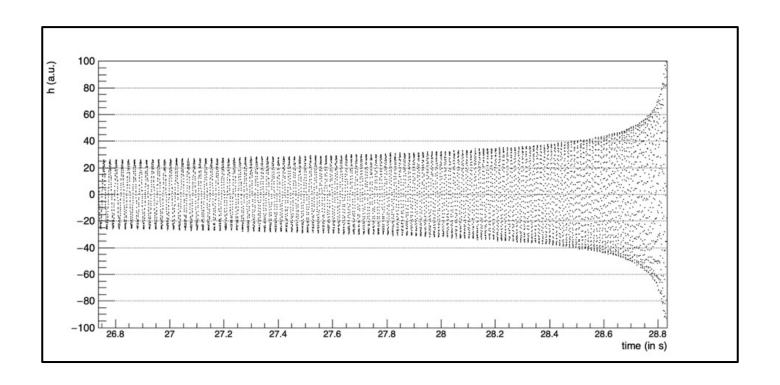
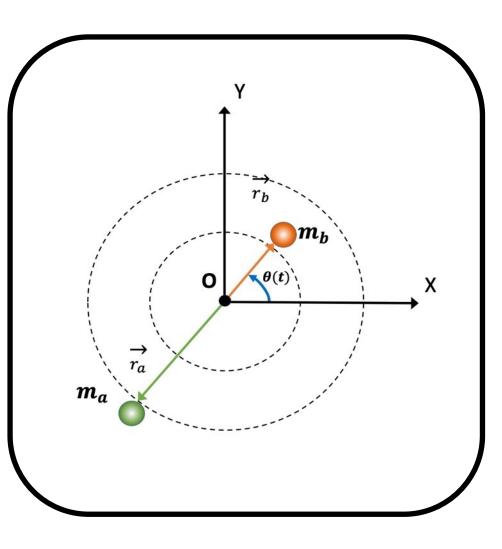
Chirp detection

An introduction



- 1. Signal shape
- 2. Detection

1. Signal shape: why a chirp?



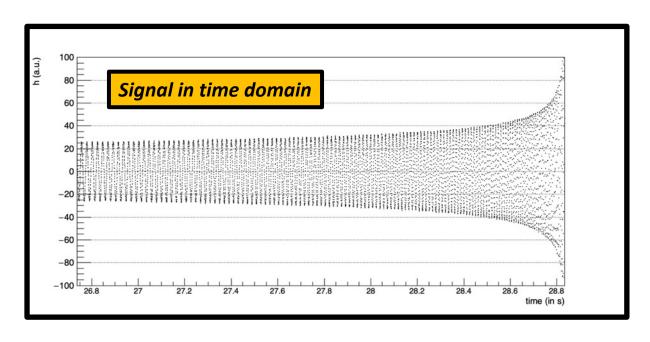
- → The system considered is made of 2 massive objects (A and B) in rotation
- → System radiates energy via gravitational waves. When objects are very dense (Neutron star, black holes) and close to each other, the amount of energy radiated becomes significant.
- \rightarrow The system is isolated, so it's total energy E_{tot} decreases.

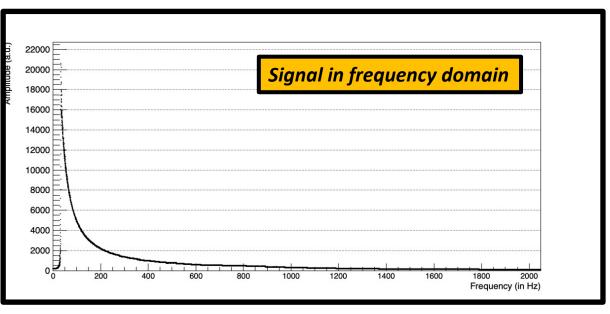
$$\frac{dE_{tot}}{dt} = G \frac{m_a m_b \dot{r}}{2r^2}$$

- → The two objects will **get closer**, rotation **speed will increase** (*Kepler's 3rd law*), and the amount of **radiated energy will increase**.
- \rightarrow The shape of the gravitational wave signal will therefore **have increasing** amplitude and frequency until $r\sim 0$ (in reality before that, at least you know that below $r=r_s$ something should happen)

1. Signal shape: the chirp

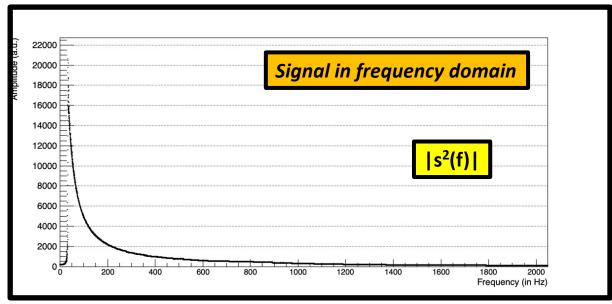
→ The signal we are looking for is the following

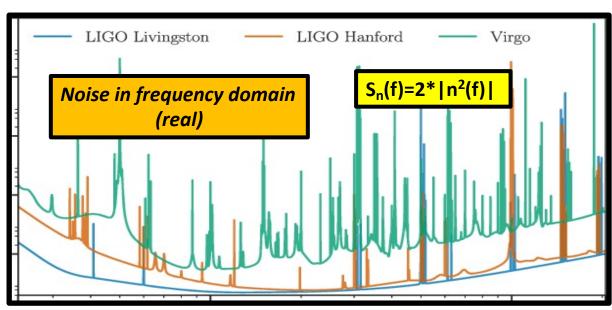




- → The time during which we can detect it depends on the sensitivity range of the detector, basically between 40 and 1000 Hz.
- → Corresponding time range depends on the signal origin, in particular on the masses of the 2 objects. It can be between 30s (low masses) and less than 1s (high masses).

2. Detection: the problem





→ Problem: we measure a strain h(t), and we need to detect when we have a signal (h(t)=s(t)+n(t)) and to stay quiet when there is nothing (h(t)=n(t)).

→ Question: how to do that?

- → Look in the frequency domain in order to reduce the number of operations
- → Idea is to transform the signal in such a way that signal to noise ratio (S/N) gets maximized.

$$x(t)=\int_{-\infty}^{+\infty}rac{ ilde{h}(f) ilde{s}^*(f)}{S_n}e^{2i\pi ft}df$$
 Matched filter

→ Matched filter is doing exactly that. When noise is larger, filter suppress the contribution, and when it gets small it enhances it.

2. Detection: the algorithm

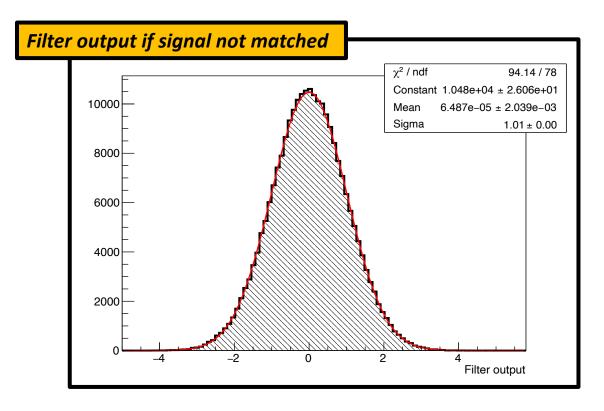
$$x(t) = \int_{-\infty}^{+\infty} \frac{\tilde{h}(f(\tilde{s}^*(f)))}{S_n} e^{2i\pi f t} df$$

- → Fourier transform of the measured strain will be continuously compared to a big collection of precomputed signals: the templates.
- → Templates are stored in a bank, which covers a given phase space (depending on what you're looking for).
- → The noise also needs to be measured precisely. It is used to compute the matched filter, but also to normalize the template amplitude. If h(t)=n(t) one has indeed:

$$\langle x^2(t) \rangle = \int_{-\infty}^{+\infty} \frac{|\tilde{s}^2(f)|}{2S_n} df = \frac{\sigma_s^2}{2}$$

 \rightarrow Normalizing each template with this factor, the output of the filter is exactly the S/N ratio $\rho(t)$

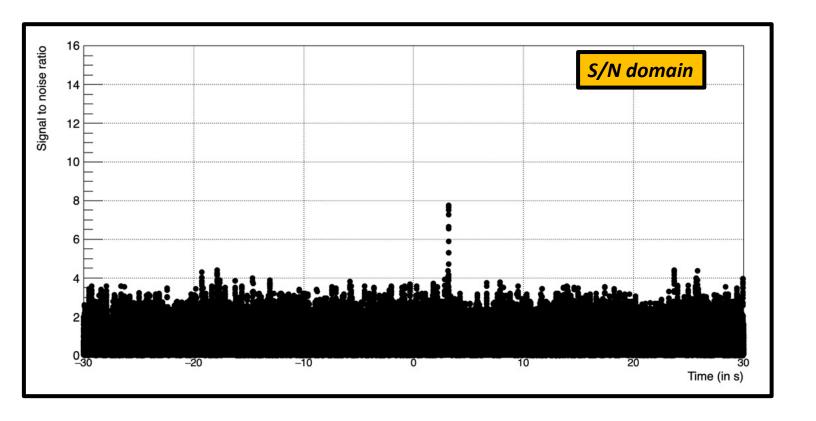
$$\rho_s(t) = \frac{\sqrt{2}}{\sigma_s} x(t)$$



 \rightarrow In the absence of signal, output is gaussian with σ =1.

2. Detection: the algorithm

- \rightarrow Whenever $|\rho(t)|$ gets larger than a given threshold, we can conclude that there is a match with a certain confidence level
- → This threshold is a tradeoff between the efficiency we want to obtain, and the maximum proportion of fake signals we can afford (false alarms)



- → For a given signal, there could be more than one template passing the threshold
- → A detailed analysis of the signal is then necessary to precisely identify its properties. **This** is the parameter estimation.

Documentation

- Document detailing the calculations: http://sviret.web.cern.ch/sviret/docs/Chirps.pdf
- Webpage with code and exercises: http://sviret.web.cern.ch/sviret/Welcome.php?n=Virgo.Ana