \frac{d\vec{v}}{dt} \hat{\vec{v}} + \vec{v} \frac{d\hat{\vec{v}}}{dt} = \frac{\vec{F}}{m} / \hat{\vec{v}}$$

$$\vec{a} \cdot \hat{\vec{v}} = \frac{d\vec{v}}{dt} \hat{\vec{v}} \cdot \hat{\vec{v}} + \vec{v} \frac{d\hat{\vec{v}}}{dt} \hat{\vec{v}} = \frac{\vec{F} \cdot \hat{\vec{v}}}{m}$$

$$\vec{a} \cdot \hat{\vec{v}} = \frac{d\vec{v}}{dt} \cdot \hat{\vec{v}} + 0 = \frac{0}{m} \rightarrow \text{iz brodoje}$$

$$\frac{d\vec{v}}{dt} = 0$$

Jednoliko kružno gibanje



$$\vec{r}(t) = R(\cos \omega t \hat{i} + \sin \omega t \hat{j})$$

$$R = |\vec{r}(t)|$$

$$x(t) = R \cos \omega t$$

$$y(t) = R \sin \omega t$$

$$\vec{v}(t) = \dot{\vec{r}}(t) = R\omega (-\sin \omega t \hat{i} + \cos \omega t \hat{j})$$

$$|\vec{v}(t)| = R\omega = \vec{v} \quad (\text{Iznos brzine se ne mijenja})$$

$$\vec{a}(t) = \ddot{\vec{r}}(t) = -\omega^2 R (\cos \omega t \hat{i} + \sin \omega t \hat{j}) = -\omega^2 \vec{r}(t)$$

$$\frac{d\vec{p}}{dt} = m\vec{a} = \vec{F}$$

$$\vec{F} = -m\omega^2 R \hat{i} \equiv \vec{F}_{cp} = \frac{m\vec{v}^2}{R}$$

$$\vec{F}_{cp} \perp \vec{v} \quad \text{dokaz: } \vec{F}_{cp} \cdot \vec{v} = (-m\omega^2 R (\cos \omega t \hat{i} + \sin \omega t \hat{j})) (R\omega (-\sin \omega t \hat{i} + \cos \omega t \hat{j})) \\ = -m\omega^3 R^2 (-\sin \omega t \cos \omega t + \sin \omega t \cos \omega t) \\ = 0$$

Kutna brzina kao vektor ($\vec{\omega}$)

$$\text{iznos: } \omega = \frac{v}{R}$$

$$\text{smjer: } \hat{\vec{\omega}} = \hat{\vec{r}} \times \hat{\vec{v}}$$

$$\hat{\vec{\omega}} \cdot \hat{\vec{r}} = (\hat{\vec{r}} \times \hat{\vec{v}}) \cdot \hat{\vec{r}} = \hat{\vec{r}} (\hat{\vec{r}} \cdot \hat{\vec{v}}) - \hat{\vec{v}} (\hat{\vec{r}} \cdot \hat{\vec{r}}) = \frac{1}{R}$$

$$\vec{r} \cdot \vec{\omega} \times \vec{r}$$

7.

Kutna količina gibanja

$$\hat{\vec{\omega}} = \vec{r} \cdot \hat{\vec{\zeta}} / rm\varphi$$

$$rm\varphi \hat{\vec{\omega}} = \vec{r} \times m\vec{\zeta} = \vec{r} \times \vec{p}$$

$$\vec{L} = \vec{r} \times \vec{p} - rm\varphi \hat{\vec{\omega}} = rm\varphi \frac{\vec{\omega}}{\varphi} = r^2 m \vec{\omega} = I \vec{\omega}$$

$I = r^2 m$ (moment tramnosti)

TEOREM: Ako se čestica giba u polju centralne sile (usmjereni u 1 točku koordinatnog sustava) kutna količina gibanja (\vec{L}) je očuvana veličina

$$\begin{aligned} \text{DOKAZ: } \frac{d\vec{L}}{dt} &= \frac{d}{dt} (\vec{r} \times \vec{p}) = \frac{d\vec{r}}{dt} \times \vec{p} + \vec{r} \times \frac{d\vec{p}}{dt} = \vec{v} \times (m\vec{v}) + \vec{r} \times \vec{F} \\ &= m\vec{v} \times \vec{v} + \vec{r} \times (\vec{F}_{\text{cent}} \cdot \vec{r}) \\ &= m\underbrace{\vec{v} \times \vec{v}}_0 + \vec{F}_{\text{cent}} \underbrace{\vec{r} \times \vec{r}}_0 = \vec{0} \end{aligned}$$

$$\frac{d\vec{L}}{dt} = \vec{r} \times \vec{F} = \vec{M}$$

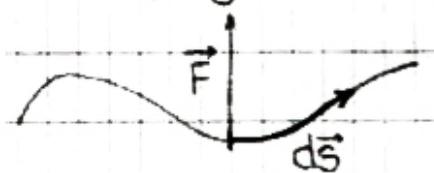
$$\frac{d\vec{L}}{dt} = I \cdot \dot{\vec{\omega}} = I \cdot \vec{\zeta} = \vec{M}$$

kutna
kutna
brzina akceleracija

Newtonova jednadžba za moment sile: $\frac{d\vec{L}}{dt} = I \cdot \vec{\zeta} = \vec{M}$

Rad

Sila \vec{F} djeluje duž diferencijala $d\vec{s}$.



Diferencijal obavljenog rada: $dW = \vec{F} \cdot d\vec{s}$

$$W = \int \vec{F} \cdot d\vec{s}$$

Rad za savijanje opnuge: $W = \frac{1}{2} kx^2$

Kinetička energija

Neka sila F djeluje na tijelo mase m duž x osi dok se tijelo giba

$$\begin{aligned} F &= \frac{d\vec{p}}{dt} = m \frac{d\vec{s}}{dt} \quad / \cdot \frac{dx}{dx} \\ &= \frac{m d\vec{s}}{dt} \cdot \frac{dx}{dx} = m \underbrace{\frac{d\vec{s}}{dx}}_{\text{brzina}} \cdot \underbrace{\frac{dx}{dt}}_{\text{prijemna ravnina brzine po prevođenom putu}} = m \vec{v} \cdot \frac{d\vec{s}}{dx} \end{aligned}$$

$$\begin{aligned} \text{Rad: } W &= \int_A^B F ds = \int_A^B \left(m \vec{v} \frac{d\vec{s}}{dx} \right) dx = m \int_A^B \vec{v} \frac{d\vec{s}}{dx} dx = \frac{m v^2}{2} \Big|_A^B \\ &= \frac{m v_B^2}{2} - \frac{m v_A^2}{2} = E_{kin.B} - E_{kin.A} = \Delta E_{kin}. \end{aligned}$$

Teorem o radu i kinetičkoj energiji: $W = \Delta E_{kin}$.

Obavljeni rad = promjeni kin. energije

Snaga

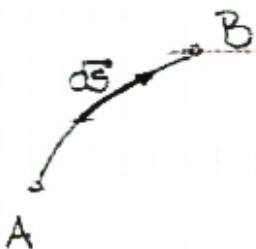
Mjera količine rada obavljenog u jedinici vremena.

$$P = \frac{dW}{dt} = \frac{d}{dt} (\vec{F} \cdot d\vec{s}) = \vec{F} \cdot \frac{d\vec{s}}{dt} = \vec{F} \cdot \vec{v}$$

Konzervativne sile i potencijalna energija

$$W = \int \vec{F} \cdot d\vec{s}$$
 (rad obavljen od A do B)

Sila za koju vrijedi da je obavljeni rad W uvek jednak neovisno o putanji zove se **KONZERVATIVNA SILA**.



$$dW_{AB} = \vec{F} \cdot d\vec{s}_{AB}$$

$$dW_{BA} = \vec{F} \cdot d\vec{s}_{BA}$$

$$= \vec{F} \cdot (-d\vec{s}_{AB})$$

$$= -\vec{F} \cdot d\vec{s}_{AB}$$

$$= -dW_{AB}$$



(krenemo od A i vratimo se u A, $W=0$)

$$\oint \vec{F} \cdot d\vec{s} = 0$$

Integral po zatvorenoj putanji

Kod konzervativne sile rad duž svake zatvorene linije je nula.

Potencijalna energija (u polju konzervativne sile) definira se kao rad koji treba obaviti da se neko tijelo dovede u neku točku.

$$U(B) = U(A) - \int_A^B \vec{F} \cdot d\vec{s}$$

$$U(x) = U(x_0) - \int_{x_0}^x \vec{F}(x) dx$$

$$\vec{F}(\vec{r}) = -\frac{dU}{dx} \vec{i} - \frac{dU}{dy} \vec{j} - \frac{dU}{dz} \vec{k} = -\frac{dU}{dr} \vec{r} = -\vec{\nabla} U$$

↓
konz.sila

$$F(x) = -\frac{d}{dx} U(x)$$

Zakon očuvanja mehaničke energije

$$W_{\text{korz}} = \Delta E_{\text{kin.}} = -\Delta U \Rightarrow \Delta(E_{\text{kin.}} + U) = 0 \quad \Delta E = 0$$

Ukupna energija: $E = E_{\text{kin.}} + U$

$$\begin{aligned} \frac{dE}{dt} &= \frac{d}{dt}(E_{\text{kin.}} + U) = \frac{d}{dt}\left(\frac{m\vec{v}\cdot\vec{v}}{2} + U(\vec{r})\right) \\ &= \frac{m}{2}(\vec{v}\cdot\ddot{\vec{v}} + \vec{v}\cdot\dot{\vec{v}}) + \frac{dU}{d\vec{r}} \cdot \frac{d\vec{r}}{dt} \\ &= m\vec{v}\cdot\ddot{\vec{v}} + \nabla U \cdot \vec{v} \\ &= \vec{v} \cdot (\vec{m}\ddot{\vec{v}} + \nabla U) = \vec{v}(\vec{m}\ddot{\vec{v}} + \nabla U) \\ &= \vec{v}(\vec{F} - \vec{F}) = 0 \end{aligned}$$

Zakon očuvanja količine gibanja

$$\vec{p} = m\vec{v}$$

Promatrasimo sistem čestica; ukupna količina gibanja: $\vec{P} = \sum_i \vec{p}_i$

$$\begin{aligned} \frac{d\vec{P}}{dt} &= \frac{d}{dt} \sum_i \vec{p}_i = \sum_i \frac{d\vec{p}_i}{dt} = \sum_i \vec{F}_i \\ \vec{F}_i &= \vec{F}_{\text{vanjska}} + \sum_{j \neq i} \vec{F}_{ij} \quad \text{Diagram: } \textcircled{i} \quad \vec{F}_i = \vec{F}_{ji} \quad \vec{F}_j = -\vec{F}_{ij} \end{aligned}$$

$$\begin{aligned} \frac{d\vec{P}}{dt} &= \sum_i \vec{F}_i = \sum_i (\vec{F}_{\text{vanjska}} + \sum_j \vec{F}_{ij}) \\ &= \sum_i \vec{F}_{\text{vanjska}} + \sum_i \sum_{j \neq i} \vec{F}_{ij} \rightarrow 0 = \sum_i \vec{F}_{\text{vanjska}} \end{aligned}$$

Ako na sistem čestica ne djeluju vanjske sile ukupna količina gibanja je očuvana.