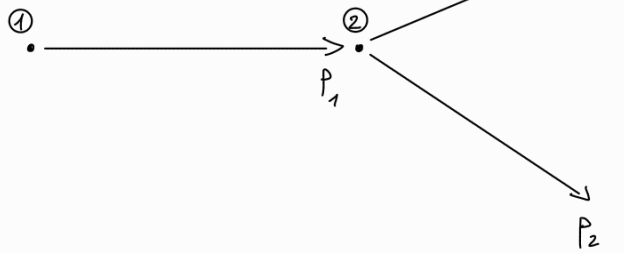


particella - nucleo = poca energia persa + grosso cambiamento di direzione
 particella - particella = viceversa

LAB



$$p_1 = (\mathcal{E}_1, 0, 0, p_1) = (\mathcal{E}_1, \vec{p}_1)$$

$$p_2 = (mc^2, 0, 0, 0) = (mc^2, \vec{0})$$

impulso trasferito

$$Q^2 = -(p_1' - p_1)^2$$

$$\stackrel{\text{CM}}{=} -((\mathcal{E}_1', p \sin \theta, 0, p \cos \theta) - (\mathcal{E}_1, 0, 0, p))^2$$

$$= -(\cancel{\mathcal{E}_1'} - \mathcal{E}_1, p \sin \theta, 0, p \cos \theta - p)^2$$

$$= -(0 - p^2 \sin^2 \theta - p^2 (\cos \theta - 1)^2)$$

$$= p^2 \sin^2 \theta + p^2 (\cos^2 \theta - 2 \cos \theta + 1)$$

$$= \underbrace{p^2 \sin^2 \theta + p^2 \cos^2 \theta}_{p^2} - 2p^2 \cos \theta + p^2$$

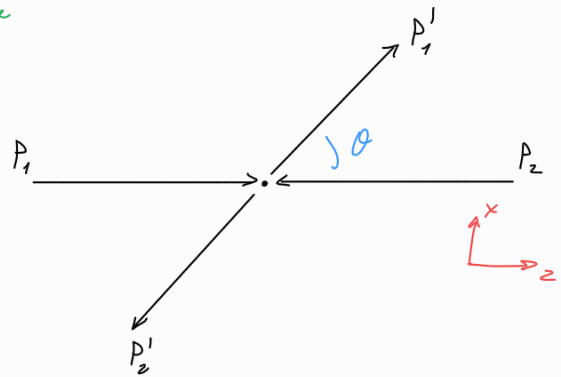
$$= 2p^2 (1 - \cos \theta)$$

$$= 4p^2 \sin^2\left(\frac{\theta}{2}\right) \quad \sin^2\left(\frac{\theta}{2}\right) = \frac{1 - \cos \theta}{2}$$

$$p_1 + p_2 = p_1' + p_2'$$

conservazione
4-impulso

CM centro di massa



$$p_1 = (\mathcal{E}_1, 0, 0, p)$$

$$p_2 = (\mathcal{E}_2, 0, 0, -p)$$

$$p_1' = (\mathcal{E}_1', p \sin \theta, 0, p \cos \theta)$$

$$p_2' = (\mathcal{E}_2', -p \sin \theta, 0, -p \cos \theta)$$

$$\mathcal{E}_1' \stackrel{\text{CM}}{=} \mathcal{E}_1 \quad \because m_1' = m_1$$

$$p^2 = (E, \vec{p})^2$$

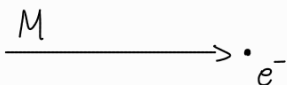
$$= E^2 - \vec{p} \cdot \vec{p}$$

$$= m^2 + |\vec{p}|^2 - |\vec{p}|^2$$

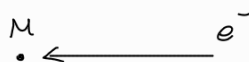
$$= m^2$$

particella grande + $e^- \Rightarrow M \gg m_e$

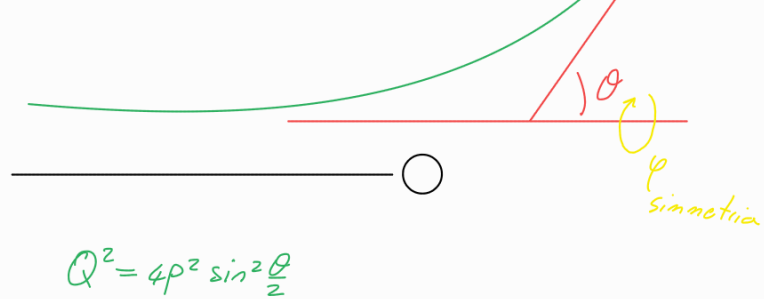
LAB



CM



$$\frac{d\sigma}{d\Omega} = \left(\frac{ze^2}{2pv}\right)^2 \sin^{-4}\left(\frac{\theta}{2}\right) \quad \text{Rutherford}$$



$$d\Omega = 2\pi \sin\theta d\theta$$

$$d\left(\sin^2\frac{\theta}{2}\right) = \underbrace{\sin\left(\frac{\theta}{2}\right)\cos\left(\frac{\theta}{2}\right)}_{\sin\theta} \frac{1}{2} d\theta = \frac{1}{2} \sin\theta d\theta \Rightarrow d\Omega = 2\pi \cdot 2 d\left(\sin^2\frac{\theta}{2}\right) = 4\pi d\left(\frac{Q^2}{4p^2}\right) = \frac{4\pi}{4p^2} d(Q^2)$$

$\leftarrow p^2 \text{ e' dato}$

$$d\Omega = \pi p^{-2} d(Q^2)$$

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma}{d(Q^2)} \frac{p^2}{\pi} = \left(\frac{ze^2}{2pv}\right)^2 \sin^{-4}\left(\frac{\theta}{2}\right) = \left(\frac{ze^2}{2pv}\right)^2 \left(\frac{4p^2}{Q^2}\right)^2 = \left(\frac{ze^2 \cancel{4p^2}}{\cancel{2p} v Q^2}\right)^2 = \left(\frac{ze^2}{\beta c Q^2}\right)^2$$

$$\Rightarrow \frac{d\sigma}{d(Q^2)} = 4\pi \left(\frac{ze^2}{\beta c Q^2}\right)^2$$

RICONTROLLA ↗

$$\begin{cases} p = (m, 0, 0, 0) & \text{iniziale elettrone} \\ p' = (E, 0, 0, p) & \text{finale elettrone} \end{cases}$$

$$\begin{aligned} Q^2 &= -(p' - p)^2 \\ &= (E - m, 0, 0, p)^2 \\ &= -((E - m)^2 - p^2) \\ &= -(E^2 + m^2 - 2Em - p^2) \\ &= -(\underbrace{E^2}_{m^2 + p^2} + m^2 - 2(\underbrace{E}_{T+m})m - p^2) \\ &= 2Tm \end{aligned}$$

energia cinetica
 $T = E - m$ persa dalla particella acquisita dall'elettrone

$$\frac{d\sigma}{d(Q^2)} = 4\pi \left(\frac{ze^2}{\beta c Q^2}\right)^2 \quad \text{con} \quad d(Q^2) = 2m dT$$

$$\frac{1}{2m} \frac{d\sigma}{dT} = 4\pi \left(\frac{ze^2}{\beta c 2mT}\right)^2 \Rightarrow \frac{d\sigma}{dT} = 2\pi \frac{1}{m} \left(\frac{ze^2}{\beta c T}\right)^2 = \frac{2\pi z^2 e^4}{mc^2 \beta^2 T^2}$$

$$T_{\min} \sim \overbrace{h\nu}^{\text{energia legame}}$$

quindi integriamo $\frac{d\sigma}{dT}$

\nwarrow riflessione totale

$$T_{\max} \rightarrow E' = \gamma mc^2, \quad p' = \gamma m \beta c \quad \swarrow \text{trasf. Lorentz}$$

$$T_{\max} = E - mc^2 = \gamma(E' + \beta c p') - mc^2$$

$$\begin{aligned}
&= \gamma^2 mc^2 + \gamma^2 m \beta^2 c^2 - mc^2 \\
&= \gamma^2 mc^2 (1 + \beta^2 - (1 - \beta^2)) \quad \gamma^{-2} = 1 - \beta^2 \\
&= \gamma^2 mc^2 (1 + \beta^2 - 1 + \beta^2) \\
&= 2\gamma^2 mc^2 \beta^2
\end{aligned}$$

Queste formule non considerano gli effetti quantistici: usando $\left. \frac{d\sigma}{dQ} \right|_{\text{Mott}}$ si ottengono delle correzioni per spin $\propto (1 - \beta^2 \sin^2 \frac{\theta}{2})$ per la sezione d'urto

$$Q^2 = 2mT = 4p^2 \sin^2 \frac{\theta}{2}$$

$$2mT_{\max} = 4p^2 \sin^2 \left(\frac{\pi}{2} \right) = 4p^2 \Rightarrow 1 - \beta^2 \sin^2 \frac{\theta}{2} = 1 - \beta^2 \frac{T}{T_{\max}}$$

A questo punto integriamo considerando la particella che entra in un materiale

$$\frac{d\sigma}{dT} = \frac{2\pi z^2 e^4}{mc^2 \beta^2 T^2} \left(1 - \beta^2 \frac{T}{T_{\max}} \right)$$

N $\frac{\text{atomi}}{\text{volume}}$ Z elettroni

stopping power:

energia persa per unità di lunghezza

$$\frac{dE}{dx} = NZ \int_{T_{\min}}^{T_{\max}} T \frac{d\sigma}{dT} dT \quad \leftarrow \text{media pesata in base all'energia di ogni urto}$$

$$\propto \int_{T_{\min}}^{T_{\max}} \frac{T}{T^2} \left(1 - \beta^2 \frac{T}{T_{\max}} \right) dT$$

$$= \left[\ln(T) - \beta^2 \frac{T}{T_{\max}} \right]_{T_{\min}}^{T_{\max}}$$

$$= \ln \left(\frac{T_{\max}}{T_{\min}} \right) - \beta^2 \left(1 - \frac{T_{\min}}{T_{\max}} \right) \quad \leftarrow 1$$

$$\approx \ln \left(\frac{T_{\max}}{T_{\min}} \right) - \beta^2$$

con costanti \Rightarrow

Formula di Bethe-Bloch

$$\frac{dE}{dx} = NZ \frac{2\pi z^2 e^4}{mc^2 \beta^2} \left(\ln \left(\frac{2\gamma^2 m \beta^2 c^2}{\hbar \langle \omega \rangle} \right) - \beta^2 \right)$$

condizioni $\left\{ \begin{array}{l} M \gg m_e \\ T_{\min} > \hbar \omega \\ \text{piccolo } \ln \text{ at } 11 \end{array} \right.$

Z nuclei materiale
 z particella incidente

piccolo bremsstrahlung

Muone in rame

$$\frac{dE}{dx}$$

