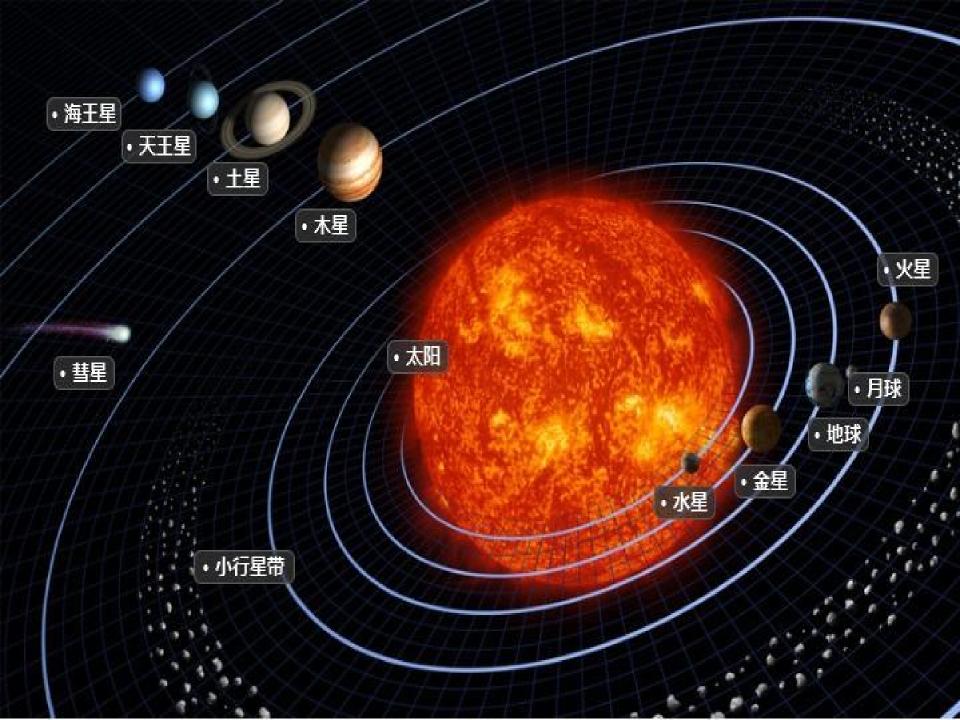
# 微为方理

# N体问题

### 内容:

- 1. N体问题的微分方程
  - ▶ 历史回顾
  - > 三维周期解演示
  - > 推荐图书
- 2. 二体问题与开普勒三大定律
- 3. N体问题的新进展



### 1. N体问题的微分方程

$$\begin{cases} m_{j} \frac{d^{2} \vec{x}_{j}}{dt^{2}} = G \sum_{k \neq j}^{N} \frac{m_{j} m_{k} (\vec{x}_{k} - \vec{x}_{j})}{|\vec{x}_{k} - \vec{x}_{j}|^{3}} \\ |\vec{x}_{j}|_{t=0} = \vec{x}_{j,0}, \quad \frac{d\vec{x}_{j}}{dt}|_{t=0} = \vec{x}_{j,1} \end{cases}$$
  $(j = 1, 2, \dots, N)$ 

其中  $m_i$ :第j个质点的质量

 $\vec{x}_i$ :第j个质点的三维空间位置向量

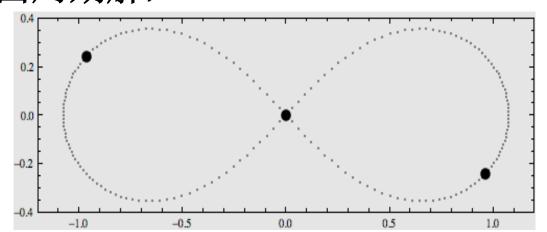
G:万有引力常数

|-|:欧氏距离

 $\vec{x}_{i,0}, \vec{x}_{i,1}$ :给定的初始条件

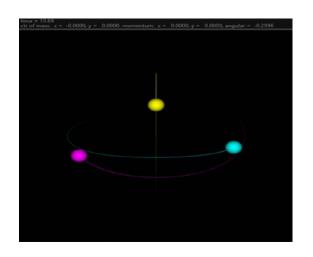
### 历史回顾

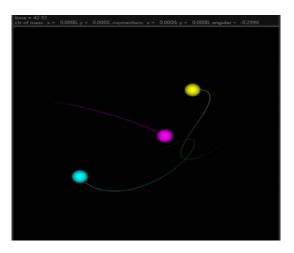
- 牛顿、1687, "自然哲学的数学原理"
- > 约翰·伯努利, 1710, 公开解决二体问题
- ▶ 布伦斯, 1887, 寻找三体问题的通解注定是无用功 (刘慈欣, 科幻小说"三体")
- ▶ 庞加莱, 1890, "关于三体问题的动态方程", 270页, Acta Mathematica (四大顶尖数学期刊之一)
- ➤ Chenciner&Montgomery, 2000, 利用变分法寻找 三体平面周期解, Annals of Mathematics, 152

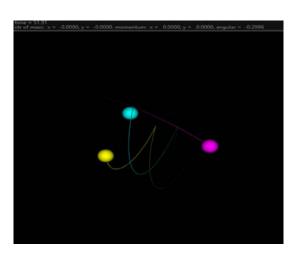


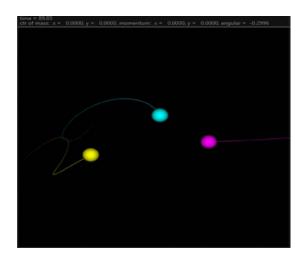
## 三维周期解演示

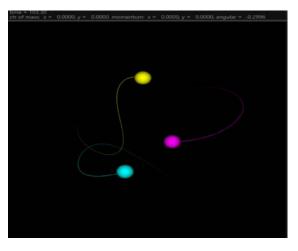
#### https://vanderbei.princeton.edu/WebGL/nBody.html

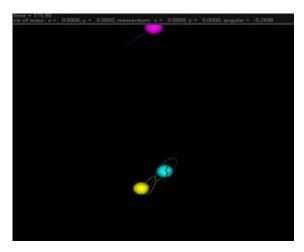






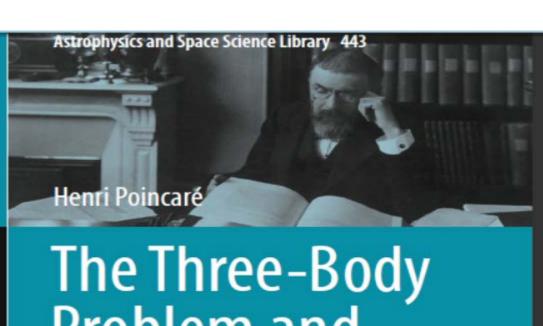






#### 推荐图书:

H. Poincaré, 1890, 法语原著的英文版



# The Three-Body Problem and the Equations of Dynamics

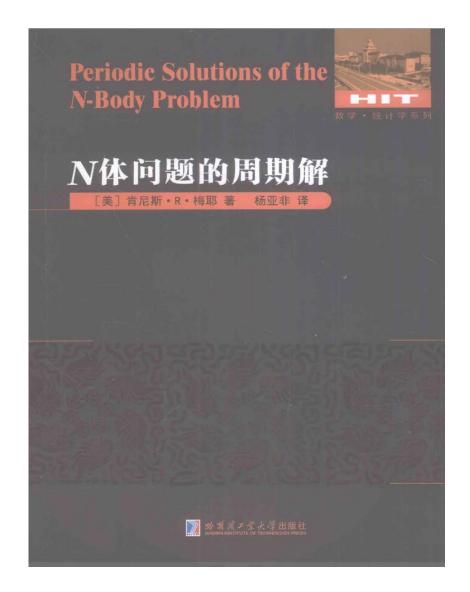
Poincaré's Foundational Work on Dynamical Systems Theory *Translated by* Bruce D. Popp

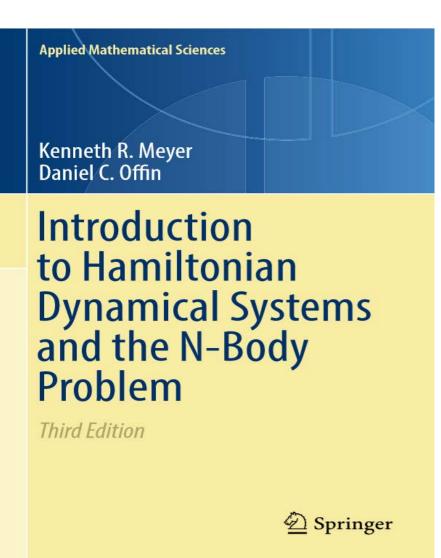




#### 推荐图书:

### K. R. Meyer (辛辛那提大学数学教授)的两本书





### 2. 二体问题与开普勒三大定律

#### 开普勒三大定律:

 $1^{st}$ .行星绕太阳的轨道为椭圆,太阳位于椭圆的一个焦点上。用极坐标表示 的椭圆轨道为

$$r = \frac{p}{1 + e\cos(\theta - \theta_0)}$$

2<sup>nd</sup>.行星向径在相等时间内扫过的面积相等

$$r^2\dot{\theta} = C$$

3rd . 行星绕太阳运动的周期平方与轨道椭圆半长径的立方成正比

$$T^2 = ka^3$$

 $T^2 = ka^3$  k对所有的行星而言是同一常数

### 二体问题的分析

对二体问题,以地日系统为例,设太阳M静止于原点,地球m的坐标向量为  $\vec{r}(t) = (x(t), y(t), z(t))$ . 由牛顿第二运动定律和万有引力定律,

$$m\frac{d^2\vec{r}(t)}{dt^2} = -\frac{GMm\vec{r}(t)}{|\vec{r}(t)|^3}$$

$$\begin{vmatrix}
\ddot{x} = -\frac{GMx}{(x^2 + y^2 + z^2)^{3/2}} \\
\Rightarrow \begin{cases}
\ddot{y} = -\frac{GMy}{(x^2 + y^2 + z^2)^{3/2}} \\
\ddot{z} = -\frac{GMz}{(x^2 + y^2 + z^2)^{3/2}}
\end{cases}
\Rightarrow
\begin{cases}
y\ddot{x} - x\ddot{y} = \frac{d}{dt}(y\dot{x} - x\dot{y}) = 0 \\
z\ddot{y} - y\ddot{z} = \frac{d}{dt}(z\dot{y} - y\dot{z}) = 0 \\
x\ddot{z} - z\ddot{x} = \frac{d}{dt}(x\dot{z} - z\dot{x}) = 0
\end{cases}$$

#### 得首次积分

$$\begin{cases} y\dot{x} - x\dot{y} = C_1 & (*) \\ z\dot{y} - y\dot{z} = A & \Rightarrow Ax + By + C_1z = 0,$$
 平面方程 
$$x\dot{z} - z\dot{x} = B \end{cases}$$

不妨设地球轨道在平面z=0上,则微分方程化为

$$\begin{cases} \ddot{x} = -\frac{\mu x}{(x^2 + y^2)^{3/2}} \\ \ddot{y} = -\frac{\mu y}{(x^2 + y^2)^{3/2}} \end{cases} (\mu = GM)$$

$$\Rightarrow \dot{x}\ddot{x} + \dot{y}\ddot{y} = -\frac{\mu(x\dot{x} + y\dot{y})}{(x^2 + y^2)^{3/2}}, \frac{d}{dt}[\dot{x}^2 + \dot{y}^2 - \frac{2\mu}{(x^2 + y^2)^{1/2}}] = 0$$

$$\Rightarrow \dot{x}^2 + \dot{y}^2 - \frac{2\mu}{(x^2 + y^2)^{1/2}} = C_2$$

引入极坐标:  $x = rcos\theta, y = rsin\theta$ , 上式化为

$$\dot{r}^2 + r^2 \dot{\theta}^2 - \frac{2\mu}{r} = C_2 \quad (**)$$

极坐标代入(\*)得开普勒第二定律:

$$-r^2\dot{\theta} = C_1 > 0$$
 (\*\*\*) (地球绕太阳顺时针转动)

再由(\*\*)(\*\*\*)得

$$\dot{r} = \pm \sqrt{C_2 + \frac{\mu^2}{C_1^2} - (\frac{C_1}{r} - \frac{\mu}{C_1})^2}$$

$$\Rightarrow \frac{dr}{d\theta} = \frac{\dot{r}}{\dot{\theta}} = \pm \frac{r^2}{C_1} \sqrt{C_2 + \frac{\mu^2}{C_1^2} - (\frac{C_1}{r} - \frac{\mu}{C_1})^2}, \quad \text{ } \Rightarrow \frac{dr}{d\theta} = \frac{\dot{r}}{\dot{\theta}} = \pm \frac{r^2}{C_1} \sqrt{C_2 + \frac{\mu^2}{C_1^2} - (\frac{C_1}{r} - \frac{\mu}{C_1})^2}, \quad \text{ } \Rightarrow \frac{dr}{d\theta} = \frac{\dot{r}}{\dot{\theta}} = \pm \frac{r^2}{C_1} \sqrt{C_2 + \frac{\mu^2}{C_1^2} - (\frac{C_1}{r} - \frac{\mu}{C_1})^2}, \quad \text{ } \Rightarrow \frac{dr}{d\theta} = \frac{\dot{r}}{\dot{\theta}} = \pm \frac{r^2}{C_1} \sqrt{C_2 + \frac{\mu^2}{C_1^2} - (\frac{C_1}{r} - \frac{\mu}{C_1})^2}, \quad \text{ } \Rightarrow \frac{dr}{d\theta} = \frac{\dot{r}}{\dot{\theta}} = \pm \frac{r^2}{C_1} \sqrt{C_2 + \frac{\mu^2}{C_1^2} - (\frac{C_1}{r} - \frac{\mu}{C_1})^2}, \quad \text{ } \Rightarrow \frac{dr}{d\theta} = \frac{\dot{r}}{\dot{\theta}} = \pm \frac{r^2}{C_1} \sqrt{C_2 + \frac{\mu^2}{C_1^2} - (\frac{C_1}{r} - \frac{\mu}{C_1})^2}, \quad \text{ } \Rightarrow \frac{dr}{d\theta} = \frac{\dot{r}}{\dot{\theta}} = \frac{\dot{r}$$

得通积分

$$\frac{\frac{C_1}{r} - \frac{\mu}{C_1}}{\sqrt{\frac{C_2 + \frac{\mu^2}{C_1^2}}{C_1^2}}} = \cos(\theta - C)$$

从而得开普勒第一定律(二次曲线中仅有椭圆不会到达无穷远):

$$r = \frac{p}{1 + e\cos(\theta - \theta_0)}, p = \frac{C_1^2}{\mu}, e = \sqrt{1 + \frac{C_2C_1^2}{\mu^2}}, \theta_0 = C$$

另外,由开普勒第二定律知单位时间向径扫过的面积为 $C_1/2$ 及椭圆性质知,

地球运行周期为

$$T = \frac{\pi ab}{C_1/2} = \frac{\pi a^2 \sqrt{1 - e^2}}{\sqrt{\mu p/2}} = \frac{2\pi a^2 \sqrt{1 - e^2}}{\sqrt{\mu a(1 - e^2)}} = \frac{2\pi}{\sqrt{\mu}} a^{3/2}$$

故得开普勒第三定律。

### 3. N体问题的新进展

> Carleo&Troyer, 2017, Science, 355, 602-606

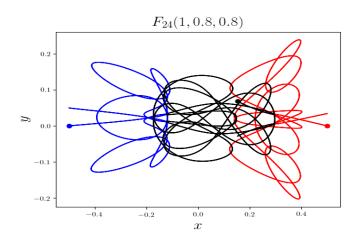
"用人工神经网络解决量子多体问题"

中文翻译: https://www.sohu.com/a/200893897\_741733

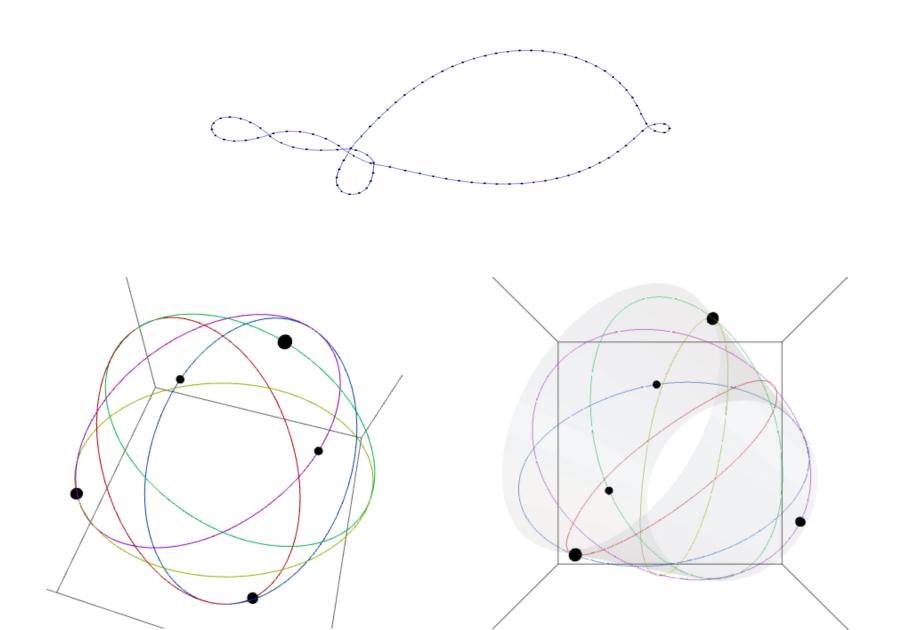
$$\Psi_{\scriptscriptstyle{M}}(\mathcal{S};\mathcal{W}) = \sum_{\{h_i\}} e^{\sum_j a_j \sigma_j^z + \sum_i b_i h_i + \sum_{ij} W_{ij} h_i \sigma_j^z}$$

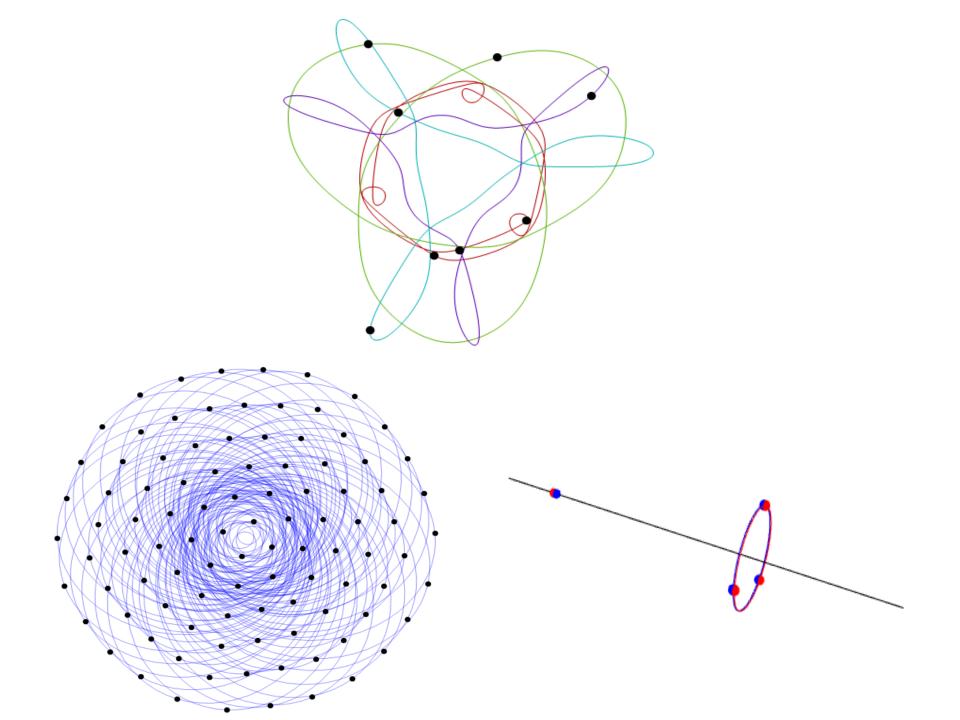
▶ 廖世俊等,2018,PASJ,70,No.4,用高精度数值 计算方法找到三体问题千种以上周期轨道

http://numericaltank.sjtu.edu.cn/three-body/three-body.htm



### >数院2017级武圣智同学的多体问题周期解算法示例:





### >最新的具有对称性的多体问题周期解:

