

# Kinks dynamics with Mass

#1D

#twokinks

#timedependigC

Simulations suggest that the model for kink dynamics developed for **slow oscillations** should be corrected by adding an **inertial term**  $+M\partial_{tt}d$  in the differential equation for the distance between two isolated kinks.

$$(M\partial_{tt}d) + \partial_t d = f(C(t))$$

$$f(C(t)) = -24\sqrt{2}C^{\frac{1}{2}}(t)[e^{-2^{\frac{1}{2}}C(t)^{\frac{1}{2}}d} - e^{-2^{\frac{1}{2}}C(t)^{\frac{1}{2}}(L-d)}]$$

Now we try to **tune by HAND** the mass **M** to fit the experimental data.

## Resume

- (1) Adding **mass** to the kink dynamics model **fits well** the data **when  $C$  is close to its minimum** value (and so when  $-\partial_t d$  is large).
- (2) Adding mass does not extend the model to cases when  $C$  is very close to zero or negative.
- (3) As time passes, the shape of the trajectory in the region of large  $C$  continues to change. An **asymmetry** in the trajectory **never stops to grow** with time. And it is not a numerical effect!

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## Simulations

We simulated a system with only two kinks, under an oscillatory  $C(t)$ .

We measured, as a function of time, the distance  $d(t)$  and the value of  $C(t)$ . So we can plot the trajectory  $(d(t), C(t))$  and compare it with the trajectory obtained by **numerically solving** the model

## Numerical integration of the model

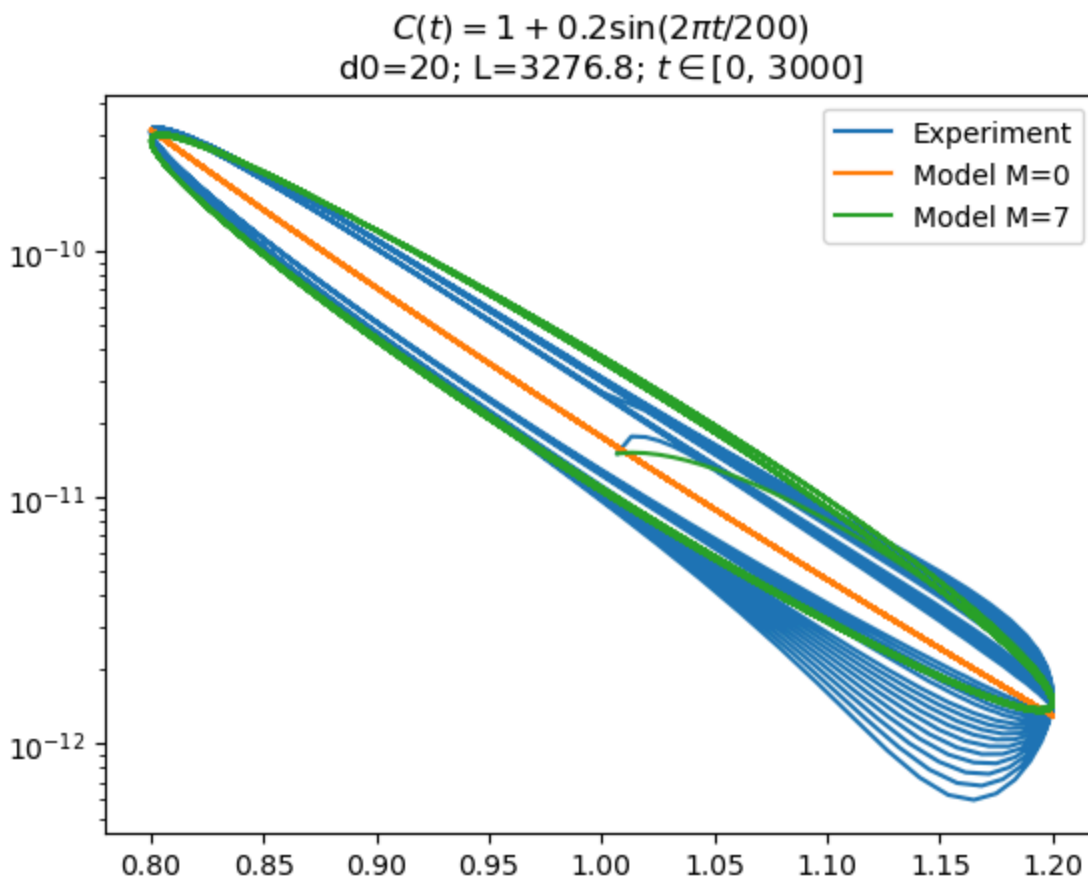
We can define  $y = \partial_t d$  so the model can be re-written as

$$\partial_t y = \frac{f(C(t)) - y}{M}$$

$$\partial_t d = y$$

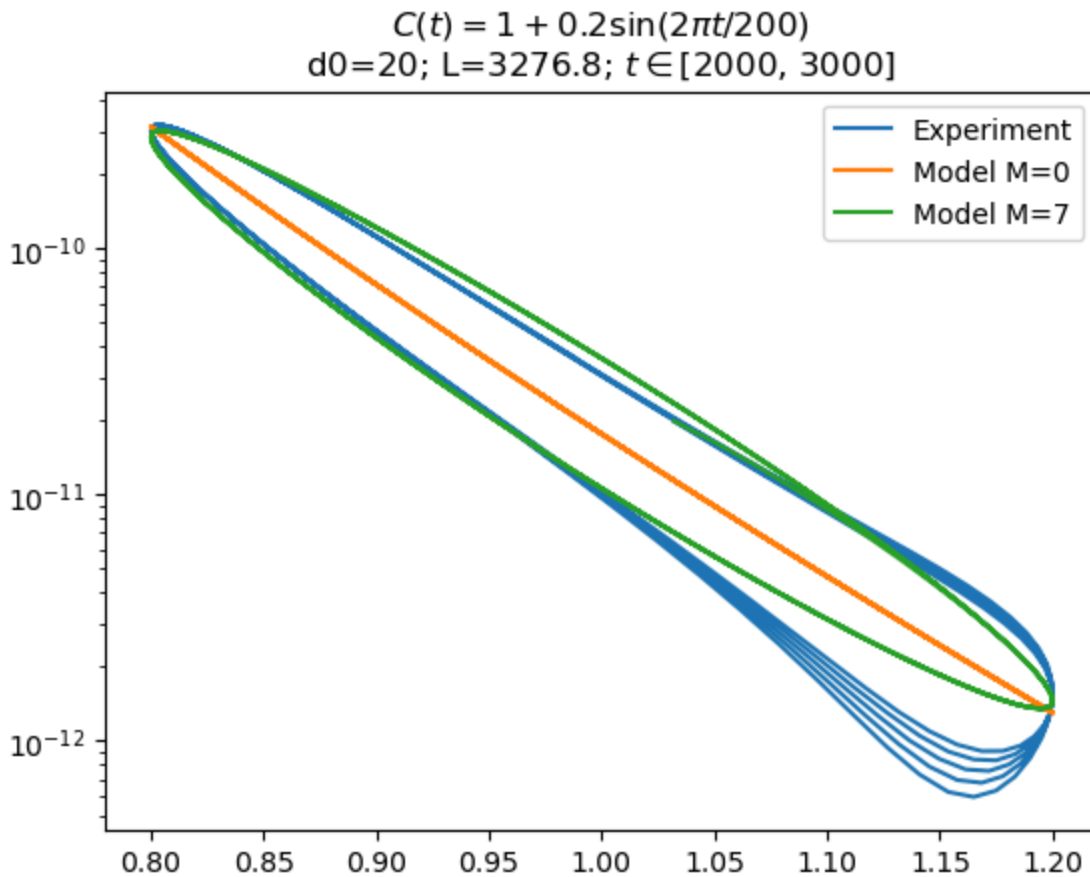
then we can integrate numerically these two coupled equations by using the **Explicit Euler scheme**, to find a solution  $(d(t), y(t))$ . The initial values  $d_0, y_0$  are taken from the simulation.

**We plot  $-\partial_t d$  v.s.  $C(t)$**



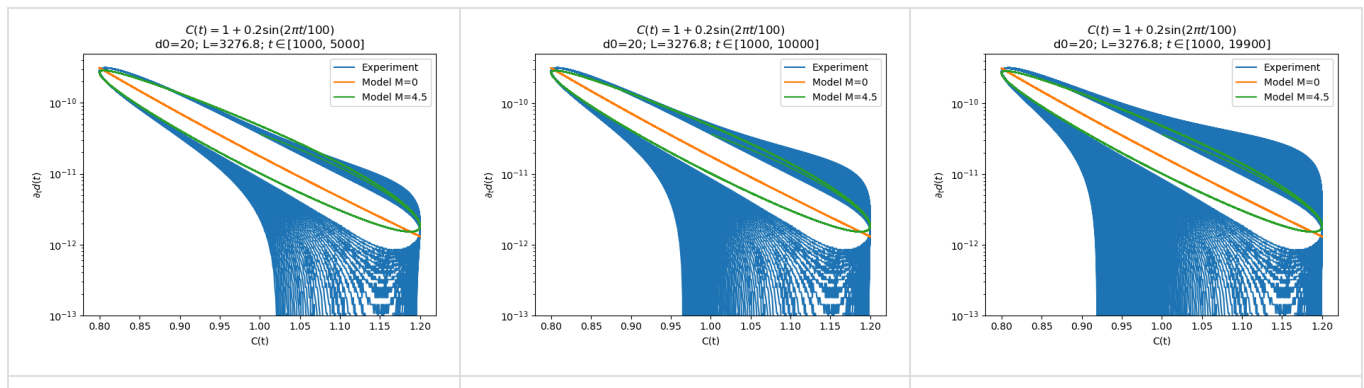
## Eliminate initial dynamics

To understand what happens as time passes, here we plot the experimental curve **after some periods from the beginning** of the simulation



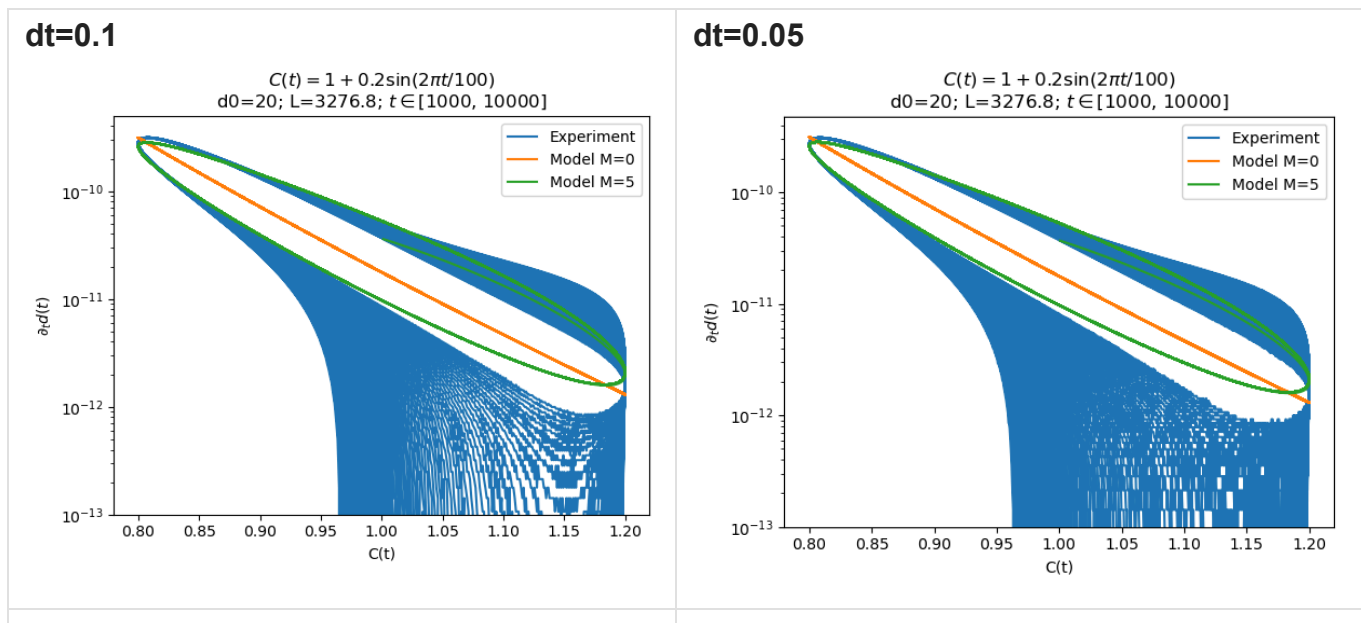
## Asymmetry growing with time

This **asymmetry** that we see at large values of  $C$  **increases** as time **passes**.



But, as the distance changes significantly when  $C$  is small (and not large), this asymmetry is responsible only for **higher order effects**.

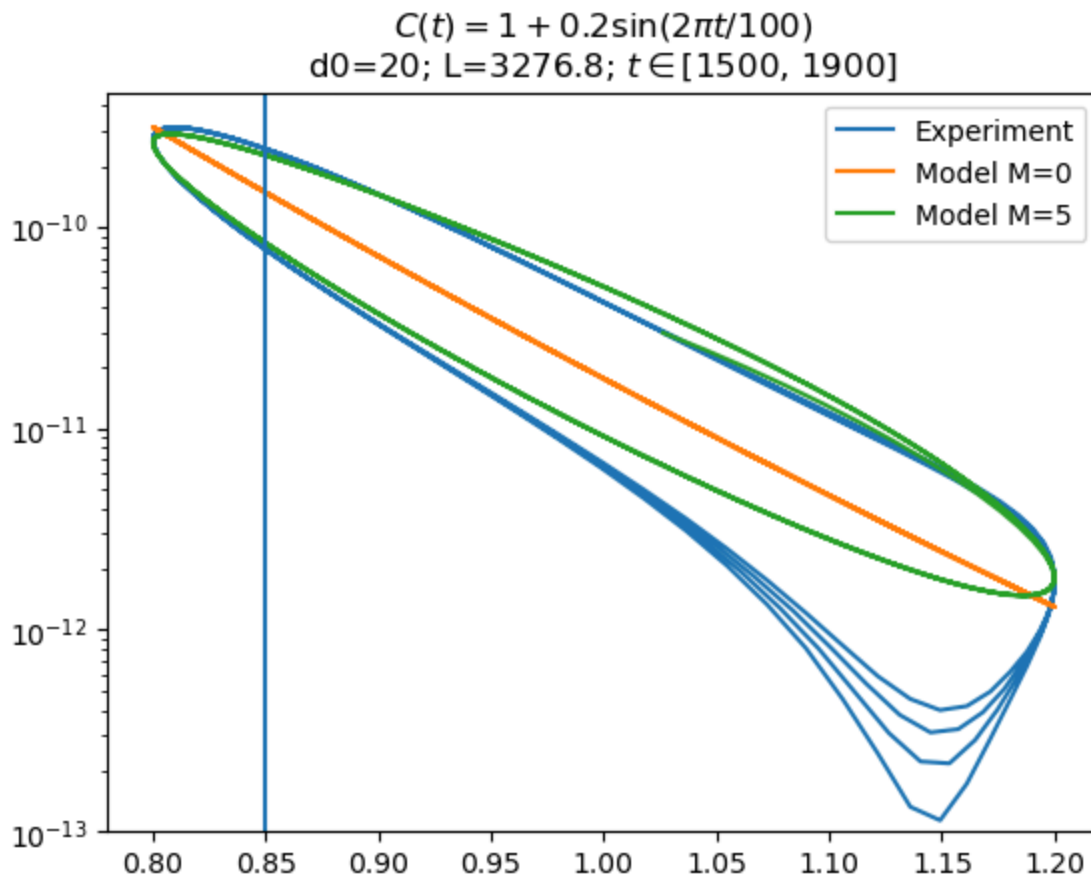
By running a simulation with a lower value of  $dt$  (0.05 instead of 0.1) the **asymmetry does NOT decrease**, so it is not a numerical error!



## Mass as a function of the period T

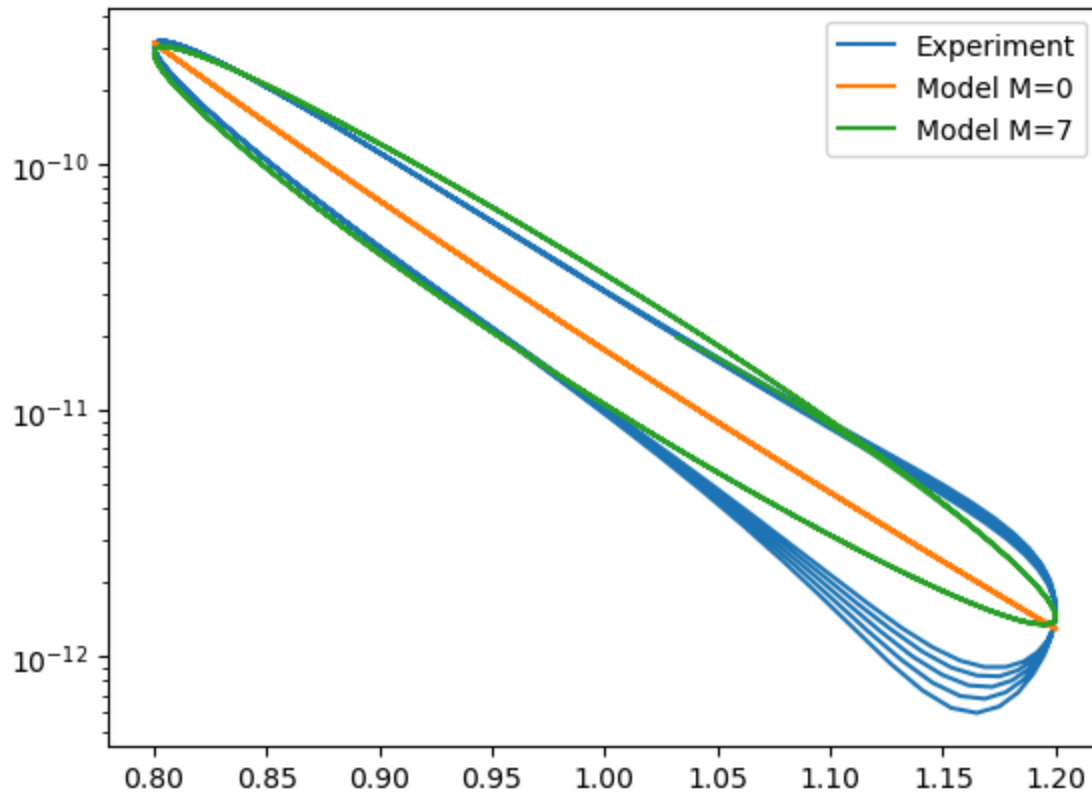
The value of the mass **M** has been tuned **by HAND**, resulting

- The mass increases with the period
- Even if the period changes of orders of magnitude, the mass remains of the same order



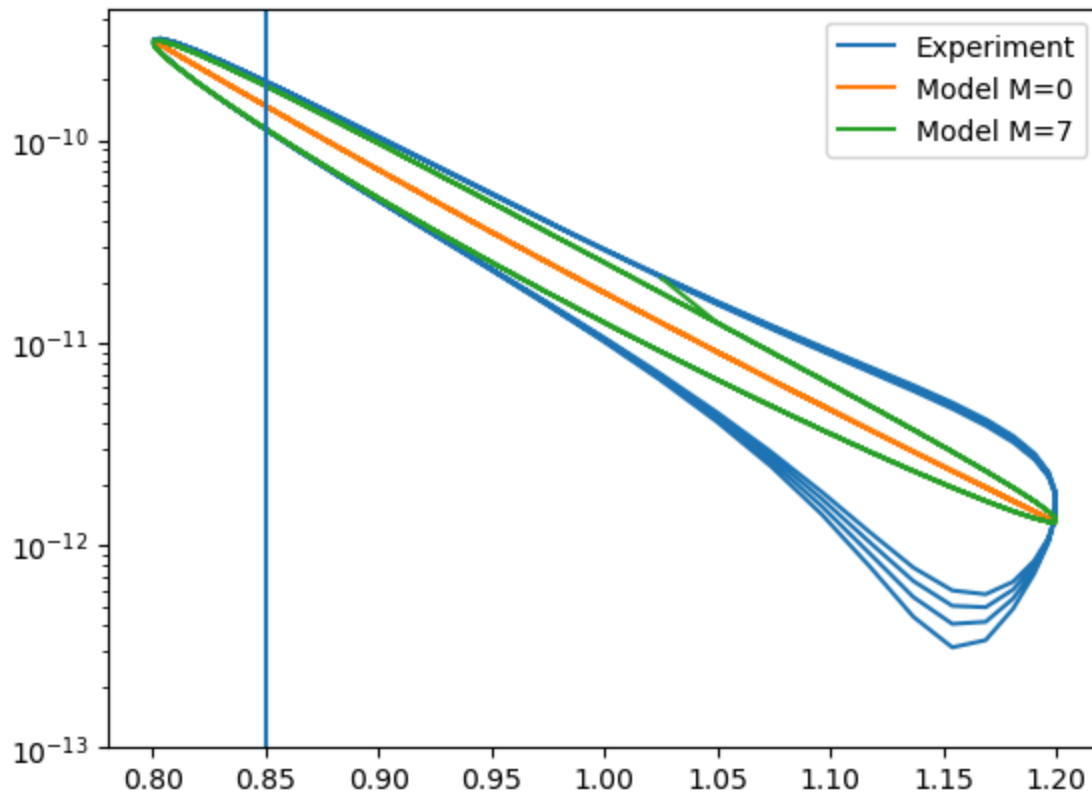
$$C(t) = 1 + 0.2\sin(2\pi t/200)$$

$d_0=20; L=3276.8; t \in [2000, 3000]$



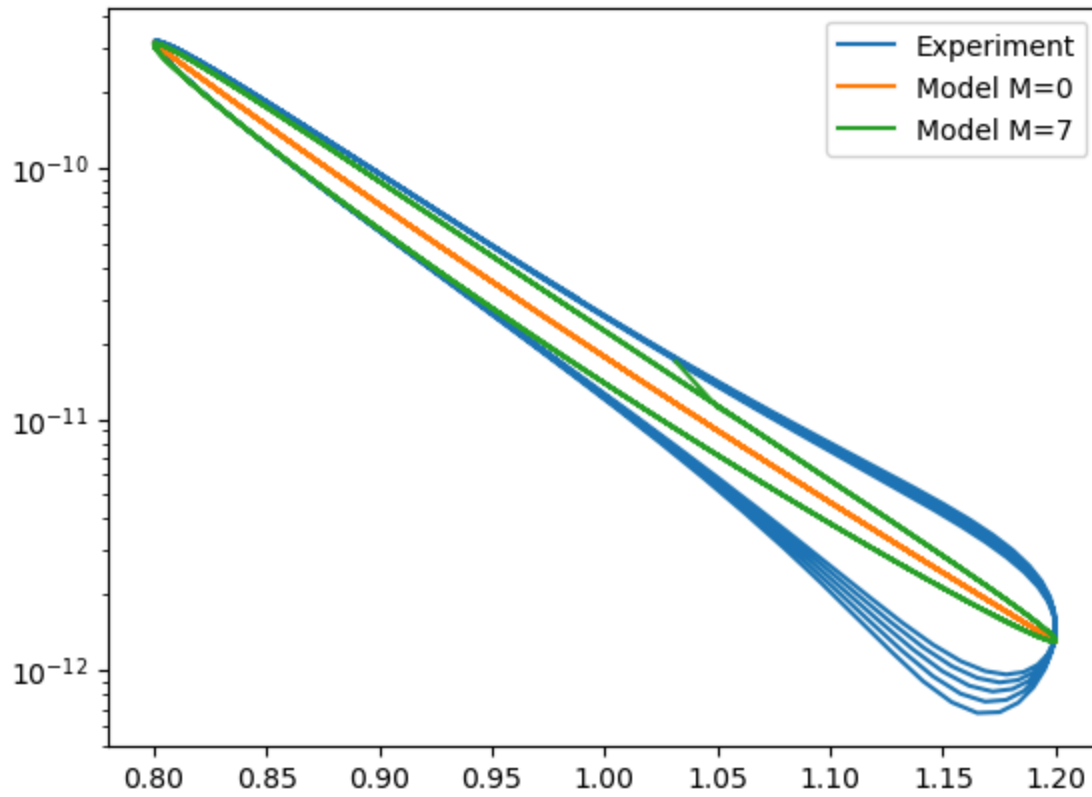
$$C(t) = 1 + 0.2\sin(2\pi t/350)$$

$d_0=20; L=3276.8; t \in [5250, 6650]$



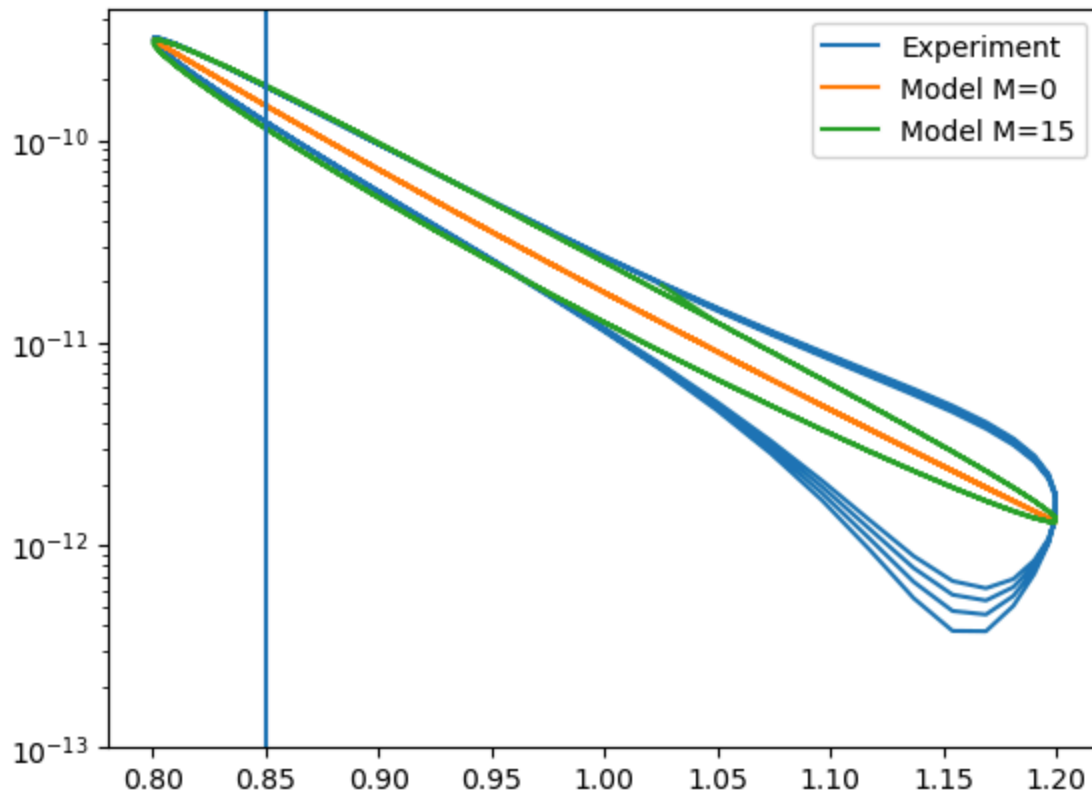
$$C(t) = 1 + 0.2\sin(2\pi t/500)$$

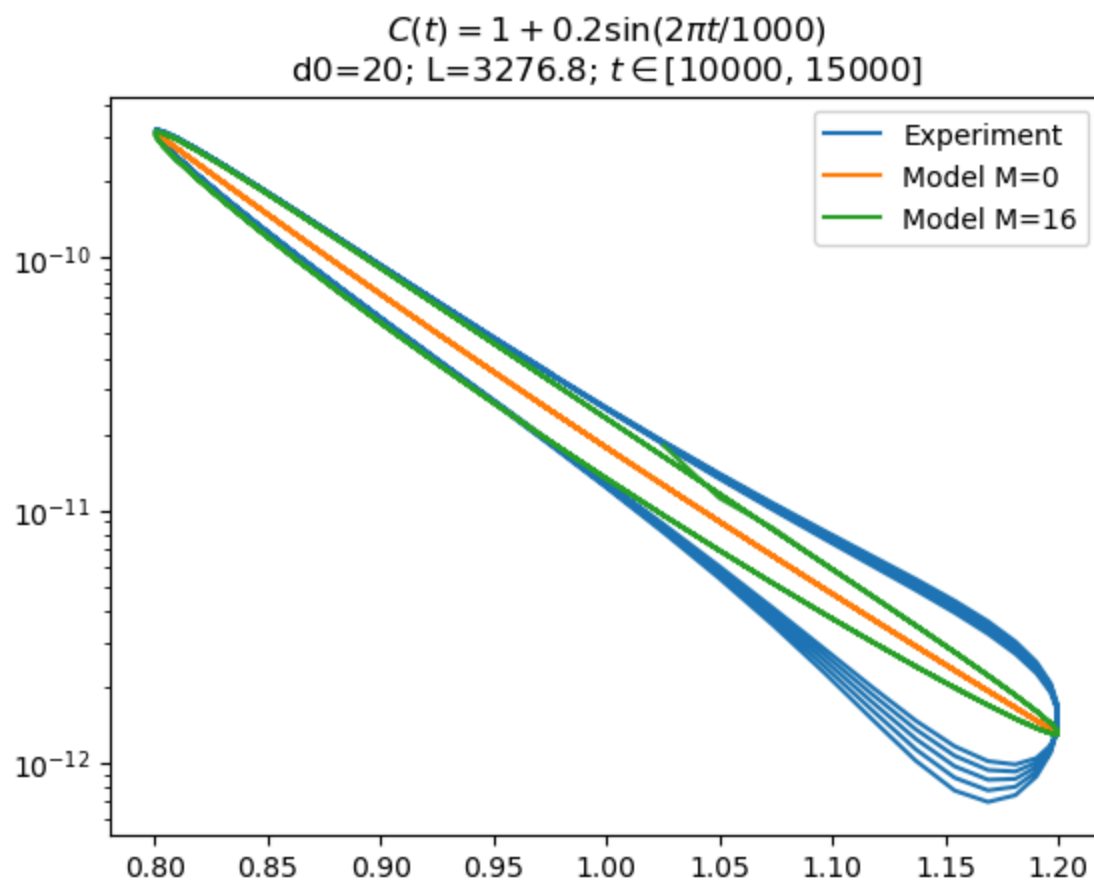
$d_0=20; L=3276.8; t \in [5000, 7500]$



$$C(t) = 1 + 0.2\sin(2\pi t/750)$$

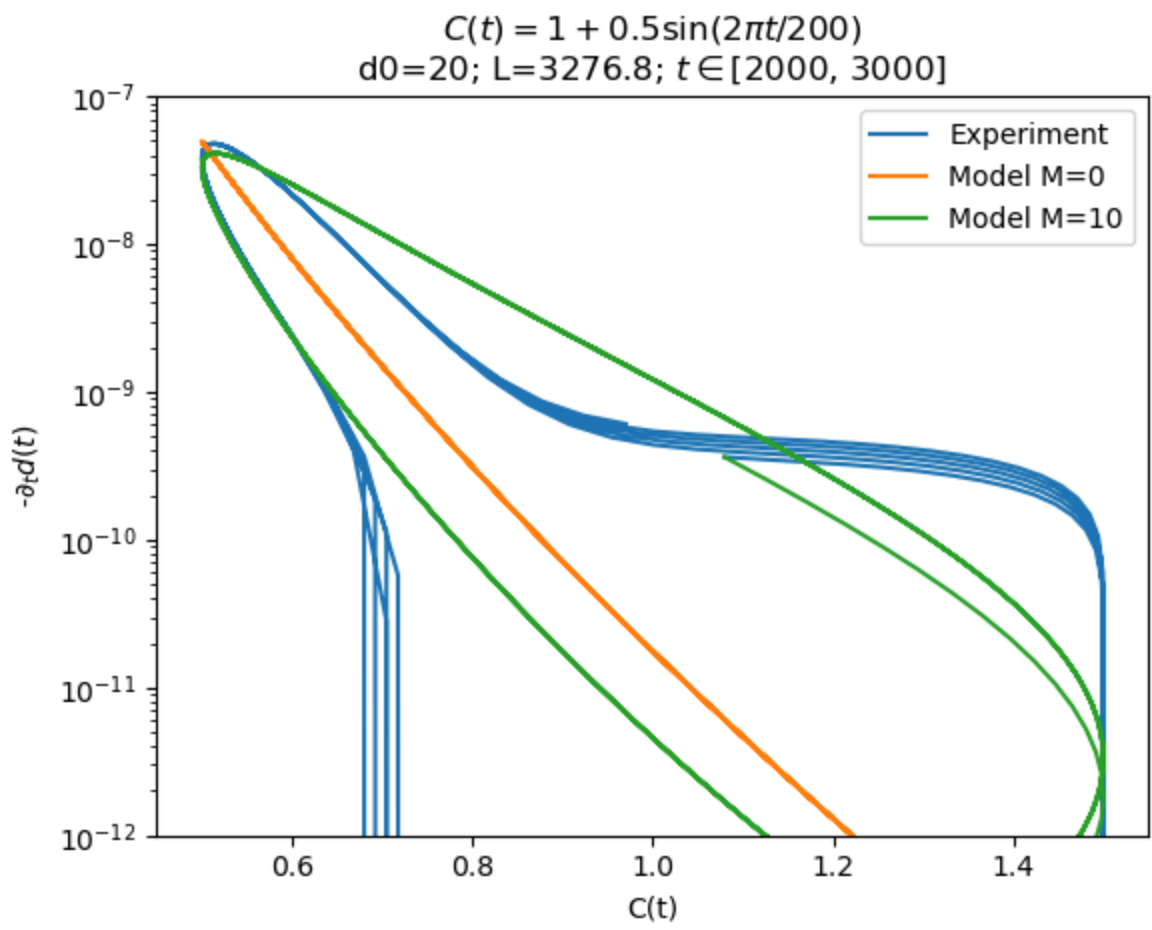
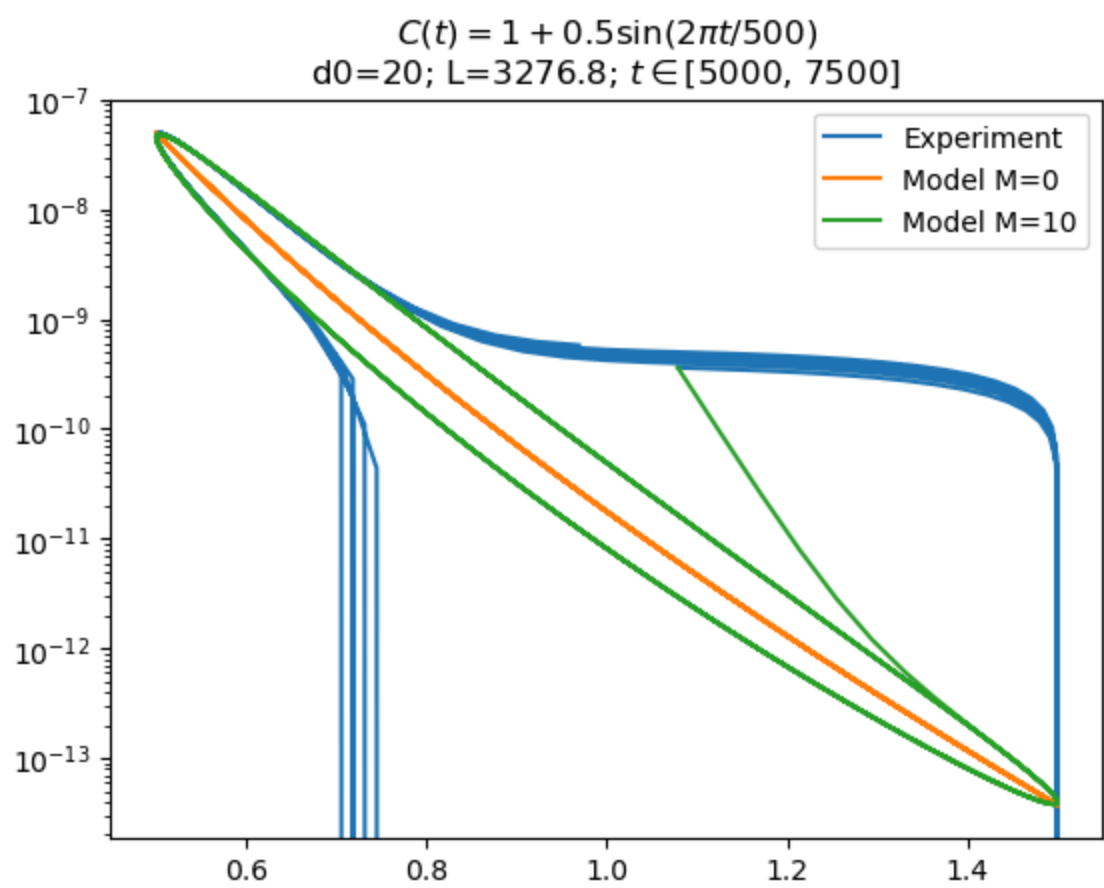
$d_0=20; L=3276.8; t \in [11250, 14250]$





**Mass as a function of the amplitude  $A$**

Also changing the amplitude, the order of magnitude of  $M$  is the same.





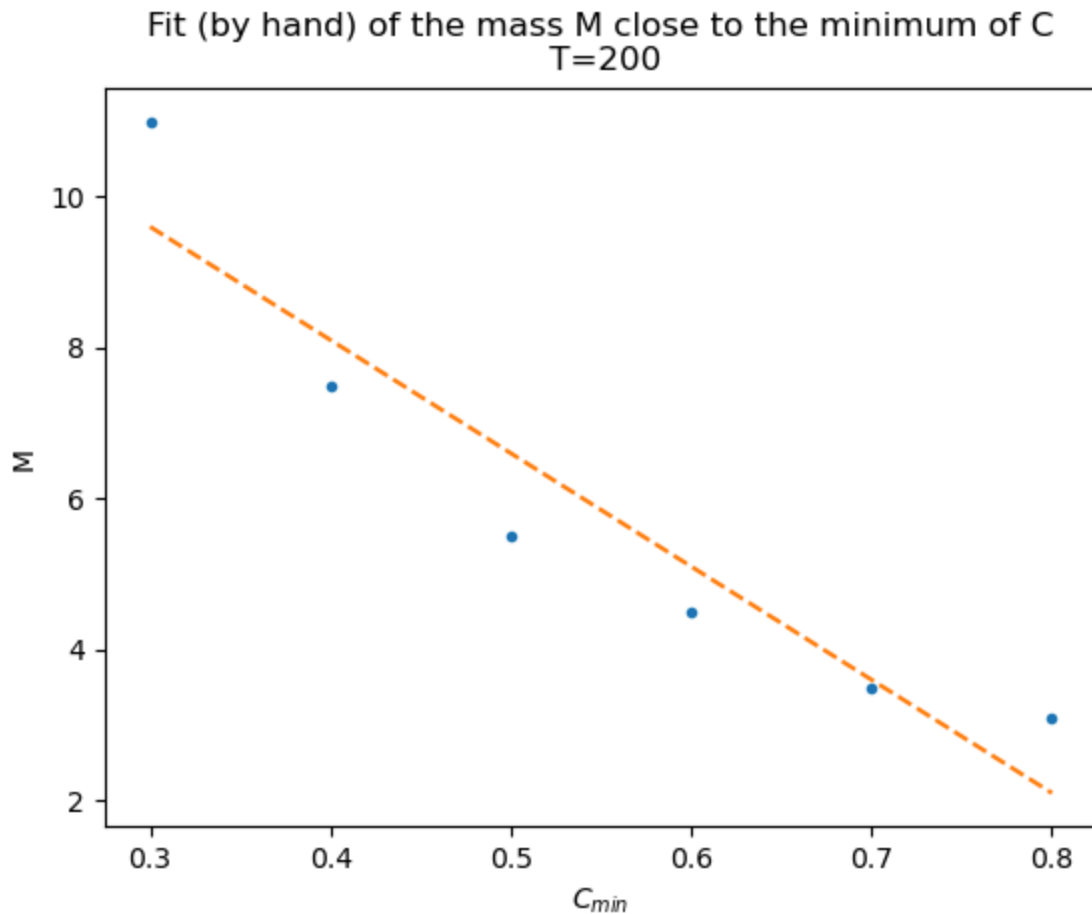
# A Mass dependent on C and Cdot

As the order  $\sim \epsilon^2$  correction depends, in principle, by  $C$  and  $\dot{C}$ , then we can try to fit the data with a mass

$$M(C, \dot{C}) = \alpha C + \beta \dot{C}$$

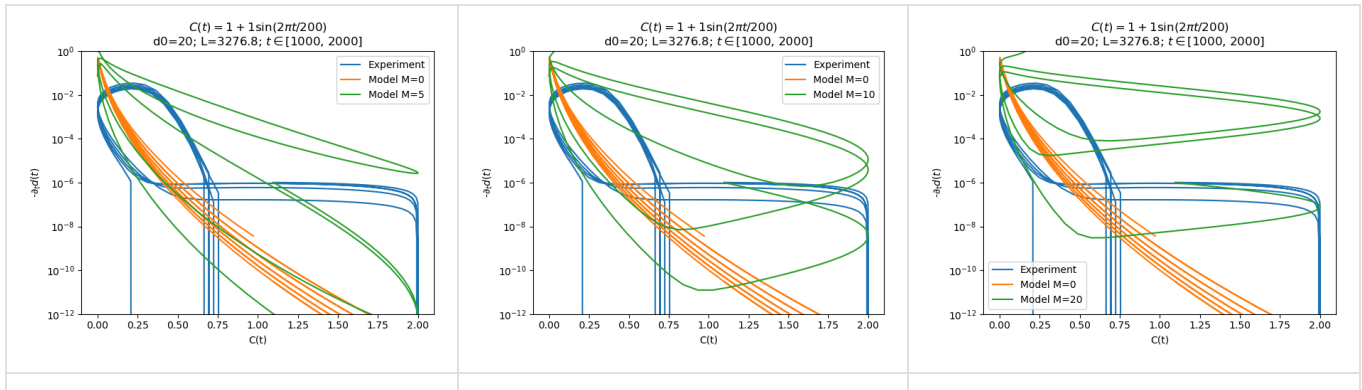
In order to find  $\alpha$ , we make a fit **by hand** by tuning  $M$  such that the fit is good close to the minimum value of  $C$ . We do this for many different values of  $C_{min}$ , while the period  $T$  is the same.

What we see below is that the data **is not distributed along a line**. Instead it fits better a power law decay  $M \sim C^{1/3}$ .



## Mass model when $A \geq \bar{C}$

In this case there is no suitable value of the mass  $M$  able to describe the deviation from the kink dynamics model.



## Conclusions

- Adding the mass  $\mathbf{M}$  to the model for "kink dynamics under slow oscillations", we have a good prediction of  $\partial_t d(C)$  **when  $C$  is small**. Although it does not work when  $C$  gets close to zero or negative.
- The order of magnitude of the fitted  $\mathbf{M}$  seems to not change with amplitude or period.
- The **asymmetry** that we see for large values of  $C$  is not predicted and it increases as time passes. Although, this asymmetry would be responsible **only for higher order** correction, as  $\partial_t d$  is order of magnitude higher when  $C$  is close to its smallest value.