(described in 18/09/17)

(modified on 3/10/17)

## Non-dimensional Equations of motion are:

$$\ddot{\mathcal{X}} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -\alpha \end{pmatrix} \left( 1 - \frac{1}{A} \right) \mathcal{V}_1 + \kappa \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -\alpha \end{pmatrix} \left( 1 - \frac{1}{B} \mathcal{L} \right) \mathcal{V}_2 - \begin{pmatrix} 0 \\ \gamma \\ 0 \end{pmatrix} \tag{1}$$

Where:

$$\mathcal{X} = \begin{pmatrix} x_p \\ y_p \\ \theta_p \end{pmatrix}$$

$$A = \sqrt{\left(h_p \sin(\theta_p) + l_p \cos(\theta_p) - x_p + x_1\right)^2 + \left(-h_p \cos(\theta_p) + l_p \sin(\theta_p) - y_p + y_1\right)^2}$$

$$B = \sqrt{\left(h_p \sin(\theta_p) - l_p \cos(\theta_p) - x_p + x_2\right)^2 + \left(-h_p \cos(\theta_p) - l_p \sin(\theta_p) - y_p + y_2\right)^2}$$

$$\mathcal{V}_{1} = \begin{pmatrix} \sin(\theta_{p}) h_{p} + \cos(\theta_{p}) l_{p} + x_{1} - x_{p} \\ -\cos(\theta_{p}) h_{p} + \sin(\theta_{p}) l_{p} + y_{1} - y_{p} \\ l_{p} \left(\cos(\theta_{p}) \left(y_{1} - y_{p}\right) - \sin(\theta_{p}) \left(x_{1} - x_{p}\right)\right) + h_{p} \left(\cos(\theta_{p}) \left(x_{1} - x_{p}\right) + \sin(\theta_{p}) \left(y_{1} - y_{p}\right)\right) \end{pmatrix}$$

$$\mathcal{V}_{2} = \begin{pmatrix} \sin(\theta_{p}) h_{p} - \cos(\theta_{p}) l_{p} + x_{2} - x_{p} \\ -\cos(\theta_{p}) h_{p} - \sin(\theta_{p}) l_{p} + y_{2} - y_{p} \end{pmatrix}$$

$$h_{p} \left(\cos(\theta_{p}) \left(x_{2} - x_{p}\right) + \sin(\theta_{p}) \left(y_{2} - y_{p}\right)\right) + l_{p} \left(\sin(\theta_{p}) \left(x_{2} - x_{p}\right) + \cos(\theta_{p}) \left(y_{p} - y_{2}\right)\right) / \alpha = \frac{m_{p} L \theta_{1}^{2}}{l_{p}} = \frac{3L \theta_{1}^{2}}{\left(w_{p}^{2} + h_{p}^{2}\right)} \quad *$$

$$\kappa = \frac{k_{2}}{k_{1}}$$

$$\mathcal{L} = \frac{L \theta_{2}}{L \theta_{1}}$$

$$\gamma = \frac{g}{L \theta_{1}} \frac{m_{p}}{k_{1}}$$

(assuming Inertia of a rectangular payload)

All dimensional quantities are originally normalized by :

$$\omega_s^2 = \frac{k_1}{m_p}$$

A symmetrical case is considered and it is:

$$k_2 \rightarrow k_1$$
 ;  $L0_2 \rightarrow L0_1$ 

\*Note: typical values that are intended to be considered are :

$$k_1 \rightarrow 200$$
,  $L0_1 \rightarrow 2$ 

$$m_p \rightarrow 2$$
,  $h_p \rightarrow 0.1$ ,  $w_p \rightarrow 1$ ,

$$g \rightarrow 9.81$$

$$x_1[t] \to 0$$

$$y_1[t] \rightarrow 0$$

$$x_2[t] \rightarrow 2w_p$$

$$y_2[t] \rightarrow y_1[t]$$

Equilibrium point is:

$$\mathcal{X} = \begin{pmatrix} w_p \\ -(\frac{1}{2}\gamma + h_p + 1) \\ 0 \end{pmatrix}$$

Linearizing about that equilibrium point:

$$x_p \to \delta x + x_{p_0}$$

$$y_p \to \delta y + y_{p_0}$$

$$\theta_p \to \delta\theta + \theta_{p_0}$$

•••

After all calculations we can describe the linearized equations as:

$$M\ddot{x} + C\dot{x} + Kx = F$$

Where:

$$M = \begin{pmatrix} (\gamma + 2) & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & (\gamma + 2) \end{pmatrix}$$

$$F = \{0\}, C = \{0\}$$

$$K = \begin{pmatrix} 2\gamma & 0 & -2\gamma h_p \\ 0 & 2 & 0 \\ -2\alpha\gamma h_p & 0 & \alpha(2\gamma h_p^2 + \gamma(\gamma + 2)h_p + 2(\gamma + 2)w_p^2) \end{pmatrix}$$

The natural frequencies can be calculated from:

$$Det[K - \omega^2 M] = 0$$

For the y component we get:

$$\omega_y^2 = 2$$

Other frequencies are:

$$= \pm \frac{\sqrt{\frac{2\gamma}{\rho} + \alpha\gamma h_p + \frac{2\alpha\gamma h_p^2}{\rho} + 2\alpha w_p^2 \pm \frac{\sqrt{-4\rho(2\alpha\gamma^2 h_p + 4\alpha\gamma w_p^2) + (-2\gamma - \alpha\gamma\rho h_p - 2\alpha\gamma h_p^2 - 2\alpha\rho w_p^2)^2}}{\rho}}{\sqrt{2}}$$

It can be shown (detailed later) that, for certain parameters, we get:

$$\omega_1 \sim \omega_x < \sqrt{2}$$
 ;  $\omega_2 \sim \omega_\theta \gg \omega_y = \sqrt{2}$ 

For the case of  $\omega_{\theta} \gg \omega_{y}$ , I would like to neglect  $\theta$  by setting  $\ddot{\theta} = \dot{\theta} \cong \mathbf{0} \rightarrow \theta = \theta_{0} + \delta\theta$ , where  $\theta_{0} = \mathbf{0}$ . It will remain true for whole the rest of the problem.

Back to the starting 3D.O.F E.O.M (1) I will isolate  $\theta = \delta\theta$  using the 3<sup>rd</sup> equation, in order to get  $\theta = f(x, y)$  and use it back in 1<sup>st</sup> two equations of x,y. and by that getting 2D.O.F dynamics to investigate asymptotically.

## The procedure is like this:

From 3<sup>rd</sup> equation of (1):

$$0=(-lpha)\left(1-rac{1}{A}
ight)\mathcal{V}_{13}+\kappa(-lpha)\left(1-rac{1}{B}\mathcal{L}
ight)\mathcal{V}_{23}$$
 , where  $\kappa=\mathcal{L}=1$ 

we write it as:

$$0 = (B(A-1))\mathcal{V}_{13} + A((B-1))\mathcal{V}_{23}$$
 (2)

In order to extract  $\theta_p$  we use first order approximation :

$$\cos(\theta_p) \to 1 \; ; \; \sin(\theta_p) \to \theta_p$$
 (3)

to get:

$$V_{13} = a_{12} + c_{11}\theta_p$$
 ;  $V_{23} = b_{12} + c_{12}\theta_p$  (4)

And use Taylor series (around  $\theta_0=0$ ) for A,B :

$$A = \sqrt{a_{11}} + \frac{a_{12}}{\sqrt{a_{11}}} \theta_p$$
 ;  $B = \sqrt{b_{11}} + \frac{b_{12}}{\sqrt{b_{11}}} \theta_p$  (5)

Where:

$$\begin{pmatrix} a_{11} \\ a_{12} \\ b_{11} \\ b_{12} \\ c_{11} \\ c_{12} \end{pmatrix} = \begin{pmatrix} (-h_p + y_1 - y_p)^2 + (w_p + x_1 - x_p)^2 \\ h_p(x_1 - x_p) + w_p(y_1 - y_p) \\ (-h_p + y_2 - y_p)^2 + (-w_p + x_2 - x_p)^2 \\ h_p(x_2 - x_p) - w_p(y_2 - y_p) \\ h_p(y_1 - y_p) - w_p(x_1 - x_p) \\ h_p(y_2 - y_p) + w_p(x_2 - x_p) \end{pmatrix} = f(h_p, w_p, y_p, y_1, y_2, x_p, x_1, x_2)$$

Setting (4,5) into (2) and eliminating  $2^{nd}$  and  $3^{rd}$  order elements of  $\theta_p$ , getting the relation :

$$\theta_p = \frac{a_{12}b_{11}(\sqrt{a_{11}} - a_{11}) + a_{11}b_{12}(\sqrt{b_{11}} - b_{11})}{a_{12}^2b_{11} + a_{11}b_{12}^2 - a_{12}b_{12}(\sqrt{a_{11}} - a_{11} + \sqrt{b_{11}} - b_{11}) - b_{11}c_{11}(\sqrt{a_{11}} - a_{11}) - a_{11}c_{12}(\sqrt{b_{11}} - b_{11})}$$
(6)

 $\theta_p$  is therefore a function of  $\{h_p, w_p, x_i, y_i\}$ ; i = 1, 2, p

Next step is:

Using (3) for the 1<sup>st</sup> two equations of (1) and getting new equations for a 2D.O.F problem, to asymptotically investigate:

## Non-dimensional 2D.O.F Equations of motion are:

$$\ddot{\mathcal{X}} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \left( 1 - \frac{1}{A} \right) \mathcal{V}_1 + \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \left( 1 - \frac{1}{B} \right) \mathcal{V}_2 - \begin{pmatrix} 0 \\ \gamma \end{pmatrix} - C \dot{\mathcal{X}} + \begin{pmatrix} Fx \\ Fy \end{pmatrix}$$
 (7)

Where:

$$\mathcal{X} = {x_p \choose y_p}$$

$$A = \sqrt{(h_p \theta_p + w_p - x_p + x_1)^2 + (-h_p + \theta_p w_p - y_p + y_1)^2}$$

$$B = \sqrt{(h_p \theta_p - w_p - x_p + x_2)^2 + (-h_p - \theta_p w_p - y_p + y_2)^2}$$

$$\mathcal{V}_1 = {w_p + x_1 - x_p + h_p \theta_p \choose -h_p + y_1 - y_p + w_p \theta_p}$$

$$\mathcal{V}_2 = {-w_p + x_2 - x_p + h_p \theta_p \choose -h_p + y_2 - y_p - w_p \theta_p}$$

$$C = {c_1 & 0 \choose 0 & c_2}$$

 $F_x$ ,  $F_y$  are the aerodynamic drag forces

 $\theta_p$  is taken from (6) above as function of  $\{h_p, w_p, x_i, y_i\}$ ; i = 1, 2, p

 $F_x$ ,  $F_y$  will be considered as result of drag on 'CONEX' cargo container payload. No sling cables drag will be considered:

$$F_{xp} = \frac{1}{2} \rho \dot{x_p}^2 (2h_p) C_{F_{x_\alpha}} \theta_p \sim \rho h_p C_D \dot{x_p}^2$$

$$F_{y_p} = \frac{1}{2} \rho \dot{y_p}^2 (2w_p) C_{F_{y_\alpha}} (\pi - \theta_p) \sim \rho w_p C_D \dot{y_p}^2$$

$$C_D = 1.1$$