

# Supplementary Material for *Conformal Symplectic Optimization for Stable Reinforcement Learning*

## S.I BASIS OF CONFORMAL SYMPLECTIC OPTIMIZATION

This section introduces some basis of conformal symplectic optimization, including symplectic integrators and splitting methods.

### A. Symplectic integrators

A one-step numerical method is called symplectic if the one-step map:

$$z_1 = \phi_h(z_0)$$

is symplectic whenever the method is applied to a smooth, conservative Hamiltonian system. The following theorems show two typical symplectic integrators with respect to the conservative Hamiltonian system  $\dot{z} = S\nabla H(z) = \mathcal{C}(z)$  [29].

**Theorem 2.** *The so-called symplectic Euler methods:*

$$p_{k+1} = p_k - h\nabla_q H(q_k, p_{k+1}), \quad q_{k+1} = q_k + h\nabla_p H(q_k, p_{k+1}), \quad (17)$$

or

$$p_{k+1} = p_k - h\nabla_q H(q_{k+1}, p_k), \quad q_{k+1} = q_k + h\nabla_p H(q_{k+1}, p_k), \quad (18)$$

are symplectic methods of order 1.

**Theorem 3.** *The Verlet schemes, i.e., leapfrog methods:*

$$\begin{aligned} p_{k+\frac{1}{2}} &= p_k - \frac{h}{2}\nabla_q H\left(q_k, p_{k+\frac{1}{2}}\right), \\ q_{k+1} &= q_k + \frac{h}{2}\left(\nabla_p H\left(q_k, p_{k+\frac{1}{2}}\right) + \nabla_p H\left(q_{k+1}, p_{k+\frac{1}{2}}\right)\right), \\ p_{k+1} &= p_{k+\frac{1}{2}} - \frac{h}{2}\nabla_q H\left(q_{k+1}, p_{k+\frac{1}{2}}\right), \end{aligned} \quad (19)$$

or

$$\begin{aligned} q_{k+\frac{1}{2}} &= q_k + \frac{h}{2}\nabla_p H\left(q_{k+\frac{1}{2}}, p_k\right), \\ p_{k+1} &= p_k - \frac{h}{2}\left(\nabla_q H\left(q_{k+\frac{1}{2}}, p_k\right) + \nabla_q H\left(q_{k+\frac{1}{2}}, p_{k+1}\right)\right), \\ q_{k+1} &= q_{k+\frac{1}{2}} + \frac{h}{2}\nabla_p H\left(q_{k+\frac{1}{2}}, p_{k+1}\right), \end{aligned} \quad (20)$$

are symplectic methods of order 2.

### B. Splitting methods

**Definition 1** (Splitting methods). *Consider an arbitrary system  $\dot{z} = \zeta(z)$  in  $\mathbb{R}^{2n}$ , and suppose that the vector field is “split” as*

$$\dot{z} = \zeta^{[1]}(z) + \zeta^{[2]}(z).$$

*If the exact flow  $\varphi_t^{[1]}$  and  $\varphi_t^{[2]}$  of the system  $\dot{z} = \zeta^{[1]}(z)$  and  $\dot{z} = \zeta^{[2]}(z)$  can be calculated explicitly, we can compose their numerical maps  $\phi_h^{[1]}$  and  $\phi_h^{[2]}$  to get the numerical approximation  $\phi_h = \phi_h^{[1]} \circ \phi_h^{[2]}$  of the system  $\dot{z} = \zeta(z)$ .*

Moreover, it is proved that splitting methods in the composition  $\phi_h^{[1]} \circ \phi_h^{[2]}$  produces a first-order integrator, while outputs a second-order integrator in the composition  $\phi_{h/2}^{[2]} \circ \phi_h^{[1]} \circ \phi_{h/2}^{[2]}$  [29].

### C. Derivative of the DLPF algorithm

This section provides a detailed analysis of the derivative of the DLPF algorithm through the discretization of conformal Hamiltonian systems comprising numerous independent microscopic one-dimensional particles. If we choose the leapfrog method (20) to integrate the conservative flow  $\varphi_t^C$  and consider the composition  $\phi_h = \phi_{h/2}^D \circ \phi_h^C \circ \phi_{h/2}^D$ , we obtain

$$\begin{aligned} q_{k+\frac{1}{2}} &= q_k + \frac{h}{2} \nabla_p H \left( q_{k+\frac{1}{2}}, e^{-\frac{1}{2}rh} p_k \right), \\ p_{k+\frac{1}{2}} &= e^{-\frac{1}{2}rh} p_k - \frac{h}{2} \left( \nabla_q H \left( q_{k+\frac{1}{2}}, e^{-\frac{1}{2}rh} p_k \right) + \nabla_q H \left( q_{k+\frac{1}{2}}, p_{k+\frac{1}{2}} \right) \right), \\ q_{k+1} &= q_{k+\frac{1}{2}} + \frac{h}{2} \nabla_p H \left( q_{k+\frac{1}{2}}, p_{k+\frac{1}{2}} \right), \\ p_{k+1} &= e^{-\frac{1}{2}rh} p_{k+\frac{1}{2}}. \end{aligned} \quad (21)$$

Then, let us consider the classical Hamiltonian with each one-dimensional particle possessing the same mass  $m$ , i.e.,  $H(\theta, p) = \sum_{i=1}^n \frac{p_i^2}{2m} + J(\theta)$ , and replace it into (21). Introducing the change of variables (7), we have

$$\begin{aligned} \theta_{k+\frac{1}{2}} &= \theta_k - \frac{1}{2} \alpha \sqrt{\beta_1} v_k, \\ v_{k+\frac{1}{2}} &= \sqrt{\beta_1} v_k + (1 - \beta_1) \nabla J \left( \theta_{k+\frac{1}{2}} \right), \\ \theta_{k+1} &= \theta_{k+\frac{1}{2}} - \frac{1}{2} \alpha v_{k+\frac{1}{2}}, \\ v_{k+1} &= \sqrt{\beta_1} v_{k+\frac{1}{2}}. \end{aligned} \quad (22)$$

Note that the algorithm remains the same if we replace successive updating rules. Thus, we rewrite them into two-step updating rules and obtain

$$\begin{aligned} v_{k+\frac{1}{2}} &= \sqrt{\beta_1} \cdot \sqrt{\beta_1} v_{k-\frac{1}{2}} + (1 - \beta_1) \nabla J \left( \theta_{k+\frac{1}{2}} \right) \\ &= \beta_1 v_{k-\frac{1}{2}} + (1 - \beta_1) \nabla J \left( \theta_{k+\frac{1}{2}} \right), \\ \theta_{k+\frac{3}{2}} &= \theta_{k+1} - \frac{1}{2} \alpha \sqrt{\beta_1} v_{k+1} \\ &= \theta_{k+\frac{1}{2}} - \frac{1}{2} \alpha v_{k+\frac{1}{2}} - \frac{1}{2} \alpha \sqrt{\beta_1} \cdot \sqrt{\beta_1} v_{k+\frac{1}{2}} \\ &= \theta_{k+\frac{1}{2}} - \frac{1}{2} \alpha (\beta_1 + 1) v_{k+\frac{1}{2}}. \end{aligned} \quad (23)$$

The DLPF algorithm is a second-order conformal symplectic integrator [25], whose pseudocode is shown in Algorithm 5. It is worth noting that  $\theta_k$  and  $v_k$  in Algorithm 5 actually means  $\theta_{k+\frac{1}{2}}$  and  $v_{k-\frac{1}{2}}$  in (23), respectively. But the only distinction is that their subscripts are different, which is written on purpose for the convenience and unification of the pseudocode.

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#### Algorithm 5 Dissipative leapfrog (DLPF) algorithm [21]

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**Input:** parameters of neural network  $\theta_0$  and their conjugate momenta  $v_0$ , learning rate  $\alpha > 0$ , first-order momentum coefficient

$0 < \beta_1 < 1$

1: **for**  $k = 0$  **to**  $N - 1$  **do**

2:    $v_{k+1} = \beta_1 v_k + (1 - \beta_1) \nabla J(\theta_k)$

3:    $g_k = \frac{1}{2}(\beta_1 + 1) v_{k+1}$

4:    $\theta_{k+1} = \theta_k - \alpha g_k$

5: **end for**

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## S.II MORE DETAILS ON RAD

### A. Relativistic Hamiltonian

Here, we derive the relativistic Hamiltonian in the form of (11). In special relativity, the well-known mass-energy equation is

$$E = m\gamma_r c^2,$$

where  $E$  is the total energy of a single particle,  $m$  is the rest mass of the single particle,  $\gamma_r = \frac{1}{\sqrt{1-s^2/c^2}}$  is the relativistic coefficient,  $s$  is the speed of the single particle, and  $c$  is the speed of light. Therefore, the rest energy, in which case the  $s = 0$ , is

$$E_0 = mc^2,$$

and the momentum of a single particle is

$$p_i = m\gamma_r s.$$

Here, it is easy to obtain the following equation:

$$c^2 p_i^2 + E_0^2 = \frac{m^2 c^4}{1 - s^2/c^2} = E^2,$$

Hence, the kinetic energy of the single particle is

$$T(p_i) = E - E_0 = \sqrt{c^2 p_i^2 + E_0^2} - E_0 = c\sqrt{p_i^2 + m^2 c^2} - E_0.$$

Since the rest energy  $E_0$  is constant, having no contribution to the system's canonical equations, we can ignore it and have the Hamiltonian

$$H(q, p) = T(p) + U(q) = \sum_i T(p_i) + U(q) = \sum_i c\sqrt{p_i^2 + m^2 c^2} + U(q).$$

Finally, we obtain (11) by replacing the potential energy  $U(q)$  with the objective  $J(\theta)$ .

### B. Derivative of the second-order RAD

Consider the relativistic Hamiltonian (11) and replace it into (21), we receive the following integrator:

$$\begin{aligned}\theta_{k+\frac{1}{2}} &= \theta_k + \frac{hc}{2} \frac{e^{-\frac{1}{2}rh}}{\sqrt{e^{-rh}p_k^2 + m^2 c^2}} p_k, \\ p_{k+\frac{1}{2}} &= e^{-\frac{1}{2}rh} p_k - h \nabla J \left( \theta_{k+\frac{1}{2}} \right), \\ \theta_{k+1} &= \theta_{k+\frac{1}{2}} + \frac{hc}{2} \frac{1}{\sqrt{p_{k+\frac{1}{2}}^2 + m^2 c^2}} p_{k+\frac{1}{2}}, \\ p_{k+1} &= e^{-\frac{1}{2}rh} p_{k+\frac{1}{2}}.\end{aligned}$$

Introducing the changes of variables (7), thus we obtain

$$\begin{aligned}\theta_{k+\frac{1}{2}} &= \theta_k - \frac{\alpha/2}{\sqrt{\delta^2 v_k^2 + \frac{1}{\beta_1}}} v_k, \\ v_{k+\frac{1}{2}} &= \sqrt{\beta_1} v_k + (1 - \beta_1) \nabla J \left( \theta_{k+\frac{1}{2}} \right), \\ \theta_{k+1} &= \theta_{k+\frac{1}{2}} - \frac{\alpha/2}{\sqrt{\delta^2 v_{k+\frac{1}{2}}^2 + 1}} v_{k+\frac{1}{2}}, \\ v_{k+1} &= \sqrt{\beta_1} v_{k+\frac{1}{2}}.\end{aligned}$$

Since replacing successive updating rules does not change the algorithm, rewriting them into two-step updating rules as

$$\begin{aligned}
v_{k+\frac{1}{2}} &= \sqrt{\beta_1} \cdot \sqrt{\beta_1} v_{k-\frac{1}{2}} + (1 - \beta_1) \nabla J \left( \theta_{k+\frac{1}{2}} \right) \\
&= \beta_1 v_{k-\frac{1}{2}} + (1 - \beta_1) \nabla J \left( \theta_{k+\frac{1}{2}} \right), \\
\theta_{k+\frac{3}{2}} &= \theta_{k+1} - \frac{\alpha/2}{\sqrt{\delta^2 v_{k+1}^2 + \frac{1}{\beta_1}}} v_{k+1} \\
&= \theta_{k+\frac{1}{2}} - \frac{\alpha/2}{\sqrt{\delta^2 v_{k+\frac{1}{2}}^2 + 1}} v_{k+\frac{1}{2}} - \frac{\alpha/2}{\sqrt{\delta^2 \beta_1 v_{k+\frac{1}{2}}^2 + \frac{1}{\beta_1}}} \sqrt{\beta_1} v_{k+\frac{1}{2}} \\
&= \theta_{k+\frac{1}{2}} - \left( \frac{1}{\sqrt{\delta^2 v_{k+\frac{1}{2}}^2 + 1}} + \frac{1}{\sqrt{\delta^2 v_{k+\frac{1}{2}}^2 + \frac{1}{\beta_1^2}}} \right) \frac{\alpha}{2} v_{k+\frac{1}{2}},
\end{aligned}$$

and we immediately receive the original form of the second-order RAD related to the relativistic Hamiltonian (11).

### S.III BASIS OF OTHER OPTIMIZATION ALGORITHMS

This section introduces some basis of the other integrators involved in this paper, including SGD, NAG and ADAM.

#### A. Stochastic gradient descent algorithm

SGD is precisely a forward Euler discretization for the differential equation  $\dot{\theta} = -\nabla J(\theta)$  [13], and its pseudocode is shown in Algorithm 6.

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**Algorithm 6** Stochastic gradient descent (SGD) algorithm [13]

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**Input:** parameters of neural network  $\theta_0$ , learning rate  $\alpha > 0$

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1: for  $k = 0$  to  $N - 1$  do
2:    $g_k = \nabla J(\theta_k)$ 
3:    $\theta_{k+1} = \theta_k - \alpha g_k$ 
4: end for

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#### B. Nesterov accelerated gradient algorithm

It is worth noting that there is a close similarity between DLPF and NAG [15]. DLPF would be exactly NAG if we replace  $\theta_{k+\frac{1}{2}}$  with  $\theta_k$  in the third equation of (22) as

$$\begin{aligned}
\theta_{k+\frac{1}{2}} &= \theta_k - \frac{1}{2}\alpha\sqrt{\beta_1}v_k, \\
v_{k+\frac{1}{2}} &= \sqrt{\beta_1}v_k + (1 - \beta_1)\nabla J\left(\theta_{k+\frac{1}{2}}\right), \\
\theta_{k+1} &= \theta_k - \frac{1}{2}\alpha v_{k+\frac{1}{2}}, \\
v_{k+1} &= \sqrt{\beta_1}v_{k+\frac{1}{2}}.
\end{aligned}$$

Rewriting them into two-step updating rules like DLPF, we have

$$\begin{aligned}
v_{k+\frac{1}{2}} &= \sqrt{\beta_1} \cdot \sqrt{\beta_1}v_{k-\frac{1}{2}} + (1 - \beta_1)\nabla J\left(\theta_{k+\frac{1}{2}}\right) \\
&= \beta_1 v_{k-\frac{1}{2}} + (1 - \beta_1)\nabla J\left(\theta_{k+\frac{1}{2}}\right), \\
\theta_{k+\frac{3}{2}} &= \theta_{k+1} - \frac{1}{2}\alpha\sqrt{\beta_1}v_{k+1} \\
&= \theta_k - \frac{1}{2}\alpha v_{k+\frac{1}{2}} - \frac{1}{2}\alpha\sqrt{\beta_1} \cdot \sqrt{\beta_1}v_{k+\frac{1}{2}} \\
&= \theta_{k+\frac{1}{2}} + \frac{1}{2}\alpha\sqrt{\beta_1}v_k - \frac{1}{2}\alpha v_{k+\frac{1}{2}} - \frac{1}{2}\alpha\beta_1 v_{k+\frac{1}{2}} \\
&= \theta_{k+\frac{1}{2}} - \frac{1}{2}\alpha(1 - \beta_1)\nabla J\left(\theta_{k+\frac{1}{2}}\right) - \frac{1}{2}\alpha\beta_1 v_{k+\frac{1}{2}} \\
&= \theta_{k+\frac{1}{2}} - \frac{1}{2}\alpha\left(\beta_1 v_{k+\frac{1}{2}} + (1 - \beta_1)\nabla J\left(\theta_{k+\frac{1}{2}}\right)\right).
\end{aligned}$$

This is the famous NAG algorithm, whose pseudocode is shown in Algorithm 7. Note that we adjust the subscripts for the convenience and unification of the pseudocode as the same in Algorithm 5. Intuitively, pulling the updating base backwards, i.e., replacing  $\theta_{k+\frac{1}{2}}$  with  $\theta_k$  in the third updating rule of (22), introduces unreal dissipation into the original dynamical system [25]. Therefore, although NAG is a first-order integrator of the classical Hamiltonian system, it is not conformal symplectic. Indeed, while DLPF exactly preserves the same dissipation of the continuous-time system, NAG introduces some extra contraction or expansion to the symplectic form, thus changing the behaviour of the original system slightly.

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**Algorithm 7** Nesterov accelerated gradient (NAG) algorithm [15]

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**Input:** parameters of neural network  $\theta_0$ , learning rate  $\alpha > 0$

```

1: for  $k = 0$  to  $N - 1$  do
2:    $v_{k+1} = \beta_1 v_k + (1 - \beta_1)\nabla J(\theta_k)$ 
3:    $g_k = \frac{1}{2}(\beta_1 v_{k+1} + (1 - \beta_1)\nabla J(\theta_k))$ 
4:    $\theta_{k+1} = \theta_k - \alpha g_k$ 
5: end for

```

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### C. Adaptive moment gradient algorithm

ADAM is hailed as the most promising algorithm for stochastic nonconvex optimization [30]. It has been utilized in many DL problems and has proven experimentally effective. This algorithm estimates the gradients' first-order and secondary raw moments, then individually adjusts the effective learning rates of different parameters. Its pseudocode from the original paper is shown in Algorithm 8, wherein every operation on vector is element-wise [17].

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**Algorithm 8** Adaptive moment estimation-original (ADAM) algorithm (original) [30]

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**Input:** parameters of neural network  $\theta_0$ , first-order momenta  $v_0$ , second-order momenta  $y_0$ , learning rate  $\alpha > 0$ , first-order momentum coefficient  $0 < \beta_1 < 1$ , second-order momentum coefficient  $0 < \beta_2 < 1$ , rational factor  $\hat{\epsilon} > 0$

```

1: for  $k = 0$  to  $N - 1$  do
2:    $v_{k+1} = \beta_1 v_k + (1 - \beta_1) \nabla J(\theta_k)$  (Update biased first-order moment estimate)
3:    $y_{k+1} = \beta_2 y_k + (1 - \beta_2) (\nabla J(\theta_k))^2$  (Update biased secondary raw moment estimate)
4:    $\hat{v}_{k+1} = v_{k+1} / (1 - \beta_1^{k+1})$  (Compute bias-corrected first-order moment estimate)
5:    $\hat{y}_{k+1} = y_{k+1} / (1 - \beta_2^{k+1})$  (Compute bias-corrected secondary raw moment estimate)
6:    $g_k = \hat{v}_{k+1}$  (Estimate gradients of the objective function)
7:    $\alpha_k = \frac{\alpha}{\sqrt{\hat{y}_{k+1} + \hat{\epsilon}}}$  (Adjust effective learning rates)
8:    $\theta_{k+1} = \theta_k - \alpha_k g_k$ 
9: end for

```

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We can compactly write the updating rules by integrating the bias-correction steps into the following steps, thus formulating ADAM in the same format as other algorithms shown in this paper (see Algorithm 4). Note that the rational factor  $\epsilon$  prevents the denominator of the effective learning rate  $\alpha_k$  from zero. ADAM has two essential properties, including its moment estimation of gradients and naturally performing self-adaption on learning rates.

## S.IV EXPERIMENTAL SETTINGS

The detailed experimental settings are shown in Table IV and Table V.

TABLE IV: Experimental settings I

Task	CartPole-v1	Hopper-v3	Other MuJoCo tasks
RL algorithm	DDPG & SAC	SAC	SAC
Discount factor	0.99	0.99	0.99
Exploration noise	$\varepsilon \sim \mathcal{N}(0, 0.1)$	/	/
Temperature coefficient	0.2	0.2	0.2
Approximate function	MLP	MLP	MLP
FC layer size	$256 \times 256$	$256 \times 256$	$256 \times 256$
Activation function	ReLU	ReLU	ReLU
Learning rate decay	/	CosineAnnealingLR	/
Critic learning rate	$5 \times 10^{-4}$	$1 \times 10^{-3} \rightarrow 1 \times 10^{-4}$	$1 \times 10^{-3}$
Actor learning rate	$5 \times 10^{-5}$	$1 \times 10^{-3} \rightarrow 1 \times 10^{-4}$	$1 \times 10^{-3}$
Target network learning rate	$5 \times 10^{-3}$	$5 \times 10^{-3}$	$5 \times 10^{-3}$
Maximum iteration	$3 \times 10^4$	$5 \times 10^5$	$1 \times 10^6$
Batch size	1024	256	256
First-order momentum coeff.	0.9	0.9	0.9
Second-order momentum coeff.	0.999	0.999	0.999

TABLE V: Experimental settings II

Task	MuJoCo tasks	Atari games
RL algorithm	TD3	DQN
Discount factor	0.99	0.99
Exploration noise	$\varepsilon \sim \mathcal{N}(0, 0.1)$	$\varepsilon$ -greedy policy
Approximate function	MLP	CNN
Number of conv. layers	/	3
FC layer size	$256 \times 256$	512
Activation function	ReLU	ReLU
Learning rate decay	/	/
Critic learning rate	$3 \times 10^{-4}$	$1 \times 10^{-4}$
Actor learning rate	$3 \times 10^{-4}$	/
Target network learning rate	$5 \times 10^{-3}$	/
Maximum iteration	$1 \times 10^6$	$1 \times 10^6$
Batch size	32	256
First-order momentum coeff.	0.9	0.9
Second-order momentum coeff.	0.999	0.999