

微分方程

二维、三维Laplace方程
的分离变量法

二维、三维Laplace方程的边值问题

■ 矩形域上的边值问题

散热片的横截面为一矩形 $[0, a] \times [0, b]$ ，它的一边 $y=b$ 保持较高温 U ，其它三边保持零度。求横截面上的稳态温度分布。

$$\text{解: } \begin{cases} u_{xx} + u_{yy} = 0, (x, y) \in (0, a) \times (0, b) \\ u|_{x=0} = u|_{x=a} = 0, y \in [0, b] \\ u|_{y=0} = 0, u|_{y=b} = U, x \in [0, a] \end{cases}$$

考虑分离解 $X(x)Y(y) \equiv 0$ 且 $X(0) = X(a) = 0$. 代入方程有

$$\frac{X''(x)}{X(x)} = -\frac{Y''(y)}{Y(y)} := -\lambda \Rightarrow \begin{cases} X''(x) + \lambda X(x) = 0, 0 < x < a \\ X(0) = X(a) = 0 \end{cases}$$

\Rightarrow 特征值 $\lambda_n = (\frac{n\pi}{a})^2$, 特征函数 $X_n(x) = \sin(\frac{n\pi}{a}x), n \geq 1$.

而 $Y''(y) - \lambda_n Y(y) = 0$ 的通解

为 $Y_n(y) = C_n \cosh(\frac{n\pi}{a}y) + D_n \sinh(\frac{n\pi}{a}y)$.

$$\begin{aligned} \text{令形式解 } u(x, y) &= \sum_{n \geq 1} X_n(x) Y_n(y) \\ &= \sum_{n \geq 1} [C_n \cosh(\frac{n\pi}{a}y) + D_n \sinh(\frac{n\pi}{a}y)] \sin(\frac{n\pi}{a}x), \end{aligned}$$

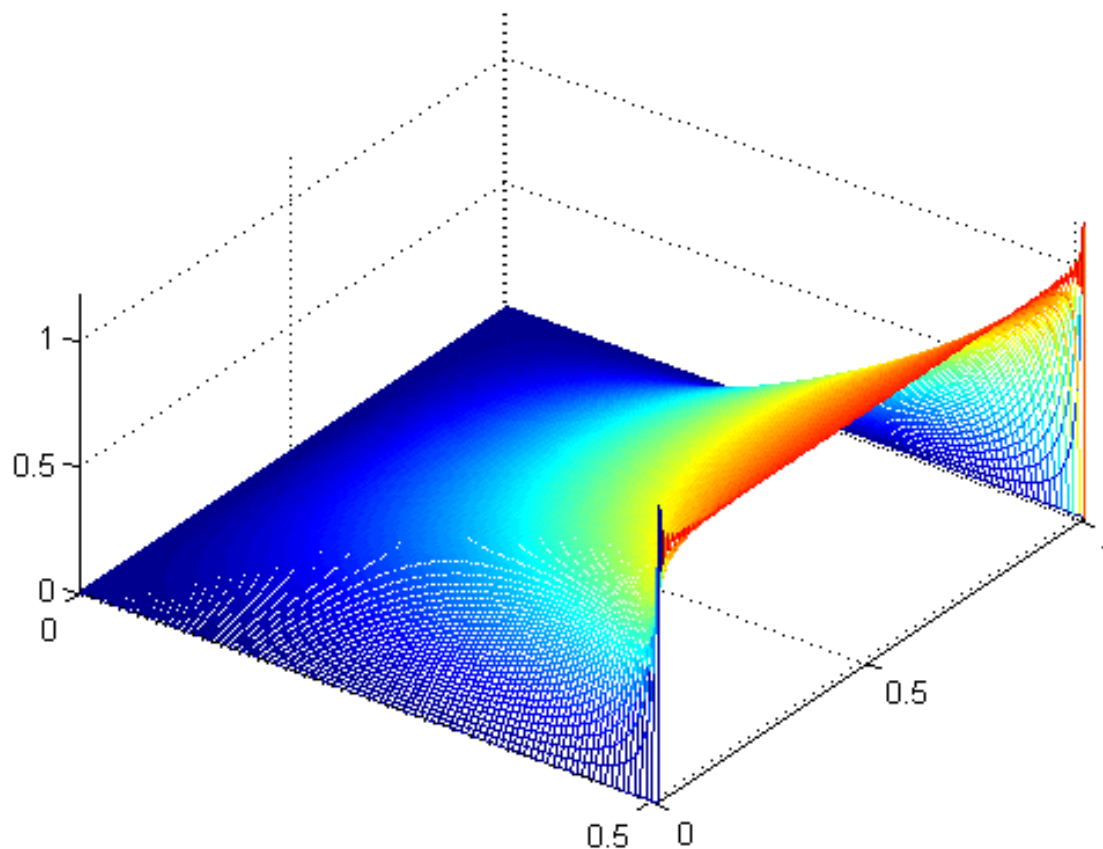
代入边界条件 $u|_{y=0} = 0, u|_{y=b} = U$, 有

$$\sum_{n \geq 1} C_n \sin(\frac{n\pi}{a}x) = 0 \Rightarrow C_n = 0,$$

$$\sum_{n \geq 1} D_n \sinh(\frac{n\pi}{a}b) \sin(\frac{n\pi}{a}x) = U \Rightarrow D_n = \frac{1}{\sinh(\frac{n\pi}{a}b)} \frac{\langle U, X_n \rangle}{\|X_n\|^2}$$

$$\begin{aligned} &= \frac{1}{\sinh(\frac{n\pi}{a}b)} \frac{U \int_0^a \sin(\frac{n\pi}{a}x) dx}{\int_0^a \sin^2(\frac{n\pi}{a}x) dx} = \frac{1}{\sinh(\frac{n\pi}{a}b)} \frac{U \frac{a}{n\pi} [1 - (-1)^n]}{\frac{a}{2}} = \frac{2U [1 - (-1)^n]}{n\pi \sinh(\frac{n\pi}{a}b)}. \end{aligned}$$

$$\therefore u(x, y) = \frac{4}{\pi} U \sum_{n=0}^{\infty} \frac{1}{2n+1} \frac{\sinh\left(\frac{2n+1}{a} \pi y\right)}{\sinh\left(\frac{2n+1}{a} \pi b\right)} \sin\left(\frac{2n+1}{a} \pi x\right)$$



参数选取

$$a = 1$$

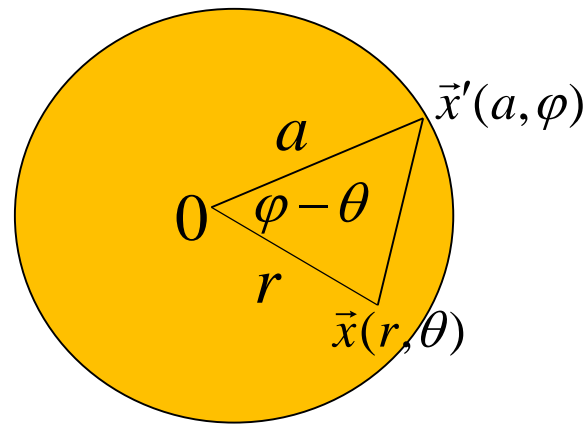
$$b = \frac{1}{2}$$

$$U = 1$$

■ 圆域内的边值问题

一个半径为 a 的薄圆盘，上下两面绝热，圆周边缘的温度分布为已知函数 $F(x,y)$ ，求稳恒状态时圆盘内的温度分布。

$$\begin{cases} u_{xx} + u_{yy} = 0, & x^2 + y^2 < a^2 \\ u|_{x^2+y^2=a^2} = F(x, y) \end{cases}$$



前面已解决且有Poisson公式

$$u(r, \theta) = \frac{a^2 - r^2}{2\pi} \int_0^{2\pi} \frac{f(\varphi)}{a^2 + r^2 - 2ar \cos(\varphi - \theta)} d\varphi.$$

注： 1. 由余弦公式易得 $u(\vec{x}) = \frac{a^2 - |\vec{x}|^2}{2\pi a} \int_{|\vec{x}'|=a} \frac{u(\vec{x}')}{|\vec{x} - \vec{x}'|^2} dS(\vec{x}')$

2. 成立平均值公式 $u(\vec{0}) = \frac{1}{2\pi} \int_0^{2\pi} f(\varphi) d\varphi = \frac{1}{2\pi a^2} \int_{|\vec{x}'|=a} u(\vec{x}') dS(\vec{x}')$

■ 球域内Laplace方程的边值问题

$$\begin{cases} \Delta_3 u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} = 0, & x^2 + y^2 + z^2 < a^2 \\ u|_{x^2+y^2+z^2=a^2} = F(x, y, z) \end{cases}$$

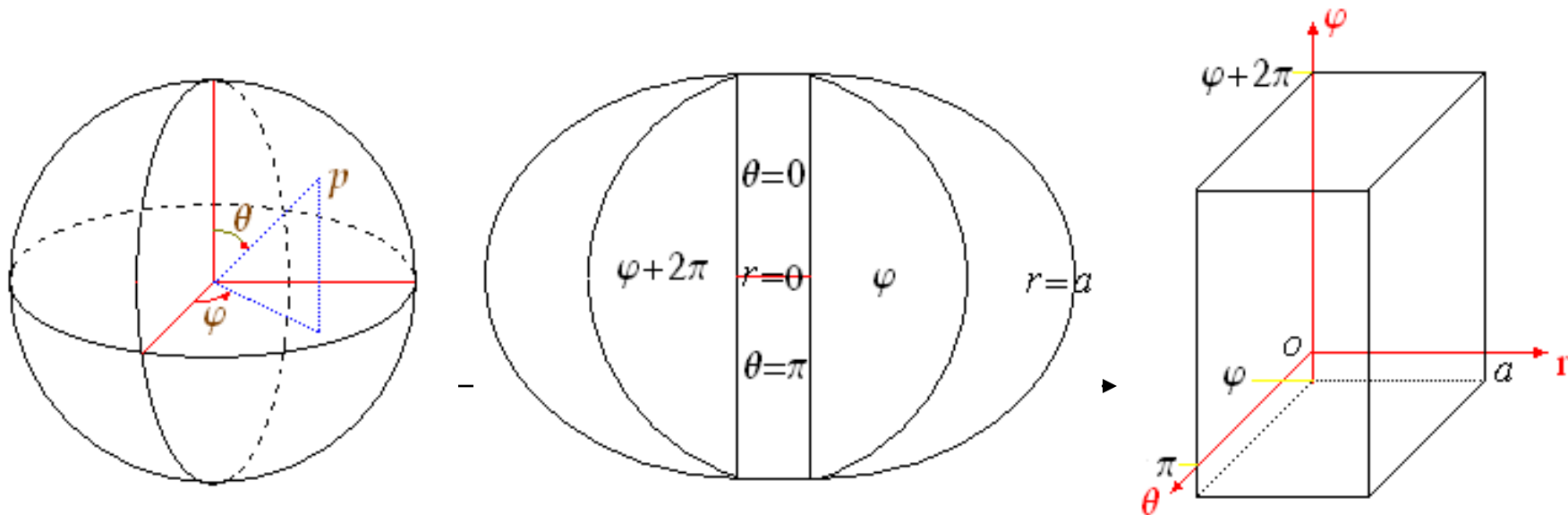
球坐标变换

$$\begin{cases} x = r \sin \theta \cos \varphi \\ y = r \sin \theta \sin \varphi \\ z = r \cos \theta \end{cases}$$

$$0 \leq r \leq a$$

$$0 \leq \theta \leq \pi$$

$$0 \leq \varphi \leq 2\pi$$



$$\left\{ \begin{array}{l} \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial u}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial u}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 u}{\partial \varphi^2} = 0, \\ u|_{r=a} = F(a \sin \theta \cos \varphi, a \sin \theta \sin \varphi, a \cos \theta) := f(\theta, \varphi) \\ u|_{r=0} = \text{有限值}, \\ u(r, \theta, \varphi) = u(r, \theta, \varphi + 2\pi), \\ u|_{\theta=0, \pi} = \text{有限值}, \end{array} \right.$$

隐含着的周期边值条件和球内约束条件

第一步： 求满足方程、周期边界条件和球内约束条件的分离解 $u(r, \theta, \varphi) = R(r)\Theta(\theta)\Phi(\varphi)$.

将分离解代入方程并对自变量逐层分离，有

$$\frac{(r^2 R')'}{R} + \frac{(\sin \theta \Theta')'}{\sin \theta \Theta} + \frac{1}{\sin^2 \theta} \frac{\Phi''}{\Phi} = 0,$$

$$R(r): \begin{cases} (r^2 R')' - \lambda R := (r^2 R')' - l(l+1)R = 0 \\ R(0) = \text{有限值} \end{cases}$$

$$\Phi(\varphi): \begin{cases} \Phi'' + \mu \Phi := \Phi'' + m^2 \Phi = 0 \\ \Phi(\varphi) = \Phi(\varphi + 2\pi) \end{cases}$$

$$\Theta(\theta): \begin{cases} \sin \theta (\sin \theta \Theta')' + [l(l+1) \sin^2 \theta - m^2] \Theta = 0 \\ \Theta|_{\theta=0, \pi} = \text{有限值} \end{cases}$$

第二步：求 $R(r), \Phi(\varphi), \Theta(\theta)$ 的具体表示式.

$$R(r): \begin{cases} (r^2 R')' - l(l+1)R = 0 \\ R(0) = \text{有限值} \end{cases}$$

欧拉方程

$$R(r) = A_l r^l + B_l \frac{1}{r^{l+1}} \Rightarrow R(r) = A_l r^l$$

$$\Phi(\varphi): \begin{cases} \Phi'' + m^2 \Phi = 0 \\ \Phi(\varphi) = \Phi(\varphi + 2\pi) \end{cases} \Rightarrow$$

$$\Phi(\varphi) = C_m \cos m\varphi + D_m \sin m\varphi, \quad m \geq 0$$

$$\Theta(\theta): \begin{cases} \sin \theta (\sin \theta \Theta')' + [l(l+1) \sin^2 \theta - m^2] \Theta = 0 \\ \Theta|_{\theta=0,\pi} = \text{有限值} \end{cases}$$

作变量变换 $x = \cos \theta$, 令 $y(x) = \Theta(\theta) = \Theta(\arccos x)$, 则易得

$$\begin{cases} (1-x^2)y'' - 2xy' + [l(l+1) - \frac{m^2}{1-x^2}]y = 0 \\ y|_{x=\pm 1} = \text{有限值} \end{cases}$$

伴随勒让德方程

$$\Theta(\theta) = y(\cos \theta) = P_l^m(\cos \theta), \quad l \geq 0, 0 \leq m \leq l$$

(详细推导可以参考季孝达“数学物理方程”第二版P104-P107)

第三步：利用边界条件求解

$$u(r, \theta, \varphi) = \sum_{l=0}^{+\infty} \sum_{m=0}^l r^l (C_{lm} \cos m\varphi + D_{lm} \sin m\varphi) P_l^m(\cos \theta)$$

$$C_{lm} = \frac{(2l+1)(l-m)!}{2\pi\delta_m a^l (l+m)!} \int_0^{2\pi} \int_0^\pi f(\theta, \varphi) P_l^m(\cos \theta) \cos m\varphi \sin \theta d\theta d\varphi$$

$$D_{lm} = \frac{(2l+1)(l-m)!}{2\pi a^l (l+m)!} \int_0^{2\pi} \int_0^\pi f(\theta, \varphi) P_l^m(\cos \theta) \sin m\varphi \sin \theta d\theta d\varphi$$

$$\delta_m = \begin{cases} 2, & m = 0 \\ 1, & m \neq 0 \end{cases}$$

例 半径为 a 的球内部没有电荷，球面上的电势为 $\sin^2 \theta \cos \varphi \sin \varphi$ ，求球形区域内部的电势分布.

解:
$$\begin{cases} \Delta u = 0, & r < a \\ u|_{r=a} = \sin^2 \theta \cos \varphi \sin \varphi \end{cases}$$
$$u(a, \theta, \varphi) = \sum_{l=0}^{+\infty} \sum_{m=0}^l a^l (C_{lm} \cos m\varphi + D_{lm} \sin m\varphi) P_l^m(\cos \theta)$$

$$= \sin^2 \theta \cos \varphi \sin \varphi = \frac{1}{6} 3 \sin^2 \theta \sin 2\varphi$$

$$= \frac{1}{6} P_2^2(\cos \theta) \sin 2\varphi \quad (\because P_2^2(x) = 3(1-x^2))$$

$$\Rightarrow D_{22} = \frac{1}{6a^2}, \text{其它系数全部为0.}$$

$$u(r, \theta, \varphi) = \frac{1}{6a^2} r^2 P_2^2(\cos \theta) \sin 2\varphi$$

注：

对于其它特殊区域上的定解问题可以利用分离变量法进行求解

例如：

半球内或球外、圆柱上的Laplace方程的边值问题