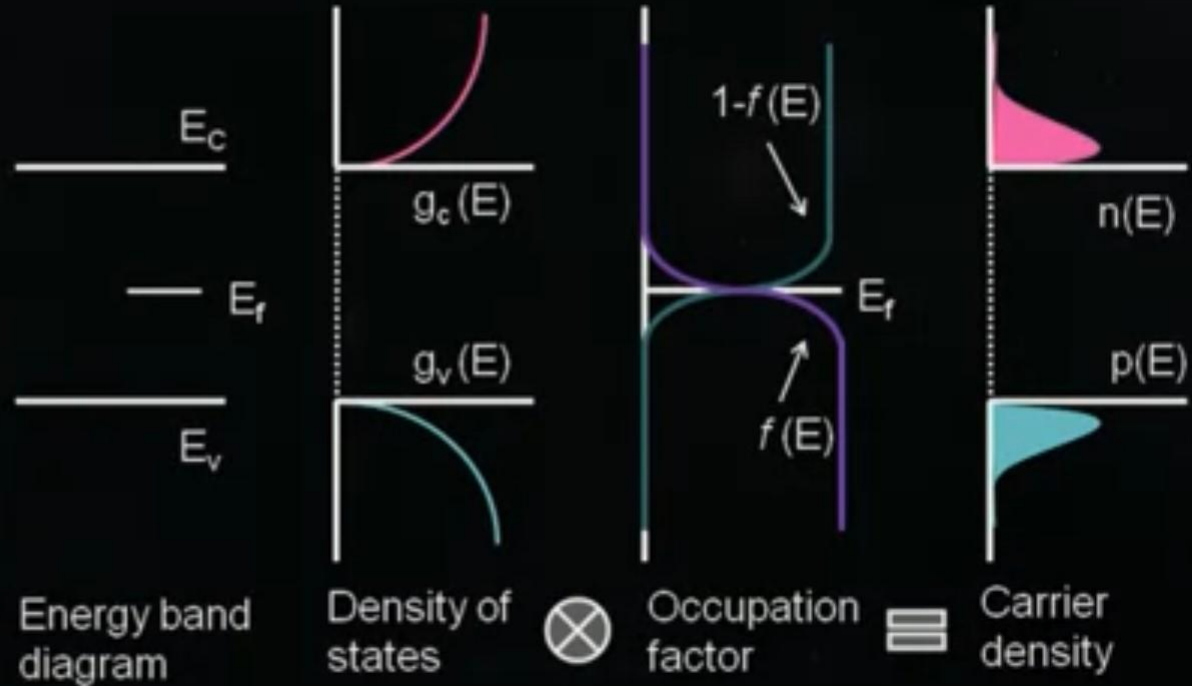


Density of states and carrier density

Press **Esc** to exit full screen

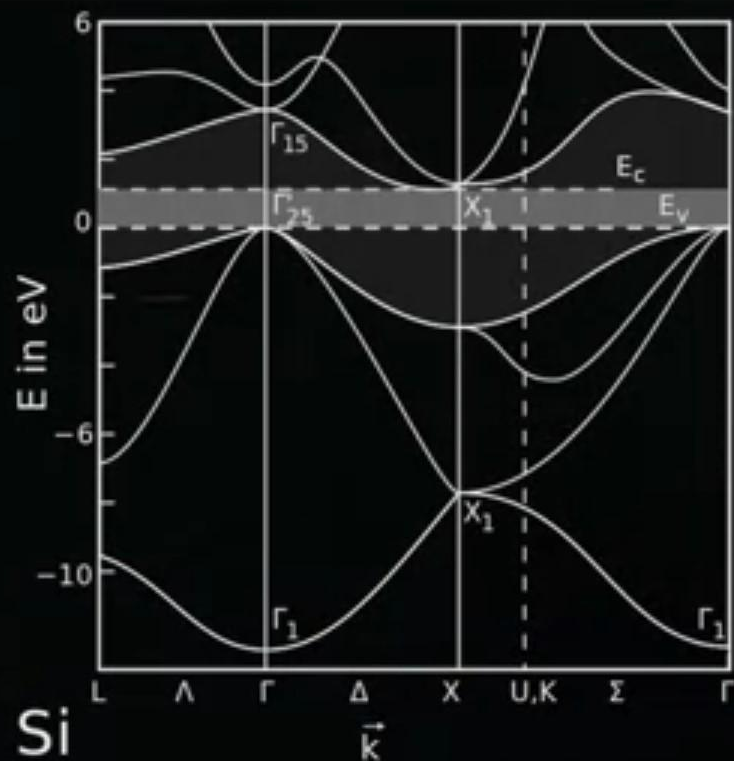


<https://nanohub.org/wiki/CarrierStatisticsPage/Image:Image.png>

Intrinsic and extrinsic semiconductors



Energy Band Structure



https://upload.wikimedia.org/wikipedia/commons/thumb/0/04/Band_structure_Si_schematic.svg/580px-Band_structure_Si_schematic.svg.png



Intrinsic semiconductors

$$n(T) = p(T) = n_i(T)$$

$$n_i(T) = \sqrt{N_v(T)N_c(T)} \exp\left(-\frac{E_g(T)}{2k_B T}\right)$$

$$E_F = E_v + \frac{1}{2}E_g + \frac{1}{2}k_B T \ln\left(\frac{N_v}{N_c}\right)$$

$$= E_v + \frac{1}{2}E_g + \frac{3}{4}k_B T \ln\left(\frac{m_v^*}{m_c^*}\right)$$

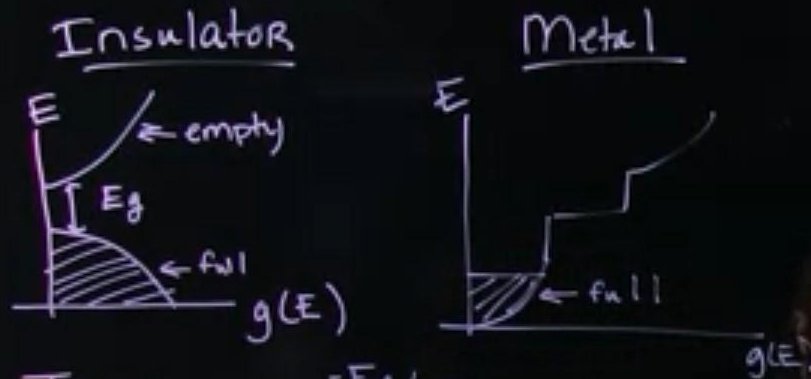
- ① $T \rightarrow 0K$ $E_F \rightarrow \text{midgap}$
- ② $\ln\left(\frac{m_v^*}{m_c^*}\right) \rightarrow 1$ E_F within $\pm k_B T$ of midgap
- ③ $k_B T \ll E_F$ $E_F \rightarrow \text{midgap}$

Extrinsic semiconductors

Treat impurities as
an additional charge

$+e = \text{donor}$

$-e = \text{acceptor}$



$$T > 0 \text{ K}, \propto e^{-E_g/k_B T}$$

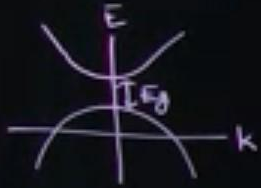
$$\text{ex) } \begin{cases} T = 300 \text{ K} \\ E_g = 4 \text{ eV} \end{cases} \text{ prob} = 10^{-55}$$

$$\begin{aligned} E_g &= 0.25 \text{ eV} \\ T &= 300 \text{ K} \end{aligned} \quad \text{" " } = 10^{-2}$$

So, you can really see from this example

Metals, Insulators & Semiconductors

Semiconductors:

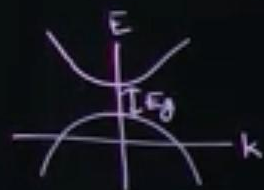


- Special kind of insulators whose band gap is small enough
- How small is small enough?
 - Enough to thermally excite electrons across band gap
 - $< 2\sim 3$ eV.
- Energy bandgaps of semiconductors (room temp.):

InSb: 0.16 eV	Ge: 0.67 eV
Si: 1.12 eV	GaAs: 1.42 eV
GaP: 2.2 eV	GaN: 3.5 eV
ZnS: 3.68 eV	

So now, let's talk about the rigorous definition of a semiconductor.

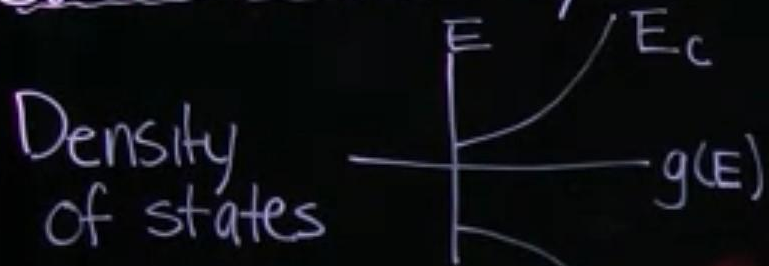
Rigorous definition of semiconductor



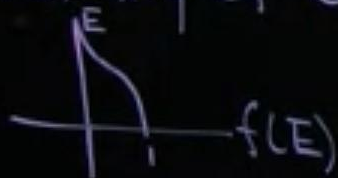
The solids that are insulators at $T = 0$ K but whose energy bandgap is of such a size that thermal excitation leads to observable conductivity at temperatures below its melting point are called the semiconductors.



Carrier Density



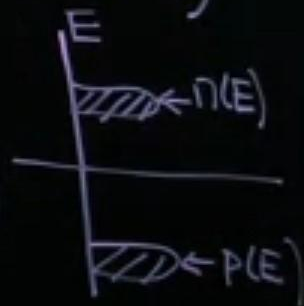
Probability of occupation



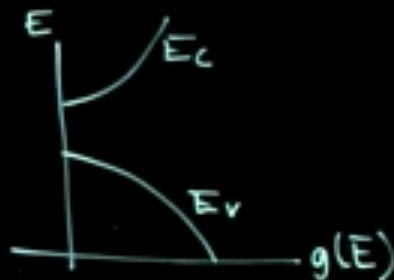
$$n(T) = \int_{E_c}^{\infty} dE g_c(E) f(E)$$

$$p(T) = \int_{-\infty}^{E_v} dE g_v(E) (1 - f(E))$$

$$f(E) = \frac{1}{1 + \exp\left(\frac{E - E_F}{k_B T}\right)}$$



Density of states



$g(E) \sim \# \text{ states/Volume/Energy}$

$g(E)dE \sim \# \text{ states/volume}$



$$E = \frac{\hbar^2 k^2}{2m^*} = \frac{\hbar^2}{2m^*} (k_x^2 + k_y^2 + k_z^2)$$

..... Ground state of N electron system

$\psi(x) = \psi(x+L)$

$E < E_F$

$$E_F = \frac{\hbar^2 k_F^2}{2m^*}$$

$$V_F = \frac{4}{3} \pi k_F^3 = \frac{4}{3} \pi \left(\frac{2m^* E_F}{\hbar^2} \right)^{3/2}$$

1 allowed state per $(2\pi)^3/V \leftarrow \text{volume of crystal}$

$$N = V_F / (2\pi)^3 / V = \frac{2V}{6\pi^2} \left(\frac{2m^* E_F}{\hbar^2} \right)^{3/2}$$

$$g(E) = \frac{1}{V} \frac{dN(E)}{dE} = \frac{1}{2\pi^2} \left(\frac{2m^*}{\hbar^2} \right)^{3/2} E^{1/2}$$

halves times E to the one half.

Nondegeneracy

$$E_c - E_F \gg k_B T$$

$$E_F - E_v \gg k_B T$$

\Rightarrow Fermi Dirac Statistics
 \approx Boltzman (classical)
 Statistics

$$f(E) = \frac{1}{1 + e^{\frac{(E - E_F)}{k_B T}}}$$

$$\approx \exp\left(-\frac{(E - E_F)}{k_B T}\right)$$

for $E > E_c$

$$1 - f(E) = \frac{1}{\exp\left(\frac{(E_F - E)}{k_B T}\right) + 1} \approx \exp\left(-\frac{(E_F - E)}{k_B T}\right)$$

$E < E_v$

us to come up with analytic expressions for our integrals.

$$n(T) = \int_{E_c}^{\infty} dE g_c(E) \exp\left(-\frac{(E-E_F)}{k_B T}\right)$$

$$p(T) = \int_{-\infty}^{E_v} dE g_v(E) \exp\left(-\frac{(E_F-E)}{k_B T}\right)$$

$$n(T) = N_c(T) \exp\left(-\frac{(E_c-E_F)}{k_B T}\right)$$

$$p(T) = N_v(T) \exp\left(-\frac{(E_F-E_v)}{k_B T}\right)$$

$$N_c(T) = \int_{E_c}^{\infty} dE g_c(E) \exp\left(-\frac{(E-E_c)}{k_B T}\right)$$

$$N_v(T) = \int_{-\infty}^{E_v} dE g_v(E) \exp\left(-\frac{(E_v-E)}{k_B T}\right)$$

Parabolic bands

$$N_c(T) = \frac{1}{4} \left(\frac{2m_c^* k_B T}{\pi \hbar^2} \right)^{3/2}$$

$$N_v(T) = \frac{1}{4} \left(\frac{2m_v^* k_B T}{\pi \hbar^2} \right)^{3/2}$$

Law of mass action

$$n(T)p(T) = N_c N_v e^{-E_g/k_B T}$$

E_g over $k_B T$.

Fermi levels

Intrinsic semiconductor

$$E_F = E_i = E_v + \frac{1}{2}E_g + \frac{3}{4}k_B T \ln\left(\frac{m_v^*}{m_c^*}\right)$$

Extrinsic semiconductor, p doped with $p = N_A$

$$E_F = E_i - k_B T \ln\left(\frac{N_A}{n_i}\right) \text{ or } E_F - E_v = k_B T \ln\left(\frac{N_v}{N_A}\right)$$

So now, let's look at the case



Fermi levels

Intrinsic semiconductor

$$E_F = E_i = E_v + \frac{1}{2}E_g + \frac{3}{4}k_B T \ln\left(\frac{m_v^*}{m_c^*}\right)$$

Extrinsic semiconductor, p doped with $p = N_A$

$$E_F = E_i - k_B T \ln\left(\frac{N_A}{n_i}\right) \text{ or } E_F - E_v = k_B T \ln\left(\frac{N_v}{N_A}\right)$$

So E_F is now equal
to E_i minus $k_B T$,

Donor Energy Level

Donor E

$$E_d = \left\{ \left(\frac{m_c^*}{m_0} \right) \left(\frac{E_0}{E_s} \right)^2 \cdot 13.6 \text{ eV} \right\}$$

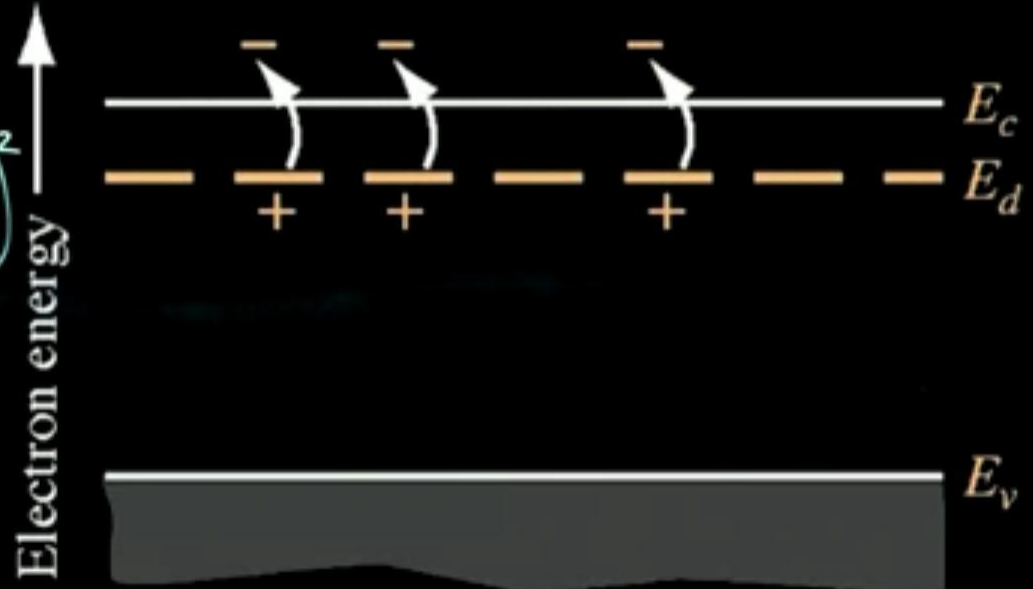


will change the energy level that appears.

Donor Energy Level

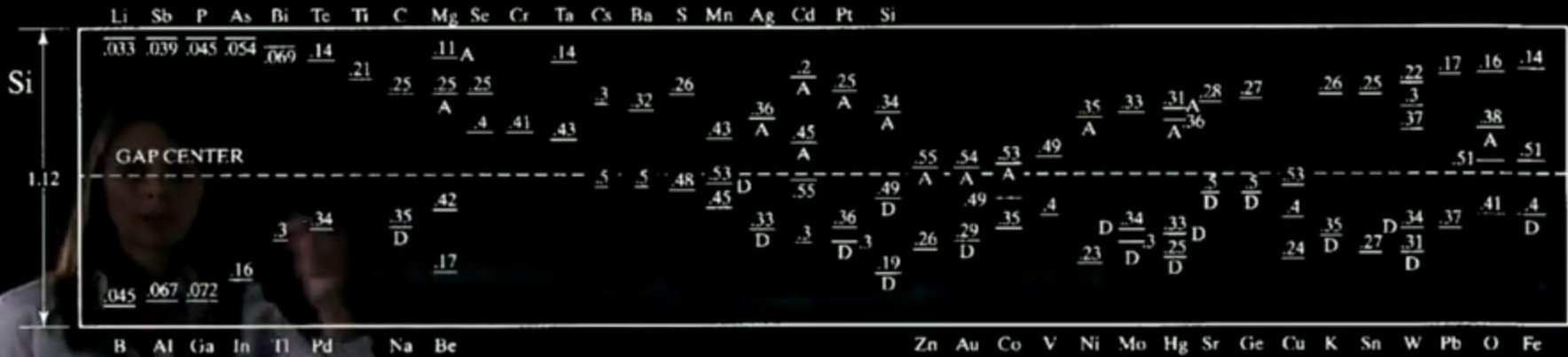
Donor E

$$E_d = \left\{ \left(\frac{m_c^*}{m_0} \right) \left(\frac{E_0}{E_s} \right)^2 \cdot 13.6 \text{ eV} \right\}$$



from being essentially an insulator to a good conductor.

Si



band.

Charge neutrality

$$n_c + n_d + N_A = N_D + p_v + p_A$$

$$n_c, p_v = \frac{1}{2} \left[\Delta n^2 + 4n_i^2 \right]^{1/2} \pm \frac{1}{2} \Delta n$$

Fully ionized impurities

$$n_d = p_A = 0$$

$$n_c + N_A = N_D + p_v$$

$$\left. \begin{array}{l} N_D > N_A \\ \text{fully ionized} \end{array} \right\} \Delta n = n_c - p_v = N_D - N_A$$

$$n_i^2 = n_c p_v$$

$$= \frac{1}{2} \left[(N_D - N_A)^2 + 4n_i^2 \right]^{1/2}$$

$$\pm \frac{1}{2} (N_D - N_A)$$

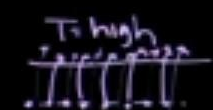
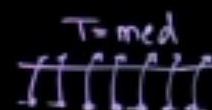
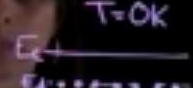
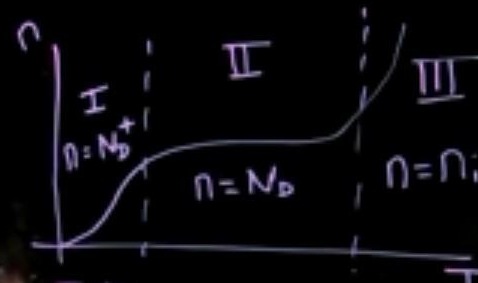
the one half plus or minus one half N_D minus N_A .

Low impurity concentration

$$\left. \begin{matrix} n \\ p \end{matrix} \right\} = n_i \pm \frac{1}{2} (N_D - N_A)$$

High impurity concentration

$$\left. \begin{matrix} n \\ p \end{matrix} \right\} \begin{matrix} N_D \gg N_A \\ n = N_D - N_A \\ p = \frac{n_i^2}{N_D - N_A} \end{matrix}$$



$$n = N_D^+$$

$$n = N_D$$

to the knobs that we have from doping and also from temperature.

Recommended References

The references below are the go-to references for semiconductor physics. These texts are completely optional, but note that the last text is freely available.

S. M. Sze, Physics of Semiconductor Devices, Wiley, 2007.

G. Streetman and S. Banerjee, Solid State Electronic Devices, Prentice Hall, 2000

K. F. Brennan, The Physics of Semiconductors, Cambridge, 1999

B. Van Zeghbroeck, [Principles of Semiconductor Devices](#)