§0.1 曲面的结构方程

利用自然标架运动方程进一步研究第一、第二基本形式系数之间的关系。关系式(Gauss方程与Codazzi方程)的导出基于一个朴素的出发点:光滑函数的二阶偏导数可交换求导次序。

回顾曲面自然标架 $(r; r_1, r_2, N)$ 的运动方程

$$\begin{cases} \frac{\partial r}{\partial u^{\alpha}} = r_{\alpha}, & \alpha = 1, 2; \quad (M_{1}) \\ \frac{\partial r_{\alpha}}{\partial u^{\beta}} = \Gamma^{\gamma}_{\beta\alpha} r_{\gamma} + b_{\alpha\beta} N, & \alpha, \beta = 1, 2; \quad (M_{2}) \\ \frac{\partial N}{\partial u^{\alpha}} = -b^{\beta}_{\alpha} r_{\beta}, & \alpha = 1, 2 \quad (M_{3}) \end{cases}$$

其中

$$\begin{split} g_{\alpha\beta} &= \langle r_{\alpha}, r_{\beta} \rangle, \quad \alpha, \beta = 1, 2; \\ b_{\alpha\beta} &= \langle r_{\alpha\beta}, N \rangle = - \langle r_{\alpha}, N_{\beta} \rangle, \quad \alpha, \beta = 1, 2; \\ b_{\alpha}^{\beta} &= b_{\alpha\gamma} g^{\gamma\beta} = g^{\beta\gamma} b_{\gamma\alpha}, \quad \alpha, \beta = 1, 2; \\ \Gamma_{\alpha\beta}^{\gamma} &= \Gamma_{\beta\alpha}^{\gamma} = \frac{1}{2} g^{\gamma\xi} (\frac{\partial g_{\beta\xi}}{\partial u^{\alpha}} + \frac{\partial g_{\alpha\xi}}{\partial u^{\beta}} - \frac{\partial g_{\alpha\beta}}{\partial u^{\xi}}), \quad \alpha, \beta, \gamma = 1, 2. \end{split}$$

 (M_1) 再求一次偏导数,即 (M_2) ,可交换次序得到

$$\Gamma^{\gamma}_{\beta\alpha} = \Gamma^{\gamma}_{\alpha\beta}, \quad b_{\alpha\beta} = b_{\beta\alpha}.$$

接下来对 (M_2) , (M_3) 再求一次偏导数,并交换次序。

先看 (M_3) :对 (M_3) 求一次偏导数并利用运动方程可得

$$\begin{split} N_{\alpha\gamma} &= & \frac{\partial}{\partial u^{\gamma}} \frac{\partial N}{\partial u^{\alpha}} = \frac{\partial}{\partial u^{\gamma}} (-b_{\alpha}^{\beta} r_{\beta}) \\ &= & -\frac{\partial b_{\alpha}^{\beta}}{\partial u^{\gamma}} r_{\beta} - b_{\alpha}^{\beta} (\Gamma_{\gamma\beta}^{\xi} r_{\xi} + b_{\beta\gamma} N). \end{split}$$

同样有

$$\begin{split} N_{\gamma\alpha} &= \frac{\partial}{\partial u^{\alpha}} \frac{\partial N}{\partial u^{\gamma}} = \frac{\partial}{\partial u^{\alpha}} (-b_{\gamma}^{\beta} r_{\beta}) \\ &= -\frac{\partial b_{\gamma}^{\beta}}{\partial u^{\alpha}} r_{\beta} - b_{\gamma}^{\beta} (\Gamma_{\alpha\beta}^{\xi} r_{\xi} + b_{\beta\alpha} N). \end{split}$$

其中法向分量的相等已知,即

$$-b_{\alpha}^{\beta}b_{\beta\gamma} = -b_{\alpha\eta}g^{\eta\beta}b_{\beta\gamma} = -b_{\alpha\eta}b_{\gamma}^{\eta} = -b_{\alpha\beta}b_{\gamma}^{\beta}.$$

由切向分量相等可得Codazzi方程

$$\frac{\partial b_{\alpha}^{\xi}}{\partial u^{\gamma}} + \Gamma_{\gamma\beta}^{\xi} b_{\alpha}^{\beta} = \frac{\partial b_{\gamma}^{\xi}}{\partial u^{\alpha}} + \Gamma_{\alpha\beta}^{\xi} b_{\gamma}^{\beta}. \quad \text{(Codazzi')}$$

对 (M_2) 求一次偏导数即坐标切向量的两阶方向导数。为分清其中只与第一基本形式有关的部分,引入如下记号:一个向量 $Y \in T_P \mathbb{R}^3$ 在切平面和法线上的投影分别记为 Y^T, Y^{\perp} ,即

$$Y^T = Y - \langle Y, N \rangle N, \quad Y^{\perp} = \langle Y, N \rangle N.$$

设 $X = X^{\beta}(u, v)r_{\beta}$ 为切向量场,定义

$$\nabla_{\frac{\partial}{\partial u^{\alpha}}} X = (\frac{\partial X}{\partial u^{\alpha}})^T := \frac{\partial X}{\partial u^{\alpha}} - \langle \frac{\partial X}{\partial u^{\alpha}}, N \rangle N.$$

特别有

$$\nabla_{\frac{\partial}{\partial u^{\beta}}} r_{\alpha} = (\frac{\partial r_{\alpha}}{\partial u^{\beta}})^{T} = \Gamma_{\beta\alpha}^{\xi} r_{\xi},$$

$$\nabla_{\frac{\partial}{\partial u^{\gamma}}} (\nabla_{\frac{\partial}{\partial u^{\beta}}} r_{\alpha}) = \left[\frac{\partial}{\partial u^{\gamma}} (\Gamma_{\beta\alpha}^{\xi} r_{\xi}) \right]^{T} = \left[\frac{\partial \Gamma_{\beta\alpha}^{\xi}}{\partial u^{\gamma}} r_{\xi} + \Gamma_{\beta\alpha}^{\eta} \frac{\partial r_{\eta}}{\partial u^{\gamma}} \right]^{T}$$
$$= \left(\frac{\partial \Gamma_{\beta\alpha}^{\xi}}{\partial u^{\gamma}} + \Gamma_{\gamma\eta}^{\xi} \Gamma_{\beta\alpha}^{\eta} \right) r_{\xi}.$$

而

$$[\frac{\partial}{\partial u^{\gamma}}(\Gamma^{\xi}_{\beta\alpha}r_{\xi})]^{\perp} = [\frac{\partial \Gamma^{\xi}_{\beta\alpha}}{\partial u^{\gamma}}r_{\xi} + \Gamma^{\eta}_{\beta\alpha}\frac{\partial r_{\eta}}{\partial u^{\gamma}}]^{\perp} = \Gamma^{\eta}_{\beta\alpha}b_{\eta\gamma}N.$$

由 (M_2) 计算

$$\begin{split} \frac{\partial}{\partial u^{\gamma}} (\frac{\partial r_{\alpha}}{\partial u^{\beta}}) &= \frac{\partial}{\partial u^{\gamma}} (\Gamma^{\xi}_{\beta\alpha} r_{\xi} + b_{\alpha\beta} N) \\ &= [\frac{\partial}{\partial u^{\gamma}} (\Gamma^{\xi}_{\beta\alpha} r_{\xi})]^{T} + [\frac{\partial}{\partial u^{\gamma}} (\Gamma^{\xi}_{\beta\alpha} r_{\xi})]^{\perp} + \frac{\partial b_{\alpha\beta}}{\partial u^{\gamma}} N + b_{\alpha\beta} N_{\gamma} \\ &= \nabla_{\frac{\partial}{\partial u^{\gamma}}} (\nabla_{\frac{\partial}{\partial u^{\beta}}} r_{\alpha}) + \Gamma^{\eta}_{\beta\alpha} b_{\eta\gamma} N + \frac{\partial b_{\alpha\beta}}{\partial u^{\gamma}} N + b_{\alpha\beta} (-b^{\xi}_{\gamma} r_{\xi}) \\ &= [\nabla_{\frac{\partial}{\partial u^{\gamma}}} (\nabla_{\frac{\partial}{\partial u^{\beta}}} r_{\alpha}) - b_{\alpha\beta} b^{\xi}_{\gamma} r_{\xi}] + (\frac{\partial b_{\beta\alpha}}{\partial u^{\gamma}} + \Gamma^{\xi}_{\beta\alpha} b_{\xi\gamma}) N. \end{split}$$

交换 β , γ 的次序得

$$\frac{\partial}{\partial u^\beta}(\frac{\partial r_\alpha}{\partial u^\gamma}) = \left[\nabla_{\frac{\partial}{\partial u^\beta}}(\nabla_{\frac{\partial}{\partial u^\gamma}}r_\alpha) - b_{\alpha\gamma}b_\beta^\xi r_\xi\right] + \left(\frac{\partial b_{\gamma\alpha}}{\partial u^\beta} + \Gamma_{\gamma\alpha}^\xi b_{\xi\beta}\right)N.$$

由它们相等得到

$$\begin{split} & \nabla_{\frac{\partial}{\partial u^{\gamma}}}(\nabla_{\frac{\partial}{\partial u^{\beta}}}r_{\alpha}) - \nabla_{\frac{\partial}{\partial u^{\beta}}}(\nabla_{\frac{\partial}{\partial u^{\gamma}}}r_{\alpha}) \\ & = & (\frac{\partial \Gamma^{\xi}_{\beta\alpha}}{\partial u^{\gamma}} + \Gamma^{\xi}_{\gamma\eta}\Gamma^{\eta}_{\beta\alpha})r_{\xi} - (\frac{\partial \Gamma^{\xi}_{\gamma\alpha}}{\partial u^{\beta}} + \Gamma^{\xi}_{\beta\eta}\Gamma^{\eta}_{\gamma\alpha})r_{\xi} \\ & = & b_{\alpha\beta}b^{\xi}_{\gamma}r_{\xi} - b_{\alpha\gamma}b^{\xi}_{\beta}r_{\xi}, \end{split}$$

以及

$$\frac{\partial b_{\beta\alpha}}{\partial u^{\gamma}} + \Gamma^{\xi}_{\beta\alpha}b_{\xi\gamma} = \frac{\partial b_{\gamma\alpha}}{\partial u^{\beta}} + \Gamma^{\xi}_{\gamma\alpha}b_{\xi\beta}.$$

分别称为Gauss方程

$$\frac{\partial \Gamma_{\beta\alpha}^{\xi}}{\partial u^{\gamma}} - \frac{\partial \Gamma_{\gamma\alpha}^{\xi}}{\partial u^{\beta}} + \Gamma_{\gamma\eta}^{\xi} \Gamma_{\beta\alpha}^{\eta} - \Gamma_{\beta\eta}^{\xi} \Gamma_{\gamma\alpha}^{\eta} = b_{\gamma}^{\xi} b_{\beta\alpha} - b_{\beta}^{\xi} b_{\gamma\alpha}; \quad (Gauss)$$

与Codazzi方程(稍后验证它与(Codazzi')等价)

$$\frac{\partial b_{\beta\alpha}}{\partial u^{\gamma}} - \Gamma^{\xi}_{\gamma\alpha}b_{\xi\beta} = \frac{\partial b_{\gamma\alpha}}{\partial u^{\beta}} - \Gamma^{\xi}_{\beta\alpha}b_{\xi\gamma}. \quad (\text{Codazzi})$$

合起来称为曲面的Gauss-Codazzi方程,或称为曲面的结构方程。

定理0.1. 曲面结构方程包括

$$\frac{\partial \Gamma_{\beta\alpha}^{\xi}}{\partial u^{\gamma}} - \frac{\partial \Gamma_{\gamma\alpha}^{\xi}}{\partial u^{\beta}} + \Gamma_{\gamma\eta}^{\xi} \Gamma_{\beta\alpha}^{\eta} - \Gamma_{\beta\eta}^{\xi} \Gamma_{\gamma\alpha}^{\eta} = b_{\gamma}^{\xi} b_{\beta\alpha} - b_{\beta}^{\xi} b_{\gamma\alpha}; \quad (Gauss)$$

和

$$\frac{\partial b_{\beta\alpha}}{\partial u^{\gamma}} - \Gamma^{\xi}_{\gamma\alpha}b_{\xi\beta} = \frac{\partial b_{\gamma\alpha}}{\partial u^{\beta}} - \Gamma^{\xi}_{\beta\alpha}b_{\xi\gamma}. \quad (Codazzi)$$

Codazzi方程的两种形式

$$\frac{\partial b_{\beta\alpha}}{\partial u^{\gamma}} - \Gamma^{\xi}_{\gamma\alpha}b_{\xi\beta} = \frac{\partial b_{\gamma\alpha}}{\partial u^{\beta}} - \Gamma^{\xi}_{\beta\alpha}b_{\xi\gamma}. \quad (\text{Codazzi})$$

$$\frac{\partial b_{\alpha}^{\xi}}{\partial u^{\gamma}} + \Gamma_{\gamma\beta}^{\xi} b_{\alpha}^{\beta} = \frac{\partial b_{\gamma}^{\xi}}{\partial u^{\alpha}} + \Gamma_{\alpha\beta}^{\xi} b_{\gamma}^{\beta}. \quad (Codazzi')$$

将验证这两种形式等价。为方便与(Codazzi)比较,将(Codazzi')改写为

$$\frac{\partial b_{\beta}^{\eta}}{\partial u^{\gamma}} + \Gamma_{\gamma p}^{\eta} b_{\beta}^{p} = \frac{\partial b_{\gamma}^{\eta}}{\partial u^{\beta}} + \Gamma_{\beta p}^{\eta} b_{\gamma}^{p}. \quad \text{(Codazzi')}$$

从而等价性由下面论断直接得到:

$$g^{\alpha\eta}(\frac{\partial b_{\beta\alpha}}{\partial u^{\gamma}} - \Gamma^{\xi}_{\gamma\alpha}b_{\xi\beta}) = \frac{\partial b^{\eta}_{\beta}}{\partial u^{\gamma}} + \Gamma^{\eta}_{\gamma p}b^{p}_{\beta}. \quad (1)$$

先证明一个常用关系式:

$$\frac{\partial g^{\alpha\beta}}{\partial u^{\gamma}} = -g^{\alpha p} g^{\beta q} \frac{\partial g_{pq}}{\partial u^{\gamma}}.$$

证明:

$$\frac{\partial (g^{\alpha p}g_{pq})}{\partial u^{\gamma}} = 0 = \frac{\partial g^{\alpha p}}{\partial u^{\gamma}}g_{pq} + g^{\alpha p}\frac{\partial g_{pq}}{\partial u^{\gamma}}$$

乘以 $g^{\beta q}$ 并对q求和可得上述关系式。

证明(1)式:

$$\begin{split} g^{\alpha\eta}(\frac{\partial b_{\beta\alpha}}{\partial u^{\gamma}} - \Gamma^{\xi}_{\gamma\alpha}b_{\xi\beta}) &= \frac{\partial b^{\eta}_{\beta}}{\partial u^{\gamma}} - \frac{\partial g^{\alpha\eta}}{\partial u^{\gamma}}b_{\beta\alpha} - g^{\alpha\eta}\frac{1}{2}g^{\xi p}(\frac{\partial g_{\alpha p}}{\partial u^{\gamma}} + \frac{\partial g_{\gamma p}}{\partial u^{\alpha}} - \frac{\partial g_{\alpha\gamma}}{\partial u^{p}})b_{\xi\beta} \\ &= \frac{\partial b^{\eta}_{\beta}}{\partial u^{\gamma}} + g^{\alpha p}g^{\eta q}\frac{\partial g_{pq}}{\partial u^{\gamma}}b_{\beta\alpha} - g^{\alpha\eta}\frac{1}{2}(\frac{\partial g_{\alpha p}}{\partial u^{\gamma}} + \frac{\partial g_{\gamma p}}{\partial u^{\alpha}} - \frac{\partial g_{\alpha\gamma}}{\partial u^{p}})b^{p}_{\beta} \\ &= \frac{\partial b^{\eta}_{\beta}}{\partial u^{\gamma}} + b^{p}_{\beta}g^{\alpha\eta}[\frac{\partial g_{p\alpha}}{\partial u^{\gamma}} - \frac{1}{2}(\frac{\partial g_{\alpha p}}{\partial u^{\gamma}} + \frac{\partial g_{\gamma p}}{\partial u^{\alpha}} - \frac{\partial g_{\alpha\gamma}}{\partial u^{p}})] \quad \text{by } q \to \alpha \\ &= \frac{\partial b^{\eta}_{\beta}}{\partial u^{\gamma}} + b^{p}_{\beta}\frac{1}{2}g^{\alpha\eta}[\frac{\partial g_{\alpha p}}{\partial u^{\gamma}} + \frac{\partial g_{\gamma \alpha}}{\partial u^{p}} - \frac{\partial g_{\gamma p}}{\partial u^{\alpha}}] \\ &= \frac{\partial b^{\eta}_{\beta}}{\partial u^{\gamma}} + \Gamma^{\eta}_{\gamma p}b^{p}_{\beta}. \end{split}$$

Codazzi方程

$$\frac{\partial b_{\beta\alpha}}{\partial u^{\gamma}} - \Gamma^{\xi}_{\gamma\alpha}b_{\xi\beta} = \frac{\partial b_{\gamma\alpha}}{\partial u^{\beta}} - \Gamma^{\xi}_{\beta\alpha}b_{\xi\gamma} \quad \text{(Codazzi)}$$

中令 $\gamma = 1, \beta = 2$,分别取 $\alpha = 1, 2$ 得到两个独立方程:

$$\frac{\partial b_{21}}{\partial u^{1}} - \Gamma_{11}^{\xi} b_{2\xi} = \frac{\partial b_{11}}{\partial u^{2}} - \Gamma_{21}^{\xi} b_{1\xi},$$
$$\frac{\partial b_{22}}{\partial u^{1}} - \Gamma_{12}^{\xi} b_{2\xi} = \frac{\partial b_{12}}{\partial u^{2}} - \Gamma_{22}^{\xi} b_{1\xi}.$$

Codazzi方程给出第二基本形式一阶导数之间的(交换)关系。

例:如果参数(u,v)使得 r_u,r_v 为互相正交的主方向,即其积分曲线为曲率线。此时F=M=0,即 $g_{12}=b_{12}=0$,由此可化简Codazzi方程。注:非脐点P,存在它的邻域以及参数(u,v)使得u,v-曲线为曲率线[do Carmo, § 3.4]。

$$\frac{\partial b_{21}}{\partial u^1} - \Gamma_{11}^{\xi} b_{2\xi} = \frac{\partial b_{11}}{\partial u^2} - \Gamma_{21}^{\xi} b_{1\xi},$$

$$(g^{\alpha\beta}) = \frac{1}{\det(g_{\alpha\beta})} \begin{pmatrix} g_{22} & -g_{12} \\ -g_{21} & g_{11} \end{pmatrix} = \frac{1}{EG - F^2} \begin{pmatrix} G & -F \\ -F & E \end{pmatrix}$$

$$H = \frac{1}{2}tr(W) = \frac{1}{2}(k_1 + k_2) = \frac{1}{2}\frac{LG - 2MF + NE}{EG - F^2},$$

可知

$$\begin{split} \frac{\partial b_{11}}{\partial u^2} &= L_v = \Gamma_{21}^{\xi} b_{1\xi} - \Gamma_{11}^{\xi} b_{2\xi} = \Gamma_{21}^1 L - \Gamma_{11}^2 N \\ &= \frac{1}{2} L g^{11} E_v - \frac{1}{2} N g^{22} (-E_v) \\ &= \frac{1}{2} E_v (L \frac{G}{EG} + N \frac{E}{EG}) \\ &= \frac{1}{2} (\frac{L}{E} + \frac{N}{G}) E_v \\ &= H E_v, \end{split}$$

类似也有

$$N_u = HG_u$$
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