ACM 100b

Properties of the Fourier transform

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Recap

• In the previous lecture we introduced the Fourier transform pair:

$$F(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) \exp(-ikx) dx,$$

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(k) \exp(ikx) dk$$

- We can transform any function which is integrable over the interval $-\infty < x < \infty$ to get the transform F(k)
- By thinking of F as a function of a complex variable k we can extend the transform to a wider class of functions.
- The original function can then be recovered by performing the second integral given F(k).



Recap

- We saw that the Fourier transform is like an extension of Fourier series except for the singular situation of an infinite interval -∞ < x < ∞
- There is a completeness relation that bears some similarity to the completeness relation for discrete sets of complete functions:

$$\delta(x - x_0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \exp(ikx) \exp(-ikx_0) dk$$

- In this lecture we will introduce some additional properties of the Fourier transform
- We will also see how the transform can be used to help solve various ODE's in the interval $-\infty < x < +\infty$
- We will then introduce some new transforms, the sine and cosine transforms appropriate for the interval $0 \le x < \infty$

Linearity of the Fourier transform

- Fourier transforms are very useful for solving constant coefficient ODE's and PDE's on infinite domains.
- The most important property is linearity.
- Let \mathcal{F} denote the operation of performing a Fourier transform:

$$\mathcal{F}[f(x)] \equiv \int_{-\infty}^{\infty} f(x) \exp(-ikx) dx.$$

Then it is easy to see that

$$\mathcal{F}[f(x)+g(x)]=\mathcal{F}[f(x)]+\mathcal{F}[g(x)].$$

• It is this property plus the way the transform behaves under differentiation that makes the transform useful for linear problems.



Transform of a derivative

- Another important property is that derivatives of functions transform simply.
- It is customary to denote the transform of a function y(x) by Y(k).
- Assume also that the function under consideration y(x) has the property that

$$y(x) \to 0$$
 as $|x| \to \infty$.

- Actually, this requirement can be relaxed by considering complex values of k and thinking of the transform as a contour integral.
- Now consider the transform of the derivative:

$$\frac{1}{\sqrt{2\pi}}\int_{-\infty}^{\infty}y'(x)\exp(-ikx)dx.$$



Transform of a derivative

Using integration by parts we have

$$\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} y'(x) \exp(-ikx) dx = \frac{1}{\sqrt{2\pi}} y(x) \exp(-ikx)|_{-\infty}^{\infty}$$
$$+ ik \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \exp(-ikx) y(x) dx.$$

As long as we have

$$y(x) \to 0$$
 as $|x| \to \infty$,

we see that

$$\mathcal{F}[y'(x)] = ikY(k).$$



Transform of a derivative

And repeating this analysis for the second derivative

$$\mathcal{F}[y''(x)] = -k^2 Y(k),$$

as long as

$$y'(x) \to 0$$
 as $|x| \to \infty$.

- This property is very important in that it can be used to transform differential equations into algebraic equations.
- More about this later.
- One thing that is sometimes forgotten is that if you write

$$\mathcal{F}[y'(x)] = ikY(k).$$

you have actually implicitly assumed

$$y(x) \to 0$$
 as $|x| \to \infty$,

 This is a useful way to impose boundary conditions but you need to be aware that you are doing this.

- Let $\mathcal{F}[f(x)]$ be the Fourier transform for f(x) and written as F(k)
- Correspondingly let $\mathcal{F}[g(x)]$ be given by G(k).
- Define the *convolution* of f and g by

$$f\star g\equiv \frac{1}{\sqrt{2\pi}}\int_{-\infty}^{\infty}f(\zeta)g(x-\zeta)d\zeta.$$

- We assume f and g behave such that this integral exists.
- Then the convolution theorem for Fourier transforms states that

$$\mathfrak{F}[f\star g]=F(k)G(k).$$



To show this consider the integral

$$\int_{-\infty}^{\infty} F(k)G(k) \exp(ikx)dk,$$

• This integral can be manipulated as follows:

$$\int_{-\infty}^{\infty} F(k)G(k)e^{ikx}dk = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} G(k)e^{ikx}dk \int_{-\infty}^{\infty} f(\zeta)e^{-ik\zeta}d\zeta$$
$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(\zeta) \int_{-\infty}^{\infty} G(k)e^{(ik(x-\zeta))}dk d\zeta$$
$$= \int_{-\infty}^{\infty} f(\zeta)g(x-\zeta)d\zeta,$$

• Multiply both sides of this relation by $1/\sqrt{2\pi}$ and this shows the result.

Consider the following convolution:

$$\frac{1}{\sqrt{2\pi}}\int_{-\infty}^{\infty}f(\zeta)\delta(x-\zeta)d\zeta.$$

- In other words we're convolving some function f(x) with a δ function.
- We can apply the convolution theorem to take the Fourier transform of this integral.
- The Fourier transform of f(x) can simply be labeled F(k).
- The Fourier transform of $\delta(x)$ has been computed already:

$$\mathcal{F}[\delta(x)] = \frac{1}{\sqrt{2\pi}}.$$



The convolution theorem tells us that

$$\frac{1}{\sqrt{2\pi}}\mathcal{F}\left[\int_{-\infty}^{\infty}f(\zeta)\delta(x-\zeta)d\zeta\right]=\frac{F(k)}{\sqrt{2\pi}}$$

- Now consider taking the inverse Fourier transform of both sides of this relationship
- From this we infer immediately that it must be that

$$\int_{-\infty}^{\infty} f(\zeta)\delta(x-\zeta)d\zeta = f(x),$$

ullet This is just the sifting property for the δ function we had derived earlier.

Parseval's theorem again

The convolution theorem says

$$\int_{-\infty}^{\infty} F(k)G(k) \exp(ikx)dk = \int_{-\infty}^{\infty} f(\zeta)g(x-\zeta)d\zeta.$$

- Suppose we set x = 0 in the convolution theorem.
- We then get

$$\int_{-\infty}^{\infty} F(k)G(k)dk = \int_{-\infty}^{\infty} f(\zeta)g(-\zeta)d\zeta.$$

Now set

$$g(-\zeta) = \overline{f(\zeta)},$$

where we now want to think of f and g as complex.

So this means

$$g(\zeta) = \overline{f(-\zeta)}$$



Parseval's theorem again

The Fourier transform of g is given by

$$G(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \overline{f(-\zeta)} \exp(-ik\zeta) d\zeta$$

 From this definition of g and the definition of the Fourier transform we can infer that

$$G(k) = \overline{F(k)}$$
.

Substitute this into the result

$$\int_{-\infty}^{\infty} F(k)G(k)dk = \int_{-\infty}^{\infty} f(\zeta)g(-\zeta)d\zeta,$$

We get

$$\int_{-\infty}^{\infty} |F(k)|^2 dk = \int_{-\infty}^{\infty} |f(x)|^2 dx.$$

Parseval's theorem again

 Recall when we were dealing with Fourier series we derived the Parseval relation:

$$\int_{-L}^{L} f(x)^{2} dx = L \left[2B_{0}^{2} + \sum_{n=1}^{\infty} \left(A_{n}^{2} + B_{n}^{2} \right) \right]$$

- In this relation the A_n and B_n are the Fourier coefficients over the interval $-L \le x \le L$.
- The Parseval relation for the Fourier transform is the extension of this result to the fully infinite interval $-\infty < x < \infty$.