

Egg. Maxwell (unità SI)

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \wedge \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\vec{\nabla} \wedge \vec{B} = \mu_0 \left(\vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t} \right)$$

ϵ_0 e μ_0 sono costanti
del vuoto

Scegliamo unità di misura per $\{m, l, t, q\}$

corrente $I = \frac{dq}{dt}$, $\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$

① elettrostatica $F = k_1 \frac{q_1 q_2}{r^2} \Rightarrow [k_1 q_1 q_2] = [F r^2] = N l^2 = m l t^{-2} l^2 = m l^3 t^{-2}$
 $\Rightarrow [k_1] = m l^3 t^{-2} q^{-2}$

② campi $E = \frac{F}{q_1} = k_1 \frac{q_2}{r^2}$

③ $\frac{dF}{dl} = 2 k_2 \frac{I_1 I_2}{J}$ *forza tra due fili in cui passa corrente* $\Rightarrow [k_2 I_1 I_2] = m l t^{-2}$ $[I] = \left[\frac{dq}{dt} \right] = q t^{-1}$
 $[k_2] = m l t^{-2} q^{-2} t^2$

Dunque $\left[\frac{k_1}{k_2} \right] = \frac{m l^3 t^{-2} q^{-2}}{m l t^{-2} q^{-2} t^2} = l^2 t^{-2} = [v^2]$

$B = 2 k_2 \alpha \frac{I}{J} \Rightarrow \left[\frac{E}{B} \right] = \frac{l}{t} \frac{1}{\alpha}$

$\vec{\nabla} \wedge \vec{E} + k_3 \frac{\partial \vec{B}}{\partial t} = \vec{0} \Rightarrow [k_3] = \frac{1}{\alpha}$

Egg Maxwell (unità generiche)

$$\vec{\nabla} \cdot \vec{E} = 4\pi k_1 \rho$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

$$\vec{\nabla} \wedge \vec{E} + k_3 \frac{\partial \vec{B}}{\partial t} = \vec{0}$$

$$\vec{\nabla} \wedge \vec{B} = 4\pi k_2 \alpha \vec{J} + \frac{k_2}{k_1} \alpha \frac{\partial \vec{E}}{\partial t}$$

senza sorgenti: $\vec{J} = \vec{0} \wedge \rho = 0$

$$\begin{cases} \nabla_\perp \bar{E} + k_3 \frac{\partial \bar{B}}{\partial t} = \bar{0} \\ \nabla_\perp \bar{B} = \frac{k_2}{k_1} \alpha \frac{\partial \bar{E}}{\partial t} \end{cases}$$

$$\nabla_\perp (\nabla_\perp \bar{E}) = \nabla (\nabla \cdot \bar{E}) - \nabla^2 \bar{E} = -\nabla^2 \bar{E}$$

$$k_3 \nabla_\perp \left(\frac{\partial \bar{B}}{\partial t} \right) = k_3 \frac{\partial}{\partial t} (\nabla_\perp \bar{B}) = k_3 \frac{k_2}{k_1} \alpha \frac{\partial^2 \bar{E}}{\partial t^2}$$

$$\nabla_\perp \left(\nabla_\perp \bar{E} + k_3 \frac{\partial \bar{B}}{\partial t} \right) = \nabla_\perp \bar{0} = \bar{0} \Rightarrow \nabla^2 \bar{E} - \frac{k_2 k_3}{k_1} \alpha \frac{\partial^2 \bar{E}}{\partial t^2} = 0$$

ma sappiamo che $\nabla^2 \bar{E} - \frac{1}{c^2} \frac{\partial^2 \bar{E}}{\partial t^2} = \bar{0} \Rightarrow \frac{k_2 k_3}{k_1} \alpha = \frac{1}{c^2}$

siccome $\frac{k_1}{k_2} = c^2 \Rightarrow \frac{1}{k_3 \alpha} = 1 \Rightarrow k_3 = \frac{1}{\alpha}$

| | | | | |
|--------|----------------------------------|------------------------|---------------------|--------------|
| SI: | $k_1 = \frac{1}{4\pi\epsilon_0}$ | $k_2 = \frac{1}{4\pi}$ | $k_3 = 1$ | $\alpha = 1$ |
| Gauss: | $k_1 = 1$ | $k_2 = \frac{1}{c^2}$ | $k_3 = \frac{1}{c}$ | $\alpha = c$ |

Trasf. galileiane sulle eqq. di Maxwell

$$\begin{cases} \bar{x}' = \bar{x} - \bar{u}t \\ t' = t \end{cases} \quad \text{trasf. galileiane} \quad \begin{matrix} \bar{v}' = \bar{v} - \bar{u} \\ \bar{a}' = \bar{a} \end{matrix} \quad \begin{matrix} \text{velocità relativa} \\ \text{tra SR} \end{matrix}$$

$$f(\bar{x}', t') = f(\bar{x}'(\bar{x}, t), t'(\bar{x}, t)) \quad \frac{\partial}{\partial t} (x^i)' + u^i t = \frac{\partial x^i}{\partial t} + u^i = (v^i)' + u^i = v^i \quad ???$$

$$\frac{\partial}{\partial t'} = \frac{\partial t}{\partial t'} \frac{\partial}{\partial t} + \frac{\partial x^i}{\partial t'} \frac{\partial}{\partial x^i} = \frac{\partial t}{\partial t'} \frac{\partial}{\partial t} + \frac{\partial t}{\partial t'} \frac{\partial x^i}{\partial t} \frac{\partial}{\partial x^i} = \frac{\partial}{\partial t} + u^i \frac{\partial}{\partial x^i} \Rightarrow \frac{\partial}{\partial t'} = \frac{\partial}{\partial t} + \bar{u} \cdot \nabla$$

$$\frac{\partial}{\partial (x^i)'} = \frac{\partial x^i}{\partial (x^i)'} \frac{\partial}{\partial x^i} + \frac{\partial t}{\partial (x^i)'} \frac{\partial}{\partial t} = \frac{\partial}{\partial x^i} \Rightarrow \nabla' = \nabla$$

volume $\begin{cases} V' = V \\ q' = q \end{cases} \Rightarrow \rho' = \rho$ principio sperimentale $\frac{\partial \rho}{\partial t} = 0$

$$\bar{J} = \rho \bar{v} \Rightarrow \bar{J}' = \rho \bar{v}' = \rho (\bar{v} - \bar{u}) = \bar{J} - \rho \bar{u}$$

$$\bar{a}' = \bar{a} \Rightarrow F' = F \quad \therefore m' = m$$

$$\vec{F} = q(\vec{E} + \vec{v} \wedge \vec{B}) \quad \text{forza di Lorentz}$$

$$\begin{cases} \vec{F}' = \vec{F} \\ q' = q \end{cases} \Rightarrow \vec{E} + \vec{v} \wedge \vec{B} = \vec{E}' + \vec{v}' \wedge \vec{B}' = \vec{E}' + (\vec{v} - \vec{u}) \wedge \vec{B}'$$

vorremmo che la trasf. dipendesse solo da \vec{u} quindi

$$\vec{E} + \vec{v} \wedge \vec{B} = \vec{E}' + (\vec{v} - \vec{u}) \wedge \vec{B}'$$

$$\vec{E} + \vec{v} \wedge (\vec{B} - \vec{B}') = \vec{E}' - \vec{u} \wedge \vec{B}'$$

$$\Rightarrow \begin{cases} \vec{B}' = \vec{B} \\ \vec{E}' = \vec{E} + \vec{u} \wedge \vec{B}' \end{cases}$$

questo non e' detto
che succeda

$$\begin{cases} \vec{\nabla}' \cdot \vec{B}' = 0 \\ \vec{\nabla}' = \vec{\nabla} \\ \vec{B}' = \vec{B} \end{cases} \Rightarrow \vec{\nabla} \cdot \vec{B} = 0 \quad \checkmark$$

$$\begin{aligned} \vec{\nabla}' \wedge \vec{E}' + \frac{\partial \vec{B}'}{\partial t'} &= \vec{\nabla} \wedge (\vec{E} + \vec{u} \wedge \vec{B}) + \left(\frac{\partial}{\partial t} + \vec{u} \cdot \vec{\nabla} \right) \vec{B} = \overbrace{\vec{\nabla} \wedge \vec{E}'}^{=0} + \frac{\partial \vec{B}}{\partial t} + \underbrace{\vec{\nabla} \wedge (\vec{u} \wedge \vec{B})}_{=0} + (\vec{u} \cdot \vec{\nabla}) \vec{B} \\ &= \underbrace{\vec{u} (\vec{\nabla} \cdot \vec{B})}_{=0} - \underbrace{\vec{B} (\vec{\nabla} \cdot \vec{u})}_{\vec{u} \text{ cost.}} + \underbrace{(\vec{B} \cdot \vec{\nabla}) \vec{u}}_{=0} - \cancel{(\vec{u} \cdot \vec{\nabla}) \vec{B}} + \cancel{(\vec{u} \cdot \vec{\nabla}) \vec{B}} = 0 \quad \checkmark \end{aligned}$$

$$\begin{aligned} \vec{\nabla}' \cdot \vec{E}' - \frac{\rho'}{\epsilon_0} &= \vec{\nabla} \cdot (\vec{E} + \vec{u} \wedge \vec{B}) - \frac{\rho}{\epsilon_0} = \overbrace{\vec{\nabla} \cdot \vec{E}}^{=0} - \frac{\rho}{\epsilon_0} + \vec{\nabla} \cdot (\vec{u} \wedge \vec{B}) = \\ &= \vec{B} \cdot \underbrace{(\vec{\nabla} \wedge \vec{u})}_{=0} - \vec{u} \cdot (\vec{\nabla} \wedge \vec{B}) = \boxed{-\vec{u} \cdot (\vec{\nabla} \wedge \vec{B})} \neq 0 \end{aligned}$$

in generale questi
non valgono

$$\begin{aligned} \vec{\nabla}' \wedge \vec{B}' - \mu_0 \vec{J}' - \frac{1}{c^2} \frac{\partial \vec{E}'}{\partial t'} &= \vec{\nabla} \wedge \vec{B} - \mu_0 (\vec{J} - \rho \vec{u}) - \frac{1}{c^2} \left(\frac{\partial}{\partial t} + \vec{u} \cdot \vec{\nabla} \right) (\vec{E} + \vec{u} \wedge \vec{B}) \\ &= \underbrace{\vec{\nabla} \wedge \vec{B} - \mu_0 \vec{J} - \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t}}_{=0} + \boxed{\mu_0 \rho \vec{u} - \frac{1}{c^2} \left(\vec{u} \wedge \frac{\partial \vec{B}}{\partial t} + (\vec{u} \cdot \vec{\nabla}) (\vec{u} \wedge \vec{B}) \right)} \neq 0 \end{aligned}$$

$$\vec{\nabla}' \cdot \vec{J}' + \frac{\partial \rho'}{\partial t'} = \text{DIMOSTRA}$$