

Variation of kink distance over a period

#twokinks

Definition

If two kinks are far from each other and $C(t)$ is a **slow and strictly positive oscillation** we know ([Kinks effective dynamics under slow POSITIVE oscillations](#)) the kink effective dynamics model states

$$\dot{d}(t) \simeq -24\sqrt{2}C^{\frac{1}{2}}(t)e^{-2^{\frac{1}{2}}C(t)^{\frac{1}{2}}d}$$

If we consider the variation of the distance over one period $\Delta d(d)$

$$\Delta d(d) = \int_0^T dt \partial_t d = -24\sqrt{2} \int_0^T dt C^{1/2}(t) e^{-\sqrt{2}dC^{1/2}(t)}$$

where this is a function of d , because you can **assume d to be CONSTANT in the integrand** as it does not change significantly over a period.

By changing variable $t \rightarrow \tau = \frac{t}{T}$

$$\Delta d(d) = -24\sqrt{2}T \int_0^1 d\tau C^{1/2}(\tau) e^{-\sqrt{2}dC^{1/2}(\tau)}$$

$$C(\tau) = \bar{C} + A \sin(2\pi\tau)$$

If we define

$$I(d) = \int_0^1 d\tau e^{-\sqrt{2}dC^{1/2}(\tau)}$$

then

$$\Delta d(d) = +24T \frac{dI(d)}{d(d)}$$

Parabola approximation ($\bar{C} \geq A$)

As simulations show that the distance changes significantly when $C(t)$ is close to its minimum value $\bar{C} - A$, then it is natural to approximate, in the integrand:

$$C(\tau) \simeq C_{min} + \alpha \left(\tau - \frac{3}{4} \right)^2$$

(as a **parabola** and for every $\tau \in [0, 1]$), then

If $\bar{C} = A$

Here $C_{min} = 0$ so

$$C^{1/2}(\tau) = \tau - \frac{3}{4}$$

$$I(d) \simeq \int_0^1 d\tau e^{-\sqrt{2}d\alpha|\tau-3/4|}$$

Now, if you change variable

$$z = (\sqrt{2}d\alpha) \left(\tau - \frac{3}{4} \right)$$

$$I(d) = (\sqrt{2}d\alpha)^{-1} \int_{z_-}^{z_+} e^{-|z|} dz$$

where

$$z_- = -\frac{3}{4}(\sqrt{2}\alpha)^{-1}d^{-1}$$

$$z_+ = +\frac{1}{4}(\sqrt{2}\alpha)^{-1}d^{-1}$$

so the integral gives an exponential correction to the power-law decay

$$I(d) \sim d^{-1}$$

TO CORRECT

Derivating the last expression respect to d , leads to

$$\Delta d \sim e^{\sqrt{2}d(\bar{C}-A)^{1/2}} d^{-1/2} \left[\sqrt{2}(\bar{C} - A) + \frac{1}{2}d^{-1} \right]$$

That is similar to the expression you have for constant C

$$\Delta d \sim e^{\sqrt{2}dC^{1/2}}$$

but now $C \rightarrow (\bar{C} - A)$ and a **power-law** term is multiplying the exponential decay.

- If $\bar{C} > A$

The exponential dominates the behaviour of $\Delta d(d)$. So the variation of d is ruled by an exponential.

$$\Delta d \sim e^{-d(\bar{C}-A)}$$

- If $\bar{C} = A$

The exponential term vanishes, but also the term in $[\dots]$ is affected, such that

$$\Delta d \sim d^{-3/2}$$

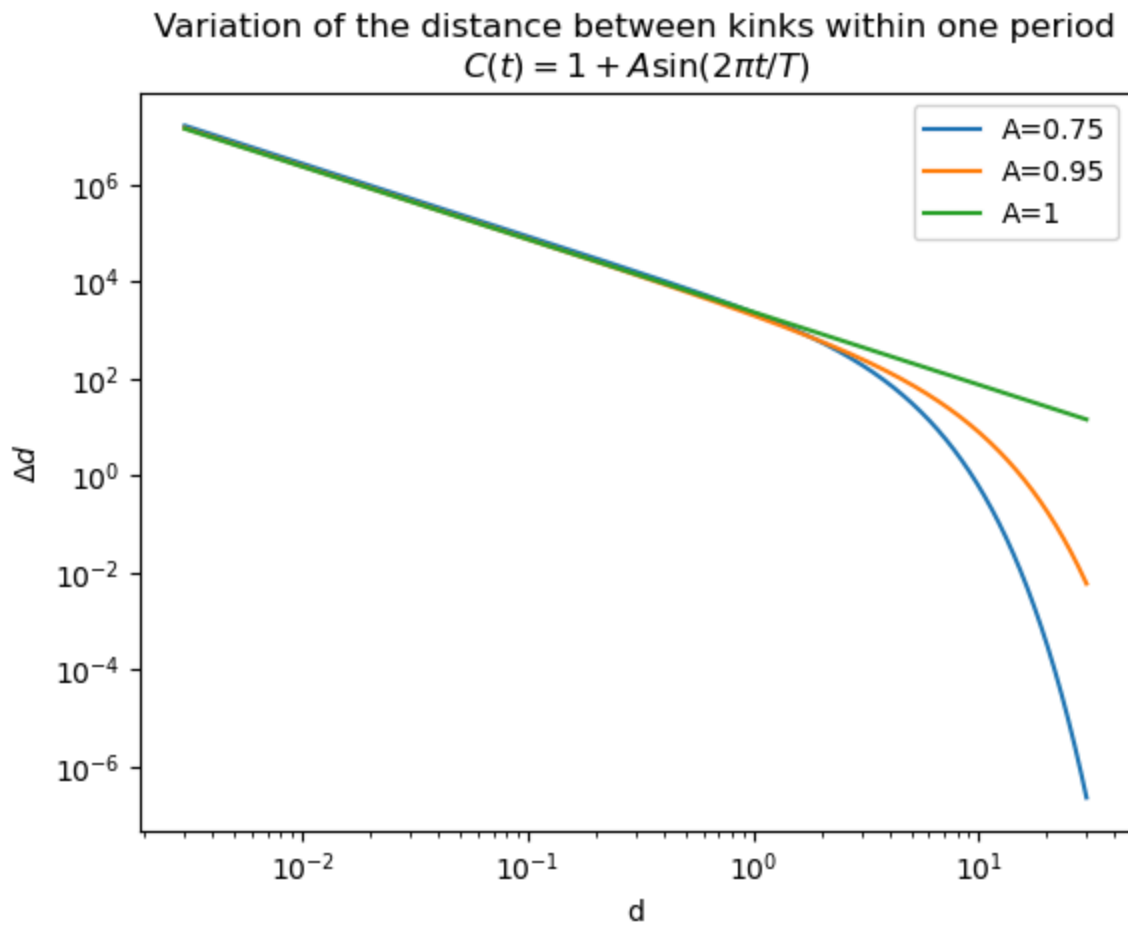
- If $\bar{C} < A$

Here, it is necessary to extend the kink effective dynamics model to cases when $C(t) < 0$ sometimes. Indeed, the model is developed considering $C(t)$ as a strictly positive oscillation. We do this **by assuming** $\partial_t d = 0$ **when** $C(t) < 0$. Then it is not possible to approximate $C^{1/2}$ to a parabola as it is not possible to calculate the square root. We should approximate $C(\tau)$ to a parabola, then take its square root. But then if you use the approximation $(1 + \epsilon)^{1/2} \simeq 1 + \frac{1}{2}\epsilon$, this approximation is not good, **so it is not clear what to expect.**

Summary

From the considerations above, we **expect** $\Delta d(d)$ to decay as

- A power-law when d is small
- Exponentially when d is large. **Unless** $A = \bar{C}$ **so the decay is power-law for any** d .
And we're interested in the behaviour far from annihilation, so at large d .



Simulations

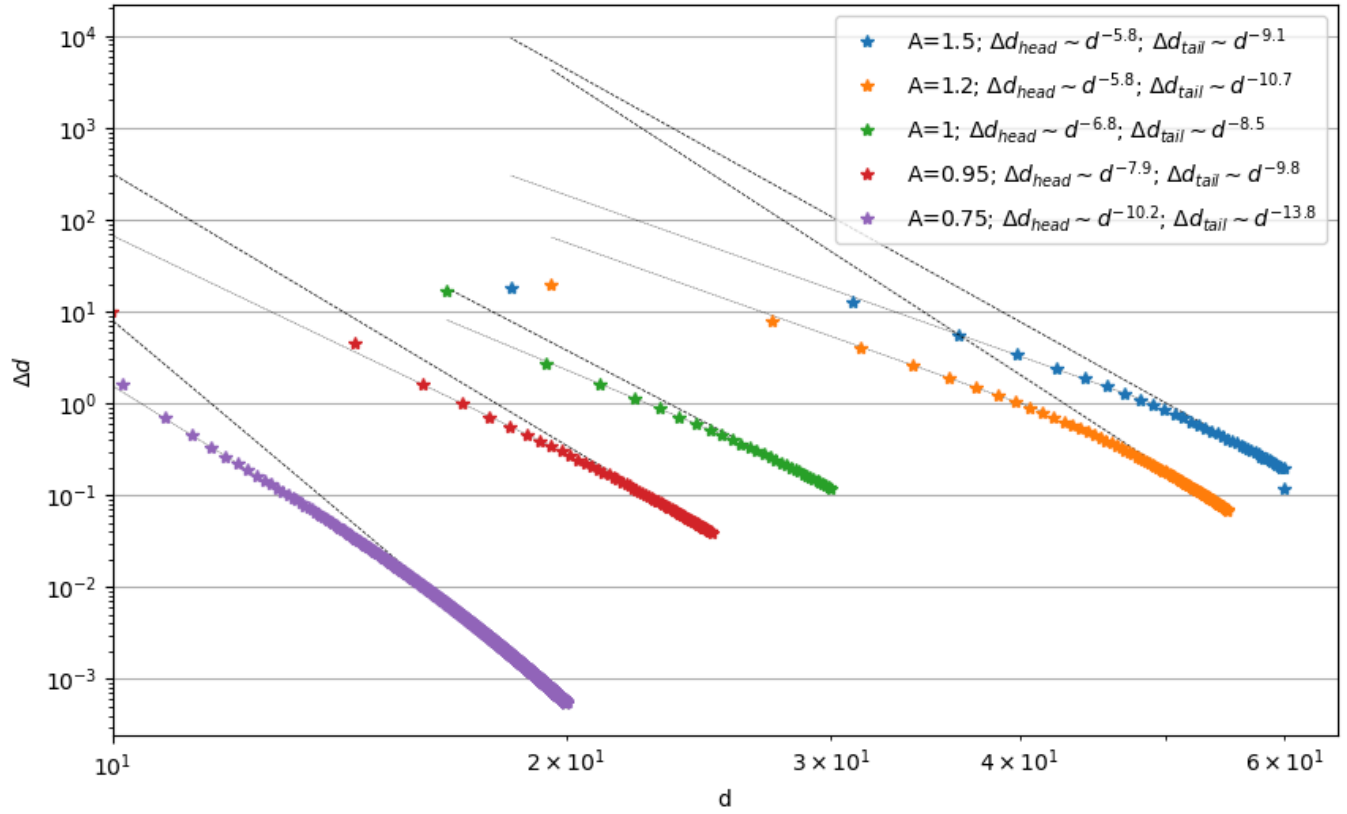
In the simulations below, the tail is fitted with a line and the slope of the line is reported in the legend. This value is way far from the expected one ($-3/2=-1.5$) and the next plot enhances that the decay is exponential (and not power-law) also when $A \geq \bar{C}$.

$$C(t) = 1 + A \sin\left(\frac{2\pi t}{500}\right)$$

Δd is the variation of the distance between two kinks in one period

d is the distance at the beginning of the period

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