

Asymptotic analysis of genealogies induced by sequential Monte Carlo algorithms

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1 Introduction

Since its introduction to the statistics literature by Gordon et al. (1993), sequential Monte Carlo (SMC) has found a huge range of applications and has become an indispensable tool to practitioners from many fields. It is particularly useful in applications such as target tracking, where the model has a natural sequential structure. However, the sequential nature of SMC is helpful not only in inherently sequential settings, but in any setting where the quantity to be inferred has highly correlated components.

In these cases the more traditional Markov chain Monte Carlo methods are essentially useless because the strong dependence structure causes them to mix extremely slowly. Conversely, SMC exploits this dependence structure to create recursive Monte Carlo algorithms that can be efficient in this difficult setting. SMC allows a much wider class of models to be solved beyond the few cases where we can find an exact analytic solution.

The method is not without its problems, and it is one of these problems that we will examine in this work.

Sections 2 and 3 explain the background material on SMC and coalescent theory respectively. With these in place we are then able to explain the contribution of this work in the context of the SMC literature. In Section 5 we present the first novel result of this work with an outline proof; the full proof is in the Appendix. Section 6 introduces the next focus of this work, for which nothing is proved yet. Finally we summarise future directions of the work, and its potential impact, in Section 7.

Throughout the document we will use the compact notation $X_{m:n}$ as shorthand for X_m, X_{m+1}, \dots, X_n , as well as $X_{-n} := X_0, \dots, X_{n-1}, X_{n+1}, \dots, X_N$. We denote falling factorial powers $(x)_a := x(x-1) \dots (x-a+1)$, with the convention $(x)_0 = 1$.

2 Sequential Monte Carlo

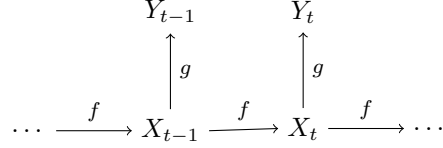
References for this section are Doucet et al. (2001), Del Moral et al. (2006), and Doucet and Johansen (2011).

2.1 Class of models

Although sequential Monte Carlo (SMC) methods can be applied in a much more general setting, they are particularly easy to motivate in the setting of state space models, where the “sequential” nature follows naturally from the discrete time steps present in the model. For the purposes of presenting the algorithm, let us consider a time-homogeneous state space model consisting of an unobservable discrete-time Markov process $X_{0:T}$ and observables $Y_{0:T}$, satisfying the conditional independence structure

$$\begin{aligned} (X_{t+1:T} \perp\!\!\!\perp X_{0:t-1}) &| X_t \\ (Y_t \perp\!\!\!\perp Y_{-t}, X_{-t}) &| X_t \end{aligned}$$

for all $t \in \{0, 1, \dots, T\}$, as represented by the graphical model below.



We assume for notational convenience that x_0, \dots, x_T take values in a common state space \mathcal{X} , and y_0, \dots, y_T in a common state space \mathcal{Y} , but these assumptions can be dropped.

Suppose we have the following model:

$$\begin{aligned}
X_0 &\sim \mu(\cdot) \\
X_{t+1} \mid (X_t = x_t) &\sim f(\cdot \mid x_t) \quad t = 0, \dots, T-1 \\
Y_t \mid (X_t = x_t) &\sim g(\cdot \mid x_t) \quad t = 0, \dots, T
\end{aligned}$$

where $(X_t)_{t=0}^T$ is an unobservable discrete-time Markov process and the observables $(Y_t)_{t=0}^T$ satisfy $Y_t \perp\!\!\!\perp \{Y_{-t}, X_{-t}\} \mid X_t$.

We assume that the *transition* and *emission* kernels have densities which are denoted by f and g respectively, but this is not necessary in general. We only require that we can sample from $\mu(\cdot)$ and $f(\cdot \mid x)$, and calculate *unnormalised* potentials $g(y \mid x)$, for all x, y .

As a concrete example, let us consider the application of target tracking. Suppose we are using radar to track the position of an aeroplane. The true trajectory of the aeroplane is unknown and is represented by $X_{0:T}$ (perhaps a sequence of positions in \mathbb{R}^3), with f encoding our model for how an aeroplane moves (perhaps some differential equations). What we observe is the output $Y_{0:T}$ of our radar equipment, which has some measurement uncertainty that is encoded in g .

2.2 Inference in state space models

Suppose we are in a Bayesian setting, where μ is our prior distribution at time 0, observations y_t arrive sequentially, and we want to infer information about the hidden states (either on- or off-line). The three main inference problems are:

Filtering (where is it now?) $p(x_t \mid y_{0:t})$

Prediction (where will it go next?) $p(x_{t+1} \mid y_{0:t})$

Smoothing (where has it been?) $p(x_{0:t} \mid y_{0:t})$

In the on-line setting, we take as our prior the posterior distribution from the previous time step $t-1$, and update it using the new observation y_t . The inference must be fast enough to keep up with the rate of arrival of observations, so in particular the complexity of the update must not increase with T . In the off-line setting, we take μ as the prior distribution, and infer the set of posteriors once all $T+1$ observations have arrived.

Prediction and filtering are essentially equivalent, because given a filtering distribution, the corresponding predictive distribution can be obtained by applying the transition kernel f . Smoothing is considered a harder task because it requires us to infer many more parameters from the same amount of information; indeed the dimension of the problem increases linearly with T .

In the case of linear Gaussian state space models, the posterior distributions of interest are available analytically, by way of the Kalman filter (Kalman, 1960) and Rauch-Tung-Striebel (RTS) smoother recursions (Rauch et al., 1965). Recursions are also available for some other conjugate models: see for example Vidoni (1999). The other analytic case occurs if the state space of $(X_t)_{t=0}^\infty$ is finite, in which case the forward-backward algorithm (Baum et al., 1970) yields the exact posteriors.

2.3 Particle approximation

In more complex models such techniques are not feasible, and we are forced to resort to Monte Carlo methods. For state space models, Markov chain Monte Carlo methods are not very effective due to the high dimension of the parameter space. But we can exploit the sequential nature of the underlying dynamics to decompose

the problem into a sequence of inferences of more manageable dimension. This is the motivation behind sequential Monte Carlo (SMC) methods.

The conditional independence structure in the model implies that the (joint) marginal distribution of the hidden states $X_{0:t}$ is given by

$$p(x_{0:t}) = \mu(x_0) \prod_{i=1}^t f(x_i | x_{i-1})$$

and that the likelihood of the observations $y_{0:t}$ given the underlying states $x_{0:t}$ takes the form

$$p(y_{0:t} | x_{0:t}) = \prod_{i=0}^t g(y_i | x_i).$$

Using the conditional independence structure, we can write

$$p(x_{0:t} | y_{0:t}) \propto g(y_t | x_t) f(x_t | x_{t-1}) p(x_{0:t-1} | y_{0:t-1}) \quad (1)$$

$$\propto \mu(x_0) g(y_0 | x_0) \prod_{i=1}^t f(x_i | x_{i-1}) g(y_i | x_i) \quad (2)$$

for $t = 0, \dots, M$, where the one-step recursion (1) is obtained using Bayes rule, and (2) is obtained by applying (1) t times. The filtering distribution $p(x_t | y_{0:t})$ can be obtained from (1) by marginalising out $x_{0:t-1}$, which is straightforward if Monte Carlo samples are available. The predictive distributions can also be derived from the smoothing distributions using

$$p(x_{t+1} | y_{0:t}) \propto f(x_{t+1} | x_t) p(x_{0:t} | y_{0:t}).$$

SMC provides a particle method to approximate to (1), given a model specification and a sequence of observations. Like the underlying process, the algorithm proceeds sequentially, returning its approximation to the smoothing distribution at each time step. This approximation is the empirical distribution of the particles:

$$\hat{p}(x_{0:t} | y_{0:t}) = \frac{1}{N} \sum_{i=1}^N \delta_{X_{0:t}^{(i)}} \quad (3)$$

The particle approximation is justified by various convergence results - see for example Del Moral (2013) for details.

A generic SMC algorithm is presented in Algorithm 1. For the state space model described above, we can take $K_{t+1}(x_t, \cdot) \equiv f(\cdot | x_t)$ and $g_{t+1}(x_t, x_{t+1}) \equiv g(y_{t+1} | x_{t+1})$.

Figure 1 illustrates the particle approximation arising from such an algorithm on a linear Gaussian model, with the exact posterior for reference.

If only the latest filtering distribution is required, we can marginalise out $\mathbf{x}_{0:t-1}$ at each step by simply throwing away the particle histories and keeping only the particle approximation \mathbf{x}_t to the filtering distribution at the current time t . The algorithm progresses in a Markovian fashion, only ever referring to the particles at the immediately previous step, so filtering distributions can be approximated with minimal memory usage. If, say, the mean and variance of $X_t | y_{0:t}$ at each time t are required, we can store just these summary statistics, plus the two most recent generations of particles, and throw away all other information about the particles at previous time steps. This is vital if one wishes to carry out filtering in an on-line fashion, as it prevents the memory requirements accumulating more than necessary.

The form of the RESAMPLE function in Algorithm 1 is discussed in Section 6.

- Add a section including theoretical justification for SMC...?

- If so, rearrange to have Sec: State Space Models \rightarrow Sub: Class of Models; Inference Problems; Methods (non-SMC) + Sec: SMC \rightarrow Sub: SMC for SSMS; Theoretical Justification for SMC.

Algorithm 1 Standard SMC

Require: $N, T, \mu, \{K_t\}, \{g_t\}, y_{0:T}$

```
1: for  $i \in \{1, \dots, N\}$  do
2:   Sample  $X_0^{(i)} \sim \mu(\cdot)$  ▷ initialise
3:    $w_0^{(i)} \leftarrow \frac{g_0(X_0^{(i)})}{\sum_{j=1}^N g_0(X_0^{(j)})}$ 
4: end for
5: for  $t \in \{0, \dots, T-1\}$  do
6:   Sample  $a_t^{(1:N)} \sim \text{RESAMPLE}(\{1, \dots, N\}, w_t^{(1:N)})$  ▷ resample particles
7:   for  $i \in \{1, \dots, N\}$  do
8:     Sample  $X_{t+1}^{(i)} \sim K_{t+1}(X_t^{(a_t^{(i)})}, \cdot)$  ▷ propagate particles
9:      $w_{t+1}^{(i)} \leftarrow g_{t+1}(X_t^{(a_t^{(i)})}, X_{t+1}^{(i)})$  ▷ calculate weights
10:  end for
11:   $W \leftarrow \sum_{j=1}^N w_{t+1}^{(j)}$ 
12:  for  $i \in \{1, \dots, N\}$  do
13:     $w_{t+1}^{(i)} \leftarrow \frac{1}{W} w_{t+1}^{(i)}$  ▷ normalise weights
14:  end for
15: end for
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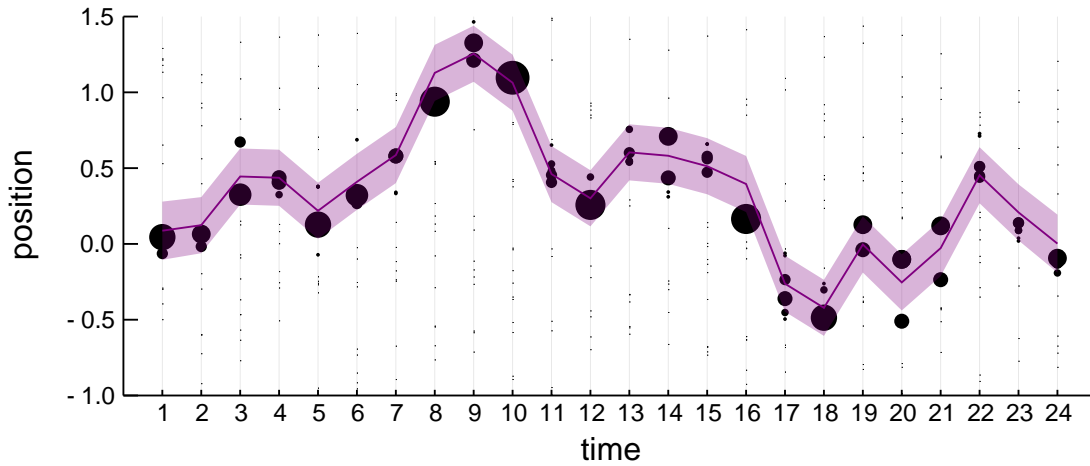


Figure 1: SMC particles before resampling for a linear Gaussian model. The purple ribbon shows the exact posterior mode and 95% credible interval, computed using the Kalman filter and RTS smoother. The black dots show the positions of the SMC particles, with size proportional to weight. After resampling all particles have equal weights but some are duplicated.

3 Coalescent Theory

- asymptotic coalescents for various population models
- time is going backwards now...
- Cannings model: exchangeability, also a key assumption for our SMC tings
- without exchangeability, can't even expect Markov process

Definition 1 (Möhle (1998)). Let \mathcal{E}_n denote the set of equivalence relations on $\{1, \dots, n\}$. A *discrete-time coalescent* is a stochastic process $(R_t)_{t \in \mathbb{N}}$ taking values in \mathcal{E}_n such that $R_0 = \{(1, 1), (2, 2), \dots, (n, n)\}$ and $\mathbb{P}[R_{t+1} = \eta \mid R_{0:t-1}, R_t = \xi] > 0$ only if $\xi \subseteq \eta$.

That is, the initial state is the trivial relation where each index is in its own equivalence class, and the only possible forward-in-time transitions are staying the same or merging some equivalence classes together. An obvious consequence of this is that the state where all of the indices are in the same equivalence class is an absorbing state for the process.

The genealogical interpretation of the equivalence relations is that $(i, j) \in R_t$ if and only if individuals i and j share a common ancestor in generation t .

3.1 Kingman's coalescent

Imagine we have a population with fixed size N over discrete generations, where each individual is descended from one randomly chosen individual of the previous generation. Then for each individual in the present generation, we can trace their *lineage* back through the generations. If we trace two lineages back in time, at some generation they may descend from the same individual, at which point we say they have *coalesced*. Once two lineages have coalesced they will stay together going backwards in time. The combined lineages of $n \leq N$ of the present individuals therefore forms a tree, or several non-overlapping trees, the entirety of which we refer to as the *ancestry* or *genealogy* of those n individuals.

Kingman's *n-coalescent* provides a model for such genealogies. Kingman showed in (Kingman, 1982a,b,c) that the *n-coalescent* is the limiting process for samples from a wide class of population models as $N \rightarrow \infty$.

The defining feature of the model is that each pair of lineages merges with unit rate. This means that many coalescences occur while there are many distinct lineages present. In particular, the *n-coalescent* can be formulated as a Poisson process where pairs of lineages coalesce independently at rate 1, with the pair to coalesce being chosen uniformly at random (Wakeley, 2009, Section 3.2).

In the notation of Wakeley (2009), let $T_i; i = 2, \dots, n$ be the i^{th} coalescence time, that is, the length of time for which there are exactly i branches in the sample genealogy. The *n-coalescent* is the process in which these times are distributed as independent Exponentials with rate $\binom{i}{2}$.

Möhle (1998) writes the same process in terms of the infinitesimal generator Q of a Markov process on the set of equivalence relations on n elements, having entries

$$q_{\xi\eta} = \begin{cases} -\binom{b}{2} & \text{if } \xi = \eta \\ 1 & \text{if } \xi \prec \eta \\ 0 & \text{otherwise} \end{cases}$$

where b is the number of equivalence classes of ξ , and $\xi \prec \eta$ means that η is a state with exactly one more pair of lineages coalesced compared to ξ .

4 SMC genealogies as coalescents

- motivate the work: see notes on back of Adam's scribbled-on copy & commented paragraph in intro
- resampling viewed backwards in time: branching process \rightarrow coalescent process
- asymptotic properties of SMC lit review: CLT, path storage, coalescence etc.
- the gap in knowledge that we aim to fill
- why the simple SSM described at start is sufficient to demonstrate coalescence
- how to deal with the difference between Kingman (time stops once all coalesced) and SMC genealogies (fixed time frame T)

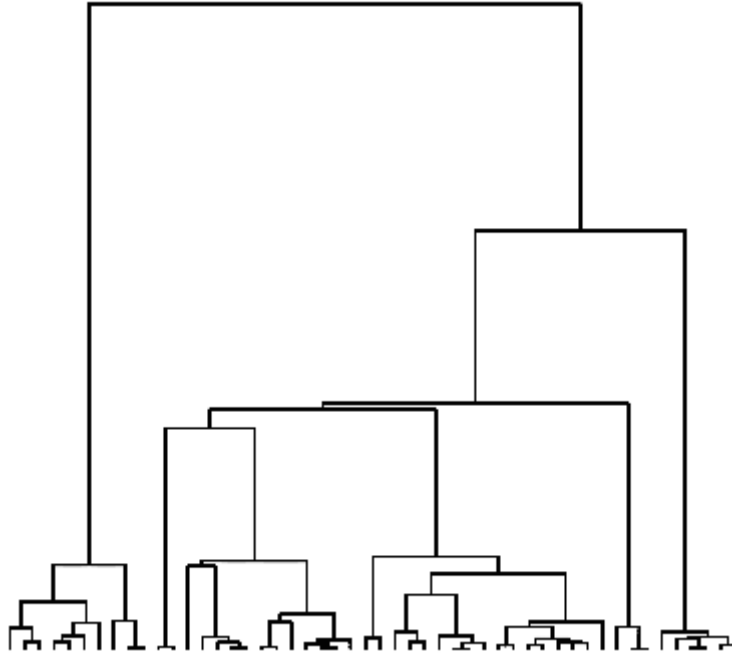


Figure 2: A realisation of Kingman’s n -coalescent for a sample of size $n = 50$. At first there are many distinct lineages, and mergers happen rapidly. Once there are fewer distinct lineages left, they take longer to merge. The process spends about half of its time with just two or three distinct lineages. (Source: Wikimedia Commons)

4.1 Ancestral degeneracy

The resampling step in Algorithm 1 induces a genealogical tree. During resampling, some particles have multiple offspring while others have none. The particles with no offspring “die out”; they are not in the lineage of any time T particle. So unless the offspring variance is low, the N time T particles are likely to originate from only a few distinct time 0 ancestors. An example of this is shown in Figure 3.

In order to estimate filtering distributions $p(x_t|y_{0:t})$, we only require a sample of particles at the current time step, so if \mathcal{X} is continuous we typically have N distinct positions given by the N particles. Then the empirical measure has mass in N locations, and the Monte Carlo error for estimating expectations under $p(x_t|y_{0:t})$ scales as $O(N^{-1/2})$ [REF].

However, we do not achieve the same performance in the case of estimating the smoothing distributions $p(x_{0:t}|y_{0:t})$. In this case we require a sample of trajectories over times $0 : t$ as opposed to a sample of particles at time t . The coalescence of lineages is an unavoidable effect of resampling, and it causes more and more of these trajectories to coincide the further into the past we look. The resulting empirical measure typically consists of N distinct masses, but the marginals at early times may just consist of a single mass repeated N times. This phenomenon, known as *ancestral degeneracy*, is illustrated in Figure 4.

So if we are really interested in the smoothed states a long way into the past, the estimation variance will be huge. This problem was identified even in the early literature (Gordon et al., 1993), where some ad hoc methods were proposed to reduce it. Since then there has been a lot of work towards mitigating ancestral degeneracy, some of which are discussed in Section 6.

- Ancestral vs. weight degeneracy

- More discussion of general techniques to mitigate it

4.2 Existing results

Rewrite this section after adding details in Sec 3...

The first results showing convergence of population models to the Kingman coalescent appear in Kingman’s

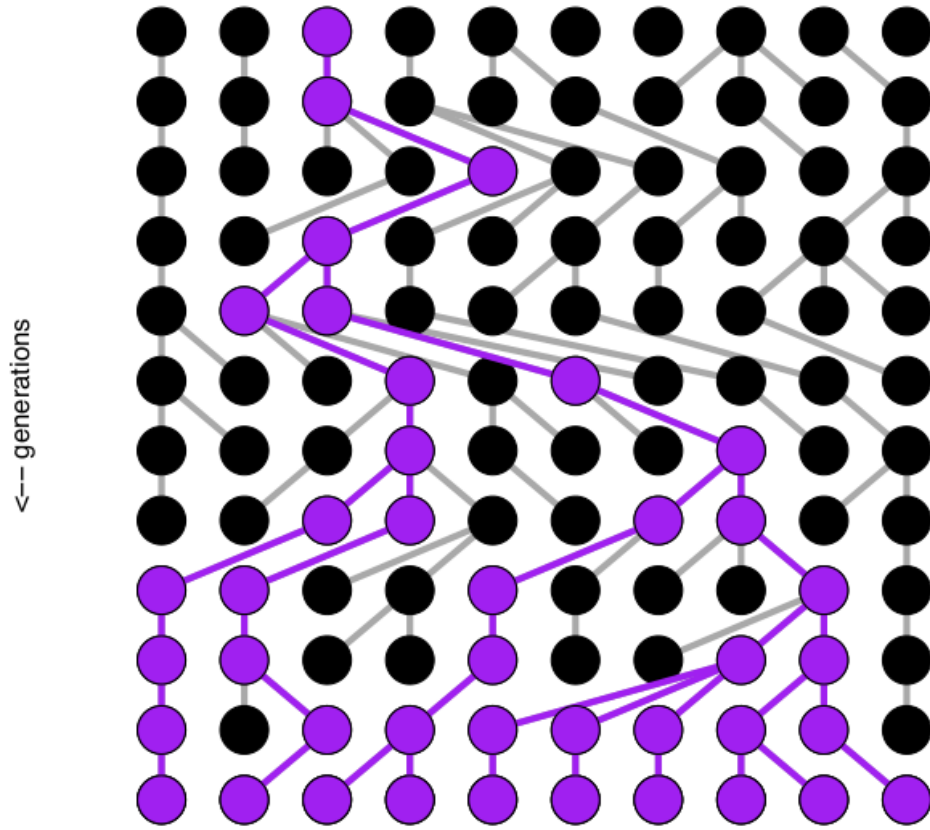


Figure 3: Genealogical tree induced by resampling over 12 generations with $N = 10$ particles. At each resampling step, any particles with no offspring “die out”; they are not in the lineage of any time T particle. In this realisation, the N particles at time T all originate from the same time 0 ancestor. “Dead” particles/lineages are coloured black/grey, while the “live” tree is highlighted in purple.

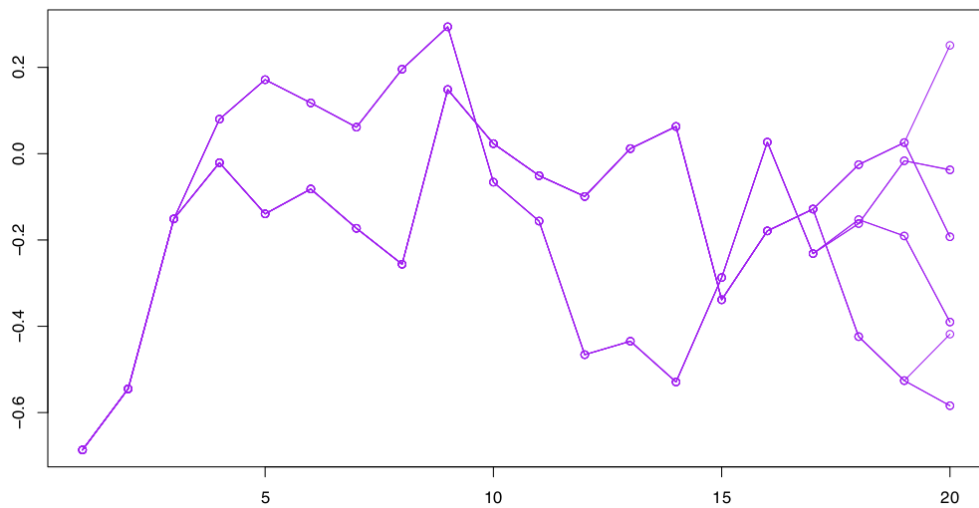


Figure 4: A sample of $N = 6$ trajectories, illustrating ancestral degeneracy. At the “present” time there are six distinct lineages, but just three steps back they have coalesced onto only two lineages, and it takes less than 20 steps back before only one lineage remains.

original paper (Kingman, 1982c) introducing the Kingman coalescent. This includes, but is not limited to, the neutral Wright-Fisher model (Fisher, 1923, 1930; Wright, 1931) and the Moran model (Moran, 1958). It is known that a general class of exchangeable models known as neutral Cannings models converge to the Kingman coalescent (Etheridge, 2011, Section 2.2). Möhle (1998) proved convergence for a larger class, including some non-exchangeable models.

Koskela et al. (2018) presents the first application of this type of analysis to SMC genealogies. Their result relies heavily on the methods introduced by Möhle (1998). They were able to prove convergence to the Kingman coalescent for genealogies induced by standard SMC algorithms with multinomial resampling. In the following sections we attempt to extend their result to cover some other SMC algorithms.

5 Conditional SMC

- Introduce particle MCMC (perhaps including basic algorithm)
- Then motivate conditional SMC as update in PMCMC

Conditional SMC differs from the standard algorithm in that one predetermined trajectory (that is, a sequence of particle positions and the corresponding ancestral line) is conditioned to survive all of the propagation and resampling steps. We will refer to this sequence as the *immortal trajectory*, following the terminology used for conditioned Galton-Watson processes, and the *immortal particle* will refer to the particle in a particular generation that is part of the immortal trajectory.

The conditional SMC algorithm was proposed by Andrieu et al. (2010) for use in the *particle Gibbs* sampler, which they introduce as part of a more general class of particle MCMC methods. In the particle Gibbs sampler, the standard SMC algorithm does not admit the desired target distribution, so this conditional version must be used instead.

When used as a component of the particle Gibbs algorithm, the immortal trajectory $x_{0:T}^*$ for each SMC run is sampled from the trajectories output from the previous run (Andrieu et al., 2010, Section 2.4.3). However, for our purposes we just consider a single SMC run for which the immortal trajectory is fixed.

A conditional SMC algorithm employing multinomial resampling is described in Algorithm 2.

In the particle Gibbs sampler, it is crucial that the conditional SMC output maintains at least two distinct trajectories. The immortal trajectory will of course be among the surviving trajectories, but additionally, the new immortal trajectory (for the next SMC run) is chosen from among the surviving trajectories. Thus if all the trajectories coalesce onto the immortal trajectory, we are forced to choose the same immortal trajectory for the next run, at least for some early time steps. One can imagine that if there was a high probability of full coalescence on each run, we could easily end up with samples from $p(x_{0:T}|y_{0:T})$ that are identical in some coordinates $0 : t$, which would not lead to good results overall.

The problem can be avoided by using a sufficiently large number of particles for the fixed time window T of the conditional SMC runs. This would require a priori knowledge of the coalescence mechanism, which is not available. However, Corollary 1 could possibly provide such knowledge. If, say, we want to ensure that the probability of all N lineages coalescing is below a certain threshold, all of the relevant information is encoded in the distribution of the time to MRCA of the genealogical process. For the Kingman coalescent this distribution is known, and Corollary 1 states that as $N \rightarrow \infty$ the genealogy is a Kingman coalescent. The remaining question is whether the Kingman coalescent provides a reasonable approximation outside of the asymptotic regime - since in reality we simulate finitely many particles. We intend to investigate this question by way of a simulation study.

5.1 Genealogies of conditional SMC algorithms

In this section we calculate various quantities related to the genealogical process induced by conditional SMC with multinomial resampling. By writing these in terms of the corresponding quantities for standard SMC with multinomial resampling, we are able to apply results from Koskela et al. (2018). In this way we will show that the genealogical process converges to the Kingman coalescent, in the sense of finite-dimensional distributions, as the number of particles $N \rightarrow \infty$.

The derivations of the expressions (4), (5), (6), along with details of the application of results from Koskela et al. (2018), are relegated to the appendix. Below is an overview of the proof. To prove convergence to the

Algorithm 2 Conditional SMC with multinomial resampling

Require: $N, T, \mu, \{K_t\}, \{g_t\}, y_{0:T}, x_{0:T}^*$

```
1: for  $i \in \{1, \dots, N\}$  do
2:   Sample  $X_0^{(i)} \sim \mu(\cdot)$  ▷ initialise
3: end for
4: Sample  $a_0^* \sim \text{Uniform}(\{1, \dots, N\})$ 
5:  $X_0^{(a_0^*)} \leftarrow x_0^*$ 
6: for  $i \in \{1, \dots, N\}$  do
7:    $w_0^{(i)} \leftarrow \frac{g_0(X_0^{(i)})}{\sum_{j=1}^N g_0(X_0^{(j)})}$ 
8: end for
9: for  $t \in \{0, \dots, T-1\}$  do
10:  Sample  $a_t^{(1:N)} \sim \text{Categorical}(\{1, \dots, N\}, w_t^{(1:N)})$  ▷ resample particles
11:  Sample  $a_{t+1}^* \sim \text{Uniform}(\{1, \dots, N\})$ 
12:   $a_t^{(a_{t+1}^*)} \leftarrow a_t^*$ 
13:  for  $i \in \{1, \dots, N\}$  do
14:    Sample  $X_{t+1}^{(i)} \sim K_{t+1}(X_t^{(a_t^{(i)})}, \cdot)$  ▷ propagate particles
15:  end for
16:   $X_{t+1}^{(a_{t+1}^*)} \leftarrow X_{t+1}^*$ 
17:  for  $i \in \{1, \dots, N\}$  do
18:     $w_{t+1}^{(i)} \leftarrow g_{t+1}(X_t^{(a_t^{(i)})}, X_{t+1}^{(i)})$  ▷ calculate weights
19:  end for
20:   $W \leftarrow \sum_{j=1}^N w_{t+1}^{(j)}$ 
21:  for  $i \in \{1, \dots, N\}$  do
22:     $w_{t+1}^{(i)} \leftarrow \frac{1}{W} w_{t+1}^{(i)}$  ▷ normalise weights
23:  end for
24: end for
```

Kingman coalescent, we must control the rates of different types of mergers. In particular, we ensure that in the large population limit (under an appropriate time-scaling), pairwise mergers happen at the correct rate, and larger mergers never occur.

Throughout the following we use tilde to indicate the conditional SMC versions of the untilded quantities relating to standard SMC, always with multinomial resampling.

Firstly, we have the expected coalescence rate:

$$\mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] = \frac{N-2}{N}\mathbb{E}[c_N(t)|\mathcal{F}_{t-1}] + \frac{2}{N}\mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \quad (4)$$

Then the expected rate of super-binary mergers (that is, more than two lineages merging simultaneously into one or more lineages) is bounded above by:

$$\begin{aligned} \mathbb{E}[\tilde{D}_N(t)|\mathcal{F}_{t-1}] &\leq \mathbb{E}[D_N(t)|\mathcal{F}_{t-1}] + \frac{3}{N}\mathbb{E}[(w_t^{(1)})^2|\mathcal{F}_{t-1}] + \frac{4}{N^2}\mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \\ &\quad + \frac{4}{N}\sum_{i=2}^N\mathbb{E}[w_t^{(1)}(w_t^{(i)})^2|\mathcal{F}_{t-1}] + \frac{2}{N^2}\sum_{i=2}^N\mathbb{E}[w_t^{(1)}w_t^{(i)}|\mathcal{F}_{t-1}] + \frac{1}{N^2}\sum_{i=2}^N\mathbb{E}[(w_t^{(i)})^2|\mathcal{F}_{t-1}] \end{aligned} \quad (5)$$

And lastly the expectation of the squared coalescence rate is bounded above by:

$$\begin{aligned} \mathbb{E}[\tilde{c}_N(t)^2|\mathcal{F}_{t-1}] &\leq \mathbb{E}[c_N(t)^2|\mathcal{F}_{t-1}] + \frac{4}{N}\mathbb{E}[(w_t^{(1)})^3|\mathcal{F}_{t-1}] + \frac{12}{N^2}\mathbb{E}[(w_t^{(1)})^2|\mathcal{F}_{t-1}] + \frac{4}{N(N-2)}\mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \\ &\quad + \frac{4}{N}\sum_{i=2}^N\mathbb{E}[w_t^{(1)}(w_t^{(i)})^2|\mathcal{F}_{t-1}] \end{aligned} \quad (6)$$

We then apply Lemma 3 of Koskela et al. (2018) to obtain the more tractable expressions

$$\begin{aligned} \frac{\varepsilon^4}{Na^4} + O(N^{-2}) &\leq \mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] \leq \frac{a^4}{N\varepsilon^4} + O(N^{-2}) \\ \mathbb{E}[\tilde{D}_N(t)|\mathcal{F}_{t-1}] &\leq \frac{C}{N}\mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] + O(N^{-3}) \\ \mathbb{E}[\tilde{c}_N(t)^2|\mathcal{F}_{t-1}] &\leq \frac{C}{N}\mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] + O(N^{-3}) \end{aligned}$$

and define the time-scaling

$$\tilde{\tau}_N(t) := \min \left\{ s \geq 1 : \sum_{r=1}^s \tilde{c}_N(r) \geq t \right\}. \quad (7)$$

Then, using Koskela et al. (2018, Lemma 2), which readily generalises to our modified quantities, we are able to verify the four conditions of Koskela et al. (2018, Theorem 1). Finally we are able to conclude the following.

Corollary 1. *Under the conditions of Koskela et al. (2018, Lemma 3), the genealogy of any n particles from a conditional SMC algorithm with multinomial resampling converges to Kingman's n -coalescent in the sense of finite-dimensional distributions, under the time-scaling defined in (7).*

6 Alternative resampling schemes

- stratified, systematic
- Gerber et al's optimal scheme
- Del Moral's toy scheme?
- any others...?
- need to be explicit about how the Unif RV is employed in the resampling

- residual resampling when weights are equal doesn't resample
- stratified/systematic resampling: dependence on order of particles
- results: theorem for residual resampling (hopefully) - haha you wish
- maybe results for other schemes (wow, so ambitious)

There is a great deal of flexibility in the function referred to as RESAMPLE in Algorithm 1. The most straightforward choice is multinomial resampling (Efron and Tibshirani, 1994), which is also the easiest to analyse. However, multinomial resampling is well known to be sub-optimal in terms of the resulting Monte Carlo variance, and is rarely used in practice. For instance, Douc et al. (2005) proves that both residual resampling and stratified resampling yield lower variance. In this section we will present some resampling schemes that claim to perform better than multinomial resampling.

Multinomial resampling. (Efron and Tibshirani, 1994) This is the simplest way to do resampling. Parental indices are chosen independently from $\{1, \dots, N\}$, each with probability given by the weight of the corresponding particle $w_t^{(i)}$. That is, $a_t^{(1:N)} \sim \text{Categorical}(\{1, \dots, N\}, w_t^{(1:N)})$. This implies the joint distribution of the offspring counts is $v_t^{(1:N)} \stackrel{d}{=} \text{Multinomial}(N, w_t^{(1:N)})$. Note that in this case the parental indices are chosen independently, but the resulting offspring counts are negatively correlated.

Residual resampling. (Liu and Chen, 1998; Whitley, 1994) Each particle $x_t^{(i)}$ is deterministically assigned $\lfloor Nw_t^{(i)} \rfloor$ offspring, and the remaining $R := N - \sum \lfloor Nw_t^{(i)} \rfloor$ offspring are assigned multinomially in proportion to the unaccounted-for weight. This yields offspring counts $v_t^{(1:N)} \stackrel{d}{=} \lfloor Nw_t^{(1:N)} \rfloor + \text{Multinomial}(R, (Nw_t^{(1:N)} - \lfloor Nw_t^{(1:N)} \rfloor)/R)$. This scheme incorporates a deterministic part which ensures that every particle with weight $> 1/N$ is guaranteed to survive. This is a desirable property as it prevents the random loss of promising particles.

7 Discussion

- results so far
- impact of this work: to practitioners, to enriching the SMC literature, interpretation within pop gen.
- future directions

A Proof of Corollary 1

In the derivation of (4) – (6) we will make extensive use of the formula for factorial moments of the multinomial distribution given in Mosimann (1962, p.67):

$$\mathbb{E}[(X_i)_a (X_j)_b] = (n)_{a+b} p_i^a p_j^b \quad (8)$$

where $(X_1, \dots, X_k) \sim \text{Multinomial}(n, \mathbf{p})$. To apply this formula we need to write everything in terms of falling factorial powers. The required conversions are summarised in Table 1.

In standard SMC with multinomial resampling, the marginal offspring distributions, conditioned on the filtration \mathcal{F}_{t-1} generated by the previous offspring counts, are

$$v_t^{(i)} \stackrel{d}{=} \text{Binomial}(N, w_t^{(i)}), \quad i = 1, \dots, N$$

where $v_t^{(i)}$ is the number of offspring in generation $t+1$ of the i th particle in generation t , N is the number of particles and $w_t^{(i)}$ is the weight associated with the i th particle in generation t .

In conditional SMC we condition on the immortal trajectory surviving each resampling step. By exchangeability we can set without loss of generality that the immortal trajectory consists of particle 1 in each generation. At each resampling step, particle 1 must therefore choose particle 1 as its parent, while the remaining $N-1$ offspring are assigned multinomially to the N possible parents. The marginal offspring

x	$=$	$(x)_1$
x^2	$=$	$(x)_2 + (x)_1$
x^3	$=$	$(x)_3 + 3(x)_2 + (x)_1$
x^4	$=$	$(x)_4 + 6(x)_3 + 7(x)_2 + (x)_1$
xy	$=$	$(x)_1(y)_1$
x^2y	$=$	$(x)_2(y)_1 + (x)_1(y)_1$
xy^2	$=$	$(x)_1(y)_2 + (x)_1(y)_1$
x^2y^2	$=$	$(x)_2(y)_2 + (x)_2(y)_1 + (x)_1(y)_2 + (x)_1(y)_1$
$(x+1)_2$	$=$	$(x)_2 + 2(x)_1$
$(x+1)^2$	$=$	$(x)_2 + 3(x)_1 + 1$
$(x+1)_2(x+1)$	$=$	$(x)_3 + 5(x)_2 + 4(x)_1$
$(x+1)_2^2$	$=$	$(x)_4 + 8(x)_3 + 14(x)_2 + 4(x)_1$

Table 1: Conversion of ordinary powers into falling factorial powers

distributions are then

$$\begin{aligned}\tilde{v}_t^{(1)} &\stackrel{d}{=} 1 + \text{Binomial}(N-1, w_t^{(1)}) \\ \tilde{v}_t^{(i)} &\stackrel{d}{=} \text{Binomial}(N-1, w_t^{(i)}), \quad i = 2, \dots, N.\end{aligned}$$

First let us consider the pair-merger rate

$$c_N(t) := \frac{1}{(N)_2} \sum_{i=1}^N (v_t^{(i)})_2.$$

For standard SMC the expected value is, using the tower rule,

$$\mathbb{E}[c_N(t)|\mathcal{F}_{t-1}] = \frac{1}{(N)_2} \sum_{i=1}^N \mathbb{E}[\mathbb{E}[(v_t^{(i)})_2]|\mathcal{F}_{t-1}] = \frac{1}{(N)_2} \sum_{i=1}^N \mathbb{E}[(N)_2(w_t^{(i)})^2|\mathcal{F}_{t-1}] = \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^2|\mathcal{F}_{t-1}]$$

as stated in Koskela et al. (2018, Remark 3). In the case of conditional SMC we separate the first term (corresponding to the immortal particle) from the sum to get

$$\begin{aligned}\mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] &= \frac{1}{(N)_2} \sum_{i=1}^N \mathbb{E}[(\tilde{v}_t^{(i)})_2|\mathcal{F}_{t-1}] = \frac{1}{(N)_2} \mathbb{E}[(\tilde{v}_t^{(1)})_2|\mathcal{F}_{t-1}] + \frac{1}{(N)_2} \sum_{i=2}^N \mathbb{E}[(\tilde{v}_t^{(i)})_2|\mathcal{F}_{t-1}] \\ &= \frac{1}{(N)_2} \left\{ (N-1)_2 \mathbb{E}[(w_t^{(1)})^2|\mathcal{F}_{t-1}] + 2(N-1) \mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] + \sum_{i=2}^N (N-1)_2 \mathbb{E}[(w_t^{(i)})^2|\mathcal{F}_{t-1}] \right\} \\ &= \frac{(N-1)_2}{(N)_2} \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^2|\mathcal{F}_{t-1}] + \frac{2(N-1)}{(N)_2} \mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \\ &= \frac{N-2}{N} \mathbb{E}[c_N(t)|\mathcal{F}_{t-1}] + \frac{2}{N} \mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}]\end{aligned}$$

which gives us (4).

An upper bound on the rate of super-binary mergers is given by

$$D_N(t) := \frac{1}{N(N)_2} \sum_{i=1}^N (v_t^{(i)})_2 \left(v_t^{(i)} + \frac{1}{N} \sum_{j \neq i} (v_t^{(j)})^2 \right).$$

In the standard case this quantity has expectation

$$\begin{aligned}\mathbb{E}[D_N(t)|\mathcal{F}_{t-1}] &= \frac{1}{N(N)_2} \sum_{i=1}^N \left\{ (N)_3 \mathbb{E}[(w_t^{(i)})^3 | \mathcal{F}_{t-1}] + 2(N)_2 \mathbb{E}[(w_t^{(i)})^2 | \mathcal{F}_{t-1}] \right\} \\ &\quad + \frac{1}{N^2(N)_2} \sum_{i=1}^N \sum_{j \neq i} \left\{ (N)_4 \mathbb{E}[(w_t^{(i)})^2 (w_t^{(j)})^2 | \mathcal{F}_{t-1}] + (N)_3 \mathbb{E}[(w_t^{(i)})^2 w_t^{(j)} | \mathcal{F}_{t-1}] \right\}\end{aligned}$$

while in the conditional case, again separating the terms involving particle 1,

$$\begin{aligned}\tilde{D}_N(t) &= \frac{1}{N(N)_2} (\tilde{v}_t^{(1)})_2 \left(\tilde{v}_t^{(1)} + \frac{1}{N} \sum_{j \neq 1} (\tilde{v}_t^{(j)})^2 \right) + \frac{1}{N(N)_2} \sum_{i \neq 1} (\tilde{v}_t^{(i)})_2 \left(\tilde{v}_t^{(i)} + \frac{1}{N} (\tilde{v}_t^{(1)})^2 + \frac{1}{N} \sum_{1 \neq j \neq i} (\tilde{v}_t^{(j)})^2 \right) \\ &= \frac{1}{N(N)_2} \left\{ (\tilde{v}_t^{(1)})_2 \tilde{v}_t^{(1)} + \frac{1}{N} \sum_{j \neq 1} (\tilde{v}_t^{(1)})_2 (\tilde{v}_t^{(j)})^2 + \frac{1}{N} \sum_{i \neq 1} (\tilde{v}_t^{(i)})_2 (\tilde{v}_t^{(1)})^2 \right\} \\ &\quad + \frac{1}{N(N)_2} \sum_{i \neq 1} \left\{ (\tilde{v}_t^{(i)})_2 \tilde{v}_t^{(i)} + \frac{1}{N} \sum_{1 \neq j \neq i} (\tilde{v}_t^{(i)})_2 (\tilde{v}_t^{(j)})^2 \right\} \\ &= \frac{1}{N(N)_2} \left\{ (\tilde{v}_t^{(1)})_2 \tilde{v}_t^{(1)} + \frac{1}{N} \sum_{i \neq 1} \left((\tilde{v}_t^{(1)})_2 (\tilde{v}_t^{(i)})^2 + (\tilde{v}_t^{(i)})_2 (\tilde{v}_t^{(1)})^2 \right) \right\} \\ &\quad + \frac{1}{N(N)_2} \sum_{i \neq 1} \left\{ (\tilde{v}_t^{(i)})_3 + 2(\tilde{v}_t^{(i)})_2 + \frac{1}{N} \sum_{1 \neq j \neq i} \left((\tilde{v}_t^{(i)})_2 (\tilde{v}_t^{(j)})_2 + (\tilde{v}_t^{(i)})_2 \tilde{v}_t^{(j)} \right) \right\}\end{aligned}$$

and so by applying the moments from (8) and Table 1 we find the expectation

$$\begin{aligned}
& \mathbb{E}[\tilde{D}_N(t)|\mathcal{F}_{t-1}] = \\
& = \frac{1}{N(N)_2} \left\{ (N-1)_3 \mathbb{E}[(w_t^{(1)})^3|\mathcal{F}_{t-1}] + 5(N-1)_2 \mathbb{E}[(w_t^{(1)})^2|\mathcal{F}_{t-1}] + 4(N-1) \mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \right\} \\
& + \frac{1}{N^2(N)_2} \sum_{i=2}^N \left\{ 2(N-1)_4 \mathbb{E}[(w_t^{(1)})^2(w_t^{(i)})^2|\mathcal{F}_{t-1}] + (N-1)_3 \mathbb{E}[(w_t^{(1)})^2 w_t^{(i)}|\mathcal{F}_{t-1}] \right. \\
& \quad \left. + 5(N-1)_3 \mathbb{E}[w_t^{(1)}(w_t^{(i)})^2|\mathcal{F}_{t-1}] + 2(N-1)_2 \mathbb{E}[w_t^{(1)} w_t^{(i)}|\mathcal{F}_{t-1}] + (N-1)_2 \mathbb{E}[(w_t^{(i)})^2|\mathcal{F}_{t-1}] \right\} \\
& + \frac{1}{N(N)_2} \sum_{i=2}^N \left\{ (N-1)_3 \mathbb{E}[(w_t^{(i)})^3|\mathcal{F}_{t-1}] + 2(N-1)_2 \mathbb{E}[(w_t^{(i)})^2|\mathcal{F}_{t-1}] \right\} \\
& + \frac{1}{N^2(N)_2} \sum_{i=2}^N \sum_{1 \neq j \neq i} \left\{ (N-1)_4 \mathbb{E}[(w_t^{(i)})^2(w_t^{(j)})^2|\mathcal{F}_{t-1}] + (N-1)_3 \mathbb{E}[(w_t^{(i)})^2 w_t^{(j)}|\mathcal{F}_{t-1}] \right\} \\
& = \frac{1}{N(N)_2} \sum_{i=1}^N \left\{ (N-1)_3 \mathbb{E}[(w_t^{(i)})^3|\mathcal{F}_{t-1}] + 2(N-1)_2 \mathbb{E}[(w_t^{(i)})^2|\mathcal{F}_{t-1}] \right\} \\
& + \frac{1}{N^2(N)_2} \sum_{i=1}^N \sum_{j \neq i} \left\{ (N-1)_4 \mathbb{E}[(w_t^{(i)})^2(w_t^{(j)})^2|\mathcal{F}_{t-1}] + (N-1)_3 \mathbb{E}[(w_t^{(i)})^2 w_t^{(j)}|\mathcal{F}_{t-1}] \right\} \\
& + \frac{1}{N(N)_2} \left\{ 3(N-1)_2 \mathbb{E}[(w_t^{(1)})^2|\mathcal{F}_{t-1}] + 4(N-1) \mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \right\} \\
& + \frac{1}{N^2(N)_2} \sum_{i=2}^N \left\{ 4(N-1)_3 \mathbb{E}[w_t^{(1)}(w_t^{(i)})^2|\mathcal{F}_{t-1}] + 2(N-1)_2 \mathbb{E}[w_t^{(1)} w_t^{(i)}|\mathcal{F}_{t-1}] + (N-1)_2 \mathbb{E}[(w_t^{(i)})^2|\mathcal{F}_{t-1}] \right\} \\
& \leq \mathbb{E}[D_N(t)|\mathcal{F}_{t-1}] + \frac{1}{N(N)_2} \left\{ 3(N-1)_2 \mathbb{E}[(w_t^{(1)})^2|\mathcal{F}_{t-1}] + 4(N-1) \mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \right\} \\
& + \frac{1}{N^2(N)_2} \sum_{i=2}^N \left\{ 4(N-1)_3 \mathbb{E}[w_t^{(1)}(w_t^{(i)})^2|\mathcal{F}_{t-1}] + 2(N-1)_2 \mathbb{E}[w_t^{(1)} w_t^{(i)}|\mathcal{F}_{t-1}] + (N-1)_2 \mathbb{E}[(w_t^{(i)})^2|\mathcal{F}_{t-1}] \right\} \\
& \leq \mathbb{E}[D_N(t)|\mathcal{F}_{t-1}] + \frac{3}{N} \mathbb{E}[(w_t^{(1)})^2|\mathcal{F}_{t-1}] + \frac{4}{N^2} \mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \\
& + \frac{4}{N} \sum_{i=2}^N \mathbb{E}[w_t^{(1)}(w_t^{(i)})^2|\mathcal{F}_{t-1}] + \frac{2}{N^2} \sum_{i=2}^N \mathbb{E}[w_t^{(1)} w_t^{(i)}|\mathcal{F}_{t-1}] + \frac{1}{N^2} \sum_{i=2}^N \mathbb{E}[(w_t^{(i)})^2|\mathcal{F}_{t-1}]
\end{aligned}$$

The second line of the first equality relies on multiplying the relevant terms in Table 1. For the second equality we recombine the terms in particle 1 into the sum. The inequalities follow by bounding e.g. $N-1$ by N , and identifying the first two lines with $\mathbb{E}[D_N(t)|\mathcal{F}_{t-1}]$. This gives us the inequality (5).

We also need control of the squared coalescence rate:

$$c_N(t)^2 = \frac{1}{(N)_2^2} \left(\sum_{i=1}^N (v_t^{(i)})_2 \right)^2 = \frac{1}{(N)_2^2} \left\{ \sum_{i=1}^N \mathbb{E}[(v_t^{(i)})_2^2] + \sum_{i=1}^N \sum_{j \neq i} \mathbb{E}[(v_t^{(i)})_2 (v_t^{(j)})_2] \right\}$$

A bound on its expected value is proved in Koskela et al. (2018), but here we will use a different, more explicit

bound to allow direct comparison between the standard and conditional cases. For standard SMC we have:

$$\begin{aligned}
\mathbb{E}[c_N(t)^2 | \mathcal{F}_{t-1}] &= \frac{1}{(N)_2^2} \left\{ (N)_4 \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^4 | \mathcal{F}_{t-1}] + 4(N)_3 \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^3 | \mathcal{F}_{t-1}] + 2(N)_2 \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^2 | \mathcal{F}_{t-1}] \right\} \\
&\quad + \frac{1}{(N)_2^2} (N)_4 \sum_{i=1}^N \sum_{j \neq i} \mathbb{E}[(w_t^{(i)})^2 (w_t^{(j)})^2 | \mathcal{F}_{t-1}] \\
&= \frac{1}{(N)_2} \left\{ (N-2)_2 \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^4 | \mathcal{F}_{t-1}] + 4(N-2) \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^3 | \mathcal{F}_{t-1}] + 2 \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^2 | \mathcal{F}_{t-1}] \right. \\
&\quad \left. + (N-2)_2 \sum_{i=1}^N \sum_{j \neq i} \mathbb{E}[(w_t^{(i)})^2 (w_t^{(j)})^2 | \mathcal{F}_{t-1}] \right\}
\end{aligned}$$

For conditional SMC, we again separate the terms involving particle 1:

$$\begin{aligned}
\tilde{c}_N(t)^2 &= \frac{1}{(N)_2^2} \left\{ \sum_{i=1}^N \mathbb{E}[(\tilde{v}_t^{(i)})_2^2] + \sum_{i=1}^N \sum_{j \neq i} \mathbb{E}[(\tilde{v}_t^{(i)})_2 (\tilde{v}_t^{(j)})_2] \right\} \\
&= \frac{1}{(N)_2^2} \left\{ \sum_{i=2}^N \mathbb{E}[(\tilde{v}_t^{(i)})_2^2] + \mathbb{E}[(\tilde{v}_t^{(1)})_2^2] + \sum_{i=2}^N \sum_{1 \neq j \neq i} \mathbb{E}[(\tilde{v}_t^{(i)})_2 (\tilde{v}_t^{(j)})_2] + 2 \sum_{i=2}^N \mathbb{E}[(\tilde{v}_t^{(1)})_2 (\tilde{v}_t^{(i)})_2] \right\}
\end{aligned}$$

and then use the same techniques as for $\tilde{D}_N(t)$ to calculate the expectation:

$$\begin{aligned}
& \mathbb{E}[\tilde{c}_N(t)^2 | \mathcal{F}_{t-1}] = \\
& = \frac{1}{(N)_2^2} \left\{ (N-1)_4 \sum_{i=2}^N \mathbb{E}[(w_t^{(i)})^4 | \mathcal{F}_{t-1}] + 4(N-1)_3 \sum_{i=2}^N \mathbb{E}[(w_t^{(i)})^3 | \mathcal{F}_{t-1}] + 2(N-1)_2 \sum_{i=2}^N \mathbb{E}[(w_t^{(i)})^2 | \mathcal{F}_{t-1}] \right\} \\
& \quad + \frac{1}{(N)_2^2} \left\{ (N-1)_4 \mathbb{E}[(w_t^{(1)})^4 | \mathcal{F}_{t-1}] + 8(N-1)_3 \mathbb{E}[(w_t^{(1)})^3 | \mathcal{F}_{t-1}] + 14(N-1)_2 \mathbb{E}[(w_t^{(1)})^2 | \mathcal{F}_{t-1}] \right\} \\
& \quad + \frac{1}{(N)_2^2} 4(N-1) \mathbb{E}[w_t^{(1)} | \mathcal{F}_{t-1}] + \frac{1}{(N)_2^2} (N-1)_4 \sum_{i=2}^N \sum_{1 \neq j \neq i} \mathbb{E}[(w_t^{(i)})^2 (w_t^{(j)})^2 | \mathcal{F}_{t-1}] \\
& \quad + \frac{2}{(N)_2^2} \sum_{i=2}^N \left((N-1)_4 \mathbb{E}[(w_t^{(1)})^2 (w_t^{(i)})^2 | \mathcal{F}_{t-1}] + 2(N-1)_3 \mathbb{E}[w_t^{(1)} (w_t^{(i)})^2 | \mathcal{F}_{t-1}] \right) \\
& = \frac{1}{(N)_2^2} \left\{ (N-1)_4 \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^4 | \mathcal{F}_{t-1}] + 4(N-1)_3 \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^3 | \mathcal{F}_{t-1}] + 2(N-1)_2 \sum_{i=1}^N \mathbb{E}[(w_t^{(i)})^2 | \mathcal{F}_{t-1}] \right\} \\
& \quad + \frac{(N-1)_4}{(N)_2^2} \sum_{i=1}^N \sum_{j \neq i} \mathbb{E}[(w_t^{(i)})^2 (w_t^{(j)})^2 | \mathcal{F}_{t-1}] \\
& \quad + \frac{1}{(N)_2^2} \left\{ 4(N-1)_3 \mathbb{E}[(w_t^{(1)})^3 | \mathcal{F}_{t-1}] + 12(N-1)_2 \mathbb{E}[(w_t^{(1)})^2 | \mathcal{F}_{t-1}] + 4(N-1) \mathbb{E}[w_t^{(1)} | \mathcal{F}_{t-1}] \right\} \\
& \quad + \frac{1}{(N)_2^2} \left\{ 4(N-1)_3 \sum_{i=2}^N \mathbb{E}[w_t^{(1)} (w_t^{(i)})^2 | \mathcal{F}_{t-1}] \right\} \\
& \leq \mathbb{E}[c_N(t)^2 | \mathcal{F}_{t-1}] + \frac{1}{(N)_2^2} \left\{ 4(N-1)_3 \mathbb{E}[(w_t^{(1)})^3 | \mathcal{F}_{t-1}] + 12(N-1)_2 \mathbb{E}[(w_t^{(1)})^2 | \mathcal{F}_{t-1}] \right\} \\
& \quad + \frac{1}{(N)_2^2} \left\{ 4(N-1) \mathbb{E}[w_t^{(1)} | \mathcal{F}_{t-1}] + 4(N-1)_3 \sum_{i=2}^N \mathbb{E}[w_t^{(1)} (w_t^{(i)})^2 | \mathcal{F}_{t-1}] \right\} \\
& \leq \mathbb{E}[c_N(t)^2 | \mathcal{F}_{t-1}] + \frac{4}{N} \mathbb{E}[(w_t^{(1)})^3 | \mathcal{F}_{t-1}] + \frac{12}{N^2} \mathbb{E}[(w_t^{(1)})^2 | \mathcal{F}_{t-1}] + \frac{4}{N(N)_2} \mathbb{E}[w_t^{(1)} | \mathcal{F}_{t-1}] \\
& \quad + \frac{4}{N} \sum_{i=2}^N \mathbb{E}[w_t^{(1)} (w_t^{(i)})^2 | \mathcal{F}_{t-1}]
\end{aligned}$$

The conditions (18) and (19) of Koskela et al. (2018, Lemma 3) give us control over the weights so that we have $w_t^{(i)} = O(1)$ for all i . Under these conditions, in the limit as $N \rightarrow \infty$, the three modified expectations derived above simplify to:

$$\begin{aligned}
& \mathbb{E}[\tilde{c}_N(t) | \mathcal{F}_{t-1}] \leq \mathbb{E}[c_N(t) | \mathcal{F}_{t-1}] + O(N^{-2}) \\
& \mathbb{E}[\tilde{D}_N(t) | \mathcal{F}_{t-1}] \leq \mathbb{E}[D_N(t) | \mathcal{F}_{t-1}] + O(N^{-3}) \\
& \mathbb{E}[\tilde{c}_N(t)^2 | \mathcal{F}_{t-1}] \leq \mathbb{E}[c_N(t)^2 | \mathcal{F}_{t-1}] + O(N^{-3})
\end{aligned}$$

This shows that each of these quantities for conditional SMC is bounded above by the corresponding standard SMC quantity, plus some vanishing error term. This will allow us to apply Koskela et al. (2018, Theorem 1), as we will show in the following.

Next we apply the result of Koskela et al. (2018, Lemma 3), so that for our modified quantity $\tilde{c}_N(t)$ we

have the upper bound:

$$\begin{aligned}
\mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] &= \frac{N-2}{N}\mathbb{E}[c_N(t)|\mathcal{F}_{t-1}] + \frac{2}{N}\mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \\
&\leq \mathbb{E}[c_N(t)|\mathcal{F}_{t-1}] + \frac{2}{N}\mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \\
&\leq \frac{a^4}{N\varepsilon^4} + \frac{2}{N}\mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \\
&= \frac{a^4}{N\varepsilon^4} + O(N^{-2})
\end{aligned} \tag{9}$$

and lower bound:

$$\begin{aligned}
\mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] &= \frac{N-2}{N}\mathbb{E}[c_N(t)|\mathcal{F}_{t-1}] + \frac{2}{N}\mathbb{E}[w_t^{(1)}|\mathcal{F}_{t-1}] \\
&\geq \frac{N-2}{N} \frac{\varepsilon^4}{Na^4} + O(N^{-2}) \\
&= \frac{\varepsilon^4}{Na^4} - \frac{2\varepsilon^4}{N^2a^4} + O(N^{-2}) \\
&= \frac{\varepsilon^4}{Na^4} + O(N^{-2})
\end{aligned} \tag{10}$$

corresponding to (22) in Koskela et al. (2018), except for the addition of a vanishing error term. Furthermore, we obtain for the other quantities (where the constant C may change from one line to the next):

$$\begin{aligned}
\mathbb{E}[\tilde{D}_N(t)|\mathcal{F}_{t-1}] &\leq \mathbb{E}[D_N(t)|\mathcal{F}_{t-1}] + O(N^{-3}) \\
&\leq \frac{C}{N}\mathbb{E}[c_N(t)|\mathcal{F}_{t-1}] + O(N^{-3}) \\
&= \frac{C}{N}\mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] + O(N^{-3})
\end{aligned} \tag{11}$$

and

$$\begin{aligned}
\mathbb{E}[\tilde{c}_N(t)^2|\mathcal{F}_{t-1}] &\leq \mathbb{E}[c_N(t)^2|\mathcal{F}_{t-1}] + O(N^{-3}) \\
&\leq \frac{C}{N}\mathbb{E}[c_N(t)|\mathcal{F}_{t-1}] + O(N^{-3}) \\
&= \frac{C}{N}\mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] + O(N^{-3})
\end{aligned} \tag{12}$$

according to equations (20) and (21) in Koskela et al. (2018), again with additional error terms.

Now let us define the time-scaling:

$$\tilde{\tau}_N(t) := \min \left\{ s \geq 1 : \sum_{r=1}^s \tilde{c}_N(r) \geq t \right\}$$

which is a generalised inverse of $\tilde{c}_N(t)$ and thus satisfies the property:

$$t - s - 1 \leq \sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \tilde{c}_N(r) \leq t - s + 1. \tag{13}$$

We are finally ready to verify the conditions of Koskela et al. (2018, Theorem 1). The conditions are the following.

(Standing Assumption) The conditional distribution of parental indices $a_t^{(1:N)}$ given offspring counts $v_t^{(1:N)}$ is uniform over all valid assignments.

$$(A) \lim_{N \rightarrow \infty} \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \tilde{D}_N(r) \right] = 0$$

$$(B) \lim_{N \rightarrow \infty} \mathbb{E}[\tilde{c}_N(t)] = 0$$

$$(C) \lim_{N \rightarrow \infty} \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \tilde{c}_N(r)^2 \right] = 0$$

$$(D) \mathbb{E}[\tilde{\tau}_N(t) - \tilde{\tau}_N(s)] \leq C_{t,s}N$$

These five conditions are verified below.

(Standing Assumption) This holds by the exchangeability of offspring assignments arising from Algorithm 2.

(B) Using (9) and applying the tower rule, we find

$$\mathbb{E}[\tilde{c}_N(t)] = \mathbb{E}[\mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}]] \leq \frac{a^4}{N\varepsilon^4} + O(N^{-2}) \xrightarrow{N \rightarrow \infty} 0$$

(C) Using Koskela et al. (2018, Lemma 2) along with (12) and the upper bound in (13),

$$\begin{aligned} \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \tilde{c}_N(r)^2 \right] &= \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \mathbb{E}[\tilde{c}_N(r)^2|\mathcal{F}_{t-1}] \right] \leq \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \left(\frac{C}{N} \mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] + O(N^{-3}) \right) \right] \\ &= \frac{C}{N} \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] \right] + O(N^{-2}) = \frac{C}{N} \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \tilde{c}_N(r) \right] + O(N^{-2}) \\ &\leq \frac{C}{N} (t-s+1) + O(N^{-2}) \xrightarrow{N \rightarrow \infty} 0 \end{aligned}$$

(A) The above calculation replacing (12) with (11) yields

$$\begin{aligned} \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \tilde{D}_N(r) \right] &= \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \mathbb{E}[\tilde{D}_N(r)|\mathcal{F}_{t-1}] \right] \\ &\leq \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \left(\frac{C}{N} \mathbb{E}[\tilde{c}_N(t)|\mathcal{F}_{t-1}] + O(N^{-3}) \right) \right] \xrightarrow{N \rightarrow \infty} 0 \end{aligned}$$

(D) Using (10), the upper bound in (13) and Koskela et al. (2018, Lemma 2),

$$\begin{aligned} \mathbb{E}[\tilde{\tau}_N(t) - \tilde{\tau}_N(s)] &= \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} 1 \right] = \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \frac{\mathbb{E}[\tilde{c}_N(r)|\mathcal{F}_{t-1}]}{\mathbb{E}[\tilde{c}_N(r)|\mathcal{F}_{t-1}]} \right] \leq \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \frac{\mathbb{E}[\tilde{c}_N(r)|\mathcal{F}_{t-1}]}{\frac{\varepsilon^4}{Na^4} + O(N^{-2})} \right] \\ &= \frac{1}{\frac{\varepsilon^4}{Na^4} + O(N^{-2})} \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \mathbb{E}[\tilde{c}_N(r)|\mathcal{F}_{t-1}] \right] = \frac{1}{\frac{\varepsilon^4}{Na^4} + O(N^{-2})} \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \tilde{c}_N(r) \right] \\ &\leq \frac{t-s+1}{\frac{\varepsilon^4}{Na^4} + O(N^{-2})} = \frac{(t-s+1)a^4N}{\varepsilon^4 + O(N^{-1})} = (t-s+1) \frac{a^4}{\varepsilon^4} N + O(1) \end{aligned}$$

where the last equality follows by a Taylor expansion of $(\frac{\varepsilon^4}{Na^4} + O(N^{-2}))^{-1}$.

Similarly we derive a lower bound using (9), the lower bound in (13) and Koskela et al. (2018, Lemma

2):

$$\begin{aligned}
\mathbb{E}[\tilde{\tau}_N(t) - \tilde{\tau}_N(s)] &= \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} 1 \right] = \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \frac{\mathbb{E}[\tilde{c}_N(r)|\mathcal{F}_{t-1}]}{\mathbb{E}[\tilde{c}_N(r)|\mathcal{F}_{t-1}]} \right] \geq \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \frac{\mathbb{E}[\tilde{c}_N(r)|\mathcal{F}_{t-1}]}{\frac{a^4}{N\varepsilon^4} + O(N^{-2})} \right] \\
&= \frac{1}{\frac{a^4}{N\varepsilon^4} + O(N^{-2})} \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \mathbb{E}[\tilde{c}_N(r)|\mathcal{F}_{t-1}] \right] = \frac{1}{\frac{a^4}{N\varepsilon^4} + O(N^{-2})} \mathbb{E} \left[\sum_{r=\tilde{\tau}_N(s)+1}^{\tilde{\tau}_N(t)} \tilde{c}_N(r) \right] \\
&\geq \frac{t-s-1}{\frac{a^4}{N\varepsilon^4} + O(N^{-2})} = \frac{(t-s-1)\varepsilon^4 N}{a^4 + O(N^{-1})} = (t-s-1) \frac{\varepsilon^4}{a^4} N + O(1)
\end{aligned}$$

Therefore we have as required

$$\mathbb{E}[\tilde{\tau}_N(t) - \tilde{\tau}_N(s)] \sim C_{t,s}N$$

as $N \rightarrow \infty$.

This concludes the proof of Corollary 1. □

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