

# ACM 100c

## Properties of the Sturm-Liouville eigenfunctions and eigenvalues

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# Properties of the solutions to the S-L problem

- The Lagrange identity remains true for this scalar product as well provided we use separable boundary conditions

$$\int_a^b \bar{u} L v dx = \int_a^b v L \bar{u} dx$$

- So far we have shown that the Sturm-Liouville operator is self adjoint
- We have also shown that for specific choices of  $p(x)$ ,  $q(x)$ , and  $r(x)$  solutions exist for special values of  $\lambda$ :

$$\frac{d}{dx} \left( p(x) \frac{dy}{dx} \right) - q(x)y(x) + \lambda r(x)y = 0, \quad a < x < b,$$

- We will call the special values of  $\lambda$  at which nontrivial solutions exist *eigenvalues*
- Correspondingly we will call the associated solutions  $y(x)$  the *eigenfunctions*

# Analogy with the linear algebra eigenvalue problem

- These names make sense because we can write the S-L problem in the form

$$-\frac{d}{dx} \left( p(x) \frac{dy}{dx} \right) + q(x)y(x) = \lambda r(x)y, \quad a < x < b,$$

which looks formally like the algebraic eigenvalue problem:

$$Ly(x) = \lambda r(x)y(x)$$

- Here  $L$  plays the role of the “matrix”,  $y(x)$  plays the role of the eigenvector, and  $\lambda$  plays the role of the eigenvalue.

# Properties of the solutions to the S-L problem

- We will show that if there are eigenfunctions for other choices of  $p(x)$ ,  $q(x)$ , and  $r(x)$ , then they all have certain common properties.
- It is actually remarkable how much we can infer about these solutions even though we cannot solve the ODE for general functions  $p(x)$ ,  $q(x)$ , and  $r(x)$ .
- We assume for now that the coefficient functions obey all the rules we discussed above

$$p(x) > 0 \quad r(x) > 0$$

- We also assume that the boundary conditions are also of the linear separable form discussed above.

$$\begin{aligned} c_1 u(a) + c_2 u'(a) &= 0, \\ d_1 u(b) + d_2 u'(b) &= 0 \end{aligned}$$

# Properties of the solutions to the S-L problem

- We will show the following results:
- All the eigenvalues of the Sturm-Liouville ODE are real.
- If  $\phi_1(x)$  and  $\phi_2(x)$  are two eigenfunctions corresponding to different eigenvalues  $\lambda_1$  and  $\lambda_2$  then the eigenfunctions are *orthogonal* in the following sense:

$$\int_a^b r(x)\phi_1(x)\phi_2(x)dx = 0.$$

- The eigenvalues of the Sturm-Liouville problem are all simple - that is there are no “multiple roots”

# More properties

- The sequence of eigenvalues  $\lambda_1, \lambda_2$  etc. can be ordered according to increasing magnitude.
- If this is done it is seen that  $\lambda_n \rightarrow +\infty$  as  $n \rightarrow \infty$ .
- In other words the eigenvalues have no point of accumulation (except at  $\infty$ )
- If an eigenfunction  $\phi_1$  has an eigenvalue  $\lambda_1$  and an eigenfunction  $\phi_2$  has an eigenvalue  $\lambda_2$  with  $\lambda_2 > \lambda_1$  then there is at least one zero of the eigenfunction  $\phi_2$  that lies between the zeroes of the eigenfunction of  $\phi_1$ .

# Expansions of arbitrary functions as series of eigenfunctions

- We say that any piecewise continuous function  $f(x)$  is “square integrable” if

$$\int_a^b |f(x)|^2 r(x) dx \text{ is finite}$$

- For all such functions there exists a series expansion in terms of Sturm-Liouville eigenfunctions that converges to the function in the sense of *mean-square*
- This means

$$\int_a^b \left| f(x) - \sum_{n=1}^N B_n \phi_n(x) \right|^2 r(x) dx \rightarrow 0 \text{ as } N \rightarrow \infty$$

- The coefficients of the series expansion are unique and can be obtained using the orthogonality properties of the eigenfunctions.

# Showing the eigenvalues are real

- We'll now show that all the eigenvalues are real.
- Suppose  $\lambda$  is an eigenvalue (possibly complex)
- In this case the corresponding eigenfunction is complex.
- Let the eigenvalue be expressed in terms of real and imaginary parts

$$\lambda = \mu + i\nu,$$

- Let the corresponding eigenfunction also be expressed in terms of real and imaginary parts:

$$\phi(x) = U(x) + iV(x).$$



# Showing the eigenvalues are real

- In Lagrange's identity set  $u = \phi$  and  $v = \phi$ .
- Then the complex form of Lagrange's identity is still

$$(L\phi, \phi) = (\phi, L\phi)$$

- Or in terms of integrals

$$\int_a^b dx \bar{\phi} L\phi = \int_a^b dx \phi L\bar{\phi}.$$

- But recall that the Sturm-Liouville ODE implies

$$L\phi = \lambda r(x)\phi,$$

- So

$$(\lambda r\phi, \phi) = (\phi, \lambda r\phi),$$

# Showing the eigenvalues are real

- So we have

$$\int_a^b \lambda r(x) \phi \bar{\phi} dx = \int_a^b \phi \bar{\lambda} r(x) \bar{\phi} dx,$$

- Another way to write this is

$$(\lambda - \bar{\lambda}) \int_a^b r(x) \phi \bar{\phi} dx = 0.$$

- But from complex numbers we know

$$\phi \bar{\phi} \geq 0$$

and can have the value zero only at isolated points

# Showing the eigenvalues are real

- So we must have

$$\int_a^b r(x) \phi \bar{\phi} dx > 0,$$

- So it must be that

$$\lambda = \bar{\lambda},$$

- This means that  $\lambda$  is real.
- This is the same kind of proof we used to show reality of eigenvalues for symmetric matrices.
- This also implies that the eigenfunction  $\phi(x)$  is real (up to some complex multiplicative constant which we could take to be real).
- We also see that we need only use the real version of the scalar product for these eigenfunctions because we've shown everything is real

# Showing the eigenvectors are orthogonal

- We next show that if  $\phi_1(x)$  and  $\phi_2(x)$  are two eigenfunctions corresponding to *differing eigenvalues*, that is

$$\lambda_1 \neq \lambda_2$$

then

$$\int_a^b \phi_1(x)\phi_2(x)r(x)dx = 0,$$

- if you think about this in terms of scalar products this means  $\phi_1$  and  $\phi_2$  are orthogonal relative to the weighting function  $r(x)$ .
- To show this, recall

$$L\phi_1 = \lambda r(x)\phi_1,$$

$$L\phi_2 = \lambda r(x)\phi_2.$$

# Showing the eigenvectors are orthogonal

- Now recall the Lagrange identity and set  $u = \phi_1$  and  $v = \phi_2$ .
- Then

$$(L\phi_1, \phi_2) - (\phi_1, L\phi_2) = 0,$$

- So this means

$$\lambda_1 \int_a^b \phi_1(x)\phi_2(x)r(x)dx - \lambda_2 \int_a^b \phi_1(x)\phi_2(x)r(x)dx = 0,$$

- This implies

$$(\lambda_1 - \lambda_2) \int_a^b \phi_1(x)\phi_2(x)r(x)dx = 0.$$

# Showing the eigenvectors are orthogonal

- But since  $\lambda_1 \neq \lambda_2$  it must be that

$$\int_a^b \phi_1(x)\phi_2(x)r(x)dx = 0.$$

and so the eigenvectors are orthogonal relative to the weighting factor  $r(x)$ .

- Note the scalar product for orthogonality is different from that used in the Lagrange identity unless  $r(x) = 1$
- It is always possible to redefine the S-L problem slightly (as well as the eigenfunctions) so that the weighting function is always 1 provided  $r(x) > 0$ .
- Note the proof above is really the same as that used to show orthogonality of eigenvectors of symmetric matrices.