Physics 106b — Classical Mechanics

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Parametrically Driven Oscillators

Outline

- Parametric drive: what is it?
- Square wave drive exact solution
- Sinusoidal drive (Mathieu equation) harmonic analysis
- Applications

Parametrically driven oscillators

Hill equation: A *parameter* of an oscillator is periodically modulated. The general form for a linear oscillator is the *Hill equation*

$$\ddot{q} + a(t)\dot{q} + b(t)q = 0$$

with a(t), b(t) periodic functions with some period T.

Floquet's theorem: For a linear ODE with coefficients that are periodic in time with period T we can find solutions in the form

$$q(t) = e^{\sigma t} P(t)$$

with σ a constant, complex in general, and P(t) a periodic function with the same period T: P(t + T) = P(t).

Stability: Re σ plays the role of a stability parameter: Re $\sigma > 0$ indicates an exponentially growing solution.

Hand and Finch write σ as $i\mu$.

Parametric drive: square wave

Equation of motion

$$\ddot{\theta} + \gamma \dot{\theta} + F(t)\theta = 0 \quad \text{with} \quad F(t) = \begin{cases} 1 + r & \text{for } 0 \le t < T/2 \\ 1 - r & \text{for } T/2 \le t < T \end{cases}$$

• General solutions for $F(t) = 1 \pm r$

$$\theta_{\pm}(t) = e^{-\gamma t/2} [A_{\pm} \cos \omega_{\pm} t + B_{\pm} \sin \omega_{\pm} t] \quad \text{with } \omega_{\pm} = \sqrt{1 \pm r - (\gamma/2)^2}$$

■ The "cosine" and "sine" like solutions over each half cycle

$$\theta_{c\pm}(t) = e^{-\gamma t/2} \left[\cos \omega_{\pm} t + \frac{\gamma}{2\omega_{\pm}} \sin \omega_{\pm} t \right]$$

$$\theta_{s\pm}(t) = e^{-\gamma t/2} \frac{1}{\omega_{+}} \sin \omega_{\pm} t$$

■ Calculate the monodromy matrix M matrix by multiplying M_+ and M_- for the two half cycles

$$M = \begin{bmatrix} \theta_{c-}(T/2) & \theta_{s-}(T/2) \\ \dot{\theta}_{c-}(T/2) & \dot{\theta}_{s-}(T/2) \end{bmatrix} \begin{bmatrix} \theta_{c+}(T/2) & \theta_{s+}(T/2) \\ \dot{\theta}_{c+}(T/2) & \dot{\theta}_{s+}(T/2) \end{bmatrix}$$

Parametric drive: sinusoidal

Sinusoidal drive gives the damped Mathieu equation

$$\ddot{\theta} + \gamma \dot{\theta} + [1 + h \cos \omega_d t]\theta = 0$$

Study primary n=1 tongue $\omega_d\approx 2$ using perturbation theory for small γ,h Floquet's theorem: look for a solution

$$\theta(t) = e^{\sigma t} \left[A e^{i\omega_d t/2} + \text{c.c.} \right] + \text{harmonics}$$

Use method of Harmonic Analysis: substitute in to the EOM and set the coefficient of $e^{i\omega_d t/2}$ term to zero.

Tricky term: $\cos \omega_d t \ \theta = \frac{1}{2} (e^{i\omega_d t} + e^{-i\omega_d t}) (Ae^{i\omega_d t/2} + A^* e^{-i\omega_d t/2}) e^{\sigma t}$

$$\left[-\left(\frac{i\omega_d}{2} + \sigma\right)^2 + \gamma \left(\frac{i\omega_d}{2} + \sigma\right) + 1 \right] A + \frac{h}{2} A^* = 0$$

For $\omega_d \approx 2$ and leaving out terms in σ^2 and $\sigma \gamma$ this can be approximated

$$\left[\left(2-\omega_{d}\right)+i\left(2\sigma+\gamma\right)\right]A=-\frac{1}{2}hA^{*}$$

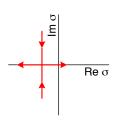
Parametric drive: sinusoidal

$$[(2 - \omega_d) + i (2\sigma + \gamma)] A = -\frac{1}{2} h A^*$$

Magnitude-squared gives the growth rate

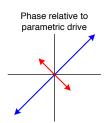
$$\sigma = -\frac{\gamma}{2} \pm \frac{1}{4} \sqrt{h^2 - (\omega_d - 2)^2}$$

For $\omega_d = 2$ instability occurs for $h > 2\gamma$

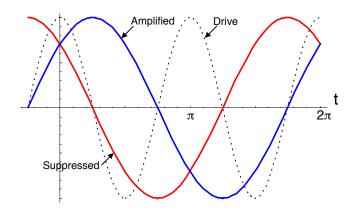


On resonance and at instability, $A = |A|e^{i\delta}$ with δ the phase relative to the drive

$$\sigma=0 \Rightarrow iA=-A^*$$
 so that $\delta=\frac{\pi}{4}$
$$\sigma=-\gamma \Rightarrow -iA=-A^*$$
 so that $\delta=-\frac{\pi}{4}$



Parametric drive: sinusoidal

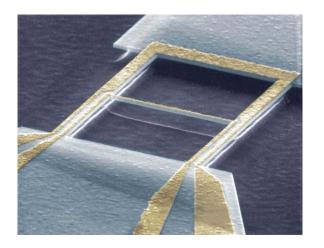


Application

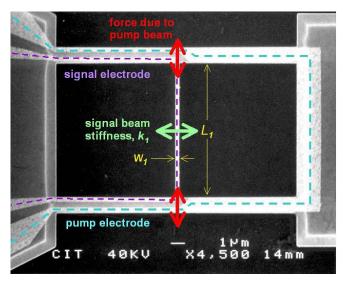
Nanoscale mechanical amplifier

- Darrell Harrington, Caltech Thesis (2002)
- Karabalin, Feng, and Roukes, Nano Letters 9, 3116 (2009)
- Karabalin, Masmanidis, and Roukes, Appl. Phys. Lett. 97, 183101 (2010)

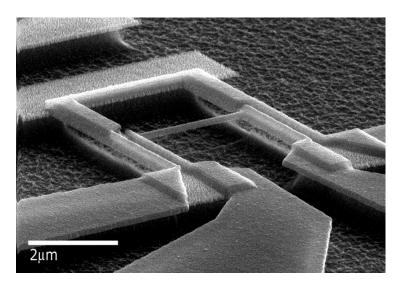
Nanomechanical Beam Resonator



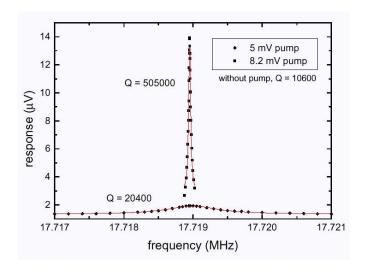
Nanomechanical Beam Resonator



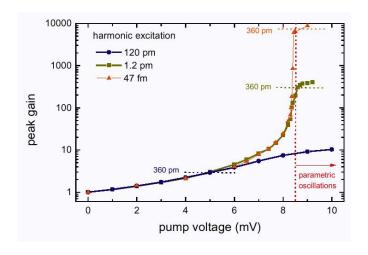
Nanomechanical Beam Resonator



Parametric Drive: Enhanced Q



Parametric Amplification



Squeezing of Thermal Noise

