Inverting kink motion with initial condition

#twokinks #1D

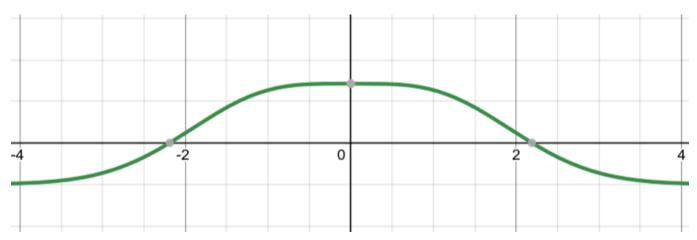
Examples

(1) Gaussian profile

We know that a Gaussian will broader in the linear regime, but what happens in general?

$$u(x,t_0)=u_0\mathcal{N}\left[g_+(x)+g_-(x)-rac{1}{2\sqrt{2\pi}\sigma}
ight]$$

$$g_{\pm}(x)=rac{e^{-(x-x_{\pm})^2/2\sigma^2}}{\sqrt{2\pi}\sigma}$$



with $\sigma=2(x_+-x_-)$ such that the second derivative in $x=rac{x_++x_-}{2}$ is zero and so $u(x,t_0)$ is almost flat between the kinks.

Second derivative at the positive kink $x_{0_{+}}$

Plots of this function and its derivatives, show that it is possible to approximate the second derivative at the right kink of $u(x,t_0)$ as the second derivative of the gaussian centered at x_+ .

$$\partial_{xx} u \simeq g_+(x) \left[-rac{1}{\sigma^2} (x-x_+)^2 + 1
ight]$$

Plots also show that $\exists x^*: \partial_{xx} u>0 \quad \forall x>x^*$, so we can estimate x^* by looking at the zero of the above function

$$x^* = \sigma + x_+$$

Estimating the right kink

For the same reason, we can approximate the position of the positive kink x_{0_+} as the zero of $g_+(x)-rac{1}{2\sqrt{2\pi}\sigma}$

$$x_{0_+} = \sqrt{\ln 4} \sigma + rac{L}{2}$$

We see that

$$x_{0_{\scriptscriptstyle \perp}}>x^*$$

so the kink always moves towards right (see example below) and so the kinks **repulse each other** for any distance and also in the non-linear regime.

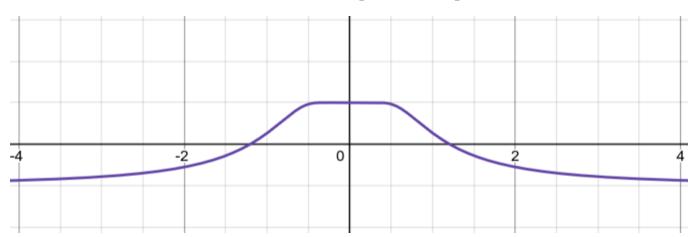
We expect the interface's shape to **relax** towards a tanh, so eventually the kinks will attract.

(2) Exponential profile

Let's consider a system with two **kinks** (a kink is **defined as a zero** of u(x)) at positions $x=\pm \frac{L}{2}$. It is possible to engine the shape of the state u(x,t=0) in proximity of the kinks such that the velocity of each kink at time t=0 corresponds to a repulsive interaction, instead of an acctractive one.

To achieve this, you prepare the initial state as

$$u(x,t=0) = 2u_0 \left[rac{1}{2} - e^{-1/(lpha x)^2}
ight]$$



here α is related to the distance L between the kinks (that are the zeroes of u(x)) as $\alpha=\frac{2}{L}\sqrt{\frac{1}{\log 2}}.$

To calculate the kink's velocity at time t=0 we consider the TDGL equation

$$\partial_t u = \partial_{xx} u + C(t)u - u^3 \quad \forall x, t$$

now, if you assume that, close enough to the right kink $x_+=+\frac{L}{2}$, you can assume that the kink propagates without changing shape, then

$$\partial_t u \simeq -\dot{x_+}\partial_x u \quad ext{if } x \simeq x_+$$

then, if you put this in the previous equation and you evaluate it at $x=x_{\pm}$ and $t=t_{0}=0$

$$-\dot{x_+}\partial_x u|_{x_+,t_0}=\partial_{xx}u|_{x_+,t_0}$$

this equation is so simple, because we used $u|_{x_+,t_0}=u^3|_{x_+,t_0}=0$.

From the last equation, we can compute the **sign** of the initial velocity of the right kink $\dot{x_+}$ by computing the sign of the first and the second derivative of the initial state at $x = x_+$.

First derivative

The sign of the first derivative is trivial, as the function u(x,t=0) is decreasing at the position of the right kink, so $\partial_x u|_{x_+,t_0} < 0$.

Second derivative

Here we should look at the concavity of the function $e^{-1/(\alpha x)^2}$.

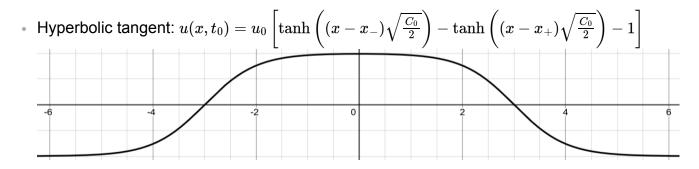
- $heta_x u(x,t=0) = -4 lpha^{-2} u_0 x^{-3} e^{-1/(lpha x)^2}$
- $\partial_{xx}u(x,t=0)|_{x_+,t_0}=-2\alpha^{-2}u_0x_+^{-4}\left[-\frac{3}{2}-\alpha^{-2}x_+^{-2}\right]>0\quad \forall x_+,\alpha$ where in the last expression I used that $u(x_+,t=0)=0$.

So we find that

- $ullet \partial_{xx} u(x,t_0)|_{x_+,t_0}>0$
- $ullet \partial_x u(x,t_0)|_{x_+,t_0}<0$

and this means that $\dot{x_+} > 0$, so **the two kinks get far apart** instead of attracting. At least **at the beginning**, because then we expect the shape of the kinks to relax towoards the steady state profile ($\sim \tanh x$).

Examples where the kinks attract



In all these cases, the kink always attract since $t = t_0$, for any value of the parameters!