A brief introduction to 2D Spectral Element Method and programing guide written in C

Zeming Su szmxjtu@163.com

April 26, 2016

1   General Introduciton

The primary goal of this introduction is to demonstrate how to implement spectral element method in 2d. We seek to produce the displacement field caused by an earthquake in a finite earth model with volume Ω. In order to simplify calculation, we just use stress-free boundary condition. We will talk about artificial absorbing boundary condition in detail later. The displacement field $\boldsymbol{u}$ is produced by an earthquake is governed by the momentum equation,

$$\rho\, \partial_t^2 \boldsymbol{u} = \boldsymbol{\nabla} \cdot \boldsymbol{T} + \boldsymbol{f} \tag{1}$$

,where $\boldsymbol{T} = (\boldsymbol{T}_1, \boldsymbol{T}_2, \boldsymbol{T}_3)^T$, and $\boldsymbol{T}_i$ is vector function. The weak form of above equation is produced by dotting the momentum equation with an arbitrary test vector $\boldsymbol{w} = (w_1(\boldsymbol{x}), w_2(\boldsymbol{x}), w_3(\boldsymbol{x}))^T$, integrating by the model volume Ω. Since we use stress-free boundary condition, which is easy to derive the weak form of wave equation, we have $\boldsymbol{T}|_{\partial\Omega} = 0$

$$
\begin{aligned}
&\int_\Omega \boldsymbol{\nabla} \cdot \boldsymbol{T} \cdot \boldsymbol{w}(\boldsymbol{x})\, \mathrm{d}^3\boldsymbol{x} \\
&= \sum_{i=1}^3 \int_\Omega w_i(\boldsymbol{x}) \nabla \cdot \boldsymbol{T}_i\, \mathrm{d}^3\boldsymbol{x} \\
&= \sum_{i=1}^3 \int_\Omega \nabla \cdot (w_i(\boldsymbol{x})\boldsymbol{T}_i) - \nabla w_i(\boldsymbol{x}) \cdot \boldsymbol{T}_i\, \mathrm{d}^3\boldsymbol{x} \\
&= \sum_{i=1}^3 \left(\int_{\partial\Omega} w_i(\boldsymbol{x})\boldsymbol{T}_i \cdot\, \mathrm{d}\boldsymbol{s} - \int_\Omega \nabla w_i(\boldsymbol{x}) \cdot \boldsymbol{T}_i\, \mathrm{d}^3\boldsymbol{x} \right) \\
&= -\sum_{i=1}^3 \int_\Omega \nabla w_i(\boldsymbol{x}) \cdot \boldsymbol{T}_i\, \mathrm{d}^3\boldsymbol{x} \\
&= -\int_\Omega \boldsymbol{\nabla}\boldsymbol{w}(\boldsymbol{x}) : \boldsymbol{T}\, \mathrm{d}^3\boldsymbol{x}
\end{aligned}
\tag{2}
$$

The Source function is

$$\boldsymbol{f} = -\boldsymbol{M} \cdot \boldsymbol{\nabla}\delta(\boldsymbol{x} - \boldsymbol{x}_s)S(t) \tag{3}$$

Using the properties of Dirac delta distribution, after integration, it transformed in the following way,

$$
\begin{aligned}
&\int_\Omega \boldsymbol{f} \cdot \boldsymbol{w}\, \mathrm{d}^3\boldsymbol{x} \\
&= -S(t) \int_\Omega \boldsymbol{M} \cdot \boldsymbol{\nabla}\delta(\boldsymbol{x} - \boldsymbol{x}_s) \cdot \boldsymbol{w}\, \mathrm{d}^3\boldsymbol{x} \\
&= \boldsymbol{M} : \boldsymbol{\nabla}\boldsymbol{w}(\boldsymbol{x}_s)S(t)
\end{aligned}
\tag{4}
$$

In 2 dimension, We can derive this using Green Formula.

$$\int_\Omega \rho\, \boldsymbol{w} \cdot \partial_t^2 \boldsymbol{u}\, \mathrm{d}^2\boldsymbol{x} = -\int_\Omega \boldsymbol{\nabla}\boldsymbol{w}(\boldsymbol{x}) : \boldsymbol{T}\, \mathrm{d}^2\boldsymbol{x} + \boldsymbol{M} : \boldsymbol{\nabla}\boldsymbol{w}(\boldsymbol{x}_s)S(t) \tag{5}$$

In the following parts, we will explain the general procedure of specfem.

First, we divide the region into a number of non-overlapping elements, Ω_e, $e = 1, \ldots, n_e$, such that $\Omega = \bigcup_e^{n_e} \Omega_e$, as shown in Fig.1. In SEM, the shape of element is restricted to quadrilateral. The integration in the whole region become the sum of integration over every element, as shown below.

$$\sum_e^{n_e} \int_{\Omega_e} \rho \, \boldsymbol{w} \cdot \partial_t^2 \boldsymbol{u} \, \mathrm{d}^2 \boldsymbol{x} = -\sum_e^{n_e} \int_{\Omega_e} \boldsymbol{\nabla w}(\boldsymbol{x}) : \boldsymbol{T} \, \mathrm{d}^2 \boldsymbol{x} + \boldsymbol{M} : \boldsymbol{\nabla w}(\boldsymbol{x}_s) S(t) \quad (6)$$

Since the equation holds for any test function $\boldsymbol{w}$, we can use $\boldsymbol{w}$ to extract the value of displacement$\boldsymbol{u}$ at one point by setting w equals 1 at this point and equals 0 at other points. It's undoubted that this kind of $\boldsymbol{w}$ did exist. Using Gauss-Lobatto-Legendre points, the integration become the forms like below:

$$M\ddot{U} + K = F, \quad (7)$$

where M denotes the global mass matrix, K the global stiffness matrix and F the source term. Let's build the mass matrix M first. According the quadrature, we can know that the mass matrix is diagonal.

$$\int_{\Omega_e} \rho \boldsymbol{w} \cdot \partial_t^2 \boldsymbol{s} \, \mathrm{d}\boldsymbol{x}$$

$$\approx \sum_{\alpha',\beta'} \omega_{\alpha'} \omega_{\beta'} \boldsymbol{J}_e^{\alpha'\beta'} \rho^{\alpha'\beta'} \sum_{i,j=1}^{2} \hat{\boldsymbol{x}}_i \cdot \hat{\boldsymbol{x}}_j$$

$$\times \sum_{\alpha,\beta} w_i^{\alpha\beta} \ell_\alpha(\xi_{\alpha'}) \ell_\beta(\eta_{\beta'}) \sum_{\sigma\tau} \ddot{s}_j^{\sigma\tau}(t) \ell_\sigma(\xi_{\alpha'}) \ell_\tau(\eta_{\beta'})$$

$$= \sum_{\alpha',\beta'} \omega_{\alpha'} \omega_{\beta'} \boldsymbol{J}_e^{\alpha'\beta'} \rho^{\alpha'\beta'} \sum_{i,j}^{3} \delta_{ij} \sum_{\alpha,\beta} w_i^{\alpha\beta} \delta_{\alpha\alpha'} \delta_{\beta\beta'} \qquad (8)$$

$$\times \sum_{\sigma\tau} \ddot{s}_j^{\sigma\tau}(t) \delta_{\sigma\alpha'} \delta_{\tau\beta'}$$

$$= \sum_{\alpha,\beta} \omega_\alpha \omega_\beta \boldsymbol{J}_e^{\alpha\beta} \rho^{\alpha\beta} \sum_{i=1}^{2} w_i^{\alpha\beta} \ddot{s}_i^{\alpha\beta}(t)$$

Now, Let's consider how to build the stiffness matrix. We can start from a local integration. The first step is to construct the stress tensor $\boldsymbol{T}$. To compute this term, we need the strain tensor first.

$$\partial_i s(\boldsymbol{x}(\xi_\alpha, \eta_\beta), t) \approx \partial_i \left[s_j^{\alpha\beta} \sum_{\alpha,\beta}^{n_\ell} \ell_\alpha(\xi_\alpha) \ell_\beta(\eta_\beta) \right]$$

$$= \left[\sum_\sigma^{n_\ell} s_j^{\sigma\beta}(t) \ell_\sigma'(\xi_\alpha) \right] \partial_i \xi(\xi_\alpha, \eta_\beta) \qquad (9)$$

$$+ \left[\sum_\sigma^{n_\ell} s_j^{\alpha\sigma}(t) \ell_\sigma'(\eta_\beta) \right] \partial_i \eta(\xi_\alpha, \eta_\beta)$$

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And the stress tensor $\boldsymbol{T}$ is obtained by the following.

$$\boldsymbol{T}(\boldsymbol{x}(\xi_\alpha, \eta_\beta), t) = \boldsymbol{c} : \boldsymbol{\nabla}\boldsymbol{s}(\boldsymbol{x}(\xi_\alpha, \eta_\beta), t). \tag{10}$$

where $\boldsymbol{c}$ denotes the stress-strain tensor, is a forth order tensor.