

Quantum Bernoulli factory

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1 Introduction

2 Quantum information

This section is based on the course by Noah Linden (University of Bristol) and discussion with Thomas Hebdige and David Jennings (Imperial College London). For a comprehensive introduction see for example Nielsen and Chuang (2002) or Wilde (2013).

The basis of quantum mechanics is simply linear algebra in a Hilbert space. The definitions in Sections 2.2 to 2.4 are exactly what you would find in a linear algebra course, except for the notation. The famously strange quantum behaviour arises from the “rules of quantum mechanics” (Section 2.5) pertaining to measurements. The majority of phenomena besides are described perfectly naturally by linear algebra, and are only counter-intuitive when one departs from the safety of mathematical abstraction and attempts to interpret quantum behaviour within the framework of the natural world, as perceived on a human scale. It is not surprising that such attempts should lead to confusion, since quantum mechanics is based on complex numbers and so diverges from human reality even at its very fundamentals.

2.1 Dirac notation

Dirac notation is a convenient way of denoting vectors such that it is easy to visually identify inner and outer products, and thus quickly recognise scalars, vectors and matrices:

- $|v\rangle$ denotes a column vector
- $\langle v|$ denotes a row vector
- $\langle u|v\rangle$ denotes an inner product (resulting in a scalar)
- $|u\rangle\langle v|$ denotes an outer product (resulting in a matrix)

Additionally, $\bar{\alpha}$ denotes the complex conjugate of a scalar α , and U^\dagger denotes the adjoint (conjugate transpose) of an operator U .

2.2 Hilbert space

A *Hilbert space* is a vector space with an inner product $\langle \cdot | \cdot \rangle$ satisfying the following:

- $\langle u | (\alpha|v\rangle + \beta|w\rangle) = \alpha\langle u|v\rangle + \beta\langle u|w\rangle$
- $\langle u|v\rangle = \overline{\langle v|u\rangle}$
- $\langle v|v\rangle \geq 0$ with equality if and only if $|v\rangle$ is the zero vector.

2.3 Orthonormal bases

An *orthonormal basis* of a Hilbert space \mathcal{H} is a set of vectors $\{v_1, \dots, v_n\}$ in \mathcal{H} such that:

- $\text{span}\{v_1, \dots, v_n\} = \mathcal{H}$
- $\langle v_i | v_j \rangle = \delta_{ij}$

We now restrict to the Hilbert space \mathbb{C}^2 , which is the state space of a single quantum bit (qubit), and is sufficient to describe the basics of quantum information. The *computational basis* $\{|0\rangle, |1\rangle\} = \{(1, 0)^T, (0, 1)^T\}$ is taken as the canonical basis for \mathbb{C}^2 and is henceforth used wherever not specified otherwise. Since it is an orthonormal basis, every vector $|v\rangle$ in \mathbb{C}^2 has a unique representation

$$|v\rangle = \alpha|0\rangle + \beta|1\rangle \equiv (\alpha, \beta)^T$$

for some $\alpha, \beta \in \mathbb{C}$. For reasons which will probably not become apparent in this treatment, we restrict ourselves to *normalised* vectors, requiring also $|\alpha|^2 + |\beta|^2 = 1$. We will also consider two vectors equivalent if they differ only by an overall phase, i.e. $|u\rangle \equiv |v\rangle$ if $|u\rangle = e^{i\theta}|v\rangle$ for some θ , since it is impossible to distinguish between two such vectors with any measurement.

To ensure coherency with the properties of the inner product, we have that

$$\langle v| = \bar{\alpha}\langle 0| + \bar{\beta}\langle 1|.$$

The inner product of $|v\rangle = \alpha|0\rangle + \beta|1\rangle$ with $|u\rangle = \gamma|0\rangle + \delta|1\rangle$ is therefore computed as

$$\begin{aligned} \langle v|u\rangle &= (\bar{\alpha}\langle 0| + \bar{\beta}\langle 1|)(\gamma|0\rangle + \delta|1\rangle) \\ &= \bar{\alpha}\gamma\langle 0|0\rangle + \bar{\alpha}\delta\langle 0|1\rangle + \bar{\beta}\gamma\langle 1|0\rangle + \bar{\beta}\delta\langle 1|1\rangle \\ &= \bar{\alpha}\gamma + \bar{\beta}\delta. \end{aligned}$$

One alternative choice of orthonormal basis which is worth mentioning is given by $\{|+\rangle, |-\rangle\}$, consisting of the states

$$\begin{aligned} |+\rangle &:= \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \\ |-\rangle &:= \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle). \end{aligned}$$

2.4 Linear operators

A linear operator is an operator with the property

$$A(\alpha|u\rangle + \beta|v\rangle) = \alpha A|u\rangle + \beta A|v\rangle.$$

It is therefore fully defined by its action on an orthonormal basis. For instance, the quantum NOT operator (usually denoted X) is defined by

$$\begin{aligned} X|0\rangle &= |1\rangle \\ X|1\rangle &= |0\rangle \end{aligned}$$

Equivalently, X can be expressed as a matrix with respect to the computational basis:

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

In Dirac notation, operators are written in terms of outer products of basis states:

$$X = |0\rangle\langle 1| + |1\rangle\langle 0|$$

This is equivalent to the matrix form:

$$X = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \end{pmatrix} + \begin{pmatrix} 0 \\ 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}.$$

Then the action of X on a general state $|v\rangle = \alpha|0\rangle + \beta|1\rangle$ can be calculated:

$$\begin{aligned} X|v\rangle &= (|0\rangle\langle 1| + |1\rangle\langle 0|)(\alpha|0\rangle + \beta|1\rangle) \\ &= \alpha|0\rangle\langle 1|0\rangle + \alpha|1\rangle\langle 0|0\rangle + \beta|0\rangle\langle 1|1\rangle + \beta|1\rangle\langle 0|1\rangle \\ &= \alpha|1\rangle + \beta|0\rangle. \end{aligned}$$

Definition 1. Let U be a linear operator.

- U is said to be *self-adjoint* (or *hermitian*) if $U^\dagger = U$.
- U is said to be *unitary* if $UU^\dagger = U^\dagger U = I$.

For example, the operator X is both self-adjoint and unitary. Crucially, unitary operators preserve normalisation, so they map states to states. It is also obvious that unitary transformations are always reversible (i.e. the inverse operator exists).

2.5 Rules of Quantum Mechanics

1. *States* of a quantum mechanical system correspond to normalised vectors in Hilbert space, up to an overall phase.
2. *Evolutions* of the system correspond to unitary operators.
3. *Measurements* on quantum states correspond to self-adjoint operators — see below.

2.5.1 Spectral theorem and measurement

The outcome of a measurement depends on the current state of the system and the type of measurement performed (i.e. which self-adjoint operator is applied). The spectral theorem states that every self-adjoint operator A can be represented by its spectral decomposition

$$A = \sum_i \lambda_i P_i \quad (1)$$

where $\{\lambda_1, \dots, \lambda_k\}$ is the set of *distinct* eigenvalues of A , and P_i is the *projection* operator onto the eigenspace corresponding to eigenvalue λ_i .

When we measure a state $|x\rangle$ using operator A , the measurement outcome we observe is one of the eigenvalues of A . In particular, we observe λ_i with probability $\langle x|P_i|x\rangle$. Making a measurement causes the system to collapse onto the eigenspace corresponding to the observed eigenvalue; that is, the state after measurement is proportional to $P_i|x\rangle$.

Example 1. To give a concrete example, let us consider again the operator X . This operator has eigenvalues ± 1 corresponding to eigenvectors (or *eigenstates* in quantum mechanics) $|+\rangle$ and $|-\rangle$ respectively. Therefore X admits the diagonal representation

$$X = |+\rangle\langle +| - |-\rangle\langle -|$$

which is of the form (1). Now suppose we make a measurement on the state $|v\rangle = |0\rangle$, using X . The outcome of the measurement will be an eigenvalue of X : either $+1$ or -1 . We observe the outcome $+1$ with probability

$$\langle v|P_{+1}|v\rangle = \langle 0|+\rangle\langle +|0\rangle = 1/2$$

in which case the state after measurement is

$$P_{+1}|v\rangle = |+\rangle\langle +|0\rangle \propto |+\rangle.$$

Similarly, with probability $1/2$ we observe the outcome -1 and the state after measurement is $|-\rangle$.

In this example, the two outcomes are equally likely because the state $|0\rangle$ is ‘equidistant’ from the two eigenstates of X . In general, outcomes that leave the state after measurement closer to the original state are more likely. This notion is formalised in Section 2.8.

A particularly important type of measurement is measurement in the computational basis. In this case the self-adjoint operator used is $0 \cdot |0\rangle\langle 0| + 1 \cdot |1\rangle\langle 1| = |1\rangle\langle 1|$. Measuring a state $|v\rangle = \alpha|0\rangle + \beta|1\rangle$ in the computational basis returns 0 with probability $|\alpha|^2$ or 1 with probability $|\beta|^2$. This is useful in quantum computing algorithms because it deterministically converts the basis states $|0\rangle$ and $|1\rangle$ to their classical equivalents 0 and 1, ready for classical output.

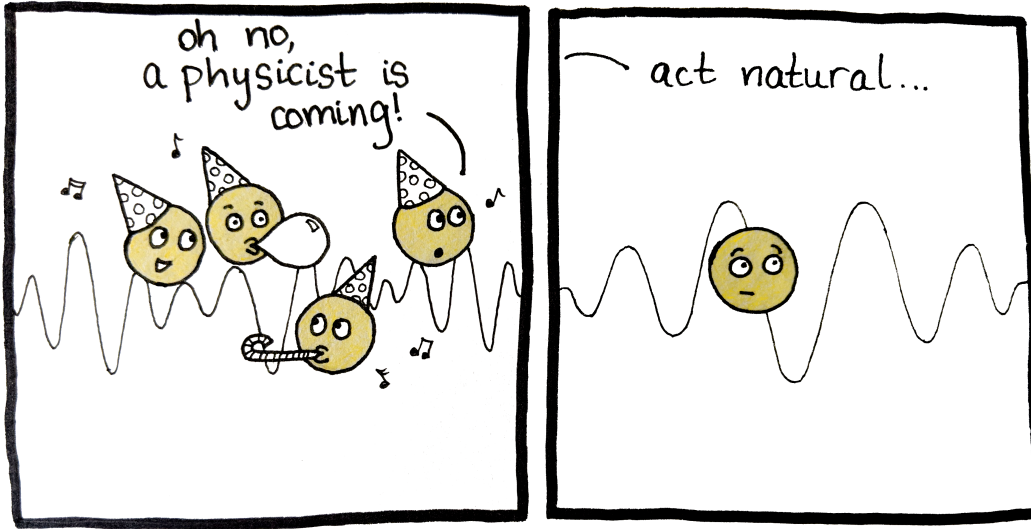


Figure 1: When a quantum system is in a superposition of states, making a measurement causes the system to collapse into the eigenstate corresponding to the observed eigenvalue.

2.6 Quantum randomness

It is worth remarking at this point on the difference between quantum and classical randomness. As statisticians we are used to dealing with randomness, but we accept that randomness is simply part of a model, and is not purported to exist in nature. We artificially introduce random variables into our models to account for a lack of information, either about the state of the system or about the (presumably deterministic) processes that govern certain phenomena.

On the contrary, our uncertainty about the outcome of a measurement on a quantum system is of a different kind. This uncertainty is not a symptom of our lack of knowledge: even when we know exactly the state of the system, as in Example 1, we are still unable to predict with certainty the outcome of a measurement on the system. The randomness here is intrinsic; the natural processes governing quantum measurement are truly stochastic.

2.7 Pure and mixed states

So far we have only encountered *pure states*, like $|0\rangle$ or $|+\rangle$. But we can also consider probabilistic mixtures of pure states, known as *mixed states*, like

$$|v\rangle = \sum_i p_i |v_i\rangle$$

where $\{p_i\}$ is a probability distribution and the $|v_i\rangle$ are pure states. You can think of this as a ‘source’ which emits a state $|v_i\rangle$ with probability p_i . A classical analogue is a bit which takes the value 1 with probability p or 0 with probability $1 - p$ (a p -coin). In this case, we can say that the expected value of the bit is p .

There is no full ordering on the set of quantum states, so it is not possible to define the expected value of a mixed state. However, we can calculate the expected value of a measurement on a mixed state, using the state’s *density operator*.

2.7.1 Density operators

The density operator of a mixed state $|v\rangle = \sum_i p_i |v_i\rangle$ is given by

$$\rho = \sum_i p_i |v_i\rangle\langle v_i|.$$

Suppose we measure this mixed state using the self-adjoint operator $A = \sum_k \lambda_k P_k$, where $\{\lambda_k\}$ are the distinct eigenvalues of A and P_k the projections onto the corresponding eigenspaces. Recall (from Section 2.5.1) that the outcome of the measurement will be an eigenvalue of A . Applying the law of total probability, the probability of

observing a particular outcome λ_j is

$$\mathbb{P}(\text{observe } \lambda_j) = \sum_i p_i \langle v_i | P_j | v_i \rangle = \text{tr}(\rho P_j)$$

where $\text{tr}(\cdot)$ denotes the trace of an operator. Note that $\text{tr}(\rho) = 1$ because the states are normalised and $\{p_i\}$ is a probability distribution. The state $\rho = I/2$ is called the *maximally mixed* state; it expresses complete ignorance, like a classical (1/2)-coin.

Different mixtures of states may have the same density operator, but in this case it is impossible to distinguish between them by any measurement.

Example 2. For example, a source $|v\rangle$ which emits $|0\rangle$ or $|1\rangle$ each with probability 1/2 has the same density operator $\rho = I/2$ as a source $|u\rangle$ which emits $|+\rangle$ or $|-\rangle$ each with probability 1/2. Suppose we measured both of these mixed states in the computational basis. The probability of observing the outcome 0 from $|u\rangle$ is

$$\langle u | P_0 | u \rangle = \frac{1}{2} \langle 0 | 0 \rangle \langle 0 | 0 \rangle + \frac{1}{2} \langle 1 | 0 \rangle \langle 0 | 1 \rangle = 1/2 \quad (2)$$

while the probability of observing 0 from $|v\rangle$ is

$$\langle v | P_0 | v \rangle = \frac{1}{2} \langle + | 0 \rangle \langle 0 | + \rangle + \frac{1}{2} \langle - | 0 \rangle \langle 0 | - \rangle = \frac{1}{2} \cdot \frac{1}{2} + \frac{1}{2} \cdot \frac{1}{2} = 1/2. \quad (3)$$

Therefore measurement in the computational basis can provide no information to distinguish between $|u\rangle$ and $|v\rangle$. And indeed nor can any other measurement — proved easily by exchanging the projection operator $|0\rangle\langle 0|$ for general P in the above — so the two sources are indistinguishable.

2.8 The Bloch sphere

Density operators have a neat geometrical representation as points in a unit sphere (Figure 2b). This is called the *Bloch sphere*, and can be thought of as the state space of a qubit.

An arbitrary pure state can be written in the form

$$|v\rangle = e^{i\psi} \left[\cos\left(\frac{\theta}{2}\right) |0\rangle + e^{i\phi} \sin\left(\frac{\theta}{2}\right) |1\rangle \right].$$

Since the overall phase $e^{i\psi}$ can be ignored, the state is parametrised by two angles θ and ϕ . Taking these as spherical coordinates with unit radius, we can think of each pure state as a point on the unit sphere (or equivalently, a unit vector). Orthogonal states lie at opposite points on the Bloch sphere.

Mixed states can also be represented in the Bloch sphere. The Bloch vector for a mixed state is found by taking the weighted average of the Bloch vectors of its component pure states. Due to the triangle inequality, mixed states therefore lie strictly inside the unit sphere. This illustrates that several different mixtures of states can have the same density operator. Quantum states are distinguishable if and only if they correspond to different points in the Bloch sphere.

There is an analogy here with classical information. The state space of a classical bit is the unit interval: its value could be 0 or 1, or we could take a probabilistic mixture of the two. The ‘pure’ states then are 0 and 1, which lie on the ‘surface’ of the interval, while ‘mixed’ states lie strictly inside the interval (Figure 2a).

Finally, unitary operators correspond to rotations of the Bloch sphere. This illustrates that unitary transformations are reversible.

2.8.1 Geometric view of measurement

Suppose we are going to measure a pure state $|v\rangle = \cos(\frac{\theta}{2}) |0\rangle + e^{i\phi} \sin(\frac{\theta}{2}) |1\rangle$ in the computational basis. There is a simple relationship between the probability of each outcome and the position of $|v\rangle$ in the Bloch sphere.

Figure 3 shows the semicircular slice of the Bloch sphere on which $|v\rangle$ lies. The value of ϕ determines which slice this is, and conditional on that the state only depends on θ — the following calculations are independent of ϕ and apply to an arbitrary pure state. The state $|v\rangle$ is projected onto the line segment between $|0\rangle$ and $|1\rangle$ as shown. Since the sphere has unit radius, the distance labelled x on Figure 3 is

$$x = 1 - \cos \theta.$$

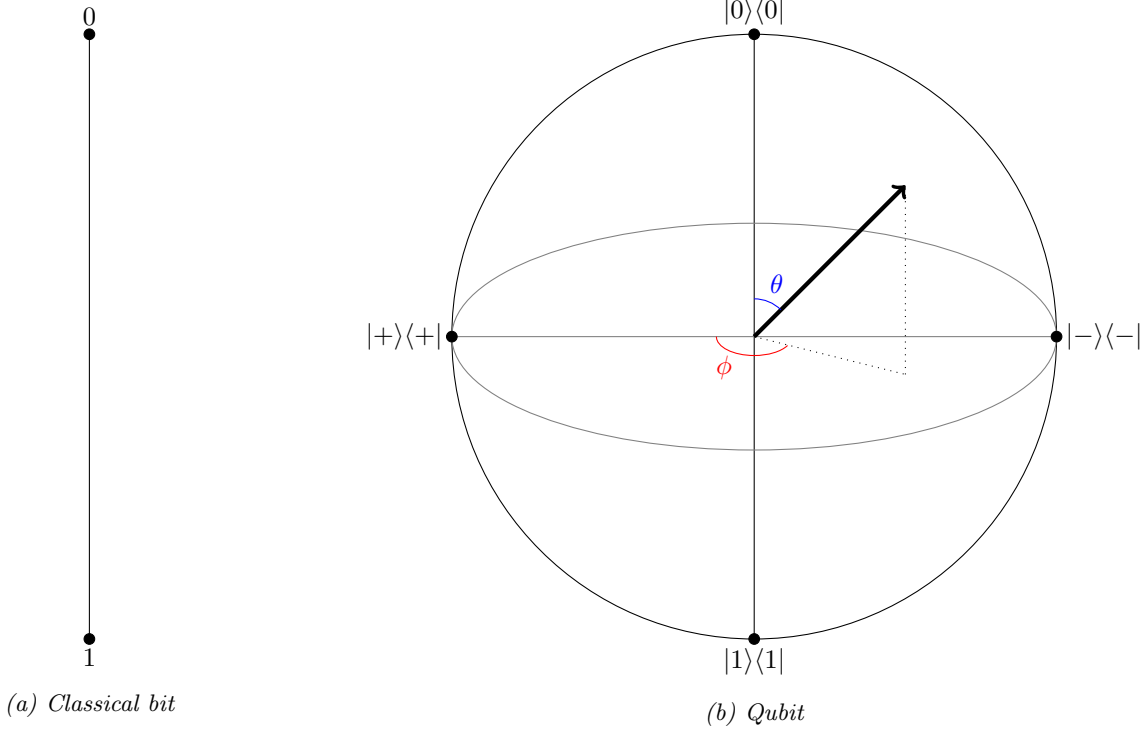


Figure 2: State space of classical bits versus qubits. The state space of a classical bit is the line segment $[0, 1]$. The ‘pure’ states 0 and 1 are at the endpoints, and probabilistic mixtures of the two lie inbetween. The state space of a qubit is the Bloch sphere. Pure states lie on the surface of the sphere, and mixed states lie in the interior. The centre of the sphere is the maximally mixed state $\rho = I/2$.

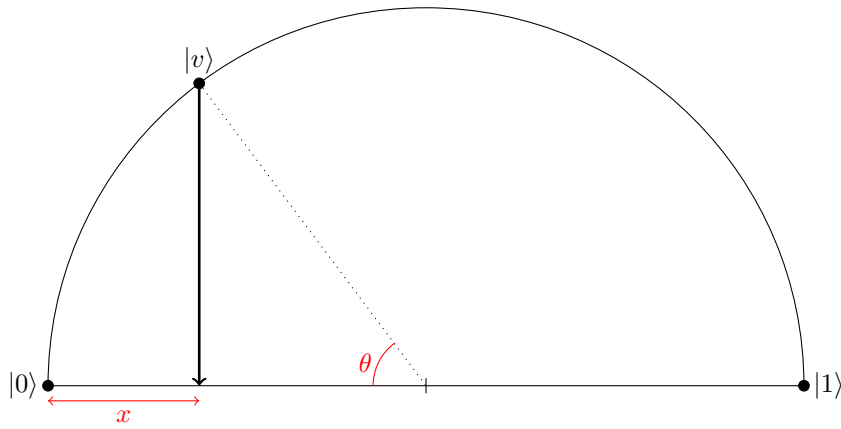


Figure 3: When measuring a state in the computational basis, the probability of observing a certain eigenvalue is related to the projected distance x from the corresponding eigenstate.

Using the expansion of $|v\rangle$ in the computational basis, we can calculate the probability of observing measurement outcome 0:

$$\mathbb{P}(\text{observe } 0) = \langle v|0\rangle \langle 0|v\rangle = \cos^2\left(\frac{\theta}{2}\right) = \frac{\cos\theta + 1}{2} = 1 - \frac{x}{2}.$$

This relationship between the projected distance and measurement probabilities also applies to mixed states. For instance, if we take any mixture of states all with the same value of θ (but different values of ϕ), the result will hold, because x and $\mathbb{P}(\text{observe } 0)$ will be the same for every component state, and hence for the mixed state too. But the result holds for general mixed states as well. Any point in the sphere whose projection onto the line segment between $|0\rangle$ and $|1\rangle$ lies a distance x from $|0\rangle$ will measure 0 with probability $1 - x/2$.

Furthermore, this illustrates a more general concept, which applies not only to measurement in the computational basis but to any quantum measurement. Any valid measurement operator on one qubit (that is non-trivial in the sense of having more than one possible outcome) has two eigenstates that lie opposite each other on the Bloch sphere. Projecting a state $|v\rangle$ onto the line segment between the two eigenstates gives the probabilities of observing the associated eigenvalues.

2.9 Entanglement

Suppose we now have two qubits and would like to describe their joint state. The joint state space will be $\mathbb{C}^2 \otimes \mathbb{C}^2$, where \otimes denotes the tensor product. This state space has canonical basis

$$\{|0\rangle \otimes |0\rangle, |0\rangle \otimes |1\rangle, |1\rangle \otimes |0\rangle, |1\rangle \otimes |1\rangle\}.$$

We also use the shorthand notations $|0\rangle \otimes |0\rangle = |0\rangle|0\rangle = |00\rangle$. The inner product on $\mathbb{C}^2 \otimes \mathbb{C}^2$ of $|u_1\rangle|u_2\rangle$ with $|v_1\rangle|v_2\rangle$ is defined as $\langle u_1|v_1\rangle \langle u_2|v_2\rangle$.

The joint state of two qubits $|u\rangle = \alpha|0\rangle + \beta|1\rangle$ and $|v\rangle = \gamma|0\rangle + \delta|1\rangle$ can be expanded in the computational basis:

$$|u\rangle \otimes |v\rangle = \alpha\gamma|00\rangle + \alpha\delta|01\rangle + \beta\gamma|10\rangle + \beta\delta|11\rangle.$$

States of this form are called *product states* as they can be written as the tensor product of two states in \mathbb{C}^2 . However, the set of all product states is only a subset of $\mathbb{C}^2 \otimes \mathbb{C}^2$. In general, a state $a|00\rangle + b|01\rangle + c|10\rangle + d|11\rangle$ may not be expressible as a product of two states in \mathbb{C}^2 . In this case, it is called an *entangled* state.

For example, consider the state $(|00\rangle + |11\rangle)/\sqrt{2}$. This state is entangled: if we try to write it in the form $|u\rangle \otimes |v\rangle$ as above, we find that either α or δ must be zero, in which case the coefficient of either $|00\rangle$ or $|11\rangle$ must be zero, producing a contradiction.

The set of four entangled states

$$\{|\phi^+\rangle, |\phi^-\rangle, |\psi^+\rangle, |\psi^-\rangle\} := \left\{ \frac{|00\rangle + |11\rangle}{\sqrt{2}}, \frac{|00\rangle - |11\rangle}{\sqrt{2}}, \frac{|01\rangle + |10\rangle}{\sqrt{2}}, \frac{|01\rangle - |10\rangle}{\sqrt{2}} \right\}$$

is another important basis for $\mathbb{C}^2 \otimes \mathbb{C}^2$, called the *Bell basis*.

3 Bernoulli factories

Roughly speaking, a Bernoulli factory is an algorithm that takes samples a Bernoulli(p) random variable as inputs and uses them to produce a Bernoulli($f(p)$) sample as output, where f is a known function but the parameter p is unknown (therefore the algorithm must not depend on p). The imaginary device for generating Bernoulli(p) random variables is sometimes referred to as a p -coin or black box; samples thereof may be called flips/tosses or queries.

Definition 2 (Bernoulli factory). Let $S \subseteq [0, 1]$ and $f : S \rightarrow [0, 1]$, and suppose we have access to a sequence of Bernoulli random variables $X_1, X_2, \dots \stackrel{iid}{\sim} \text{Bernoulli}(p)$ with unknown parameter $p \in S$. Let $U \in \mathcal{U}$ denote an auxiliary random variable with known distribution, independent of p . Let τ be a stopping time with respect to the natural filtration such that $\mathbb{P}(\tau < \infty) = 1$.

A *Bernoulli factory* is a computable function $\mathcal{A}(U, X_1, \dots, X_\tau)$ with the property that $Y := \mathcal{A}(U, X_1, \dots, X_\tau) \sim \text{Bernoulli}(f(p))$ for all $p \in S$.

The stopping time τ , which may depend on the observed values of U and X_1, X_2, \dots , is the first time at which the value of \mathcal{A} can be computed, and is called the *running time* of the Bernoulli factory.

Definition 3 (simulable). Following the terminology of Keane and O’Brien (1994), a function $f : S \rightarrow [0, 1]$ is said to be *simulable* if there is a Bernoulli factory that simulates $f(p)$ for all $p \in S$. Furthermore, it is *strongly simulable* if there is a Bernoulli factory that simulates $f(p)$ without the need for an auxiliary random variable U .

The first formal treatment of a Bernoulli factory problem was presented by Von Neumann (1951), who described a Bernoulli factory for the function $f(p) \equiv 1/2$:

To cite a human example, for simplicity, in tossing a coin it is probably easier to make two consecutive tosses independent than to toss heads with probability exactly one-half. If independence of successive tosses is assumed, we can reconstruct a 50-50 chance out of even a badly biased coin by tossing twice. If we get heads-heads or tails-tails, we reject the tosses and try again. If we get head-tails (or tails-heads), we accept the result as heads (or tails). The resulting process is rigorously unbiased, although the amended process is at most 25 percent as efficient as ordinary coin-tossing.

Most authors now use the slightly more convenient convention with the outcomes inverted, so that once 10 or 01 is observed, the output of the Bernoulli factory is equal to the last coin flip (0 or 1 respectively). It is clear that, as claimed, this procedure produces a $(1/2)$ -coin, because 01 and 10 appear with the same probability, so conditional on observing one or the other, both outcomes have probability $1/2$. The procedure is formalised in Example 3.

Example 3. $f(p) \equiv 1/2$.

This constant function is strongly simulable, using the procedure:

$$\tau = \min\{t \in \{2, 4, \dots\} : X_{t-1} \neq X_t\}$$

$$\mathcal{A}(X_1, X_2, \dots) = X_\tau.$$

The running time of this algorithm is random, and indeed unbounded (but it is finite almost surely, making it a valid Bernoulli factory). In particular, $\tau \stackrel{d}{=} 2 \times \text{Geom}(2p(1-p))$, with expectation $\mathbb{E}(\tau) = \frac{1}{p(1-p)}$. Figure 4 shows the expected running time for various values of p . The expected running time achieves its minimum value of four when $p = 1/2$ (that is, the coin was already unbiased), hence von Neumann’s claim that “the amended process is at most 25 percent as efficient as ordinary coin-tossing” (which has deterministic running time one). The expected running time approaches infinity when p is close to 0 or 1, but is moderate across a large range non-extreme values.

The algorithm given in Example 3 can be extended to simulate other constant functions of the form $f(p) \equiv 1/k$ for $k \in \{2, 3, \dots\}$. The generalised procedure is:

$$\tau = \min \left\{ t \in \{k, 2k, 3k, \dots\} : \sum_{i=t-k+1}^t X_i = 1 \right\}$$

$$\mathcal{A}(X_1, X_2, \dots) = X_\tau.$$

However, the running time increases badly with k , even under certain early stopping strategies. Of course, any constant $f(p) \equiv c$ is simulable by employing an auxiliary random variable $U \sim \text{Bernoulli}(c)$, but von Neumann’s example shows that $c = 1/2$ is also strongly simulable.

Another important example is the case of $f(p) = p^2$, which is strongly simulable using exactly two queries, returning heads if and only if both inputs are heads. Example 4 presents a generalisation of this procedure.

Example 4. $f(p) = p^k, p \in [0, 1], k \in \{1, 2, \dots\}$.

This problem is solved easily by flipping the p -coin k times and outputting heads if and only if all k flips return heads. No auxiliary random variable is required.

$$\tau = k$$

$$\mathcal{A}(X_1, \dots, X_T) = \mathbb{I}\{X_1 = 1, X_2 = 1, \dots, X_\tau = 1\}$$

In this case the running time is deterministically equal to k . It is clear that this procedure produces heads with probability p^k , since the flips are independent and each gives heads with probability p .

The running time of this algorithm can be decreased by using an early stopping strategy, whereby we stop the process as soon as a tails input appears (in which case we know immediately that the output will be tails). In this case the running time is random, but still bounded above by k :

$$\tau = k \wedge \min\{t \in \{1, 2, \dots\} : X_t = 0\} \stackrel{d}{=} \min\{k, \text{Geom}(1-p)\}.$$

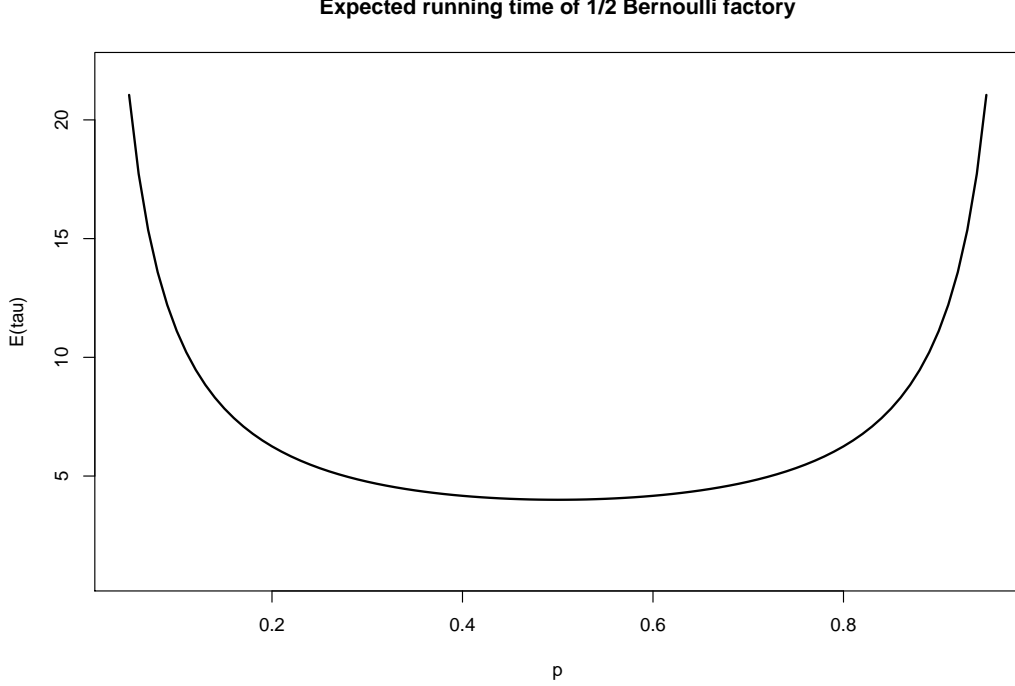


Figure 4: Expected running time $\mathbb{E}(\tau)$ of von Neumann's Bernoulli factory for $f(p) \equiv 1/2$ (Von Neumann, 1951). The expectation attains its minimum $\mathbb{E}(\tau) = 4$ at $p = 1/2$, and is moderate over a large range of p , but grows unboundedly as p approaches 0 and 1.

Figure 5 shows the expected running time of this procedure for a few values of k . We see that early stopping can significantly reduce the expected running time when p is not close to 1, and that the gain in efficiency is greater for larger k . In the case $k = 1$, where f is the identity, the number of coin flips required is deterministically equal to one, for any p .

Example 5. $f(p) = 2p(1 - p)$.

This is the probability of obtaining either 01 or 10 with two coin flips. Therefore the Bernoulli factory is simply:

$$\tau = 2$$

$$\mathcal{A}(X_1, X_2) = \mathbb{I}\{(X_1, X_2) \in \{(0, 1), (1, 0)\}\}$$

Of course, many other functions can be similarly obtained by varying the number of coin flips and the set of outcomes resulting in heads. One example, mentioned by Patel et al. (2018), is the function $f(p) = 3p(1 - p)$. Noticing that $3p(1 - p) = 3p^2(1 - p) + 3p(1 - p)^2$, this function is the probability of obtaining any outcome other than 000 or 111 with three coin flips.

Examples 3, 4 and 5 cast some light on the sort of manipulations that are possible when constructing Bernoulli factories. Lemma 1 provides some general properties, which are trivial to prove.

Lemma 1. Suppose we have a Bernoulli factory \mathcal{A}_f for the function $f : S_f \rightarrow [0, 1]$ and a Bernoulli factory \mathcal{A}_g for $g : S_g \rightarrow R_g \subseteq [0, 1]$. Then

- (a) There exists a Bernoulli factory for $1 - f : S_f \rightarrow [0, 1]$, given by inverting the output of \mathcal{A}_f .
- (b) For any $\alpha \in [0, 1]$ there exists a Bernoulli factory for $\alpha f + (1 - \alpha)g$, given by first sampling from $\text{Bernoulli}(\alpha)$ and, depending on the outcome, using either f or g .
- (c) If $R_g \subseteq S_f$, then there exists a Bernoulli factory for the composition $f \circ g : S_g \rightarrow [0, 1]$, given by the composition $\mathcal{A}_f \circ \mathcal{A}_g$.
- (d) There exists a Bernoulli factory for the product $fg : S_f \cap S_g \rightarrow [0, 1]$, given by $\mathcal{A}_f \wedge \mathcal{A}_g$.

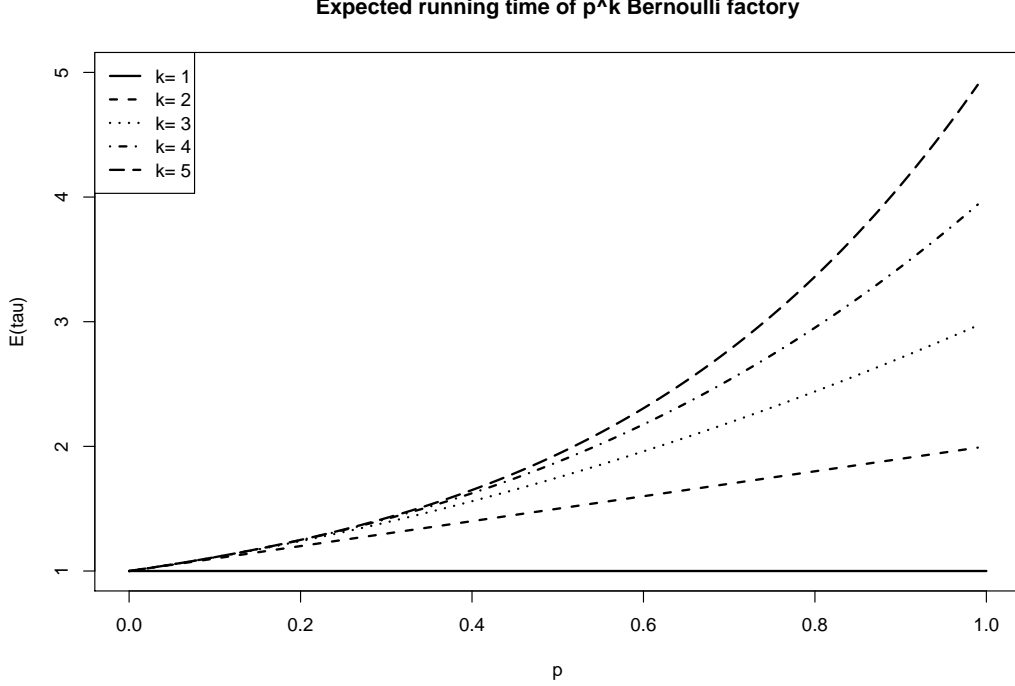


Figure 5: Expected running time for the $f(p) = p^k$ Bernoulli factory described, with early stopping. Without early stopping, the running time is deterministically equal to k . The early stopping procedure significantly increases the efficiency, particularly when p is small or k is large. The minimum number of queries is one for all k and p , and of course the $k = 1$ (identity) Bernoulli factory requires exactly one query for any p .

In the case of (b), (c) and (d) the overall running time is bounded if and only if the running time of both A_f and A_g is bounded.

Note that (b) is a special case of (c). Additionally, since there is a trivial Bernoulli factory for $f(p) \equiv 0$, (b) implies that αf is simulable for any $\alpha \in [0, 1]$.

It is not immediately obvious that not all functions are simulable. However, one counterexample is the following.

Example 6. $f(p) = 2p, p \in (0, 1/2)$.

This function turns out not to be simulable, although the truncated function $f(p) = 2p \wedge (1 - \varepsilon)$ — or equivalently, restricting the domain to $p \in (0, (1 - \varepsilon)/2)$ — is simulable for any $\varepsilon > 0$. In this case the number of queries required is $O(\varepsilon^{-1})$ (Huber, 2016).

The problem with the unrestricted function is that $f(p)$ approaches 1 within $p \in (0, 1)$, which violates one of the necessary and sufficient conditions given in the Keane-O’Brien theorem [REF].

Lemma 2 (Łatuszyński et al. (2011, Corollary 3.3)). *It is not possible to construct a Bernoulli factory for $f(p) = 2p$ where $p \in (0, 1/2)$.*

The following proof is an intuitive one, adapted from Łatuszyński et al. (2011). The intuition behind the proof is that there is no sharp change between say a $(1/4)$ -coin and a $(1/2)$ -coin; but there is a qualitative difference between the resulting $(1/2)$ -coin and 1-coin respectively. The latter is no longer random, but deterministically returns heads every time, regardless of the outcome on the $(1/2)$ -coin that seeds it.

Proof. Let \mathbb{P}_p denote the probability measure when the seeding coin has probability p of showing heads. We will compare the cases $p = 1/4$ and $p = 1/2 - \varepsilon$. First,

$$\lim_{t \rightarrow \infty} \mathbb{P}_{1/4}[\mathcal{A}(X_1, \dots, X_T) = 0, T < t] = \frac{1}{2}$$

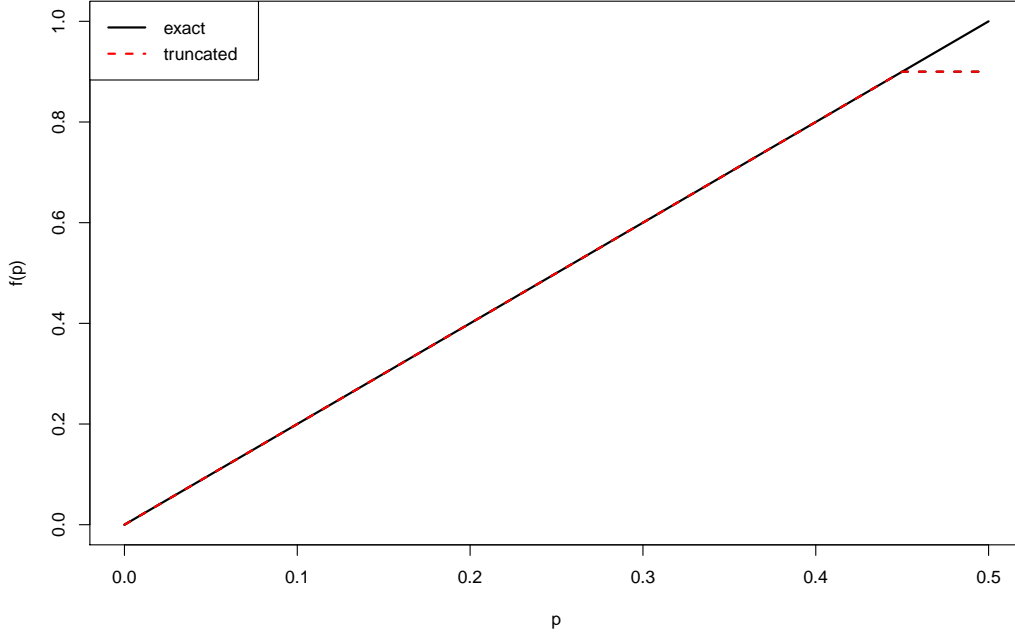


Figure 6: The function $f(p) = 2p$ (black) is not simulable; whereas the truncated function $f(p) = 2p \wedge (1 - \varepsilon)$ (red) is simulable for any $\varepsilon > 0$, using $O(\varepsilon^{-1})$ queries.

Thus there exists a $t_0 \in \mathbb{N}$ such that

$$\mathbb{P}_{1/4}[\mathcal{A}(X_1, \dots, X_T) = 0, T \leq t_0] > \frac{1}{4}$$

We have for any $p \in [1/4, 1/2]$:

$$\mathbb{P}_p[\mathcal{A}(X_1, \dots, X_T) = 0, T \leq t_0] \geq \left(\frac{1}{2}\right)^{t_0} \mathbb{P}_{1/4}[\mathcal{A}(X_1, \dots, X_T) = 0, T \leq t_0] \geq \left(\frac{1}{2}\right)^{t_0+2}$$

Since, restricted to $T \leq t_0$, \mathbb{P}_p has the same support for any $p \in (0, 1)$, we can consider the ratio of probabilities of any sequence in $\{0, 1\}^{t_0}$. In particular, we consider

$$\inf_{p \in [1/4, 1/2]} \inf_{A \subseteq \{0, 1\}^{t_0}} \frac{\mathbb{P}_p(A)}{\mathbb{P}_{1/4}(A)}$$

Now

$$\mathbb{P}_{1/2-\varepsilon}[\text{output } 1] = 1 - \mathbb{P}_{1/2-\varepsilon}[\mathcal{A}(X_1, X_2, \dots) = 0] \leq 1 - \mathbb{P}_{1/2-\varepsilon}[\mathcal{A}(X_1, \dots, X_T) = 0, T \leq t_0] \leq 1 - \left(\frac{1}{2}\right)^{t_0+2}$$

Taking the limit as $\varepsilon \rightarrow 0$, the LHS is $f(1/2) = 1$, while the RHS is independent of ε and remains strictly less than 1. This provides the contradiction, and hence such an algorithm \mathcal{A} cannot exist. \square

Theorem 1 (Keane and O'Brien (1994)). *A function $f : S \rightarrow [0, 1]$ is simulable if and only if f is continuous on S and either*

- (a) f is constant on S , or
- (b) there exists $n \in \mathbb{N}$ such that $\min\{f(p), 1 - f(p)\} \geq \min\{p^n, (1 - p)^n\}$ for all $p \in S$.

Roughly speaking, the second condition is “ $f(p)$ is polynomially bounded away from 0 and 1 for all $p \in S \cap (0, 1)$ ”. A proof is given in Keane and O'Brien (1994), but is not reproduced here. We now see that Lemma 2 follows as a corollary from the Keane-O'Brien theorem.

Corollary 1. *It is not possible to construct a Bernoulli factory for $f(p) = 2p$ where $p \in (0, 1/2)$.*

Proof. Firstly, f is continuous on $S = (0, 1/2)$, so the initial condition of Theorem 1 is satisfied. Since f is not constant, we require that there exists $n \in \mathbb{N}$ such that $\min\{f(p), 1 - f(p)\} \geq \min\{p^n, (1 - p)^n\}$ for all $p \in S$. Take $p = \frac{1}{2} - \varepsilon$. Then for all $\varepsilon > 0$, $p \in S$. We calculate the relevant quantities and, since the inequality should hold for all $\varepsilon > 0$, we can take the limit $\varepsilon \rightarrow 0$:

$$\begin{aligned} \min\{f(p), 1 - f(p)\} &= 2\varepsilon \xrightarrow{\varepsilon \rightarrow 0} 0 \\ \min\{p^n, (1 - p)^n\} &= \left(\frac{1}{2} - \varepsilon\right)^n \xrightarrow{\varepsilon \rightarrow 0} \frac{1}{2^n}. \end{aligned}$$

Since $1/2^n > 0$ for all $n \in \mathbb{N}$, we have that for all n there exists ε such that $\min\{f(p), 1 - f(p)\} < \min\{p^n, (1 - p)^n\}$; therefore f violates the second condition of Theorem 1. \square

4 The quantum advantage

It is known that for certain computational problems, quantum computers provide a theoretical advantage over classical computers. For example, the Deutsch-Jozsa algorithm (Deutsch and Jozsa, 1992) is an example of a problem (if not a particularly practical one) which a quantum computer can solve in exponentially less time than the best possible classical algorithm.

It is widely believed (at least within the field) that quantum computing really is more powerful than its classical counterpart. Quantum computers are able to do anything that classical computers can, and quantum speed-ups have been proved for a handful of carefully constructed problems (see for example Grover’s algorithm for unstructured search (Grover, 1997)), so from a theoretical viewpoint there is no denying the so-called ‘quantum supremacy’.

There is, however, a lack of rigorous results for any computational problems with significant practical applications. One such problem was apparently solved by Kerenidis and Prakash (2016), who came up with a quantum speed-up for the recommendation problem — that is, how Amazon and Netflix estimate which products you will like given yours and other customers’ ratings of all products. However, the authors were unable to prove that no comparably fast classical algorithm exists, and one has now been found by Tang (2018).

4.1 A quantum Bernoulli factory

The Bernoulli factory problem is an example of a practically applicable problem that is nevertheless simple to state and analyse. It was therefore a good candidate in the quest for a practical and provable quantum advantage, and indeed Dale et al. (2015) prove this advantage.

We will use the term *quantum Bernoulli factory* to refer to an algorithm that takes p -quoins as inputs, and outputs a Bernoulli($f(p)$) sample, where as usual $p \in S$ is unknown and $f : S \rightarrow [0, 1]$ is known. A p -quoin is the quantum analogue of a classical p -coin, given by the quantum state

$$|p\rangle = \sqrt{1 - p}|0\rangle + \sqrt{p}|1\rangle.$$

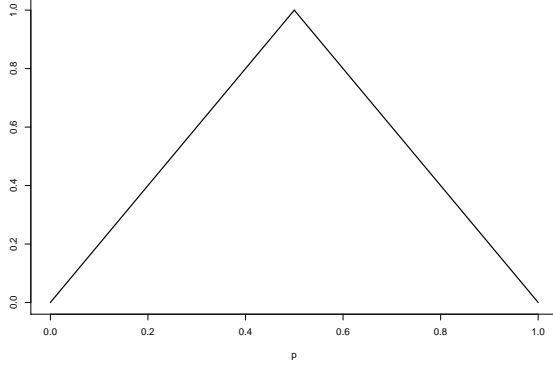
This is a sensible definition since measuring $|p\rangle$ in the computational basis gives output 0 with probability $1 - p$ or 1 with probability p , like a classical p -coin. Note that the quantum Bernoulli factory has a quantum input (quoins) and a classical output (measurement outcome). We will call a function *quantum simulable* if there exists a quantum Bernoulli factory for it.

Dale et al. (2015) prove that quantum computation allows a strictly larger class of Bernoulli factories to be constructed, taking as an example the notorious $f(p) = 2p$. While this function is not simulable in the classical setting, the authors show that it is quantum simulable, and provide an explicit algorithm using two entangled qubits. They also determine necessary and sufficient conditions (Theorem 2) for a function to be quantum simulable, which are weaker than those for classical simulability provided by the Keane-O’Brien theorem (Theorem 1).

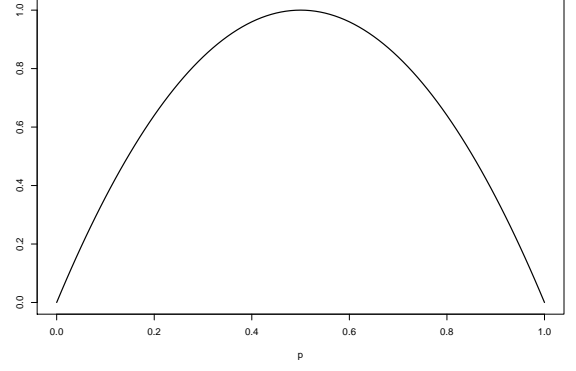
Example 7 gives the quantum algorithm of Dale et al. (2015) for the function $f_{\wedge}(p) := 2 \min\{p, 1 - p\}$, $p \in [0, 1]$. This function (Figure 7a) coincides with that of Example 6 when the domain is restricted to $p \in (0, 1/2)$, so it is clearly sufficient to find a Bernoulli factory for f_{\wedge} .

Example 7. First note that

$$f_{\wedge}(p) = 2 \min\{p, 1 - p\} = 1 - \sqrt{1 - 4p(1 - p)} = \sum_{k=1}^{\infty} \binom{2k}{k} \frac{1}{(2k - 1)2^{2k}} (4p(1 - p))^k =: \sum_{k=1}^{\infty} q_k (4p(1 - p))^k.$$



(a) The function $f_{\wedge}(p)$, which cannot be simulated classically since it touches 1 within its domain.



(b) Simulation of $f_{\wedge}(p)$ reduces to simulation of this function $4p(1-p)$, which is also classically unobtainable.

where q_k are constants not depending on p , and are hence simulable. Due to property (b) of Lemma 1, f_{wedge} is therefore simulable if there exists a Bernoulli factory for $(4p(1-p))^k$. Since this is a case of Example 4, we need only construct a Bernoulli factory for $4p(1-p)$. This function is not classically simulable since it still touches 1 at $p = 0.5$.

However, Dale et al. (2015) describe the following quantum Bernoulli factory for this function. We first prepare two quoin; that is, we have two qubits jointly in the product state $|p\rangle|p\rangle$. We then measure the joint state in the Bell basis, say with the operator

$$0 \cdot |\phi^+\rangle\langle\phi^+| + 1 \cdot |\phi^-\rangle\langle\phi^-| + 2 \cdot |\psi^+\rangle\langle\psi^+| + 3 \cdot |\psi^-\rangle\langle\psi^-|.$$

Then the measurement outcomes are:

$$\begin{aligned} \mathbb{P}(\text{observe } 0) &= \langle v|\phi^+\rangle\langle\phi^+|v\rangle = 1/2 \\ \mathbb{P}(\text{observe } 1) &= \langle v|\phi^-\rangle\langle\phi^-|v\rangle = \frac{(1-2p)^2}{2} \\ \mathbb{P}(\text{observe } 2) &= \langle v|\psi^+\rangle\langle\psi^+|v\rangle = 2p(1-p) \\ \mathbb{P}(\text{observe } 3) &= \langle v|\psi^-\rangle\langle\psi^-|v\rangle = 0 \end{aligned}$$

This means that $\mathbb{P}(\text{observe } 1 \text{ or } 2) = 1/2$, and $\mathbb{P}(\text{observe } 2 \mid \text{observe } 1 \text{ or } 2) = 4p(1-p)$, which is the required function. Thus our Bernoulli factory is as follows: measure $|p\rangle|p\rangle$ in the Bell basis until the observed outcome is either 1 or 2, then return heads if the outcome is 2, or tails if the outcome is 1.

Theorem 2 (Dale et al. (2015, Theorem 1)). *A function $f : [0, 1] \rightarrow [0, 1]$ is quantum simulable (using only operations on single qubits) if and only if:*

- (a) f is continuous on S
- (b) $Z := \{z : f(z) = 0\}$ and $W := \{w : f(w) = 1\}$ are both finite sets
- (c) $\forall z \in Z$ there exist constants $c \in \mathbb{R}$, $\delta > 0$, integer $k < \infty$ such that $c(p-z)^{2k} \leq f(p)$ for all $p \in [z-\delta, z+\delta]$
- (d) $\forall w \in W$ there exist constants $c \in \mathbb{R}$, $\delta > 0$, integer $k < \infty$ such that $1 - c(p-w)^{2k} \leq f(p)$ for all $p \in [w-\delta, w+\delta]$.

In the quantum case, the function f is allowed to touch 0 and 1 at finitely many points in its domain, as opposed to the classical case where $f(p)$ may only take the values 0 and 1 at $p \in \{0, 1\}$ or in the cases of the constant functions $f \equiv 0$ and $f \equiv 1$. The conditions (c) and (d) ensure that f does not go to 0 or 1 too quickly at these points.

Dale et al. (2015) also claim that the same necessary and sufficient conditions apply even when operations on more than one qubit are allowed. That is, taking advantage of quantum entanglement (as in Example 7) does not expand the set of simulable functions. It can, however, substantially reduce the consumption of p -quoin, as demonstrated in Patel et al. (2018).

5 Discussion

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