

Program Reference

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Overview of libcint usage

Preparing args

...

Interface

C routine

```
dim = CINTgto_cart(bas_id, bas);
```

```

dim = CINTgto_spheric(bas_id, bas);
dim = CINTgto_spinor(bas_id, bas);
fle(buf, shls, atm, natm, bas, nbas, env);
f2e(buf, shls, atm, natm, bas, nbas, env, opt);
f2e_optimizer(&opt, atm, nbas, bas, nbas, env);
CINTdel_optimizer(&opt);

```

- buf: column-major double precision array.
 - for 1e integrals of shells (i,j), data are stored as [i1j1 i2j1 ...]
 - for 2e integrals of shells (i,j|k,l), data are stored as
[i1j1k1l1 i2j1k1l1 ... i1j2k1l1 ... i1j1k2l1 ...]
 - complex data are stored as two double elements, first is real, followed by imaginary, e.g. [Re Im Re Im ...]
- shls: 0-based basis/shell indices.
 - int[2] for 1e integrals
 - int[4] for 2e integrals
- atm: int[natm*6], list of atoms. For ith atom, the 6 slots of atm[i] are
 - atm[i*6+0] nuclear charge of atom i
 - atm[i*6+1] env offset to save coordinates (env[atm[i*6+1]], env[atm[i*6+1]+1], env[atm[i*6+1]+2]) are (x,y,z)
 - atm[i*6+2] nuclear model of atom i, = 2 indicates gaussian nuclear model
 - atm[i*6+3] unused
 - atm[i*6+4] unused
 - atm[i*6+5] unused
- natm: int, number of atoms, natm has no effect **except nuclear attraction** integrals
- bas: int[nbas*8], list of basis. For ith basis, the 8 slots of bas[i] are
 - bas[i*8+0] 0-based index of corresponding atom
 - bas[i*8+1] angular momentum
 - bas[i*8+2] number of primitive GTO in basis i
 - bas[i*8+3] number of contracted GTO in basis i
 - bas[i*8+4] kappa for spinor GTO.
 - < 0 the basis $\sim j = l + 1/2$.
 - > 0 the basis $\sim j = l - 1/2$.
 - = 0 the basis includes both $j = l + 1/2$ and $j = l - 1/2$

- `bas[i*8+5]` env offset to save exponents of primitive GTOs. e.g. 10 exponents `env[bas[i*8+5]] ... env[bas[i*8+5]+9]`
- `bas[i*8+6]` env offset to save column-major contraction coefficients. e.g. 10 primitive -> 5 contraction needs a 10×5 array

```

env[bas[i*8+6] ] | env[bas[i*8+6]+10] |      | env[bas[i*8+6]+40]
env[bas[i*8+6]+1] | env[bas[i*8+6]+11] |      | env[bas[i*8+6]+41]
.                | .                  | ... | .
.                | .                  |      | .
env[bas[i*8+6]+9] | env[bas[i*8+6]+19] |      | env[bas[i*8+6]+49]

```

– ‘`bas[i*8+7]`’ unused

- `nbas`: int, number of bases, `nbas` has no effect, can be set to 0
- `env`: double[], save the value of coordinates, exponents, contraction coefficients
- struct `CINTOpt *opt`: so called “optimizer”, it needs to be initialized
`CINTOpt *opt = NULL; intname_optimizer(&opt, atm, natm, bas, nbas, env);`

every integral type has its own optimizer with the suffix `_optimizer` in its name, e.g. the optimizer for `cint2e_sph` is `cint2e_sph_optimizer`. “optimizer” is an optional argument for the integrals. It can roughly speed the integration by 10% without affecting the value of integrals. If no optimizer is wanted, set it to `NULL`.

optimizer needs to be released after using.

```
CINTdel_optimizer(&opt);
```

- if the return value equals 0, every element of the integral is 0
- short example

```

#include "cint.h"
...
CINTOpt *opt = NULL;
cint2e_sph_optimizer(&opt, atm, natm, bas, nbas, env);
for (i = 0; i < nbas; i++) {
    shls[0] = i;
    di = CINTcgto_spheric(i, bas);
    ...
    for (l = 0; l < nbas; l++) {

```

```

        shls[3] = 1;
        dl = CINTcgto_spheric(1, bas);
        buf = malloc(sizeof(double) * di * dj * dk * dl);
        cint2e_cart(buf, shls, atm, natm, bas, nbas, env, opt);
        free(buf);
    }
}
CINTdel_optimizer(&opt);

```

Fortran routine

```

dim = CINTgto_cart(bas_id, bas)
dim = CINTgto_spheric(bas_id, bas)
dim = CINTgto_spinor(bas_id, bas)
call f1e(buf, shls, atm, natm, bas, nbas, env)
call f2e(buf, shls, atm, natm, bas, nbas, env, opt)
call f2e_optimizer(opt, atm, natm, bas, nbas, env)
call CINTdel_optimizer(opt)

```

- atm and bas are 2D integer array
 - atm(1:6,i) is the (charge, offset_coord, nuclear_model, unused, unused, unused) of the ith atom
 - bas(1:8,i) is the (atom_index, angular, num_primitive_GTO, num_contract_GTO, kappa, offset_exponent, offset_coeff, unused) of the ith basis
- parameters are the same to the C function. Note that those offsets atm(2,i) bas(6,i) bas(7,i) are 0-based.
- buf is 2D/4D double precision/double complex array
- opt: an integer(8) to hold the address of so called “optimizer”, it needs to be initialized by


```
integer(8) opt call f2e_optimizer(opt, atm, natm, bas, nbas, env)
```

The optimizer can be banned by setting the “optimizer” to 0_8

```
call f2e(buf, atm, natm, bas, nbas, env, 0_8)
```

To release optimizer, execute

```
call CINTdel_optimizer(opt);
```

- short example

```

...
integer,external CINTcgto_spheric
integer(8) opt
call cint2e_sph_optimizer(opt, atm, natm, bas, nbas, env)
do i = 1, nbas
  shls(1) = i - 1
  di = CINTcgto_spheric(i-1, bas)
  ...
  do l = 1, nbas
    shls(4) = l - 1
    dl = CINTcgto_spheric(l-1, bas)
    allocate(buf(di,dj,dk,dl))
    call cint2e_sph(buf, shls, atm, natm, bas, nbas, env, opt)
    deallocate(buf)
  end do
end do
call CINTdel_optimizer(opt)

```

Supported angular momentum

$$l_{max} = 6$$

Data ordering

- for Cartesian GTO, the output data in buf are sorted as

s shell	p shell	d shell	...
...	
s	p <i>x</i>	d <i>xx</i>	
s	p <i>y</i>	d <i>xy</i>	
...	p <i>z</i>	d <i>xz</i>	
	p <i>x</i>	d <i>yy</i>	
	p <i>y</i>	d <i>yz</i>	
	p <i>z</i>	d <i>zz</i>	
	

- for real spheric GTO, the output data in buf are sorted as

s shell	p shell	d shell	f shell	...
...	
s	p x	d xy	f $y(3x^2 - y^2)$	
s	p y	d yz	f xyz	
...	p z	d z^2	f yz^2	
	p x	d xz	f z^3	
	p y	d $x^2 - y^2$	f xz^2	
	p z	...	f $z(x^2 - y^2)$	
	...		f $x(x^2 - 3y^2)$	
			...	

- for spinor GTO, the output data in buf correspond to

...	kappa=0,p shell	kappa=1,p shell	kappa=0,d shell	...
	
	$p_{1/2}(-1/2)$	$p_{1/2}(-1/2)$	$d_{3/2}(-3/2)$	
	$p_{1/2}(1/2)$	$p_{1/2}(1/2)$	$d_{3/2}(-1/2)$	
	$p_{3/2}(-3/2)$	$p_{1/2}(-1/2)$	$d_{3/2}(1/2)$	
	$p_{3/2}(-1/2)$	$p_{1/2}(1/2)$	$d_{3/2}(3/2)$	
	$p_{3/2}(1/2)$	$p_{1/2}(-1/2)$	$d_{5/2}(-5/2)$	
	$p_{3/2}(3/2)$	$p_{1/2}(1/2)$	$d_{5/2}(-3/2)$	
	$p_{1/2}(-1/2)$...	$d_{5/2}(-1/2)$	
	$p_{1/2}(1/2)$		$d_{3/2}(-3/2)$	
	$p_{3/2}(-3/2)$		$d_{3/2}(-1/2)$	
	$p_{3/2}(-1/2)$...	
	...			

Tensor

Integrals like Gradients have more than one components. The output array is ordered in Fortran-contiguous. The tensor component takes the biggest strides.

- 3-component tensor
 - X buf(:,0)
 - Y buf(:,1)
 - Z buf(:,2)
- 9-component tensor
 - XX buf(:,0)
 - XY buf(:,1)
 - XZ buf(:,2)
 - YX buf(:,3)
 - YY buf(:,4)

- YZ buf(:,5)
- ZX buf(:,6)
- ZY buf(:,7)
- ZZ buf(:,8)

Built-in function list

- Cartesian GTO integrals
 - CINTcgto_cart(int shell_id, int bas[]): Number of cartesian functions of the given shell
 - cint1e_ovlp_cart $\langle i|j \rangle$
 - cint1e_nuc_cart $\langle i|V_{nuc}|j \rangle$
 - cint1e_kin_cart $.5\langle i|\vec{p} \cdot \vec{p}|j \rangle$
 - cint1e_ia01p_cart $\langle i|\frac{\vec{r}}{r^3}|\times\vec{\nabla}|j \rangle$
 - cint1e_irixp_cart $\langle i|(\vec{r}-\vec{R}_i)\times\vec{\nabla}|j \rangle$
 - cint1e_ircxp_cart $\langle i|(\vec{r}-\vec{R}_o)\times\vec{\nabla}|j \rangle$
 - cint1e_iking_cart $0.5i\langle \vec{p} \cdot \vec{p}|U_g|j \rangle$
 - cint1e_iovlp_cart $i\langle i|U_g|j \rangle$
 - cint1e_inucg_cart $i\langle i|V_{nuc}|U_g|j \rangle$
 - cint1e_ipovlp_cart $\langle \vec{\nabla}_i|i|j \rangle$
 - cint1e_ipkin_cart $0.5\langle \vec{\nabla}_i|\vec{p} \cdot \vec{p}|j \rangle$
 - cint1e_ipnuc_cart $\langle \vec{\nabla}_i|V_{nuc}|j \rangle$

- cint1e_iprinv_cart	$\langle \vec{\nabla} i r^{-1} j \rangle$
- cint1e_rinv_cart	$\langle i r^{-1} j \rangle$
- cint2e_cart	$(ij kl)$
- cint2e_ig1_cart	$i(U_{gj} kl)$
- cint2e_ip1_cart	$(\vec{\nabla} ij kl)$

• Spheric GTO integrals

- CINTcgto_spheric(int shell_id, int bas[]):	Number of spheric functions of the given shell
- cint1e_ovlp_sph	$\langle i j \rangle$
- cint1e_nuc_sph	$\langle i V_{nuc} j \rangle$
- cint1e_kin_sph	$0.5 \langle i \vec{p} \cdot \vec{p} j \rangle$
- cint1e_ia01p_sph	$\langle i \frac{\vec{r}}{r^3} \times \vec{\nabla} j \rangle$
- cint1e_irixp_sph	$\langle i (\vec{r}_c - \vec{R}_i) \times \vec{\nabla} j \rangle$
- cint1e_ircxp_sph	$\langle i (\vec{r}_c - \vec{R}_o) \times \vec{\nabla} j \rangle$
- cint1e_iking_sph	$0.5i \langle \vec{p} \cdot \vec{p} U_{gj} \rangle$
- cint1e_iovlp_sph	$i \langle i U_{gj} \rangle$
- cint1e_inucg_sph	$i \langle i V_{nuc} U_{gj} \rangle$
- cint1e_ipovlp_sph	$\langle \vec{\nabla} i j \rangle$
- cint1e_ipkin_sph	$0.5 \langle \vec{\nabla} i \vec{p} \cdot \vec{p} j \rangle$

- `cint1e_ipnuc_sph` $\langle \vec{\nabla}_i | V_{nuc} | j \rangle$
- `cint1e_iprinv_sph` $\langle \vec{\nabla}_i | r^{-1} | j \rangle$
- `cint1e_rinv_sph` $\langle i | r^{-1} | j \rangle$
- `cint2e_sph` $(ij|kl)$
- `cint2e_ig1_sph` $i(iU_gj|kl)$
- `cint2e_ip1_sph` $(\vec{\nabla}_{ij}|kl)$
- Spinor GTO integrals
 - `CINTcgto_spinor(int shell_id, int bas[])`: Number of spinor functions of the given shell
 - `cint1e_ovlp` $\langle i | j \rangle$
 - `cint1e_nuc` $\langle i | V_{nuc} | j \rangle$
 - `cint1e_nucg` $\langle i | V_{nuc} | U_gj \rangle$
 - `cint1e_srsr` $\langle \vec{\sigma} \cdot \vec{r}_i | \vec{\sigma} \cdot \vec{r}_j \rangle$
 - `cint1e_sr` $\langle \vec{\sigma} \cdot \vec{r}_i | j \rangle$
 - `cint1e_srsp` $\langle \vec{\sigma} \cdot \vec{r}_i | \vec{\sigma} \cdot \vec{p}_j \rangle$
 - `cint1e_spsp` $\langle \vec{\sigma} \cdot \vec{p}_i | \vec{\sigma} \cdot \vec{p}_j \rangle$
 - `cint1e_sp` $\langle \vec{\sigma} \cdot \vec{p}_i | j \rangle$
 - `cint1e_spspsp` $\langle \vec{\sigma} \cdot \vec{p}_i | \vec{\sigma} \cdot \vec{p}_j \cdot \vec{p}_j \rangle$
 - `cint1e_spnuc` $\langle \vec{\sigma} \cdot \vec{p}_i | V_{nuc} | j \rangle$

– cint1e_spnucsp	$\langle \vec{\sigma} \cdot \vec{p}i V_{nuc} \vec{\sigma} \cdot \vec{p}j \rangle$
– cint1e_srnucsr	$\langle \vec{\sigma} \cdot \vec{r}i V_{nuc} \vec{\sigma} \cdot \vec{r}j \rangle$
– cint1e_sa10sa01	$0.5 \langle \vec{\sigma} \times \vec{r}_c i \vec{\sigma} \times \frac{\vec{r}}{r^3} j \rangle$
– cint1e_ovlpg	$\langle i U_g j \rangle$
– cint1e_sa10sp	$0.5 \langle \vec{r}_c \times \vec{\sigma} i \vec{\sigma} \cdot \vec{p}j \rangle$
– cint1e_sa10nucsp	$0.5 \langle \vec{r}_c \times \vec{\sigma} i V_{nuc} \vec{\sigma} \cdot \vec{p}j \rangle$
– cint1e_sa01sp	$\langle i \frac{\vec{r}}{r^3} \times \vec{\sigma} \vec{\sigma} \cdot \vec{p}j \rangle$
– cint1e_spgsp	$\langle U_g \vec{\sigma} \cdot \vec{p}i \vec{\sigma} \cdot \vec{p}j \rangle$
– cint1e_spgnucsp	$\langle U_g \vec{\sigma} \cdot \vec{p}i V_{nuc} \vec{\sigma} \cdot \vec{p}j \rangle$
– cint1e_spgsa01	$\langle U_g \vec{\sigma} \cdot \vec{p}i \frac{\vec{r}}{r^3} \times \vec{\sigma} j \rangle$
– cint1e_ipovlp	$\langle \vec{\nabla} i j \rangle$
– cint1e_ipkin	$0.5 \langle \vec{\nabla} i p \cdot p j \rangle$
– cint1e_ipnuc	$\langle \vec{\nabla} i V_{nuc} j \rangle$
– cint1e_iprinv	$\langle \vec{\nabla} i r^{-1} j \rangle$
– cint1e_ipspnucsp	$\langle \vec{\nabla} \vec{\sigma} \cdot \vec{p}i V_{nuc} \vec{\sigma} \cdot \vec{p}j \rangle$
– cint1e_ipsprinvsp	$\langle \vec{\nabla} \vec{\sigma} \cdot \vec{p}i r^{-1} \vec{\sigma} \cdot \vec{p}j \rangle$

- cint2e
 $(ij|kl)$

- cint2e_spsp1
 $(\vec{\sigma} \cdot \vec{p}i\vec{\sigma} \cdot \vec{p}j|kl)$

- cint2e_spsp1spsp2
 $(\vec{\sigma} \cdot \vec{p}i\vec{\sigma} \cdot \vec{p}j|\vec{\sigma} \cdot \vec{p}k\vec{\sigma} \cdot \vec{p}l)$

- cint2e_srsr1
 $(\vec{\sigma} \cdot \vec{r}i\vec{\sigma} \cdot \vec{r}j|kl)$

- cint2e_srsr1srsr2
 $(\vec{\sigma} \cdot \vec{r}i\vec{\sigma} \cdot \vec{r}j|\vec{\sigma} \cdot \vec{r}k\vec{\sigma} \cdot \vec{r}l)$

- cint2e_sa10sp1
 $0.5(\vec{r}_c \times \vec{\sigma}i\vec{\sigma} \cdot \vec{p}j|kl)$

- cint2e_sa10sp1spsp2
 $0.5(\vec{r}_c \times \vec{\sigma}i\vec{\sigma} \cdot \vec{p}j|\vec{\sigma} \cdot \vec{p}k\vec{\sigma} \cdot \vec{p}l)$

- cint2e_g1
 $(iU_{gj}|kl)$

- cint2e_spgsp1
 $(\vec{\sigma} \cdot \vec{p}iU_g\vec{\sigma} \cdot \vec{p}j|kl)$

- cint2e_g1spsp2
 $(iU_{gj}|\vec{\sigma} \cdot \vec{p}k\vec{\sigma} \cdot \vec{p}l)$

- cint2e_spgsp1spsp2
 $(\vec{\sigma} \cdot \vec{p}iU_g\vec{\sigma} \cdot \vec{p}j|\vec{\sigma} \cdot \vec{p}k\vec{\sigma} \cdot \vec{p}l)$

- cint2e_ip1
 $(\vec{\nabla}ij|kl)$

- cint2e_ipsp1
 $(\vec{\nabla}\vec{\sigma} \cdot \vec{p}i\vec{\sigma} \cdot \vec{p}j|kl)$

- cint2e_ip1spsp2
 $(\vec{\nabla}ij|\vec{\sigma} \cdot \vec{p}k\vec{\sigma} \cdot \vec{p}l)$

- cint2e_ipsp1spsp2
 $(\vec{\nabla}\vec{\sigma} \cdot \vec{p}i\vec{\sigma} \cdot \vec{p}j|\vec{\sigma} \cdot \vec{p}k\vec{\sigma} \cdot \vec{p}l)$

- cint2e_ssp1ssp2
 $(i\vec{\sigma}\vec{\sigma} \cdot \vec{p}j|k\vec{\sigma}\vec{\sigma} \cdot \vec{p}l)$